# **Bubble-Point Pressures and Liquid Densities of Mixtures Blended** with Difluoromethane, Pentafluoroethane, and 1,1,1-Trifluoroethane

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Measurements of the bubble-point pressures and the liquid densities of the binary difluoromethane + 1,1,1-trifluoroethane (R-32 + R-143a) and the ternary difluoromethane + pentafluoroethane + 1,1,1-trifluoroethane (R-32 + R-125 + R-143a) refrigerant mixtures at temperatures from 280 to 320 K in 10 K intervals are reported. The experimental apparatus used in the measurement is a magnetic densimeter coupled with a variable volume cell. The estimated uncertainties are ±12 mK in temperature, ±10 kPa in pressure, and ±0.2% in density, whereas the uncertainty in the determined composition of mixture is estimated to be ±0.3 mass %. On the basis of the measured bubble-point pressures, the optimization of the binary interaction parameter of the Peng–Robinson equation, which reproduces the available data within ±1%, is discussed. Empirical correlations for the saturated-liquid and compressed-liquid densities of mixtures composed of R-32, R-125, and R-143a with an accuracy within ±0.3% are also developed.

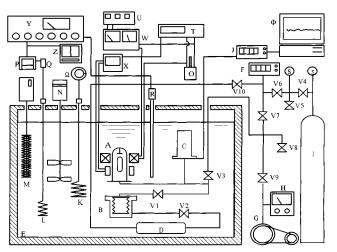
## Introduction

Binary and/or ternary mixtures composed of difluoromethane (R-32), pentafluoroethane (R-125), and 1,1,1-trifluoroethane (R-143a) are expected to be potential alternatives to chlorodifluoromethane (R-22) and R-502, an azeotropic blend of 48.8 mass % R-22 and 51.2 mass % R-115 (pentafluorochloroethane). This is due to their nearazeotropic characteristics that promote behavior similar to that of a pure compound. In an extensive research program on the promising HFC (hydrofluorocarbon) refrigerant mixtures, we have already published measurements of bubble-point pressures and liquid densities of the binary mixtures R-32 + R-125 (Widiatmo et al., 1993) and R-125 + R-143a (Widiatmo et al., 1995; Fujimine et al., 1999). The aim of the present study is to conduct similar measurements of the binary mixture R-32 + R-143a and the ternary mixture R-32 + R-125 + R-143a in wide ranges of composition and temperature.

### **Experimental Section**

A magnetic densimeter coupled with a variable-volume cell has been used in the present work to measure the liquid phase PVTx properties of the binary R-32 + R-143a and the ternary R-32 + R-125 + R-143a mixtures. The experimental apparatus is shown in Figure 1. The apparatus consists of a temperature control and measurement section, a pressure control and measurement section, and a density measurement section. The temperature is determined by means of a platinum resistance thermometer placed in the vicinity of the sample cell immersed in a thermostated fluid bath. The pressure is measured with a quartz pressure transducer, while the density is obtained

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**Figure 1.** Experimental apparatus: (A) magnetic densimeter; (B) variable-volume cell; (C) digital quartz-pressure transducer; (D) damper; (E) thermostat; (F) digital-quartz pressure gauge; (G) vacuum pump; (H) vacuum gauge; (I) nitrogen gas; (J) digital quartz pressure computer; (K) main heater; (L) subheater; (M) cooler; (N) stirrer; (O) standard resistor; (P) PID controller; (R) standard platinum resistance thermometer; (S) pressure gauge; (T) digital multimeter; (U) current controller; (V1–10) valves; (W) DC power supply; (X) galvanometer; (Y) thermometer bridge; (Z) pen recorder; ( $\Omega$ ) transformer; ( $\Phi$ ) personal computer.

by using a magnetic densimeter. By combining the magnetic densimeter with a variable-volume cell, it is possible to create a vapor—liquid coexisting state or a compressedliquid phase in the liquid sample. The bubble-point condition is determined from careful visual observation of the appearance and disappearance of a bubble in a liquid sample. In the case of mixtures, the composition of the mixture is determined through weighing the supply vessels before and after sample condensation into the variablevolume cell. Detailed explanations of the present apparatus as well as experimental procedures have been reported elsewhere (Widiatmo et al., 1993).

