# PVT Properties of Dimethyl(isopropenylethynyl)carbinol at Temperatures from 293 to 363 K and Pressures from 0.1 to 350 MPa, Studied Using the Micro-PVT System 

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#### Abstract

The effect of pressure on the volume of dimethyl(isopropenylethynyl)carbinol, $\left(\mathrm{CH}_{3}\right)_{2} \mathrm{C}(\mathrm{OH}) \mathrm{C} \equiv \mathrm{CC}\left(\mathrm{CH}_{3}\right)=$ $\mathrm{CH}_{2}$, has been measured over the pressure range ( 0.1 to 350 ) MPa and the temperature range ( 293.15 to 363.15) K using the Micro-PVT system, a new universal instrument for measuring the thermodynamic characteristics of materials. Experimental data on the variation of the volume ratio $\mathrm{k}=\mathrm{V}_{\mathrm{p}} / \mathrm{V}_{0}\left(\mathrm{~V}_{\mathrm{p}}=\right.$ sample volume at given pressure $\mathrm{P} ; \mathrm{V}_{0}=$ sample volume at 0.1 MPa ) with pressure have been approximated by two equations allowing the interpolation and extrapolation of experimental data with an acceptable uncertainty. The temperature dependence of the coefficients of the Tait equation has been studied, enabling the calculation of the volume ratio with an acceptable uncertainty for a wider temperature and pressure range. The isothermal compressibility, the isobaric thermal expansivity, and the increment of isobaric molar heat capacity from its value at 0.1 MPa have been calculated using the experimental data obtained.


## Introduction

The research of foams and colloid media and processes in these media is a promising research area. ${ }^{1-3}$ Flotation reagents used in the process of oreenrichment should have a high sorption to floatable materials. The application of materials composed of molecules having a succession of double and triple bonds as flotation reagents is promising for increased sorption in sulfide media having defective polycarbon structures. Specialists from Evrofinchermetholding Ltd., Moscow, were the first to synthesize dimethyl(isopropenylethynyl)carbinol (DMIPEC), a chemical product which meets these requirements. Its chemical formula is $\left(\mathrm{CH}_{3}\right)_{2} \mathrm{C}(\mathrm{OH}) \mathrm{C} \equiv \mathrm{CC}\left(\mathrm{CH}_{3}\right)=\mathrm{CH}_{2}$. DMIPEC was successfully used as a flotation reagent for the enrichment of ores of precious and nonferrous metals and coals. The results achieved have formed the basis for broad research on the physical and chemical properties of this synthesized material to extend its field of application. The results of PVT measurement of DMIPEC at eight temperatures over the temperature range ( 293.15 to 363.15 ) K and the pressure range ( 0.1 to 350 ) MPa are presented in this article.

## Experimental Section

DMIPEC was synthesized by means of direct interaction of acetylene with acetone and subsequent dehydration of the reaction product at thestage of rectification. The purity of the synthesized samples was $98.5 \%$ and was determined with gas-liquid chromatography using a conventional type gas-liquid chromatograph of "Chrom-5" model with 3.5 m glass column. Hydroquinone, a polymerization inhibitor,

[^0]with a concentration of $0.1 \%$ was added to the material to prevent polymerization. The synthesized material was stored at room temperature in closed canisters.

The Micro-PVT system, ${ }^{4}$ developed at VNIIGAZ and intended for the measurement of PVT properties of liquid and solid materials, was used for the study on the DMIPEC samples. The system indudes an electromechanical press and an electronic control unit connected to a personal computer. The Micro-PVT system operates with a high degree of automation and accuracy.

Measurements of the thermodynamic properties of materials using the Micro-PVT system are based on the socalled "piston in a bottle" method. Samples of the material are placed in a cylindrical high-pressure chamber made of high-strength stainless steel; the pressure is obtained by means of a tightened piston inserted from one of the flanks of the chamber and moved by a stepping motor. Friction problems in sealing are overcome by having a rotating piston, with varying rotational speeds available. The sample volume is determined on the basis of a preliminary calibration and known piston displacement, and the pressure is cal culated from the extension of a spring connected with the piston and calibrated with high accuracy; the piston movement and spring extension are controlled with electronic displacement transducers accurate to $0.1 \mu \mathrm{~m}$. A temperature sensor placed at the other end of the highpressure chamber measures the temperature in the chamber, and this sensor simultaneously seals the chamber. Data on the sample volume V , its pressure P , and temperature $T$ are recorded by the computer and used for further treatment of the results. The system allows experiments to be carried out in different regimes (isothermal, isochoric, isobaric, quasiadiabatic) and allows the experimental results obtained to be treated in real time, providing

Table 1. Experimental Data for Volume Ratio $k=V_{p} / V_{0}$ for DMIPEC at Temperatures from (293.15 to 363.15) $K$

