Conclusion

Accurate low-pressure VLE of isobutyraldehyde/ethyl acetate/DMF and the constituent binaries have been obtained. The lowering of the relative volatility of isobutyraldehyde/ethyl acetate by DMF in their ternary mixture can be understood in terms of the chemical interaction between DMF and the aldehyde group. All models considered, except the van Laar model, predict the behavior of the isobutyraldehyde/ethyl acetate/DMF ternary mixture accurately with only binary information. The DMF-aldehyde group interaction in the isobutyraldehyde/DMF and DMF/benzaldehyde mixtures are significantly different as discussed above, and the UNIFAC group contribution model cannot describe both mixtures accurately with one set of DMF-aldehyde parameters.

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Registry No. DMF, 68-12-2; isobutyraldehyde, 78-84-2; ethyl acetate, 141-78-6.

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Thermodynamic Properties of Dichloromethane Gas

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Virial coefficients, derived from an analysis of precise P-V-T data, form a basis for the determination of thermodynamic properties and the force constants of appropriate intermolecular potential models. An earlier work reports experimental P-V-T measurements for dichloromethane gas from about 350 to 510 K along with compression factors and second virial coefficients. This study presents revised values of second virial coefficients and the following other derived quantities (the figure in parentheses indicates the estimated percent uncertainty): (I) third virial coefficients (2-20); (II) Stockmayer (12-6-3) force constants; (iii) thermodynamic properties (0.2-0.5 up to 10 atm, increasing to 5 at 30 atm and above); and (iv) fugacity coefficients. Stockmayer force constants of this work, evaluated from the revised second virial coefficients, are in close agreement with literature values. The experimental third virial coefficients are reported for the first time here. Though the uncertainty in these values is large, the usefulness of the values is considerable as is evident from the excellent agreement of the present entropy values with those available in the literature.

Introduction

The present investigation is a part of an extensive research program for evaluating the chemical thermodynamic properties of selected halogenated alkanes over a wide range of temperatures and pressures. The volumetric properties of dichloromethane gas from about 350 to 510 K at pressures below the saturation pressure at each temperature were determined by using a Kay-type apparatus which was fabricated in a previous investigation (1). The compression factors and second virial coefficients computed from the measured P, V, Tvalues, with the estimated percent uncertainties of 0.02-0.5 and 0.5, respectively, have been reported elswhere (1).

Virial coefficients, derived from an analysis of precise P-V-Tdata, form a basis for the determination of thermodynamic properties and the force constants of suitable potential models. No experimental values are available for the third and higher virial coefficients of dichloromethane. Using the second virial coefficient data of Perez Masia and Diaz Pena (2) in the range 349.53-510.0 K and those of Fogg et al. (3) in the range 319.4-382.6 K as compiled by Dymond and Smith (4), Polak and Lu (5) evaluated the parameters of the Stockmayer (12-6-3) potential function. Rätzsch and Freydank (6) report the parameters of the Stockmayer (18-6-3) potential. Dzung (7) calculated the thermodynamic properties of dichloromethane up to a temperature of 423 K and a pressure of about 2.6 atm, using published values of the specific heat capacity of the saturated liquid and other pertinent data. Seshadri et al. (8) extended the range of Dzung's work by computing the thermodynamic properties of dichloromethane up to a temperature of 750 K and a pressure of 200 atm by using the 11-constant Martin-Hou equation of state (9).

This work reports revised values of second virial coefficients and the following other quantities derived from the experimental P-V-T data (1) (the figure in parentheses indicates the estimated percent uncertainty): (i) third virial coefficients (2-20); (ii) Stockmayer (12-6-3) force constants; (iii) thermodynamic properties (0.2-0.5 up to 10 atm, increasing to 5 at 30 atm and above); and (iv) fugacity coefficients. The experimental third virial coefficients are reported for the first time in the present investigation. The revised second virial coefficients are of a higher precision than those reported earlier (1), these values and the present third virial coefficients being very nearly consistent with each other.