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$x_{R-32}/mol \cdot mol^{-1}$	<i>T</i> /K	<i>P</i> /MPa	$ ho/{ m kg}{\cdot}{ m m}^{-3}$	$x_{R-32}/mol \cdot mol^{-1}$	<i>T</i> /K	P/MPa	$ ho/{ m kg}{ m \cdot}{ m m}^{-1}$
0	$249.98_1$	0.281	1097.3	$0.502_{4}$	279.98 <sub>8</sub>	0.945	1005.9
0	$259.98_4$	0.401	1065.4	$0.502_4$	$289.98_8$	1.257	968.2
0	$269.98_7$	0.562	1032.7	$0.502_4$	299.987	1.642	925.3
0	$279.98_{8}$	0.764	998.4	$0.502_{4}$	309.987	2.108	878.0
0	$279.98_{8}$	0.762	996.9	$0.502_{4}$	319.987	2.667	822.7
0	$289.98_8$	1.015	959.4	$0.749_{7}$	$279.98_{8}$	0.992	1015.2
0	$289.98_8$	1.017	961.3	$0.749_{7}$	$289.98_8$	1.324	977.1
0	$299.98_{6}$	1.326	919.9	$0.749_{7}$	$299.98_{7}$	1.734	935.0
0	$299.98_8$	1.324	918.9	$0.749_{7}$	$309.98_7$	2.231	888.7
0	$309.98_5$	1.702	873.5	$0.749_{7}$	$319.98_{7}$	2.830	833.9
0	$309.98_{6}$	1.700	872.8	$0.749_{7}$	$319.98_7$	2.832	834.1
0	$319.98_{6}$	2.149	817.7	1.0	$279.98_{7}$	1.007	1030.0
0	$319.98_{6}$	2.150	818.6	1.0	$279.98_{7}$	1.007	1030.7
0	$329.98_{6}$	2.689	750.0	1.0	$289.98_{7}$	1.347	993.8
$0.248_3$	$279.98_{8}$	0.868	1001.4	1.0	$289.98_{7}$	1.349	993.2
$0.248_3$	$289.98_8$	1.152	963.6	1.0	$299.98_{6}$	1.773	952.6
$0.248_{3}$	299.987	1.502	920.9	1.0	$299.98_{6}$	1.772	951.9
$0.248_{3}$	309.98 <sub>7</sub>	1.925	873.8	1.0	$309.98_{6}$	2.293	907.5
$0.248_{3}$	319.987	2.433	818.5	1.0	$319.98_{4}$	2.918	856.5

Through the procedure adopted in the present study, the temperature, pressure, and density measurements can be performed within the uncertainties of  $\pm 12$  mK,  $\pm 10$  kPa, and  $\pm 0.2\%$ , respectively. The uncertainty in the composition of the mixture is estimated to be not greater than  $\pm 0.3$  mass %. Analyses of our sample purities by the manufacturers resulted in 99.99 mol % for R-32, 99.96 mol % for R-125, and 99.998 mol % for R-143a, respectively. No further purification has been done in the present work.

#### **Results and Discussion**

The bubble-point pressures and the liquid-phase PVTx properties of the binary R-32 + R-143a system have been measured for temperatures from 280 to 320 K at 10 K intervals and for compositions from 0 to 100 mass % in 25 mass % intervals. The measured bubble-point pressures and the saturated-liquid densities are listed in Table 1, while the liquid-phase PVTx properties are tabulated in Table 2.

Measurements of the bubble-point pressures and liquidphase PVTx properties of the ternary R-32 + R-125 + R-143a mixture at a composition of 33.56/33.28/33.16 mol % were completed for temperatures from 280 to 320 K at 10 K intervals. The measured property data are given in Tables 3 and 4.

As we have reported previously (Widiatmo et al., 1993), the binary R-32 + R-125 system is an azeotropic mixture for which a minimum temperature exists in the vicinity of a composition with 80 mass% R-32. The binary R-125 + R-143a system is also considered an azeotropic blend with a maximum temperature, especially at lower temperatures. For the ternary R-32 + R-125 + R-143a system, the present measurements cover only a single composition so that the existence of any azeotrope in this ternary system cannot be identified experimentally. However, from our calculation of the bubble-point and dew-point pressures of the ternary R-32 + R-143a system, it seems likely that the present ternary system is a nonazeotropic blend.

The bubble-point pressure representation in terms of the Peng–Robinson (PR) equation (Peng and Robinson, 1976) has been evaluated in the present work. Table 5 tabulates the parameters used in the original PR equation. The binary interaction parameters used in the calculation of the bubble-point pressures of the binary R-32 + R-143a and R-125 + R-143a systems were determined to be constant with temperature, while those of R-32 + R-125 were determined to be a function of temperature, as already reported (Widiatmo et al., 1997). The optimized binary

interaction parameters  $k_{ij}$  are given in Table 6, where *T* denotes the temperature in K.

The bubble-point pressure deviation from the PR equation is depicted in Figure 2. Also shown in Figure 2 are the available bubble-point pressures reported by Kleemiss (1997), Nagel and Bier (1995), Zhelezny et al. (1995), and Higashi (1993). Except for some data by Higashi which show strong temperature dependence, and a single datum by Nagel and Bier, the optimized PR equation represents the available measured data within  $\pm 1\%$ , as shown in Figure 2. Thus, the optimized PR equation can be utilized to calculate the temperature glide of the refrigerant mixtures. From such calculations it has been confirmed that the temperature glide of mixtures consisting of R-32, R-125, and R-143a is smaller than 1 K, as typically shown in Figure 3. This small temperature glide may indicate the mixtures are nearly azeotropic.