| $\mathrm{P} / \mathrm{MPa}$ | k | P/MPa | k | P/ MPa | k | $\mathrm{P} / \mathrm{MPa}$ | k | P/ MPa | k | P/MPa | k | $\mathrm{P} / \mathrm{MPa}$ | k | P/ MPa | k |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{T}=293.15 \mathrm{~K}$ |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 2.146 | 0.9983 | 19.774 | 0.9849 | 49.730 | 0.9661 | 100.155 | 0.9416 | 160.438 | 0.9194 | 210.803 | 0.9041 | 259.523 | 0.8916 | 309.664 | 0.8801 |
| 4.229 | 0.9966 | 24.848 | 0.9814 | 55.832 | 0.9628 | 109.925 | 0.9377 | 170.631 | 0.9161 | 219.885 | 0.9019 | 269.671 | 0.8892 | 320.826 | 0.8777 |
| 6.027 | 0.9952 | 30.303 | 0.9777 | 59.977 | 0.9606 | 120.514 | 0.9334 | 179.749 | 0.9132 | 230.442 | 0.8989 | 280.119 | 0.8868 | 330.350 | 0.8757 |
| 7.941 | 0.9938 | 35.371 | 0.9746 | 69.751 | 0.9556 | 130.102 | 0.9299 | 190.530 | 0.9100 | 239.711 | 0.8965 | 290.619 | 0.8844 | 340.004 | 0.8737 |
| 9.883 | 0.9923 | 40.563 | 0.9713 | 79.867 | 0.9506 | 139.722 | 0.9264 | 199.134 | 0.9074 | 249.547 | 0.8940 | 299.343 | 0.8823 | 349.479 | 0.8718 |
| 15.527 | 0.9881 | 44.973 | 0.9687 | 90.354 | 0.9458 | 150.491 | 0.9227 |  |  |  |  |  |  |  |  |
| $\mathrm{T}=303.15 \mathrm{~K}$ |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 1.995 | 0.9984 | 19.971 | 0.9844 | 49.926 | 0.9647 | 100.353 | 0.9395 | 159.288 | 0.9171 | 210.136 | 0.9015 | 260.575 | 0.8882 | 310.215 | 0.8767 |
| 4.150 | 0.9966 | 24.856 | 0.9810 | 54.734 | 0.9620 | 110.273 | 0.9353 | 170.473 | 0.9134 | 220.529 | 0.8986 | 270.633 | 0.8857 | 319.596 | 0.8745 |
| 5.960 | 0.9950 | 30.169 | 0.9771 | 59.757 | 0.9592 | 120.626 | 0.9311 | 179.747 | 0.9105 | 229.709 | 0.8961 | 280.979 | 0.8833 | 330.622 | 0.8721 |
| 7.862 | 0.9935 | 35.009 | 0.9739 | 70.490 | 0.9536 | 130.010 | 0.9277 | 189.667 | 0.9074 | 239.354 | 0.8937 | 289.981 | 0.8811 | 339.911 | 0.8703 |
| 9.789 | 0.9919 | 39.957 | 0.9708 | 80.306 | 0.9487 | 140.143 | 0.9238 | 200.090 | 0.9043 | 250.688 | 0.8906 | 300.945 | 0.8787 | 349.593 | 0.8683 |
| 15.190 | 0.9879 | 45.162 | 0.9675 | 89.427 | 0.9443 | 150.800 | 0.9202 |  |  |  |  |  |  |  |  |
| $\mathrm{T}=313.15 \mathrm{~K}$ |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 1.960 | 0.9983 | 19.774 | 0.9836 | 49.527 | 0.9632 | 99.909 | 0.9370 | 160.205 | 0.9131 | 209.523 | 0.8974 | 260.892 | 0.8835 | 310.192 | 0.8718 |
| 3.783 | 0.9967 | 25.258 | 0.9796 | 55.073 | 0.9597 | 109.462 | 0.9326 | 170.138 | 0.9097 | 219.998 | 0.8943 | 270.914 | 0.8810 | 319.385 | 0.8698 |
| 6.039 | 0.9947 | 30.476 | 0.9758 | 59.884 | 0.9570 | 120.530 | 0.9279 | 179.331 | 0.9066 | 230.741 | 0.8914 | 279.665 | 0.8788 | 330.883 | 0.8672 |
| 8.333 | 0.9928 | 35.132 | 0.9725 | 69.895 | 0.9516 | 130.280 | 0.9241 | 190.029 | 0.9032 | 240.054 | 0.8887 | 290.272 | 0.8763 | 340.839 | 0.8651 |
| 10.166 | 0.9913 | 39.871 | 0.9693 | 80.291 | 0.9464 | 140.376 | 0.9202 | 199.599 | 0.9003 | 249.697 | 0.8864 | 299.316 | 0.8743 | 350.844 | 0.8631 |
| 15.102 | 0.9872 | 44.971 | 0.9660 | 90.217 | 0.9415 | 149.557 | 0.9169 |  |  |  |  |  |  |  |  |
| $\mathrm{T}=323.15 \mathrm{~K}$ |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 2.107 | 0.9980 | 20.022 | 0.9820 | 49.902 | 0.9606 | 100.505 | 0.9331 | 160.880 | 0.9087 | 210.032 | 0.8925 | 260.486 | 0.8786 | 309.725 | 0.8669 |
| 3.904 | 0.9963 | 25.037 | 0.9779 | 55.233 | 0.9572 | 109.762 | 0.9289 | 170.014 | 0.9055 | 220.489 | 0.8895 | 270.184 | 0.8763 | 319.185 | 0.8647 |
| 5.922 | 0.9944 | 29.963 | 0.9741 | 59.715 | 0.9544 | 120.415 | 0.9242 | 180.178 | 0.9021 | 229.477 | 0.8870 | 280.061 | 0.8738 | 329.890 | 0.8625 |
| 8.097 | 0.9924 | 35.163 | 0.9704 | 69.949 | 0.9485 | 130.952 | 0.9199 | 190.654 | 0.8985 | 239.801 | 0.8841 | 290.329 | 0.8711 | 340.003 | 0.8605 |
| 9.819 | 0.9909 | 39.716 | 0.9674 | 79.978 | 0.9433 | 140.651 | 0.9162 | 200.060 | 0.8957 | 250.581 | 0.8813 | 300.709 | 0.8688 | 349.196 | 0.8586 |
