

# Isochoric $PVT_x$ Measurements for the $C_2H_6 + N_2O$ Binary System

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In the search for fluids potentially suitable for low-temperature refrigeration applications, we have turned our attention to systems consisting of natural fluids, nitrous oxide ( $N_2O$ ) and ethane ( $C_2H_6$ ) in particular. Isochoric  $PVT_x$  measurements are reported for nitrous oxide + ethane for both the two-phase and the superheated vapor regions at temperatures ranging from (216 to 358) K and pressures from (494 to 5261) kPa along 11 isochores. The data obtained in the two-phase region were used to derive VLE parameters using a flash method with the Carnahan–Starling–De Santis equation of state (CSD EOS). The dew point was also found for each isochore from the intersection of the  $P$ – $T$  sequences. The dew points were then used to derive VLE parameters from CSD EOS. Results from the superheated region were compared with the predicted from the virial EOS and CSD EOS.

## Introduction

One of the chemicals potentially suitable for use in low-temperature refrigeration applications is nitrous oxide, whose critical parameters and saturated pressures are very similar to those of carbon dioxide. The main advantage of nitrous oxide over carbon dioxide lies in its very low melting temperature, whereas its global warming potential (GWP, which is 310) is low, but not low enough to be considered environmentally friendly.

The thermophysical properties of ethane are well-known. Although the ethane + nitrous oxide system has been used as a system for low-temperature refrigeration applications, its thermophysical properties have yet to be thoroughly elucidated. Isochoric measurements were consequently taken, covering temperatures from (216 to 358) K to make up for the lack of VLE and  $PVT_x$  data. A summary of the properties of nitrous oxide, carbon dioxide, and ethane is presented in Table 1. VLE parameters were derived from data in the two-phase region, applying the Carnahan–Starling–De Santis equation of state (CSD EOS),<sup>1</sup> while data obtained from the superheated region were compared with the predictions obtained with the virial equation of state.

## Experimental Section

**Chemicals.** Ethane and nitrous oxide were supplied by Sol SpA. Their purity was checked by gas chromatography, using a thermal conductivity detector, and found to be 99.99 % for both fluids, basing all estimations on an area response.

**Apparatus.** The new experimental setup is illustrated in Figure 1. The basic experimental setup has already been described elsewhere,<sup>2</sup> so it is only briefly outlined here. The main changes made to the original apparatus<sup>3,4</sup> concerned the twin thermostatic baths (7) filled with different silicone oils (Baysilone M10 and Baysilone M100, Bayer). After charging with the sample mixture, the setup could be operated over two temperature

**Table 1. Global Warming Potential (GWP), Triple Point Temperature ( $T_t$ ), Normal Boiling Temperature ( $T_b$ ), Critical Temperature ( $T_c$ ), and Critical Pressure ( $P_c$ ) of the Investigated Fluids**

	GWP	$T_t$ /K	$T_b$ /K	critical point	
				$T_c$ /K	$P_c$ /kPa
R170	20	90.35	184.55	305.33	4871
R744A	310	182.33	184.67	309.57	7245
R744	1	216.59	194.75 <sup>a</sup>	304.13	7377

<sup>a</sup> Sublimation temperature at atmospheric pressure.

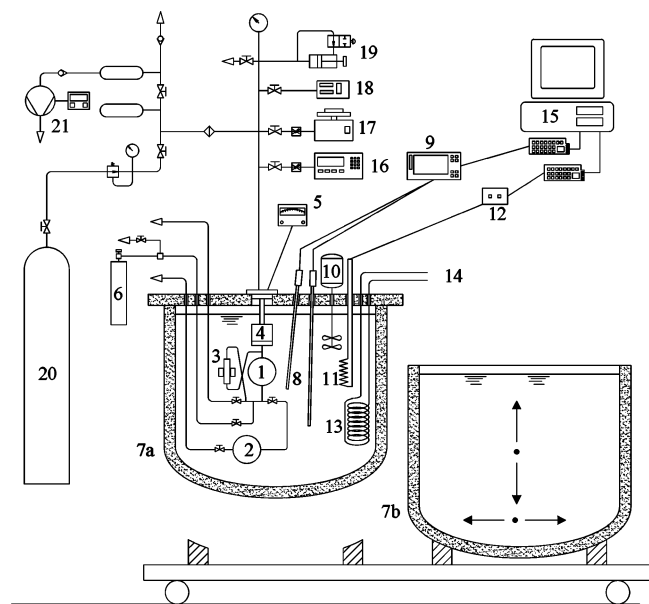
**Table 2. Measurements at Bulk Compositions ( $z_1$ ) for the  $C_2H_6$  (1) +  $N_2O$  (2) System over the Temperature Range ( $\Delta T$ ) and Pressure Range ( $\Delta P$ ) with Resultant Dew Temperatures ( $T_d$ ) and Dew Pressures ( $P_d$ )<sup>a</sup>**