Figure 2 also illustrates the deviation of the measured bubble-point pressures of the ternary R-32 + R-125 + R-143a system from the PR equation, which merely requires the binary interaction parameters of the binary R-32 + R-125, R-32 + R-143a, and R-125 + R-143a systems given in Table 6. Figure 2 indicates that the PR equation optimized with the proposed procedure represents satisfactorily the bubble-point pressure of the ternary system. The deviation of the measured bubble-point pressures from the optimized PR equation does not exceed  $\pm 5$  kPa.

The measured saturated liquid densities of refrigerant mixtures composed of R-32, R-125, and R-143a are represented using the saturated-liquid density correlation given in eq 1. This type of equation has also successfully been applied to binary refrigerant mixtures composed of R-32, R-125, and R-134a (Widiatmo et al., 1995).

$$\frac{\rho_{\rm s}}{\rho_{\rm Cm}} - 1 = A_{\rm m} \left(1 - \frac{T}{T_{\rm Cm}}\right)^{a_{\rm m}} + B_{\rm m} \left(1 - \frac{T}{T_{\rm Cm}}\right)^{b_{\rm m}} + C_{\rm m} \left(1 - \frac{T}{T_{\rm Cm}}\right)^{c_{\rm m}} (1)$$

In eq 1,  $\rho_{\rm s}$  denotes the saturated-liquid density (in kg·m<sup>-3</sup>),  $\rho_{\rm Cm}$  and  $T_{\rm Cm}$  denote the critical density (in kg·m<sup>-3</sup>) and temperature (in K), respectively, and the subscript m denotes the mixture. The critical temperatures and critical densities for refrigerant mixtures are correlated by eqs 2 and 3, while the compositional dependence of the parameters  $a_{\rm m}$ ,  $b_{\rm m}$ ,  $c_{\rm m}$ ,  $A_{\rm m}$ ,  $B_{\rm m}$ , and  $C_{\rm m}$  was simply correlated by eq 4, where  $L_i$  denotes any arbitrary coefficient. The critical