| 15.140 | 0.9861 | 44.739 | 0.9640 | 89.866 | 0.9383 | 149.859 | 0.9128 |  |  |  |  |  |  |  |  |
| $\mathrm{T}=333.15 \mathrm{~K}$ |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 1.995 | 0.9981 | 20.282 | 0.9809 | 50.132 | 0.9585 | 100.474 | 0.9304 | 159.965 | 0.9055 | 209.962 | 0.8887 | 260.672 | 0.8744 | 310.845 | 0.8621 |
| 4.109 | 0.9958 | 25.310 | 0.9767 | 54.972 | 0.9553 | 109.529 | 0.9260 | 169.740 | 0.9020 | 220.009 | 0.8857 | 270.629 | 0.8718 | 320.090 | 0.8598 |
| 6.119 | 0.9939 | 29.991 | 0.9728 | 59.722 | 0.9523 | 120.396 | 0.9212 | 179.787 | 0.8985 | 230.365 | 0.8827 | 280.742 | 0.8691 | 329.268 | 0.8578 |
| 8.169 | 0.9919 | 34.827 | 0.9689 | 70.493 | 0.9458 | 130.670 | 0.9168 | 190.515 | 0.8949 | 239.808 | 0.8801 | 289.265 | 0.8672 | 340.468 | 0.8552 |
| 10.257 | 0.9899 | 40.250 | 0.9652 | 80.338 | 0.9403 | 140.203 | 0.9129 | 200.159 | 0.8918 | 250.921 | 0.8771 | 300.037 | 0.8646 | 350.531 | 0.8531 |
| 15.108 | 0.9856 | 45.146 | 0.9619 | 90.180 | 0.9354 | 150.470 | 0.9091 |  |  |  |  |  |  |  |  |
| $\mathrm{T}=343.15 \mathrm{~K}$ |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 1.995 | 0.9979 | 19.707 | 0.9803 | 50.169 | 0.9560 | 100.360 | 0.9269 | 160.153 | 0.9010 | 209.509 | 0.8841 | 259.652 | 0.8697 | 309.625 | 0.8573 |
| 4.098 | 0.9957 | 25.151 | 0.9756 | 55.376 | 0.9525 | 110.585 | 0.9220 | 170.995 | 0.8972 | 219.705 | 0.8809 | 269.648 | 0.8669 | 320.675 | 0.8547 |
| 6.312 | 0.9933 | 30.229 | 0.9712 | 59.958 | 0.9496 | 120.199 | 0.9176 | 179.963 | 0.8940 | 230.065 | 0.8779 | 280.066 | 0.8643 | 330.081 | 0.8526 |
| 8.209 | 0.9913 | 34.750 | 0.9674 | 70.309 | 0.9433 | 130.116 | 0.9132 | 190.495 | 0.8903 | 239.884 | 0.8751 | 290.096 | 0.8619 | 339.961 | 0.8505 |
| 10.154 | 0.9894 | 39.942 | 0.9634 | 79.980 | 0.9375 | 139.831 | 0.9090 | 199.934 | 0.8872 | 250.253 | 0.8722 | 300.692 | 0.8594 | 349.982 | 0.8484 |
| 15.340 | 0.9843 | 45.379 | 0.9593 | 90.329 | 0.9319 | 149.819 | 0.9050 |  |  |  |  |  |  |  |  |
| $\mathrm{T}=353.15 \mathrm{~K}$ |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 2.096 | 0.9976 | 20.247 | 0.9784 | 49.988 | 0.9536 | 100.490 | 0.9230 | 160.359 | 0.8966 | 210.064 | 0.8791 | 259.756 | 0.8645 | 309.585 | 0.8517 |
| 4.139 | 0.9955 | 24.915 | 0.9741 | 55.572 | 0.9496 | 109.635 | 0.9185 | 169.864 | 0.8930 | 220.050 | 0.8759 | 269.742 | 0.8617 | 320.398 | 0.8493 |
| 6.255 | 0.9933 | 30.368 | 0.9693 | 59.977 | 0.9467 | 120.041 | 0.9137 | 179.884 | 0.8893 | 230.382 | 0.8728 | 279.658 | 0.8591 | 329.943 | 0.8470 |
| 8.030 | 0.9912 | 34.882 | 0.9655 | 70.046 | 0.9403 | 130.683 | 0.9088 | 190.168 | 0.8859 | 240.996 | 0.8696 | 290.063 | 0.8565 | 339.715 | 0.8448 |
| 9.873 | 0.9893 | 39.803 | 0.9613 | 80.405 | 0.9341 | 140.481 | 0.9047 | 200.536 | 0.8823 | 250.209 | 0.8671 | 300.620 | 0.8538 | 349.621 | 0.8427 |
| 15.090 | 0.9836 | 44.906 | 0.9574 | 90.393 | 0.9284 | 150.296 | 0.9006 |  |  |  |  |  |  |  |  |
| $\mathrm{T}=363.15 \mathrm{~K}$ |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 1.971 | 0.9976 | 20.123 | 0.9765 | 50.618 | 0.9501 | 100.120 | 0.9190 | 160.764 | 0.8914 | 210.325 | 0.8736 | 259.699 | 0.8586 | 310.736 | 0.8453 |
| 4.253 | 0.9947 | 25.535 | 0.9711 | 55.299 | 0.9465 | 110.520 | 0.9136 | 170.184 | 0.8878 | 220.154 | 0.8704 | 269.535 | 0.8559 | 319.515 | 0.8431 |
| 6.198 | 0.9922 | 30.211 | 0.9668 | 60.149 | 0.9431 | 120.641 | 0.9088 | 179.930 | 0.8840 | 230.199 | 0.8671 | 281.200 | 0.8528 | 330.746 | 0.8407 |
| 8.222 | 0.9898 | 34.745 | 0.9629 | 70.272 | 0.9364 | 130.605 | 0.9042 | 189.997 | 0.8806 | 240.545 | 0.8641 | 289.699 | 0.8506 | 340.298 | 0.8387 |
| 10.012 | 0.9876 | 40.213 | 0.9583 | 80.301 | 0.9303 | 140.872 | 0.8996 | 200.625 | 0.8769 | 250.048 | 0.8613 | 300.305 | 0.8479 | 350.168 | 0.8364 |
| 58 | 0.9820 | 45.425 | 0.9540 | 90.443 | 0.9243 | 150.455 | 0.8956 |  |  |  |  |  |  |  |  |