Theoretical Background

1. The Virial Equation of State. The virial equation of state expresses the deviations from ideality as an infinite power series in the density, 1/V (Leiden series)

$$Z = \frac{PV}{RT} = 1 + \frac{B}{V} + \frac{C}{V^2} + \dots$$
(1)

or, in the pressure, P (Berlin series)

$$Z = \frac{PV}{RT} = 1 + B'P + C'P^2 + \dots$$
(2)

The coefficients of the two expansions are uniquely interrelated as follows (10):

$$B' = B/RT; \quad B = B'RT \tag{3}$$

$$C' = (C - B^2) / (RT)^2; \quad C = (C' + B'^2)(RT)^2$$
 (4)

2. Thermodynamic Properties. The thermodynamic properties of a real gas are commonly expressed in the form of departure functions. For example, $[H(T) - H^{\circ}(T)]$ is the enthalpy of the gas in the real state less that of the gas in the standard state (ideal gas at 1 atm) at the same temperature. Departure functions are related to the second and third virial coefficients as follows:

$$[U(T) - U^{\circ}(T)] = -RT \left[\frac{T}{V} \frac{dB}{dT} + \frac{T}{2V^2} \frac{dC}{dT} \right]$$
(5)

$$\begin{bmatrix} H(T) - H^{\circ}(T) \end{bmatrix} = RT \left[\frac{1}{V} \left(B - T \frac{dB}{dT} \right) + \frac{1}{V^2} \left(C - \frac{T}{2} \frac{dC}{dT} \right) \right]$$
(6)

$$[S(T) - S^{\circ}(T)] = -R \left[\ln P + \frac{T}{V} \frac{dB}{dT} + \frac{B^{2}}{2V^{2}} - \frac{C}{2V^{2}} + \frac{T}{2V^{2}} \frac{dC}{dT} \right]$$
(7)

$$[G(T) - G^{\circ}(T)] = RT(B'P + C'(P^2/2)) + RT \ln P \quad (8)$$

3. Fugacity Coefficient. The fugacity coefficient, φ , is related to B' and C' as

$$\ln \varphi = B'P + C'(P^2/2)$$
 (9)

Results and Discussion

1. Virial Coefficients. The quantities, B_v and C_v , defined by eq 10 and 11, were calculated from the isothermal P-V-T

$$B_{v} = (Z - 1)V; \quad B = \lim_{(1/V) \to 0} B_{v}$$
 (10)

$$C_{v} = [(Z - 1)V - B]V; \quad C = \lim_{(1/V) \to 0} C_{v}$$
 (11)

data reported in an earlier work (1) and extrapolated to zero density to obtain *B* and *C* given in Table I. The shapes of the C_v curves are so sensitive to the value chosen for *B* that they provide excellent criteria for extrapolating the B_v curves to zero density (11). A small deviation from the "correct" values of

Table I. Second and Third Virial Coefficients for Dichloromethane

T/K	$B/(cm^3 mol^{-1})$	$C \times 10^{-5}/$ (cm ⁶ mol ⁻²)	
240.52	-546.0	(0111 11101)	
368.07	-476 8	0.29	
389.71	-413.8	0.27	
411.65	-365.4	0.26	
431.68	-322.9	0.24	
451.71	-292.0	0.23	
471.86	-262.5	0.20	
492.25	-239.7	0.18	
502.44	-227.1	0.17	
510.00	-220.9		

B results in unreasonably large positive or negative curvature in the plots of C_v at low densities. Thus, the values of *B* were selected so that the C_v plots had a small, nearly constant curvature at low densities. Not only are the extrapolated values for *B* and *C* very nearly consistent with each other, but more precise values can be assigned to *B* than previously reported (1), e.g., in the order of ± 0.15 cm³ mol⁻¹ at 492.25 K in comparison to ± 1.4 cm³ mol⁻¹ (1). The uncertainty in *C* values is estimated to be about 20% at 368.07 K and decreases to about 2% at 471.86 K and above (up to near T_c). The *C* values at the two extreme temperatures, 349.53 and 510 K (T_c), had much larger uncertainty as compared to those at intermediate temperatures and hence are not reported. The present *C* values are the only ones available so far.

Computation of Z using B values only is recommended up to about 1.5 atm at each temperature. Inclusion of C values results in the prediction of Z to within 0.5% of the experimental compression factors at high pressures. However, at the highest pressures additional correction is necessary implying the necessity of at least the fourth virial coefficient, D. Because of a large scatter in $\{[(Z - 1)V - B]V - C\}V$ vs. 1/V plots at most of the isotherms, it was not possible to get reasonable estimates for D.