_ <sub>32</sub> /mol∙mol <sup>−1</sup>	<i>T</i> /K	<i>P</i> /MPa	$ ho/{ m kg}{ m \cdot}{ m m}^{-3}$	$x_{R-32}/mol \cdot mol^{-1}$	<i>T</i> /K	P/MPa	ρ/kg•m⁻
0 <sup>a</sup>	$249.98_1$	0.281	1097.3	0 <sup>a</sup>	$299.98_{6}$	1.326	919.9
0	$249.98_{1}$	0.504	1097.9	0 <sup>a</sup>	$299.98_{6}$	1.324	918.9
0	$249.98_{1}$	1.001	1099.9	0	$299.98_{6}$	1.506	922.1
0	$249.98_{1}$	1.500	1101.3	0	$299.98_{6}$	2.002	927.4
0	$249.98_{1}$	1.993	1103.0	0	$299.98_{6}$	2.500	932.6
0	$249.98_{1}$	2.529	1104.6	0	$299.98_{6}$	3.003	937.2
0	$249.98_{1}$	3.006	1106.1	0 <sup>a</sup>	$309.98_{6}$	1.702	873.5
0 <sup>a</sup>	$259.98_4$	0.401	1065.4	0 <sup>a</sup>	$309.98_{6}$	1.700	872.8
0	$259.98_4$	1.015	1068.4	0	$309.98_{6}$	2.003	878.1
0	$259.98_4$	1.538	1070.4	0	$309.98_{6}$	2.505	885.5
0	$259.98_4$	2.014	1072.2	0	$309.98_{6}$	3.002	891.8
0	$259.98_4$	2.513	1074.1	0 <sup>a</sup>	$319.98_{6}$	2.149	817.
0	$259.98_4$	3.021	1076.2	0 <sup>a</sup>	$319.98_{6}$	2.150	818.0
0 <sup>a</sup>	$269.98_{7}$	0.562	1032.7	0	$319.98_{6}$	2.510	827.0
0	$269.98_{7}$	1.031	1035.7	0	$319.98_{6}$	2.998	837.9
0	$269.98_7$	1.532	1038.0	$0^a$	$329.98_{6}$	2.689	750.0
0	$269.98_7$	2.055	1040.7	0	$329.98_{6}$	3.013	764.0
0	$269.98_7$	2.524	1043.0	$0.248_3^a$	$279.98_{8}$	0.868	1001.4
0	269.987	3.013	1045.2	$0.248_{3}$	$279.98_{8}$	1.501	1005.
0 <sup>a</sup>	$279.98_{8}$	0.764	998.4	$0.248_{3}$	$279.98_{8}$	2.021	1008.
0 <sup>a</sup>	$279.98_{8}$	0.762	996.9	0.2483	279.98 <sub>8</sub>	2.513	1011.4
0	279.98 <sub>8</sub>	1.003	1000.1	0.2483	279.98 <sub>8</sub>	3.008	1013.
0	$279.98_{8}$	1.303	1002.6	$0.248_3^a$	289.98 <sub>8</sub>	1.152	963.
0	279.988	1.602	1004.3	0.2483	289.98 <sub>8</sub>	1.550	966.
0	279.98 <sub>8</sub>	1.901	1006.1	0.2483	289.98 <sub>8</sub>	2.015	970.2
0	279.988	2.202	1007.8	0.2483	289.98 <sub>8</sub>	2.511	974.
0	279.988	2.503	1009.7	0.2483	289.98 <sub>8</sub>	3.002	977.
0	279.98 <sub>8</sub>	3.007	1013.7	$0.248_3^a$	299.987	1.502	920.9
0 <sup>a</sup>	289.98 <sub>8</sub>	1.015	959.4	0.2483	299.987	1.997	926.
0 <sup>a</sup>	289.98 <sub>8</sub>	1.017	961.3	0.2483	299.987	2.505	931.
0	289.98 <sub>8</sub>	1.218	962.9	0.2483	299.987	3.004	935.0
0	289.98 <sub>8</sub>	1.501	965.2	$0.248_3^a$	309.98 <sub>7</sub>	1.925	873.8
0	289.98 <sub>8</sub>	1.804	967.7	0.2483	309.98 <sub>7</sub>	2.513	882.
0	289.98 <sub>8</sub>	2.203	970.8	$0.248_3$ $0.248_3$	309.987	3.006	889.0
0	289.98 <sub>8</sub>	2.603	973.7	$0.248_3^a$	319.98 <sub>7</sub>	2.433	818.
0	289.98 <sub>8</sub>	3.003	976.8	0.2483	319.987	2.995	831.
$0.502_4^a$	279.98 <sub>8</sub>	0.945	1005.9	0.7497	299.987	3.000	945.
$0.502_4$ $0.502_4$	279.98 <sub>8</sub>	1.507	1003.9	$0.749_7^a$	309.98 <sub>7</sub>	2.231	888.
$0.502_4$ $0.502_4$	279.98 <sub>8</sub>	2.017	1011.9	0.749 <sub>7</sub> 0.749 <sub>7</sub>	309.98 <sub>7</sub>	2.488	891.4
$0.502_4$ $0.502_4$	279.98 <sub>8</sub>	2.509	1011.5	$0.749_7$ $0.749_7$	309.98 <sub>7</sub>	3.004	897.3
$0.502_4$ $0.502_4$	279.988	3.015	1014.1	$0.7497^{a}$	319.98 <sub>7</sub>	2.830	833.
$0.502_4^{a}$	289.98 <sub>8</sub>	1.257	968.2	$0.7497^{a}$	319.987 319.987	2.830	833.
$0.502_4$ $0.502_4$	289.98 <sub>8</sub>	1.535	970.1	0.749 <sub>7</sub> 0.749 <sub>7</sub>	319.98 <sub>7</sub> 319.98 <sub>7</sub>	3.000	834.
	0		974.0	$1.0^{a}$	279.98 <sub>7</sub>	1.007	1030.0
$0.502_4 \\ 0.502_4$	$289.98_8$ $289.98_8$	$2.013 \\ 2.514$	974.0	1.0 <sup>a</sup>	$279.98_7$ $279.98_7$	1.007	1030.
	0						
0.5024	289.98 <sub>8</sub>	3.015	980.2	1.0	279.98 <sub>7</sub>	1.532	1033.
$0.502_4^a$	299.98 <sub>7</sub>	1.642	925.3	1.0	279.98 <sub>7</sub>	2.029	1035.
0.5024	299.987	2.029	929.0	1.0	279.987	2.523	1037.
0.5024	299.987	2.502	933.8	1.0	279.987	3.017	1039.
0.5024	299.987	3.005	937.8	1.0 <sup>a</sup>	289.987	1.347	993.
$0.502_4^a$	309.98 <sub>7</sub>	2.108	878.0	1.0 <sup>a</sup>	289.987	1.349	993.
$0.502_4$	309.98 <sub>7</sub>	2.530	883.6	1.0	289.98 <sub>7</sub>	1.515	994.
0.5024	309.98 <sub>7</sub>	3.004	889.7	1.0	289.987	2.005	996.
$0.502_4^a$	319.987	2.667	822.7	1.0	289.98 <sub>7</sub>	2.503	999.
0.5024	319.987	3.000	829.7	1.0	289.98 <sub>7</sub>	3.005	1002.2
0.749 <sub>7</sub> <sup>a</sup>	279.98 <sub>8</sub>	0.992	1015.2	1.0 <sup>a</sup>	299.98 <sub>6</sub>	1.773	952.0
0.7497	279.98 <sub>8</sub>	1.511	1017.1	1.0 <sup>a</sup>	$299.98_{6}$	1.772	951.9
0.7497	279.98 <sub>8</sub>	2.009	1020.6	1.0	$299.98_{6}$	2.015	953.
0.7497	279.98 <sub>8</sub>	2.504	1023.1	1.0	$299.98_{6}$	2.508	957.
0.7497	$279.98_{8}$	3.005	1025.3	1.0	$299.98_{6}$	3.018	960.
$0.749_7^a$	$289.98_8$	1.324	977.1	$1.0^{a}$	$309.98_{6}$	2.293	907.
$0.749_{7}$	$289.98_8$	2.003	982.0	1.0	$309.98_{6}$	2.519	909.5
$0.749_{7}$	$289.98_8$	2.504	984.8	1.0	$309.98_{6}$	3.016	914.
$0.749_{7}$	$289.98_8$	3.004	988.0	1.0 <sup>a</sup>	$319.98_4$	2.918	856.5
$0.749_7^a$	299.987	1.734	935.0	1.0	$319.98_4$	3.007	857.5
$0.749_{7}$	$299.98_{7}$	2.514	941.5				