Table 2. Densities of DMIPEC at Atmospheric Pressure ( 0.1 MPa ) and Temperatures from (293.15 to 363.15) K

| $\mathrm{T} / \mathrm{K}$ | $\rho / \mathrm{kg} \cdot \mathrm{m}^{-3}$ | $\mathrm{~T} / \mathrm{K}$ | $\rho / \mathrm{kg}^{-3} \mathrm{~m}^{-3}$ |
| :---: | :---: | :---: | :---: |
| 293.15 | 884.1 | 333.15 | 851.8 |
| 303.15 | 876.9 | 343.15 | 842.9 |
| 313.15 | 869.0 | 353.15 | 833.6 |
| 323.15 | 860.6 | 363.15 | 823.7 |

complex information on the thermodynamic parameters of the sample and allowing one to adjust the measurement regime and to improve the accuracy of the experimental results.

The volume of the chamber can be varied from ( 0.3 to $0.7) \mathrm{cm}^{3}$; the pressure can be varied from ( 0.1 to 500 ) MPa with an adjustable rate of compression and decompression over the range $0.001 \mathrm{MPa} \cdot \mathrm{s}^{-1}$ to $12.5 \mathrm{MPa} \cdot \mathrm{s}^{-1}$. The pos-
sibility of obtaining a large number of experimental PVT points is an advantage of the Micro-PVT system. The number of data points and time intervals are set by the operator or are selected automatically by the system, depending on the requirements of the experiment. The maximum sampling frequency of experimental points is 10 Hz .
The Micro-PVT system has been calibrated at the National Engineering Laboratory (Glasgow, Scotland). The uncertainty of temperature measurement, based on temperature sensor calibration on IPTS-90 against a standard platinum resistance thermometer, was estimated to be $\pm 0.03 \mathrm{~K}$ over the temperature range (273 to 373) K; the maximum relative error for the measurement of the volume and pressure was estimated to be $\pm 0.2 \%$ over the ranges of values used in the experiments.

Table 3. Coefficients for the Dependence of the Volume Ratio on Pressure for DMIPEC at Temperatures from (293.15 to 363.15) K

| T/K | C | B/ MPa | $\mathrm{a}_{0} / \mathrm{MPa}$ | $\mathrm{a}_{1}$ | $-\mathrm{a}_{2} / \mathrm{GPa}^{-1}$ | $\mathrm{a}_{3} / \mathrm{GPa}^{-2}$ | B*/MPa |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 293.15 | 0.2049 | 108.40 | 1215.74 | 5.254 | 3.06 | 1.13 | 110.13 |
| 303.15 | 0.2051 | 103.39 | 1159.60 | 5.296 | 3.54 | 1.72 | 104.89 |
| 313.15 | 0.2071 | 98.03 | 1092.95 | 5.297 | 4.48 | 3.77 | 97.87 |
| 323.15 | 0.2052 | 89.67 | 1005.31 | 5.398 | 5.22 | 5.03 | 90.92 |
| 333.15 | 0.2077 | 85.93 | 953.18 | 5.221 | 4.31 | 3.06 | 85.37 |
| 343.15 | 0.2076 | 79.95 | 890.74 | 5.248 | 5.25 | 5.23 | 79.48 |
| 353.15 | 0.2095 | 75.44 | 832.31 | 5.210 | 5.46 | 5.52 | 73.81 |
| 363.15 | 0.2083 | 68.79 | 759.14 | 5.239 | 5.61 | 5.43 | 67.96 |

The densities of the samples at atmospheric pressure were measured at different temperatures with an Anton Paar DMA 58 densimeter. The uncertainty of the density measurements was estimated to be $\pm 0.5 \mathrm{~kg} \cdot \mathrm{~m}^{-3}$. PVT curve measurements were carried out on samples using a PETROTEST VC5/5 liquid thermostat. The measurements were carried out under isothermal conditions; the initial volume of the samples was about $0.5 \mathrm{~cm}^{3}$, and the pressure varied over the range ( 0.1 to 350 ) MPa. A thousand experimental points per measurement were obtained for each PVT curve.

## Results and Discussion

The pressure dependence of the volume ratio $\mathrm{k}=\mathrm{V}_{\mathrm{p}} / \mathrm{V}_{0}$ (the ratio of the volume $\mathrm{V}_{\mathrm{p}}$ at the current pressure P to the initial volume $\mathrm{V}_{0}$ at a pressure of 0.1 MPa ) is a convenient form for the presentation of the effect of pressure on the sample volume. The experimental data obtained for the volume ratio, $k$, of DMIPEC are given in Table 1 over the temperature range (293.15 to 363.15 ) K and the pressure range ( 0.1 to 350 ) MPa. The total estimated uncertainty of the $k$ measurement was found to be $\pm 0.25 \%$.

The densities of DMIPEC at atmospheric pressure (0.1 MPa ) and at the same temperatures are given in Table 2. The densities from this table can be used for the calculation of the pressure dependence of density over the range of measured pressures and temperatures by means of dividing the density at atmospheric pressure from Table 2 by the volume ratio, $k$, from Table 1 at the required temperature and pressure. The total estimated uncertainty of the density calculation was found to be $\pm 0.3 \%$.