$z_1$	$\Delta T$ /K	$\Delta P$ /kPa	$N$ /mol	$N$ exp. points			$T_d$ /K	$P_d$ /kPa
				2ph	Vap	tot		
0.125	216–353	541–3940	0.42784	11	13	24	263.50	2524.3
0.208	218–353	603–3808	0.41356	6	11	17	261.00	2411.4
0.234	218–358	613–4525	0.49524	11	11	22	267.60	2834.5
0.345	217–353	599–3415	0.36583	9	12	21	256.32	2132.1
0.361	217–353	610–4063	0.45046	10	12	22	263.12	2551.5
0.483	220–358	666–3592	0.38211	9	13	22	259.09	2208.0
0.521	216–323	566–5261	0.78153	15	8	23	282.33	3893.2
0.628	217–353	567–3352	0.36250	9	13	22	259.48	2105.9
0.707	224–358	722–4093	0.45871	10	11	21	268.32	2547.7
0.749	218–353	572–3560	0.40631	10	12	22	264.65	2258.9
0.876	217–353	494–3457	0.38433	10	12	22	267.15	2218.0

<sup>a</sup> 2ph and Vap denote data within the VLE boundary and superheated region, respectively.

ranges, approximately from (210 to 290) K and from (290 to 360) K, depending on which bath was used. The two silicone oils have different kinematic viscosity values (10 and 100 cSt at room temperature, respectively). The one with lower kinematic viscosity, due to its higher volatility, was applied only for the low-temperature range, while that with a greater viscosity was applied only at high temperatures. The thermostatic baths were easy to move thanks to the new system configuration. The spherical cells and pressure transducer are immersed in one of

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**Figure 1.** Experimental apparatus: 1, constant volume spherical cell; 2, auxiliary cell; 3, magnetic pump; 4, differential pressure transducer; 5, electronic null indicator; 6, charging system; 7, thermostatic baths; 8, platinum thermoresistances; 9, thermometric bridge; 10, stirrer; 11, heater; 12, power system; 13, cooling coil; 14, connections to auxiliary thermostatic bath; 15, acquisition system; 16, Bourdon gage; 17, dead weight gage; 18, vibrating cylinder; 19, pressure gage; 20, precision pressure controller; 21, nitrogen reservoir; 22, vacuum pump system.

the two thermostatic baths (7). An auxiliary thermostat (14) was used to reach below-ambient temperatures. The cell volume was estimated (as explained elsewhere<sup>4</sup>) to be  $(273.5 \pm 0.3) \text{ cm}^3$  at room temperature.

**Table 3.** Experimental Data within the VLE Boundary for the  $\text{C}_2\text{H}_6$  (1) +  $\text{N}_2\text{O}$  (2) System