 Table 2. Liquid Phase PVTx Properties of R-32 + R-143a

<sup>*a*</sup> Experimentally determined bubble-point.

temperatures and critical densities of single-component refrigerants are listed in Table 5; the numerical values for the parameters *a*, *b*, *c*, *A*, *B*, and *C* for pure substances are given in Table 7.

$$\nu_{\rm Cm} = \sum x_i \nu_{\rm Ci} + \sum \sum x_i x_j \Delta \nu_{ij} (\nu_{\rm Ci}^{1/3} + \nu_{\rm Cj}^{1/3}) \nu_{\rm Ci}^{1/3} \nu_{\rm Cj}^{1/3}$$
(2)

$$\nu_{\rm Cm} T_{\rm Cm} = \sum \sum x_i x_j \Delta T_{ij} \sqrt{\nu_{\rm Ci} T_{\rm Ci} \nu_{\rm Cj} T_{\rm Cj}}$$
(3)  
$$L_{\rm m} = \sum L_i x_i + \sum \sum x_i x_j \Delta L_{ij}$$
(4)

 $\nu$  with any subscript throughout the present paper denotes the molar volume (in cm<sup>3</sup>·mol<sup>-1</sup>). The subscript i

Table 3. Bubble-Point Pressures and Saturated-Liquid Densities of  $R\mbox{-}32\mbox{ + }R\mbox{-}125\mbox{ + }R\mbox{-}143a$ 

$x_{ m R-32}/$ mol·mol <sup>-1</sup>	10 02 10 120		<i>P</i> /MPa	$ ho/{ m kg}{\cdot}{ m m}^{-3}$
$0.335_{6}$	0.3328	$279.98_{8}$	0.900	1121.7
$0.335_{6}$	$0.332_{8}$	$289.98_8$	1.201	1078.8
$0.335_{6}$	$0.332_{8}$	$299.98_{7}$	1.574	1030.1
$0.335_{6}$	$0.332_{8}$	$309.98_{7}$	2.023	975.9
$0.335_{6}$	0.3328	$319.98_{7}$	2.563	911.2

Table 4. Liquid-Phase PVTx Properties of R-32 + R-125 +R-143a

$x_{R-32}/$ mol·mol <sup>-1</sup>	$x_{ m R-125}/$ mol·mol <sup>-1</sup>	<i>T</i> /K	<i>P</i> /MPa	$ ho/{ m kg}{ m \cdot}{ m m}^{-3}$
$0.335_6{}^a$	0.3328	$279.98_{8}$	0.900	1121.7
$0.335_{6}$	$0.332_{8}$	$279.98_{8}$	1.506	1125.9
$0.335_{6}$	$0.332_{8}$	$279.98_{8}$	2.014	1129.2
$0.335_{6}$	$0.332_{8}$	$279.98_{8}$	2.507	1132.4
$0.335_{6}$	$0.332_{8}$	$279.98_{8}$	3.015	1135.3
$0.335_6^a$	$0.332_{8}$	$289.98_8$	1.201	1078.8
$0.335_{6}$	$0.332_{8}$	$289.98_8$	1.511	1081.5
$0.335_{6}$	$0.332_{8}$	$289.98_8$	2.012	1086.0
$0.335_{6}$	$0.332_{8}$	$289.98_8$	2.502	1090.1
$0.335_{6}$	$0.332_{8}$	$289.98_8$	3.011	1094.3
$0.335_6^a$	$0.332_{8}$	$299.98_7$	1.574	1030.1
$0.335_{6}$	$0.332_{8}$	$299.98_7$	2.002	1035.3
$0.335_{6}$	$0.332_{8}$	$299.98_{7}$	2.498	1040.7
$0.335_{6}$	$0.332_{8}$	$299.98_{7}$	3.005	1046.0
$0.335_6^a$	$0.332_{8}$	$309.98_{7}$	2.023	975.9
$0.335_{6}$	$0.332_{8}$	$309.98_7$	2.498	983.6
$0.335_{6}$	$0.332_{8}$	$309.98_7$	3.000	991.3
$0.335_6^a$	0.3328	$319.98_7$	2.563	911.2
$0.335_{6}$	0.3328	319.987	2.997	922.8

<sup>a</sup> Experimentally determined bubble-point.