For ease of interpolation and limited extrapolation and for calculating gas-phase nonideality, the *B* and *C* values given in Table I were fitted to a polynomial in 1/T.

$$B = 207.72 - 289.35 \times 10^{3} T^{-1} + 961.54 \times 10^{5} T^{-2} - 304.1 \times 10^{8} T^{-3}$$
(12)

$$C = -95.6 \times 10^3 + 87.1 \times 10^8 T^{-1} - 15.2 \times 10^9 T^{-2}$$
(13)

The standard deviations of the fit for *B* and *C* were 1.48 cm³ mol⁻¹ and 0.043×10^4 cm⁶ mol⁻², respectively.

2. Thermodynamic Properties and Fugacity Coefficients. Thermodynamic properties $[U(T) - U^{\circ}(T)]$, $[H(T) - H^{\circ}(T)]$, and $[S(T) - S^{\circ}(T)]$ were calculated at the experimental P, V values for a particular isotherm from eq 5, 6, and 7, respectively. In evaluating $[G(T) - G^{\circ}(T)]$ from eq 8, B' and C' were obtained from eq 3 and 4, respectively. In a similar manner, the fugacity coefficients, φ , were calculated from eq 9. Appropriate derivatives of the virial coefficients were obtained from eq 12 and 13. The computed thermodynamic properties and fugacity coefficients are reported in Table II. The error in the derived thermodynamic properties is estimated to be 0.2-0.5% up to a pressure of about 10 atm, increasing to 5% at 30 atm and above.

The agreement of the present entropy values with those of Seshadri et al. (8) is excellent as shown in Table III. Ideal gas entropy values, S° , were taken from ref 12. Comparison of the *H* values was not possible because of the difference in the choice of standard states.