$T/\text{K}$	$P/\text{kPa}$	$V/\text{dm}^3 \cdot \text{mol}^{-1}$	$T/\text{K}$	$P/\text{kPa}$	$V/\text{dm}^3 \cdot \text{mol}^{-1}$	$T/\text{K}$	$P/\text{kPa}$	$V/\text{dm}^3 \cdot \text{mol}^{-1}$	$T/\text{K}$	$P/\text{kPa}$	$V/\text{dm}^3 \cdot \text{mol}^{-1}$
$z_1 = 0.125$						$z_1 = 0.521$					
215.98	541.2	0.637	242.91	1397.9	0.638	215.64	565.9	0.349	252.92	1891.3	0.350
218.13	589.3	0.637	247.89	1628.3	0.638	218.13	621.7	0.349	258.00	2167.0	0.350
223.09	711.5	0.638	252.90	1878.4	0.638	223.08	744.6	0.349	263.05	2469.2	0.350
228.05	851.7	0.638	257.90	2174.1	0.639	228.02	884.4	0.349	268.09	2799.1	0.350
232.95	1010.1	0.638	262.89	2477.6	0.639	232.96	1042.4	0.349	273.07	3154.2	0.350
237.92	1191.9	0.638				237.91	1218.7	0.349	277.87	3528.8	0.350
$z_1 = 0.208$						$z_1 = 0.628$					
217.73	603.1	0.659	242.92	1434.8	0.660	242.87	1418.5	0.349	282.83	3899.4	0.350
223.08	736.4	0.660	252.88	1928.5	0.660	247.85	1640.6	0.350			
232.94	1040.5	0.660	257.98	2221.3	0.661	216.52	567.0	0.752	243.05	1356.1	0.753
$z_1 = 0.234$						$z_1 = 0.753$					
218.03	613.2	0.551	247.98	1678.7	0.551	223.10	719.4	0.752	248.09	1564.1	0.753
223.09	741.3	0.551	252.96	1939.2	0.551	228.11	852.8	0.753	253.11	1792.4	0.753
228.04	885.6	0.551	258.07	2236.1	0.552	233.08	1012.4	0.753	258.09	2036.2	0.754
232.98	1047.6	0.551	263.05	2555.7	0.552	238.07	1170.1	0.753			
237.98	1233.9	0.551	267.98	2834.7	0.552	223.81	721.9	0.595	248.05	1516.1	0.595
243.00	1444.2	0.551				228.12	836.5	0.595	253.04	1736.1	0.595
$z_1 = 0.345$						$z_1 = 0.707$					
216.90	598.8	0.745	243.11	1460.4	0.746	233.03	973.7	0.595	257.98	1975.4	0.596
223.09	753.8	0.746	248.09	1692.8	0.746	238.07	1136.3	0.595	263.06	2242.6	0.596
228.04	897.4	0.746	253.08	1950.5	0.747	243.04	1315.7	0.595	268.05	2511.2	0.596
233.04	1062.7	0.746	258.10	2152.4	0.747	218.14	571.8	0.671	243.11	1280.7	0.672
238.07	1249.6	0.746				223.10	682.0	0.671	249.01	1511.0	0.672
$z_1 = 0.31$						$z_1 = 0.749$					
217.35	609.5	0.605	243.13	1461.2	0.606	228.08	806.2	0.671	253.10	1687.1	0.672
223.13	755.1	0.606	248.12	1693.3	0.606	233.14	948.5	0.672	258.08	1920.4	0.672
228.12	899.9	0.606	253.11	1951.6	0.606	238.12	1105.7	0.672	263.09	2171.5	0.672
233.14	1065.7	0.606	258.10	2236.2	0.606	216.75	494.4	0.709	243.10	1173.7	0.710
238.13	1251.6	0.606	263.09	2534.2	0.607	223.08	621.6	0.710	248.09	1353.0	0.710
$z_1 = 0.483$						$z_1 = 0.876$					
219.89	665.7	0.714	243.09	1433.3	0.714	228.05	736.1	0.710	253.11	1553.3	0.711
223.12	748.6	0.714	248.07	1653.5	0.715	233.03	864.6	0.710	258.10	1771.3	0.711
228.08	890.4	0.714	253.09	1899.6	0.715	238.10	1011.6	0.710	263.14	2012.6	0.711
233.04	1048.7	0.714	258.09	2146.0	0.715						
238.04	1228.5	0.714									