Table 5. Parameters Used in the PR Equation

properties	R-32	R-125	R-143a
T <sub>C</sub> /K	$351.26^a \\ 5.780^d$	$339.17^b$	345.88 <sup>c</sup>
P <sub>C</sub> /MPa		$3.620^b$	3.764 <sup>c</sup>
$ ho_{\rm C}/{ m kg}\cdot{ m m}^{-3}$	427 <sup>a</sup>	577 <sup>b</sup>	431 <sup>c</sup>
$M/{ m g}\cdot{ m mol}^{-1}$	52.024	120.02	84.04
ω	0.277	0.306	0.263

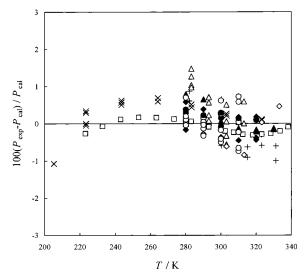
 $^a$  Higashi (1996).  $^b$  Higashi (1994).  $^c$  Higashi and Ikeda (1995).  $^d$  Qian et al. (1993).

 Table 6. Binary Interaction Parameters for the Binary
 Systems

binary system	$k_{\rm ij}$ (T in K)
R-32 + R-125 R-32 + R-143a R-125 + R-143a	$egin{array}{c} -0.0466+0.000169T\ 0.0116\ -0.0126 \end{array}$

or *j* denotes a property of single components, while the subscript *ij* denotes a binary mixture.  $\Delta v_{ij}$  and  $\Delta T_{ij}$  in eqs 2 and 3 are cross-parameters for the critical molar volume and critical temperature of the mixtures, respectively.  $\Delta L_{ij}$  in eq 4 represents the cross-parameter for  $a_{\rm m}$ ,  $b_{\rm m}$ ,  $c_{\rm m}$ ,  $A_{\rm m}$ ,  $B_{\rm m}$ , and  $C_{\rm m}$  in eq 1. Any cross-parameter  $l_{ij}$  in the present paper fulfills the conditions  $l_{ij} = l_{ji}$  and  $l_{ij,i=j} = 0$ .

Higashi (1996), Nagel and Bier (1995, 1996), and Kishizawa et al. (1999a,b) reported the critical temperatures and densities for the binary R-32 + R-125 and R-125 + R-143a systems. Mizukawa et al. (1999) of our laboratory have recently completed their measurements on the critical temperatures and densities of the binary R-32 + R-143a system. Those reported and/or referred data were adopted to develop eq 2 for the respective properties. Equations 2 and 3 are modifications of those used by Thomson et al. (1982), by introducing cross-parameters,  $\Delta v_{ij}$  and  $\Delta T_{ij}$ , that play a role to indicate whether the properties deviate to



**Figure 2.** Bubble-point pressure deviation from the PR equation: (◆) this work (R-32 + R-143a); (●) this work (R-32 + R-125 + R-143a); (▲) Fujimine et al. (1999) (R-125 + R-143a); (□) Kleemiss (1997) (R-125 + R-143a); (×) Nagel and Bier (1996) (R-125 + R-143a); (+) Zhelezny et al. (1995) (R-125 + R-143a); (○) Widiatmo et al. (1993) (R-32 + R-125); (△) Higashi (1993) (R-32 + R-125); (◇) Kleemiss (1997) (R-32 + R-125).

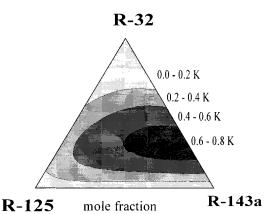


Figure 3. Temperature glide of mixtures blended with R-32, R-125, and R-143a at 500 kPa.

positive or negative against the averaged values. In relation to the critical temperature and density of the ternary R-32 + R-125 + R-143a system, there are no reported data available at present. They are calculated in the present work from eqs 2 and 3 by introducing the cross-parameters of its constituent binary systems. The cross-parameters for the critical properties of refrigerant mixtures are summarized in Table 8.

Equation 4 is used by replacing *L* with  $a_m$ ,  $b_m$ ,  $c_m$ ,  $A_m$ ,  $B_m$ , and  $C_m$  in eq 1. The subscripts 1, 2, and 3 in eq 2 represent R-32, R-125, and R-143a, respectively. The basic data sets used to optimize  $a_m$ ,  $b_m$ ,  $c_m$ ,  $A_m$ ,  $B_m$ , and  $C_m$  include the present work and those by Widiatmo et al. (1993, 1995) for the binary R-32 + R-125 and R-125 + R-143a systems and by Fujimine et al. (1999) for the binary R-125 + R-143a system. The cross-parameters used in eq 4 are given in Table 9.