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T/K	P/atm	$-[U(T) - U^{\circ}(T)]/$ (cal mol ⁻¹)	$-[H(T) - H^{\circ}(T)]/$ (cal mol ⁻¹)	$[G(T) - G^{\circ}(T)]/$ (cal mol ⁻¹)	$[S(T) - S^{\circ}(T)]/$ (cal mol ⁻¹ K ⁻¹)	φ
349.53	0.14922	5.04	7.01	-1323.3	3.7659	0.9972
• • • • • •	0.22311	7.54	10.5	-1044.9	2.9594	0.9958
	0.26727	9.04	12.6	-920.0	2.5962	0.9949
	0.46091	15.7	21.8	-544.1	1.4943	0.9912
	0.85752	29.4	40.8	-118.2	0.2212	0.9837
	2.4088	85.2	118.4	578.2	-1.9927	0.9543
	2.4283	85.9	119.4	583.5	-2.0108	0.9540
368.07	0.14852	4.31	6.04	-1396.6	3.7779	0.9976
	0.17485	5.08	7.11	-1277.5	3.4010	0.9972
	0.26232	(+03 199	10.7	-501.0	2.0380	0.9908
	0.40207	13.2	35.7		0.2182	0.9863
	0.00001	20.0	98.5	594.0	-1.8812	0.9630
	2.3330	84.4	117.9	715.2	-2.2633	0.9561
	3.5368	108.8	152.0	882.0	-2.8090	0.9442
	4.1191	128.1	179.0	986.3	-3.1655	0.9351
389.71	0.14781	3.65	5.14	-1482.1	3.7898	0.9981
	0.18433	4.55	6.41	-1311.4	3.3487	0.9976
	0.25397	6.28	8.84	-1063.9	2.7074	0.9967
	0.42128	10.4	14.7	-673.7	1.6910	0.9945
	0.83148	20.7	29.2	-151.3	0.3135	0.9892
	2.4457	62.2	87.6	667.7	-1.9378	0.9683
	2.9259	75.0	105.4	801.5	-2.3271	0.9621
	3.5766	92.6	130.1	950.3	-2.7718	0.9537
	4.3025	112.5	158.0	1085.7	-3.1912	0.9444
411.05	7.4752	205.6	200.3	14/9.0 1509 9	~4.0049 2 sean	0.9039
411.00	0.14282	3.UJ 3.71	4.20 5.90	-1093.3	3 4439	0.9981
	0.17034	4 73	6.70	-1230.8	2.9736	0.9976
	0.33816	7.20	10.2	-889.9	2.1371	0.9964
	0.81545	17.4	24.7	-174.1	0.3630	0.9912
	2.4459	53.3	75.3	709.9	1.9073	0.9737
	2.6428	57.7	81.5	771.4	-2.0720	0.9716
	2.9533	64.7	91.4	859.5	-2.3098	0.9683
	3.4014	74.9	105.8	971.0	-2.6157	0.9635
	4.0173	89.0	125.8	1101.6	-2.9813	0.9569
	4.5860	102.3	144.5	1204.7	-3.2772	0.9509
	7.3457	169.6	239.1	1564.5	-4.3807	0.9216
401.00	11.681	287.0	403.5	1902.0	-0.0987	0.8762
431.00	0.10742	2.55	4.15	-1340.1	3.0072	0.9981
	0.21008	4 85	6.88	-1162.6	2 6773	0.9976
	0.37229	6.99	9.92	-850.5	1.9473	0.9966
	0.72729	13.7	19.4	-278.9	0.6010	0.9933
	1.0390	19.6	27.8	24.6	-0.1215	0.9905
	2.4231	46.3	65.7	740.0	-1.8664	0.9779
	2.6383	50.5	71.6	811.3	-2.0453	0.9759
	2.8946	55.5	78.8	888.8	-2.2413	0.9736
	3.2174	61.9	87.8	976.9	-2.4663	0.9706
	3.5886	69.3	98.2	1067.5	-2.7005	0.9672
	3.9484	76.5	108.4	1146.6	-2.9073	0.9640
	4.2901	83.4	118.2	1215.0	-3.0884	0.9609
451 71	11.595	243.2	343.4 9 6 9	2007.3	-0.4441 3.7050	0.0203
401./1	0.15298	2.00	0.00 1 1 A	-1000.0	3.1202	0.5500
	0.17000	2.52	5 19	-1367.3	3.0155	0.9983
	0.26678	4 45	6.34	-1188.0	2.6159	0.9979
	0.41286	6.90	9.82	-797.0	1.7427	0.9968
	0.81627	13.7	19.5	-188.0	0.3731	0.9936
	2.3262	39.4	56.1	741.3	-1.7651	0.9817
	2.5060	42.5	60.5	806.8	-1.9200	0.9803
	2.7153	46.1	65.6	877.3	-2.0875	0.9787
	2.9475	50.2	71.3	949.3	-2.2595	0.9769
	3.1378	53.5	76.1	1004.1	-2.3913	0.9754
	3.3631	57.4	81.6	1064.7	-2.5378 -9.7649	0.9736
	3.7410	64.U 71.0	91.1 101.0	1107.0 1976 0	-2.1043	0.9707
	4.1308	71.0	101.0	1240.U 1339.6	-2.9190 -3.9194	0.9010
	26 485	19.4 570 0	814 5	2739 9	-7.8534	0.7991
471.86	0.14605	2.19	3.12	-1804.8	3.8183	0.9990
	0.18682	2.80	3.99	-1574.2	3.3278	0.9987
	0.21994	3.30	4.70	-1421.4	3.0024	0.9985
	0.26920	4.03	5.75	-1232.2	2.5992	0.9982
	0.37112	5.56	7.93	-931.8	1.9580	0.9975
	0.59788	8.98	12.8	-486.1	1.0031	0.9959
	0.98174	14.8	21.1	-23.5	0.0053	0.9933