A great number of equations have been developed for the presentation of experimental results of PVT measurements. A number of approximation coefficients are used in many of these equations. Malhotra and Woolf 5 have shown that the dependence

$$
\begin{equation*}
1-k=C \log \left(\frac{B+P}{B+0.1}\right) \tag{1}
\end{equation*}
$$

which is a Tait integral equation, ${ }^{6}$ allows the effect of pressure on the volume of liquids at a given temperature to be approximated with high accuracy and over a wide range of pressures using only two coefficients: C and B. In addition, the experimental results can be extrapolated outside the range of the measurements with an acceptable accuracy using eq 1 , including a gas-liquid transition zone. The coefficient $C$ as a rule is practically independent of temperature, and having studied the temperature dependence of the coefficient $B$, it is possible to obtain an equation which gives the volume ratio to a high accuracy over a large range of temperatures and pressures.

Information on the volume ratio, k, allows one to calculate the secant bulk modulus, K , the pressure dependence of which is usually approximated as follows by a polynomial of third degree:


Figure 1. Variation of the coefficient $\mathrm{B}^{*}$ with temperature.

$$
\begin{equation*}
K=\frac{P-P_{0}}{1-k}=a_{0}+a_{1} P+a_{2} P^{2}+a_{3} P^{3} \tag{2}
\end{equation*}
$$

where $P_{0}$ refers to atmospheric pressure ( 0.1 MPa ). The coefficients $a_{0}, a_{1}, a_{2}$, and $a_{3}$ are also a convenient method for the presentation of experimental data, and they may be used for the interpolation of experimental results obtained over the measurement range. Unlike eq 1, eq 2 is in general unsuitable for extrapolation outside the pressure range of the experimental data. However, eq 2 is usually a more accurate representation of the experimental data than eq 1.

Experimental results of the measurement of $k$, presented in Table 1, have been approximated by a least-squares fit using eqs 1 and 2 , and the coefficients for such an approximation are given in Table 3 for all temperatures for which measurements have been carried out. The coefficients C and B presented in Table 3 allow an approximation of the experimental data from Table 1 accurate to within $\pm 0.1 \%$. The coefficients $a_{0}, a_{1}, a_{2}$, and $a_{3}$ presented in Table 3 allow an approximation of the experimental data from Table 1 accurate to within $\pm 0.1 \%$ for pressures $P \geq$ 50 MPa and accurate to within $\pm 0.15 \%$ for lower pressures.

Analysis of the temperature dependence of the coefficient C presented in Table 3 shows that this coefficient depends weakly on temperature, and its average value is $0.2069 \pm$ 0.0017. The experimental data obtained from Table 1 were then approximated by a least-squares fit using eq 1 for a fixed value of $C=0.2069$ for all temperatures studied, and the values obtained for $\mathrm{B} *$ are presented in Table 3 as well. Using the coefficients $C=0.2069$ and corresponding $B^{*}$ values enables the approximation of experimental data from Table 1 accurate to within $\pm 0.15 \%$. The variation of B* with temperature is presented in Figure 1. It is seen from this figure that $B^{*}$ decreases linearly with increasing temperature. The approximation coefficients in the equation

$$
\begin{equation*}
B^{*}=B_{0}+B_{1} T \tag{3}
\end{equation*}
$$

Table 4. Isothermal Compressibility, $\mathbf{k}_{T}$, Secant Isothermal Compressibility, $\mathbf{k}_{\mathbf{T}}{ }^{*}$, Isobaric Thermal Expansivity, $\alpha$, and Increment of Isobaric Molar Heat Capacity, $\Delta C_{p}$, from Its Value at 0.1 MPa for DMIPEC at Temperatures from (293.15 to 363.15) K