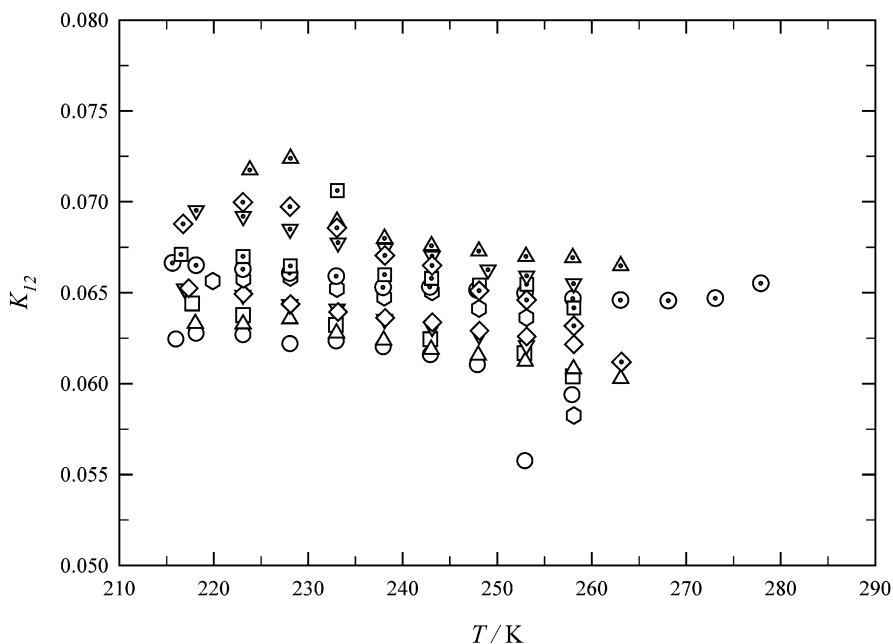
The pressure and temperature data acquisition systems were identical to those of the previous apparatus.<sup>3,4</sup> A PID device was used to control the temperature, which was measured using a calibrated resistance thermometer; the total uncertainty of the temperature measurements was  $\pm 0.02 \text{ K}$ . The uncertainty in the pressure measurements stems from the uncertainty of the transducer and null indicator system and the pressure gauges. The uncertainty of the digital pressure indicator (Ruska, model 7000) is  $\pm 0.003 \%$  of its full scale. Temperature fluctuations due to bath instability can also affect the total uncertainty in the pressure measurement, which was nonetheless found to be less than  $\pm 1 \text{ kPa}$ .

**Experimental Procedure.** Mixtures were prepared using the gravimetric method. First of all, the pure samples were charged in different bottles, degassed to remove noncondensable gases and air, and then weighed with an analytical balance (uncertainty  $\pm 0.3 \text{ mg}$ ). After evacuating the cell, the bottles were then emptied into the cell immersed in the bath. Then the bottles were weighed again, and the mass of the charge was calculated from the difference between the two masses. The mass of material in the dead volumes was estimated and subtracted from the total mass of the charge. The uncertainty in mixture preparation was estimated to be constantly lower than 0.001 in mole fraction. After reaching the experimental temperature, the mixing pump was activated for about 15 min, and next, the mixture was allowed to stabilize for about 20 min before the data recording.

## Results and Discussion

The temperature and pressure ranges are shown in Table 2, along with the mixture's composition and the number of moles charged. On the basis of the analysis of the slope of each  $T-P$