		$-[U(T) - U^{\circ}(T)]/$	$-[H(T) - H^{\circ}(T)]/$	$[G(T) - G^{\circ}(T)]/$	$[S(T) - S^{\circ}(T)]/$	
T/K	P/atm	(cal mol ⁻¹)	(cal mol ⁻¹)	(cal mol ⁻¹)	$(cal mol^{-1} K^{-1})$	arphi
	2.5127	38.2	54.4	847.8	-1.9120	0.9830
	2.8422	43.3	61.7	961.2	-2.1678	0.9807
	3.2594	49.7	70.8	1087.0	-2.4537	0.9779
	3.7977	58.1	82.8	1226.8	-2.7754	0.9743
	4.5293	69.6	99.1	1387.3	-3.1501	0.9694
	5.7655	89.4	127.2	1606.1	-3.6734	0.9610
	7.3117	114.5	162.8	1818.1	-4.1980	0.9507
	28.219	528.8	745.3	2939.9	-7.8014	0.8149
	32.577	645.8	907.5	3042.8	-8.3571	0.7877
492.25	0.12885	1.75	2.49	-2005.2	4.0684	0.9992
102.20	0.16452	2.23	3.18	-1766.3	3.5818	0.9990
	0 20614	2.79	3.98	-1546.0	3.1325	0.9988
	0.27663	3 75	5 35	-1258.6	2 5461	0.9984
	0.37345	5.06	7 22	-965.6	1 9470	0.9978
	0.53822	7 31	10.4	-609.1	1 2162	0.9968
	1.0257	14.0	19.9	18.9	-0.0788	0.9940
	2 4641	22.8	19.5	867.9	-1 8608	0.0040
	2.4041	37.9	53.0	974.9	-2 0899	0.0000
	2.1001	57.6 41.6	50.2	1066.6	_2.0000	0.3030
	3.0234	41.0	0 0 .0	1100.0	-2.2013	0.9822
	0.4400	47.0	07.0	1190.3	-2,5554	0.9797
	3.9481	04.0	11.1	1320.3	-2.8400	0.9766
	4.0200	62.7	09.3	1400.0	-3.1201	0.9734
	5.5292	77.0	109.6	1640.5	-3.0003	0.9070
	31.280	517.9	731.3	3176.2	-7.9314	0.8220
	35.351	608.7	807.9	3269.3	-8.3/42	0.8000
502.44	0.11231	1.45	2.07	-2183.7	4.3421	0.9994
	0.14883	1.92	2.74	-1902.8	3.7817	0.9992
	0.19299	2.49	3.56	-1643.6	3.2642	0.9989
	0.24565	3.17	4.53	-1403.0	2.7834	0.9986
	0.33907	4.39	6.25	-1081.7	2.1405	0.9981
	0.48430	6.27	8.94	-726.6	1.4283	0.9973
	0.99810	12.9	18.5	-7.40	-0.0220	0.9945
	2.4476	32.0	45.5	880.2	-1.8425	0.9866
	2.7645	36.1	51.5	1000.0	-2.0928	0.9848
	3.2276	42.3	60.3	1152.1	-2.4129	0.9823
	3.7798	49.6	70.7	1306.7	-2.7414	0.9793
	4.3471	57.2	81.5	1443.1	-3.0345	0.9762
	4.9637	65.6	93.4	1572.1	-3.3148	0.9728
	5.6894	75.4	107.4	1704.3	-3.6057	0.9688
	32.547	505.0	713.8	3287.7	-7.9595	0.8271
	37.401	603.6	851.5	3396.4	-8.4476	0.8025
	47.023	837.1	1175.3	3563.9	-9.4182	0.7549
510.00	0.12122	1.51	2.16	-2139.2	4.1903	0.9994
	0.16283	2.04	2.90	-1840.4	3.6029	0.9991
	0.21889	2.74	3.90	-1540.8	3.0135	0.9989
	0.28768	3.60	5.13	-1264.2	2.4688	0.9985
	0.44904	5.62	8.01	-813.8	1.5800	0.9976
	0.66814	8.37	11.9	-412.2	0.7849	0.9965
	1.0613	13.3	19.0	54.7	-0.1444	0.9944
Note: Mor	eignificent fig	ures are given to retain	internal consistency	.		

Table III.	Comparison of Real	Gas Entropy	Values for
Dichlorom	ethane at 431.68 and	502.44 K	

Table II (Continued)

		S/(cal m	nol ⁻¹ K ⁻¹)	
$T/{ m K}$	P/atm	this work	Seshadri et al. $(8)^a$	
431.68	0.20		72.814	
101100	0.21	72.624		
	1.00		69.570	
	1.04	69.412		
	2.89	67.293		
	3.00		67.275	
	10.0		64.430	
	11.6	64.089		
502.44	0.10		76.527	
	0.11	76.225		
	0.19	75.147		
	0.20		75.147	
	0.48	73.311		
	0.50		73.317	
	1.00	71.861	71.927	
	4.96	68.568		
	5.00		68.609	
	30.0		64.043	
	32.5	63.923		
	47.0	62.465		
	50.0		61.402	
^a Values interg	polated in ten	nperatures.		