| P/MPa | $\mathrm{k}_{\mathrm{T}} / 10^{-4} \mathrm{MPa}^{-1}$ | $\mathrm{k}_{\mathrm{T}}{ }^{*} / 10^{-4} \mathrm{MPa}^{-1}$ | $\alpha / 10^{-3} \mathrm{~K}^{-1}$ | $\Delta \mathrm{C}_{\mathrm{p}} / \mathrm{J} \cdot \mathrm{mol}^{-1} \cdot \mathrm{~K}^{-1}$ | P/MPa | $\mathrm{k}_{\mathrm{T}} / 10^{-4} \mathrm{MPa}^{-1}$ | $\mathrm{k}_{\mathrm{T}}{ }^{*} / 10^{-4} \mathrm{MPa}^{-1}$ | $\alpha / 10^{-3} \mathrm{~K}^{-1}$ | $\Delta \mathrm{C}_{\mathrm{p}} / \mathrm{J} \cdot \mathrm{mol}^{-1} \cdot \mathrm{~K}^{-1}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  | $\mathrm{T}=2$ | 3.15 K |  |  |  |  |
| 0.1 | 8.22 | 8.22 | 0.84 | 0 | 125 | 4.03 | 5.47 | 0.57 | -18 |
| 5 | 7.87 | 8.05 | 0.82 | -1 | 150 | 3.65 | 5.16 | 0.54 | -21 |
| 10 | 7.57 | 7.89 | 0.81 | -3 | 175 | 3.35 | 4.88 | 0.52 | -23 |
| 15 | 7.29 | 7.73 | 0.79 | -4 | 200 | 3.08 | 4.64 | 0.50 | -24 |
| 20 | 7.04 | 7.58 | 0.77 | -5 | 225 | 2.86 | 4.43 | 0.49 | -26 |
| 30 | 6.57 | 7.30 | 0.74 | -6 | 250 | 2.66 | 4.25 | 0.47 | -27 |
| 40 | 6.16 | 7.04 | 0.72 | -8 | 275 | 2.49 | 4.08 | 0.46 | -29 |
| 60 | 5.49 | 6.58 | 0.67 | -11 | 300 | 2.35 | 3.93 | 0.45 | -30 |
| 80 | 4.94 | 6.18 | 0.63 | -13 | 325 | 2.21 | 3.79 | 0.44 | -31 |
| 100 | 4.49 | 5.84 | 0.60 | -16 | 350 | 2.09 | 3.67 | 0.43 | -32 |
| $\mathrm{T}=303.15 \mathrm{~K}$ |  |  |  |  |  |  |  |  |  |
| 0.1 | 8.62 | 8.62 | 0.91 | 0 | 125 | 4.19 | 5.65 | 0.58 | -19 |
| 5 | 8.24 | 8.43 | 0.88 | -1 | 150 | 3.80 | 5.32 | 0.56 | -22 |
| 10 | 7.92 | 8.25 | 0.86 | -3 | 175 | 3.47 | 5.03 | 0.53 | -24 |
| 15 | 7.63 | 8.08 | 0.83 | -4 | 200 | 3.19 | 4.78 | 0.51 | -26 |
| 20 | 7.35 | 7.91 | 0.81 | -5 | 225 | 2.95 | 4.56 | 0.49 | -27 |
| 30 | 6.86 | 7.60 | 0.78 | -7 | 250 | 2.75 | 4.37 | 0.48 | -29 |
| 40 | 6.43 | 7.32 | 0.75 | -8 | 275 | 2.58 | 4.19 | 0.47 | -30 |
| 60 | 5.71 | 6.83 | 0.70 | -11 | 300 | 2.42 | 4.04 | 0.46 | -31 |
| 80 | 5.14 | 6.40 | 0.66 | -14 | 325 | 2.28 | 3.90 | 0.45 | -32 |
| 100 | 4.67 | 6.04 | 0.62 | -16 | 350 | 2.16 | 3.77 | 0.44 | -33 |
| $\mathrm{T}=313.15 \mathrm{~K}$ |  |  |  |  |  |  |  |  |  |
| 0.1 | 9.15 | 9.15 | 0.98 | 0 | 125 | 4.38 | 5.91 | 0.60 | -20 |
| 5 | 8.77 | 8.93 | 0.94 | -2 | 150 | 3.97 | 5.56 | 0.57 | -23 |
| 10 | 8.42 | 8.73 | 0.90 | -3 | 175 | 3.62 | 5.26 | 0.54 | -25 |
| 15 | 8.10 | 8.54 | 0.87 | -4 | 200 | 3.33 | 4.99 | 0.52 | -27 |
| 20 | 7.80 | 8.35 | 0.85 | -5 | 225 | 3.08 | 4.76 | 0.51 | -29 |
| 30 | 7.27 | 8.01 | 0.81 | -7 | 250 | 2.86 | 4.55 | 0.49 | -30 |
| 40 | 6.79 | 7.70 | 0.78 | -9 | 275 | 2.68 | 4.37 | 0.48 | -32 |
| 60 | 6.02 | 7.17 | 0.73 | -12 | 300 | 2.51 | 4.20 | 0.46 | -33 |
| 80 | 5.40 | 6.71 | 0.68 | -15 | 325 | 2.36 | 4.05 | 0.45 | -34 |
| 100 | 4.90 | 6.32 | 0.64 | -17 | 350 | 2.23 | 3.91 | 0.45 | -35 |
| $\mathrm{T}=323.15 \mathrm{~K}$ |  |  |  |  |  |  |  |  |  |
| 0.1 | 9.94 | 9.94 | 1.04 | 0 | 125 | 4.55 | 6.22 | 0.63 | -21 |
| 5 | 9.45 | 9.69 | 0.99 | -2 | 150 | 4.10 | 5.83 | 0.59 | -24 |
| 10 | 9.04 | 9.44 | 0.96 | -3 | 175 | 3.74 | 5.50 | 0.56 | -26 |
| 15 | 8.67 | 9.22 | 0.92 | -4 | 200 | 3.43 | 5.22 | 0.54 | -28 |
| 20 | 8.32 | 9.00 | 0.90 | -6 | 225 | 3.18 | 4.97 | 0.52 | -30 |
| 30 | 7.71 | 8.60 | 0.85 | -8 | 250 | 2.95 | 4.75 | 0.50 | -32 |
| 40 | 7.18 | 8.24 | 0.81 | -9 | 275 | 2.76 | 4.55 | 0.49 | -33 |
| 60 | 6.32 | 7.62 | 0.76 | -13 | 300 | 2.59 | 4.37 | 0.47 | -35 |
| 80 | 5.64 | 7.11 | 0.71 | -16 | 325 | 2.44 | 4.20 | 0.46 | -36 |
| 100 | 5.10 | 6.68 | 0.67 | -18 | 350 | 2.31 | 4.05 | 0.45 | -37 |
| $T=333.15 \mathrm{~K}$ |  |  |  |  |  |  |  |  |  |
| 0.1 | 10.49 | 10.49 | 1.11 | 0 | 125 | 4.76 | 6.47 | 0.65 | -22 |
| 5 | 10.06 | 10.21 | 1.05 | -2 | 150 | 4.28 | 6.06 | 0.61 | -25 |
| 10 | 9.61 | 9.95 | 1.01 | -3 | 175 | 3.90 | 5.71 | 0.58 | -28 |
| 15 | 9.20 | 9.70 | 0.97 | -5 | 200 | 3.58 | 5.41 | 0.56 | -30 |
| 20 | 8.83 | 9.47 | 0.94 | -6 | 225 | 3.30 | 5.14 | 0.53 | -32 |
| 30 | 8.15 | 9.04 | 0.89 | -8 | 250 | 3.07 | 4.91 | 0.51 | -34 |
| 40 | 7.58 | 8.66 | 0.85 | -10 | 275 | 2.86 | 4.70 | 0.50 | -35 |
| 60 | 6.65 | 7.99 | 0.79 | -13 | 300 | 2.69 | 4.52 | 0.48 | -36 |
| 80 | 5.93 | 7.44 | 0.74 | -17 | 325 | 2.52 | 4.35 | 0.47 | -38 |
| 100 | 5.34 | 6.97 | 0.69 | -19 | 350 | 2.39 | 4.20 | 0.46 | -39 |
| $T=343.15 \mathrm{~K}$ |  |  |  |  |  |  |  |  |  |
| 0.1 | 11.22 | 11.22 | 1.18 | 0 | 125 | 4.94 | 6.78 | 0.68 | -24 |
| 5 | 10.72 | 10.91 | 1.11 | -2 | 150 | 4.45 | 6.34 | 0.64 | -27 |
| 10 | 10.22 | 10.61 | 1.06 | -3 | 175 | 4.05 | 5.97 | 0.60 | -29 |
| 15 | 9.78 | 10.33 | 1.01 | -5 | 200 | 3.70 | 5.64 | 0.58 | -32 |
| 20 | 9.36 | 10.06 | 0.98 | -6 | 225 | 3.42 | 5.36 | 0.55 | -34 |
| 30 | 8.62 | 9.58 | 0.92 | -9 | 250 | 3.18 | 5.11 | 0.53 | -35 |
| 40 | 7.99 | 9.15 | 0.88 | -11 | 275 | 2.97 | 4.89 | 0.51 | -37 |
| 60 | 6.97 | 8.42 | 0.82 | -14 | 300 | 2.78 | 4.69 | 0.49 | -38 |
| 80 | 6.18 | 7.81 | 0.76 | -17 | 325 | 2.62 | 4.50 | 0.48 | -39 |
| 100 | 5.56 | 7.31 | 0.72 | -20 | 350 | 2.48 | 4.33 | 0.47 | -41 |
| $T=353.15 \mathrm{~K}$ |  |  |  |  |  |  |  |  |  |
| 0.1 | 12.01 | 12.01 | 1.25 | 0 | 125 | 5.06 | 7.10 | 0.70 | -25 |
| 5 | 11.40 | 11.65 | 1.16 | -2 | 150 | 4.54 | 6.62 | 0.66 | -28 |
| 10 | 10.90 | 11.31 | 1.10 | -4 | 175 | 4.14 | 6.22 | 0.62 | -31 |
| 15 | 10.41 | 11.00 | 1.06 | -5 | 200 | 3.79 | 5.88 | 0.59 | -33 |
| 20 | 9.90 | 10.70 | 1.02 | -7 | 225 | 3.51 | 5.58 | 0.57 | -35 |
| 30 | 9.05 | 10.16 | 0.96 | -9 | 250 | 3.27 | 5.32 | 0.54 | -37 |
| 40 | 8.35 | 9.69 | 0.91 | -11 | 275 | 3.07 | 5.08 | 0.52 | -39 |
| 60 | 7.23 | 8.88 | 0.84 | -15 | 300 | 2.89 | 4.87 | 0.50 | -40 |
| 80 | 6.38 | 8.22 | 0.79 | -18 | 325 | 2.73 | 4.68 | 0.49 | -41 |
| 100 | 5.72 | 7.67 | 0.75 | -21 | 350 | 2.60 | 4.50 | 0.48 | -43 |