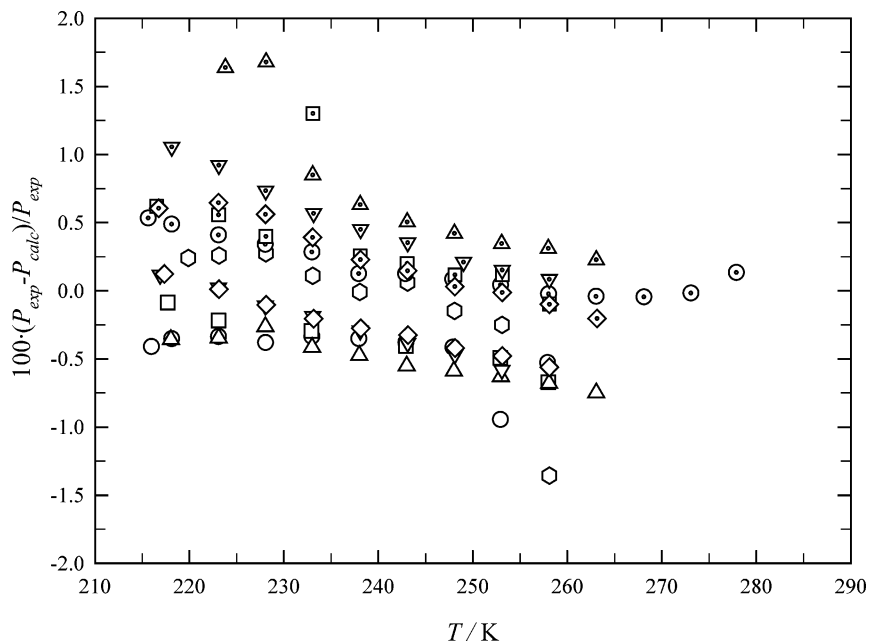
**Figure 2.**  $K_{12}$  values found by the flash method for the  $\text{C}_2\text{H}_6$  (1) +  $\text{N}_2\text{O}$  (2) system.

○ $z_j=0.1248$	▽ $z_j=0.3453$	⊙ $z_j=0.5208$	▽ $z_j=0.7490$
□ $z_j=0.2078$	◇ $z_j=0.3608$	⊠ $z_j=0.6277$	◇ $z_j=0.8762$
△ $z_j=0.2339$	⬡ $z_j=0.4827$	△ $z_j=0.7066$	

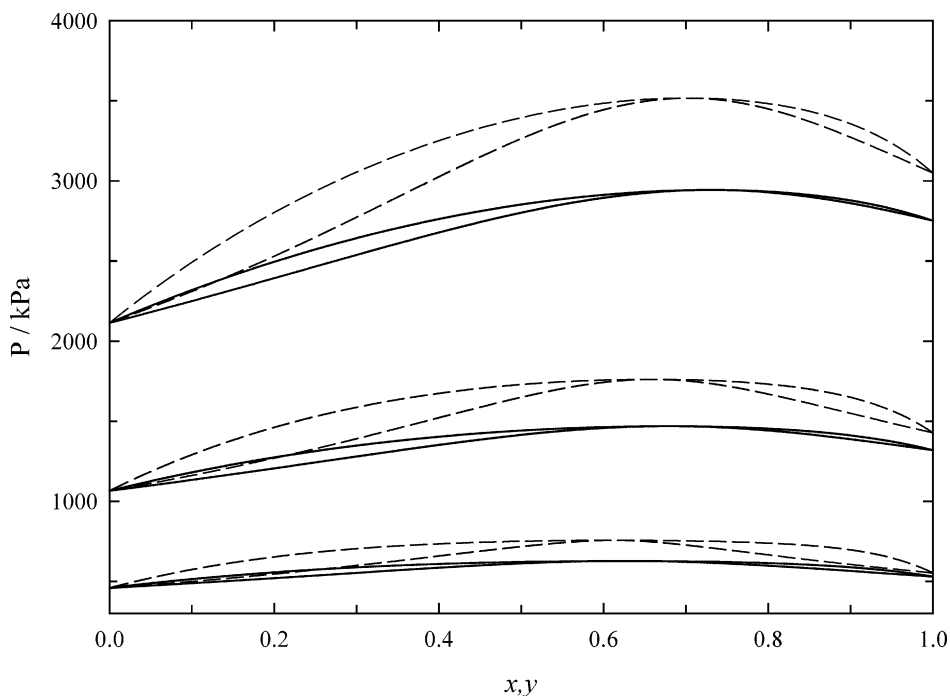
isochoric data, we also need the volumetric properties of both phases, which we calculated from the CSD EOS.  $T$ ,  $P$ ,  $z_i$ , and  $n$  (number of moles charged) were kept constant for each experimental point during the fitting procedure. The isochoric cell volume was known from the gravimetric calibration, so the binary interaction parameter ( $K_{12}$ ) and the composition at the bubble and dew points were found, considering them as dependent variables. Figure 2 shows the  $K_{12}$  values found for each data point in the two-phase region, while Figure 3 shows the scatter diagram of the relative pressure deviations. The

behavior of the binary interaction parameters was almost temperature-independent, and the pressure deviations amounted to around  $\pm 1\%$ .