3. Force Constants for the Stockmayer (12-6-3) Potential (13). Parameters of the Stockmayer (12-6-3) potential, ϵ/k and $\sigma,$ given in Table IV were evaluated from the B values by the trial and error procedure described in detail by Hirschfelder et al. (14). The values of k and N_A used were those given by Cohen and Taylor (15), and the dipole moment was taken from McClellan (16). Table IV also compares the present values with those evaluated by Polak and Lu (5) by a least-squares regression analysis, and with the values of the Stockmayer (18-6-3) potential parameters given by Rätzsch and Freydank (6). The agreement between the present values of ϵ/k and those of Polak and Lu based on the B data of Perez Masia and Diaz Pena (2) is excellent, whereas our b_0 , t^{\bullet} , and σ values are closer to those of Polak and Lu based on the B data of Fogg et al. (3). This discrepancy may be due to, in part, the difference in the technique of evaluation and also the inclusion of B values at higher temperatures in this work. However, the reliability of the present force constants, ϵ/k and $\sigma,$ was test checked by computing B values at 431.68 K using the elaborate procedure and newer tables of Gallagher and Klein (17). Agreement between the experimental and the calculated B was within 0.5%. The present values of ϵ/k and σ can therefore be used in computing second virial coefficients from the Stockmayer (12-6-3) potential within reasonable accuracy.

Table IV. Comparison of the Stockmayer Force Constants for Dichloromethane

potential model	N	μ/D	$\sigma/{ m \AA}$	$(\epsilon/k)/{ m K}$	${{b_0}/{{({ m cm}^3 \ { m mol}^{-1})}}}$	t*	temp. range/ K	ref
(12-6-3)	7	1.631	2.764	672.6	26.63	0.48	319.4-382.6	5
(12-6-3)	5	1.631	3.150	716.6	39.42	0.304	323.16 - 423.16	5
(12-6-3)	10	1.631	2.648	715.29	23.423	0.513	349.53-510	this work
(18-6-3)	7	1.631	3.80	642.62			323.16 - 423.16	6

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Glossary

b ₀	second virial coefficient for a rigid sphere of diame-
R	second virial coefficient in eq. 1 cm ³ mol ⁻¹
D D/	second virial coefficient in eq. 2, etm ⁻¹
6	third virial coefficient in eq. 2, atm
	third virial coefficient in eq. 1, cm ² mol -
C ^r	third virial coefficient in eq 2, atm ²
G	Gibbs energy, cal mol ⁻¹
Н	enthalpy, cal mol ⁻¹
k	Boltzmann constant, J K ⁻¹
N	number of experimental points
N _A	Avogadro number, mol ⁻¹
Р	pressure, atm
R	gas constant, cm ³ atm K ⁻¹ mol ⁻¹
S	entropy, cal mol ⁻¹ K ⁻¹
t *	parameter in the Stockmayer potential
T	temperature, K
T _c	critical temperature, K
U	internal energy, cal mol ⁻¹
V	molar volume, cm ³ mol ⁻¹
Ζ	compression factor
e	parameter in the Stockmayer potential, J
μ	dipole moment, D
φ	fugacity coefficient

- σ
- parameter in the Stockmayer potential, Å

Superscript

٥ property of the substance in its standard state

Registry No. Dichloromethane, 75-09-2.

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Ion-Exchange Equilibria: Study of the Na⁺/Mn²⁺, Na⁺/Ni²⁺, and Na⁺/Cu²⁺ Exchanges on Several Lewatit Cation Exchangers

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Experimental equilibrium data of several Lewatit + Na⁺ + B²⁺ ion-exchange binary systems are given under various conditions: (I) with three exchangers of different cross-linking degree (S-100, S-115, and SP-120); (II) with two sizes of particles (SP-120: D1 = 0.74 mm, and D2 = 0.96 mm); (III) at three total ionic concentration levels of the solution phase (1, 0.1, and 0.01 N); and (IV) with three different counterions ($B^{2+} = Mn^{2+}$, Ni^{2+} , and Cu²⁺). By means of a previously developed method, where the Wilson equation is employed to evaluate the activity coefficients in the resin phase, the equilibrium constants are determined. The influence of the above variables is studied by using these equilibrium constants.

Introduction

The variables that can influence the equilibrium in a ion-exchange system can be classified in three groups: (1) those that depend on the exchanger such as the type of exchanger, cross-linking degree, and size of the particle; (2) variables that depend on the solution phase such as the kind of solvent, the set of ions involved, both interchangeable and uninterchangeable, and the total ionic concentration; and (3) those that depend on the working conditions such as pressure or temperature.

In order to carry out a systematic study of ion-exchange equilibrium it seems necessary to establish previously a parameter or a set of parameters that characterize in the best possible way such equilibrium. In a previous paper (1), the