Table 4 (Continued)

| P/MPa | $\mathrm{k}_{\mathrm{T}} / 10^{-4} \mathrm{MPa}^{-1}$ | $\mathrm{k}^{*} / 10^{-4} \mathrm{MPa}^{-1}$ | $\alpha / 10^{-3} \mathrm{~K}^{-1}$ | $\Delta \mathrm{C}_{\mathrm{p}} / \mathrm{J} \cdot \mathrm{mol}^{-1} \cdot \mathrm{~K}^{-1}$ | P/MPa | $\mathrm{k}_{\mathrm{T}} / 10^{-4} \mathrm{MPa}^{-1}$ | $\mathrm{k}^{*} / 10^{-4} \mathrm{MPa}^{-1}$ | $\alpha / 10^{-3} \mathrm{~K}^{-1}$ | $\Delta \mathrm{C}_{\mathrm{p}} / \mathrm{J} \cdot \mathrm{mol}^{-1} \cdot \mathrm{~K}^{-1}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{T}=363.15 \mathrm{~K}$ |  |  |  |  |  |  |  |  |  |
| 0.1 | 13.16 | 13.16 | 1.31 | 0 | 125 | 5.28 | 7.48 | 0.72 | -26 |
| 5 | 12.41 | 12.74 | 1.21 | -2 | 150 | 4.73 | 6.96 | 0.68 | -29 |
| 10 | 11.73 | 12.33 | 1.14 | -4 | 175 | 4.30 | 6.52 | 0.64 | -32 |
| 15 | 11.15 | 11.95 | 1.10 | -5 | 200 | 3.95 | 6.15 | 0.61 | -35 |
| 20 | 10.55 | 11.60 | 1.06 | -7 | 225 | 3.65 | 5.83 | 0.58 | -37 |
| 30 | 9.64 | 10.97 | 0.99 | -10 | 250 | 3.41 | 5.55 | 0.56 | -39 |
| 40 | 8.85 | 10.42 | 0.94 | -12 | 275 | 3.20 | 5.30 | 0.53 | -41 |
| 60 | 7.61 | 9.48 | 0.87 | -16 | 300 | 3.03 | 5.07 | 0.51 | -42 |
| 80 | 6.69 | 8.73 | 0.82 | -19 | 325 | 2.87 | 4.86 | 0.50 | -43 |
| 100 | 5.97 | 8.11 | 0.77 | -23 | 350 | 2.73 | 4.68 | 0.49 | -45 |

were found to be 288.55 MPa and $-0.6087 \mathrm{MPa} \cdot \mathrm{K}^{-1}$, respectively, for $B_{0}$ and $B_{1}$.