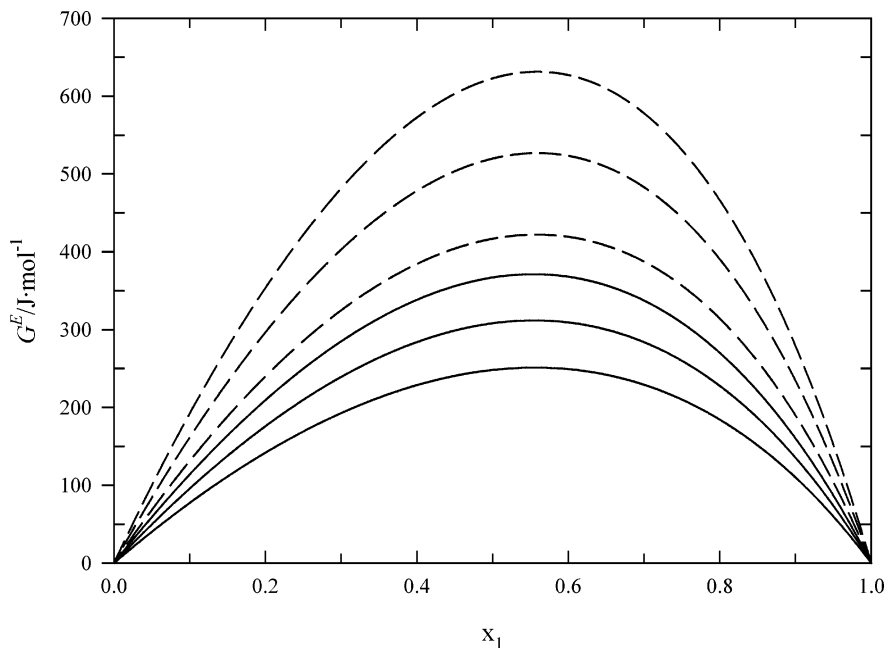
Using the  $K_{12}$  values averaged from our measurements in the two-phase region, we calculated the VLE at three different temperatures. The results are given in Figure 4. Since the properties of  $\text{N}_2\text{O}$  and  $\text{CO}_2$  are similar in terms of critical temperature, critical pressure, saturated pressure, and molecular mass, one of our goals was to compare the nitrous oxide and carbon dioxide binary systems. The VLE behavior of  $\text{C}_2\text{H}_6$  (1)



**Figure 3.** Pressure deviations between experimental values and those calculated with the  $K_{12}$  coefficients for the  $\text{C}_2\text{H}_6$  (1) +  $\text{N}_2\text{O}$  (2) system. Symbols as denoted in Figure 2.



**Figure 4.** VLE representation from the CSD EOS for the  $\text{C}_2\text{H}_6$  (1) +  $\text{N}_2\text{O}$  (2) system (solid lines) and for the  $\text{C}_2\text{H}_6$  (1) +  $\text{CO}_2$  (2) system (dashed lines) at three temperatures:  $T = 218.15$  K (lowest lines),  $T = 243.15$  K (middle lines), and  $T = 268.15$  K (upper lines).



**Figure 5.** Excess Gibbs energy for the  $\text{C}_2\text{H}_6$  (1) +  $\text{N}_2\text{O}$  (2) system (solid lines) and for the  $\text{C}_2\text{H}_6$  (1) +  $\text{CO}_2$  (2) system (dashed lines):  $T = 218.15$  K (upper lines),  $T = 243.15$  K (middle lines), and  $T = 268.15$  K (lower lines).

+  $\text{CO}_2$  (2) calculated with the binary interaction parameter ( $K_{12} = 0.1205$ ) derived from the literature data<sup>6-11</sup> is included in Figure 4 for the comparison.