Figure 4 illustrates the deviation of measured saturatedliquid densities from eq 1. Also shown in Figure 4 are our saturated liquid density data of the binary R-32 + R-125and R-125 + R-143a systems reported earlier (Widiatmo et al., 1993, 1995) and the data by Ikeda et al. (1995) in a reduced density range similar to that for the present study. Ikeda and co-workers reported saturated-liquid density

Table 7. Parameters Used in	in Eq 1
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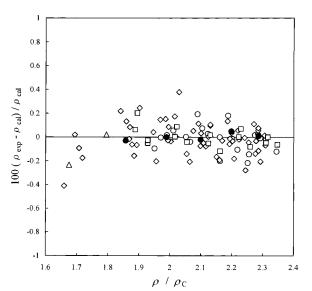
		-		
coefficient	R-32	R-125	R-143a 1.046 23 1.699 42	
$A_i$	0.991 467	2.228 14		
$B_i$	1.805 72	0.451 997		
$C_i$	2.343 30	0	0	
$a_i$	0.280	0.370	0.257	
$b_i$	$b_i = 0.530 \ c_i = 4.450$		0.599	
Ci			0	

Table 8. Cross-Parameters in Eqs 2, 3, 9, and 10

		-		
binary system	ry system $2\Delta v_{ij}$		$2\Delta P_{ij}$	$2\Delta\omega_{ij}$
R-32 + R-125	0	2.0446	-0.7290	0.042
R-32 + R-143a	-0.0228	1.9769	-0.3777	0.037
R-125 + R-143a	-0.0055	1.9984	0.0050	0.009

Table 9. Cross-Parameters in Eq 4

binary system	$2\Delta a_{ij}$	$2\Delta b_{ij}$	$2\Delta c_{ij}$	$2\Delta A_{ij}$	$2\Delta B_{ij}$	$2\Delta C_{ij}$
R-32 + R-125	0.030	-0.283	18.472	0.961	-0.972	-2.450
R-32 + R-143a	0.204	-0.019	11.420	1.135	-1.423	-2.266
R-125 + R-143a	0.027	-0.303	0	0.587	-0.770	0



**Figure 4.** Saturated-liquid density deviation from eq 1: ( $\Box$ ) this work (R-32 + R-143a); ( $\bullet$ ) this work (R-32 + R-125 + R-143a); ( $\bigcirc$ ) Widiatmo et al. (1993) (R-32 + R-125); ( $\diamondsuit$ ) Fujimine et al. (1999) (R-125 + R-143a); ( $\triangle$ ) Ikeda and Higashi (1995) (R-125 + R-143a).

data near the critical temperature of the binary R-125 + R-143a system. Their measurements were performed to obtain direct measurements of the critical temperature and density. The present saturated-liquid densities of the binary R-32 + R-143a system are underestimated by eq 1, as indicated in Figure 4. This is due to the procedure of determining the critical temperature, pressure, density, and all coefficients on the basis of the measured compressed-liquid densities instead of the saturated ones. However, the deviation falls within the claimed uncertainty of the density measurement,  $\pm 0.2\%$ .

Figure 4 shows that eq 1 excellently represents most of the saturated-liquid density data within  $\pm 0.3\%$  except those very near the critical point (0.989 <  $T_{\rm r}$  < 1) by Ikeda et al. (1995). The deviations of the data by Ikeda et al. from eq 1 range from +0.66 to -1.5%. The measured saturated-liquid densities of the ternary R-32 + R-125 + R-143a system are also satisfactorily represented. It should be noted that the representation for the ternary mixture was developed from only the cross-parameters of the corresponding binary mixtures in eqs 1–4.

The measured compressed-liquid densities of the refrigerant mixtures in the present study are represented by eq 5, which is a modification of the Sato equation (1981). Equation 5 has also been applied successfully in our previous work concerning the binary and ternary mixtures of R-32, R-125, and R-134a (Widiatmo et al., 1997).

$$\ln\frac{\rho_{\rm s}}{\rho} = E\ln\left(\frac{P+D}{P_{\rm s}+D}\right) \tag{5}$$

In eq 5,  $\rho_s$  and  $P_s$  are the saturated-liquid density (in kg·m<sup>-3</sup>) and the bubble-point pressure (in MPa), respectively.  $\rho_s$  is obtained from eq 1, while  $P_s$  is from eq 6.

$$\ln \frac{P_{\rm s}}{P_{\rm Cm}} = 5.92715 - \frac{6.09649}{T_{\rm r}} - 1.28863 \ln T_{\rm r} + 0.169347 T_{\rm r}^6 + \omega_{\rm m} \Big( 13.4902 - \frac{13.8756}{T_{\rm r}} - 11.1704 \ln T_{\rm r} + 0.385433 T_{\rm r}^6 \Big)$$
(6)