The estimate of the accuracy of the value of $k$ using the dependence obtained for $\mathrm{B}^{*}(\mathrm{~T})$ (see eq 3) and $\mathrm{C}=0.2069$ allows the calculation of the volume ratio, $k$, over the pressure range ( 0.1 to 350 ) MPa and the temperature range ( 293.15 to 363.15 ) K accurate to $\pm 0.25 \%$, and the extrapolation of data outside the measurement range over ( 0.1 to 700) MPa and (273.15 to 373.15 ) K accurate to $\pm 0.75 \%$. Outside this range extrapolation probably becomes unreliable.

The values of isothermal compressibility, $\mathrm{k}_{\mathrm{T}}$, presented in Table 4 were calculated using the following expression:
$\mathrm{k}_{\mathrm{T}}=-\left\{1 /\left[\left(\mathrm{P}-\mathrm{P}_{0}\right)-\mathrm{K}\right]\right\}\left\{1-\left(\mathrm{P}-\mathrm{P}_{0}\right) / \mathrm{K}(\partial \mathrm{K} / \partial \mathrm{P})_{T}\right\}$
The differentiation was carried out analytically, using data from Table 3 on the approximation of the secant bulk modulus by a cubic polynomial. The data presented in Table 4 show that the isothermal compressibility of DMIPEC increases significantly with increasing temperature. So, $k_{T}$ equals $8.22 \times 10^{-4} \mathrm{MPa}^{-1}$ at a temperature of 293.15 K and a pressure of 0.1 MPa , and it increases about one and a half times, becoming equal to $13.16 \times 10^{-4} \mathrm{MPa}^{-1}$ for an increase of temperature up to 363.15 K . The accuracy of the determination of $k_{T}$ is estimated as $\pm 0.4 \%$ for pressures $P \geq 50 \mathrm{MPa}$, and is perhaps greater for lower pressures.

Strictly speaking, the secant bulk modulus is not a thermodynamic function, since its determination is based on a ratio of finite increments (see eq 2), not on a differentiation. It is used because of its simplicity of calculation and because of the small number of experimental points usually available. The tangent bulk modulus $\mathrm{K}^{*}$ $=-\mathrm{V}(\partial \mathrm{P} / \partial \mathrm{V})_{\mathrm{T}}$ (or simply bulk modulus) and its inverse value, the isothermal compressibility $\mathrm{k}_{\mathrm{T}}$ (see also eq 4), which are calculated by the differentiation of the dependence of volume on pressure, are real thermodynamic values. The Micro-PVT system allows the determination of these derivatives with high accuracy due to the large number of experimental points and high accuracy of measurements. The coefficient of secant isothermal compressibility, $\mathrm{k}_{T} *$ (the inverse of the secant bulk modulus, K, calculated according to eq 2), is given in Table 4 for comparison. The values of $k_{T}$ and $k_{T} *$ coincide only at low pressures, and the difference between them increases with increasing pressure. This is stipulated by the fact that the ratio of finite increments is close to the derivative only for small increments.

Values of isobaric thermal expansivity, $\alpha$, presented in Table 4, were calculated according to formula

$$
\begin{equation*}
\alpha=\left(\partial \ln V_{m} / \partial T\right)_{P} \tag{5}
\end{equation*}
$$

where $\mathrm{V}_{\mathrm{m}}$ is the molar volume, which was determined on the basis of densities taken from Table 2 and volumeratios, k, cal culated according to eq 1 using coefficients from Table 3. The derivative was calculated analytically, for which purpose the dependence of $\ln \mathrm{V}_{\mathrm{m}}$ on temperature was approximated by a cubic polynomial. The accuracy of the determination of $\alpha$ is estimated as $\pm 0.5 \%$ for pressures $P$ $\geq 50 \mathrm{MPa}$ and is perhaps greater for lower pressures.

The calculated values of isobaric thermal expansivity, $\alpha$, allowed the calculation of the increment of isobaric molar heat capacity, $\Delta C_{p}$, from its value at 0.1 MPa according to the following formula:

$$
\begin{align*}
\Delta C_{p}=C_{p}-C_{p}(0.1 \mathrm{MPa}) & = \\
& -\int_{0.1}^{P} T M / \rho\left\{(\partial \alpha / \partial T)_{P}+\alpha^{2}\right\} d P \tag{6}
\end{align*}
$$

where $M$ is the molar mass and $\rho$ is the density at $P$. The data obtained are presented in Table 4. The values of $\alpha$ were represented by a quadratic in $T$ to enable analytic differentiation. $\Delta \mathrm{C}_{\mathrm{p}}$ has an error of about $\pm 2 \mathrm{~J} \cdot \mathrm{~mol}^{-1} \cdot \mathrm{~K}^{-1}$ for the studied range of pressures.

## Conclusion

The PVT properties of DMIPEC have been measured for eight temperatures over the temperature range (293.15 to 363.15) K and the pressure range ( 0.1 to 350 ) MPa. The data obtained on the volume ratio were approximated by the Tait equation, and the cal culated secant bulk modulus was approximated by a cubic polynomial. A dependence for the calculation of DMIPEC volumeratio has been proposed based on the analysis of the temperature dependence of coefficients in the Tait equation for different temperatures and the studied range of pressures. The isothermal compressibility, the isobaric thermal expansivity, and the increment of isobaric molar heat capacity from its value at 0.1 MPa were calculated.

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