The  $\text{C}_2\text{H}_6$  (1) +  $\text{N}_2\text{O}$  (2) system reveals strong positive deviations from Raoult's law, forming a positive azeotrope with an estimated composition around  $x_1 = 0.63, 0.68,$  and  $0.73$  at  $T = (218, 243,$  and  $268)$  K, respectively. The  $\text{C}_2\text{H}_6$  +  $\text{CO}_2$  system's behavior is qualitatively similar but with distinctly greater deviations from Raoult's law.

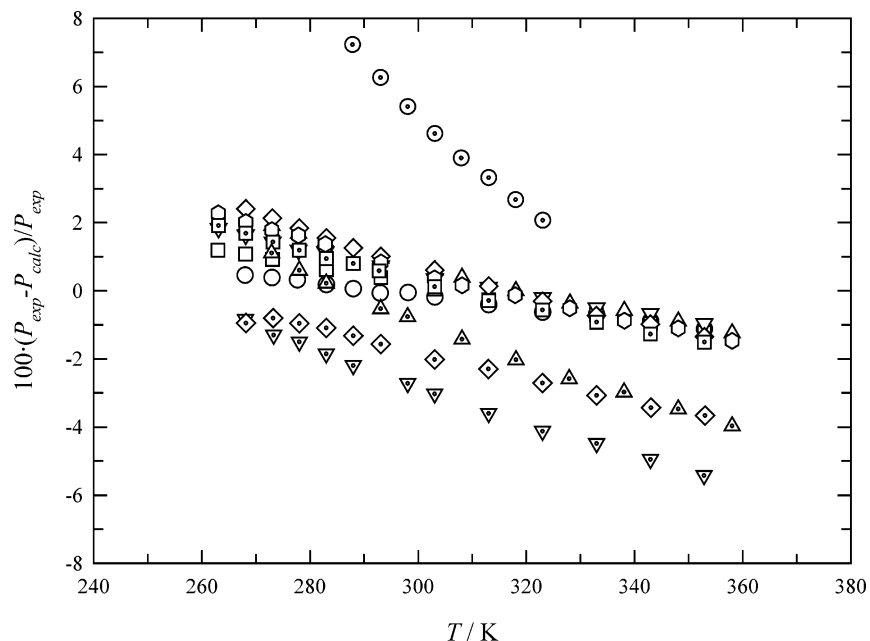
The Gibbs excess energy ( $G^E$ ) calculated using the equation:

$$G^E = RT(\ln \phi - \sum x_i \phi_i) \quad (1)$$

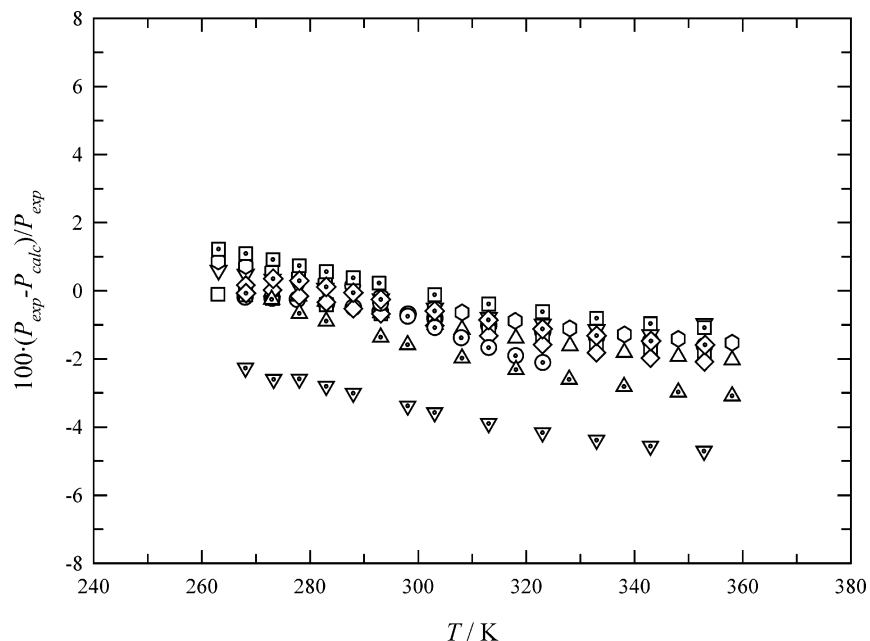
where  $\phi$  and  $\phi_i$  stand for the fugacity coefficients of the mixtures

and pure components, respectively, at the same temperature ( $T$ ) and pressure ( $P$ ), and  $R$  is the gas constant showed a maximum at around  $x_1 = 0.55$ . The trend of the  $G^E(x)$  at the three different temperatures is shown in Figure 5, again for both systems, where the latter's greater deviation from ideality is evident.