D and E are given in eqs 7 and 8

$$\frac{D}{P_{\rm Cm}} + 1 = 13.406\tau + 66.7486\tau^2 - 97.1726\tau^3 + \tau^4 \exp(-22.4997 + 226.792\omega_{\rm m} - 451.956\omega_{\rm m}^{-2})$$
(7)  
$$E = -0.147397 + 0.192097\omega_{\rm m}$$
(8)

 $T_{\rm r}$  in eq 6 is  $T/T_{\rm Cm}$ , and  $\tau$  is  $1 - T_{\rm r}$ ;  $T_{\rm Cm}$  is calculated from eqs 2 and 3. Equation 6 is the presently adopted bubblepoint pressure correlation that is different from that proposed in our previous paper (Widiatmo et al., 1997). This improvement became feasible thanks to the existence of the reported critical pressure data by Nagel and Bier (1995, 1996) for the binary R-32 + R-125 and R-125 + R-143a systems. The mixing rule for the critical pressure (in MPa)  $P_{\rm Cm}$  is given by eq 9, whereas that for the acentric factor  $\omega_{\rm m}$  is given by eq 10.

$$P_{\rm Cm} = \sum x_i P_{\rm Ci} + \sum \sum x_i x_j \Delta P_{ij} \tag{9}$$

$$\omega_{\rm m} = \sum x_i \omega_i + \sum \sum x_i x_j \Delta \omega_{ij} \tag{10}$$

The cross-parameters  $\Delta P_{ij}$  in eq 9 for the binary R-32 + R-125 and R-125 + R-143a systems were optimized on the basis of the critical pressure data by Nagel and Bier (1995, 1996). Since, for the binary R-32 + R-143a system, neither critical temperature nor critical pressure data are available, the critical pressure of this binary mixture is derived from the bubble-point pressure data with the aid of eq 6. The optimized values of  $\Delta P_{ij}$  for all the mixtures of present interest are tabulated in Table 8. Similar to that of the binary R-32 + R-143a system, there is no report of a critical pressure measurement of the ternary R-32 + R-125 + R-143a system. Therefore, in the present work, the critical pressure of the ternary system is calculated from eq 9 by introducing the cross-parameter  $\Delta P_{ij}$  of its constituent binary mixtures.

Equation 6 takes a functional form similar to that used by Lee and Kesler (1975), introducing the mixing rule for the acentric factor  $\omega_m$  given in eq 10. The coefficients in eq 6 have been optimized so as to fulfill the criteria given below for a vapor pressure correlation:

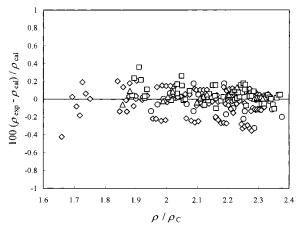


Figure 5. Liquid density deviation from eq 5: (□) this work (R-32 + R-143a); ( $\triangle$ ) this work (R-32 + R-125 + R-143a); ( $\bigcirc$ ) Widiatmo et al. (1993) (R-32 + R-125); (◊) Fujimine et al. (1999) (R-125 + R-143a).

(1) at  $T = T_{\rm Cm}$ ,  $P_{\rm s} = P_{\rm Cm}$ (2) at  $T = T_{\rm Cm}$ ,  $d\alpha/dT_{\rm r} = 0$ , where  $\alpha$  is the Riedel factor defined as  $\alpha = T_{\rm r} d \ln P_{\rm r}/dT_{\rm r}$ 

(3)  $\omega = -\log(P_{\rm s}/P_{\rm Cm})_{T_{\rm r}=0.7} - 1$ 

The basic input data are the present bubble-point pressure data in addition to those by Widiatmo et al. (1993, 1995, 1997) to cover the binary and ternary mixtures of R-32, R-125, and R-134a and those by Fujimine et al. (1999). Using eq 6, the present bubble-point pressures are reproduced within  $\pm 2\%$ .

Equations 7 and 8 are also revised versions from our previous representations (Widiatmo et al., 1997), by applying the functional forms similar to those used by Thomson et al. (1982). The coefficients of these equations are determined on the basis of the present compressedliquid densities complimented with those of our previous work (Widiatmo et al., 1997).

By incorporating eqs 6-8 with eq 5, the result leads to a compressed-liquid density representation within  $\pm 0.3\%$ , as shown in Figure 5 except a single datum by Fujimine et al. (1999) for the binary R-125 + R-143a system. Figure 5 shows not only the deviation of the compressed-liquid densities but also that of the saturated-liquid densities of the binary and ternary mixtures of R-32, R-125, and R-143a. It is important to note that eq 5 is developed mostly on the basis of compressed-liquid densities at pressures below 4 MPa. Since engineering applications of those refrigerants in air conditioning and refrigeration are limited to that pressure range, it seems worthwhile to claim the importance of eq 5 in engineering applications. Since eq 5 is also developed on the basis of the compressed-liquid densities of the mixtures blended with R-32, R-125, and/ or R-134a, it also works effectively to reproduce the corresponding experimental data for the binary and ternary mixtures of R-32, R-125, and R-134a.

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