**PVTx.** Since there are no published data on the superheated vapor region for the binary systems considered, we compared our experimental  $PVTx$  findings with data calculated using the virial EOS in the Leiden form. A temperature dependence of the third-degree polynomial was found for the second and third virial coefficients. Then the pressure values were calculated from the virial equation of state in the Leiden form, using the isochoric experimental temperatures and molar volumes. The



**Figure 6.** Relative pressure deviations for the  $C_2H_6$  (1) +  $N_2O$  (2) system between experimental values from isochoric and data calculated with the virial EOS. Symbols as denoted in Figure 2.



**Figure 7.** Relative pressure deviations for the  $C_2H_6$  (1) +  $N_2O$  (2) system between experimental values from isochoric and data calculated with the CSD EOS. Symbols as denoted in Figure 2.

second and third virial coefficients for the two pure fluids ( $B_{ii}$ ,  $B_{jj}$ ,  $C_{iii}$ , and  $C_{jjj}$ ) were drawn from our previous Burnett experiments<sup>12</sup> for nitrous oxide and from the literature<sup>13</sup> for ethane. The following expressions were adopted for the cross virial coefficients ( $B_{ij}$ ,  $C_{ijj}$ , and  $C_{jij}$ ) because the binary system revealed an almost ideal behavior:

$$B_{ij} = (B_{ii} + B_{jj})/2 \quad (2)$$

$$C_{ijj} = (2C_{iii} + C_{jjj})/3 \quad (3)$$

$$C_{jij} = (C_{iii} + 2C_{jjj})/3 \quad (4)$$

The results are given in Figure 6, showing the relative pressure deviations. Bearing in mind that the comparison is based on cross virial coefficients estimated through eqs 2, 3, and 4 and

that the temperature dependence of the virial coefficients was represented by means of third-degree polynomial expressions. The results obtained for one composition ( $z_1 = 0.5208$ ) are clearly out of the deviation trend. For this sample, the mass charged was about twice greater than for all other compositions. In our opinion, the virial equation of state truncated after the third term is not able of accurate representation of the data at greater density. This is why we compared density at superheated region also with the CSD EOS prediction. In this case, the coefficients of the CSD EOS fitted to data along saturation were extrapolated out of range, and also, the  $K_{12}$  value, tuned to the low-temperature two-phase data, was assumed to be temperature independent. Deviations between experiments and CSD EOS are presented in Figure 7. The greatest deviation for the sample  $z_1 = 0.5208$ , observed with the virial EOS prediction, was not



evident when comparing data with the CSD EOS prediction. Moreover, overall deviations in pressures from the CSD EOS prediction showed an AAD = 1.2 %, smaller than the AAD = 1.5 % obtained from the virial EOS comparison. Eventually, it is worth noting that both these models, which are based on different theoretical approaches, showed a similar trend of deviations in terms of temperature.

## Conclusions

An isochoric apparatus has been used to obtain  $PVT_x$  measurements on  $C_2H_6 + N_2O$ . The binary interaction parameters were derived from experimental data in the two-phase region, applying the flash method and the CSD EOS. The dew point parameters were found by interpolating the  $P-T$  isochoric sequences, again applying the CSD EOS. The calculated binary interaction parameter were used to derive the VLE, which revealed strong positive deviations from Raoult's law with a positive azeotrope at around  $x_1 = 0.65$  and  $G^E$  below 300 J/mol at  $T = 243.15$  K. The  $PVT_x$  data are consistent with the values predicted by the virial equation of state and the CSD EOS, using its coefficients for the system constituents derived from independent measurements. In general, the  $PVT_x$  are slightly better represented by the CSD EOS.

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