# Equilibrium Adsorption Measurements of Pure Nitrogen, Carbon Dioxide, and Methane on a Carbon Molecular Sieve at Cryogenic Temperatures and High Pressures<sup>†</sup>

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A detailed experimental study of the adsorption behavior at equilibrium of pure nitrogen, methane, and carbon dioxide gases on a commercial carbon molecular sieve (Shirasagi MSC 3K-161) is reported at temperatures between 115 K to 323 K and pressures up to 5 MPa. A volumetric-type apparatus was used to obtain over 200 excess (Gibbs) adsorption capacity data over this range of pressure and temperature with an estimated uncertainty of 4 %. The absolute adsorption isotherms were type I in the IUPAC classification with the adsorption capacity at constant pressure increasing significantly with decreasing temperature. For each gas, the adsorption data were regressed to a four parameter Toth equation to represent the temperature and pressure dependence of the data with a relative standard uncertainty of 4 %. The optimized parameters from the Toth equation included the isosteric enthalpies of adsorption that were 17 kJ·mol<sup>-1</sup>, 27 kJ·mol<sup>-1</sup>, and 18 kJ·mol<sup>-1</sup> for N<sub>2</sub>, CO<sub>2</sub>, and CH<sub>4</sub>, respectively.

# 1. Introduction

Natural gas is the most environmentally friendly of the fossil fuels, and in contrast to many other clean fuels the infrastructure necessary for its widespread use already exists in many parts of the world. However, before a field can be developed the natural gas must be processed to remove impurities, and the cost of the required processing can be critical to the field's economic viability. Conventionally, CO<sub>2</sub> removal requires the use of a water-based amine solution and large absorber and regeneration towers. The capital and ongoing operational costs for this process are substantial, and the use (regeneration and disposal) of the amine solution can be problematic from the perspectives of both safety and the environment. Furthermore this conventional CO<sub>2</sub> removal process saturates the gas with water vapor and can require the duplication of expensive gas dehydration equipment; gas produced offshore is often dehydrated prior to being piped to an onshore processing facility to avoid hydrate formation. Once onshore, the acid gas removal processes must be followed by a repeat of the dehydration process using glycol contactors and, in cryogenic plants, molecular sieve towers.1

Nitrogen can also be a problematic impurity in natural gas. It is estimated that approximately 16 % of U.S. gas reserves are contaminated with higher amounts of  $N_2$  than the 4 % by mole typically allowed by pipeline specifications.<sup>2</sup> In some liquified natural gas (LNG) plants, disposal of the end-flash vapor (excess to fuel gas requirements) can be hindered if it

contains large concentrations of  $N_2$  as the construction and operation of expensive cryogenic distillation towers to remove the nitrogen may be required. If a gas field contains  $CO_2$  and  $N_2$  at combined levels in excess of 10 % by mole, then the cost of conventional separation processes can make development of the field unviable.

Many workers are investigating the use of alternative gas separation technologies such as membranes and adsorption targeted at removing nitrogen or carbon dioxide from natural gas more efficiently.<sup>2,3</sup> Most of these investigations have focused predominantly on ambient temperatures or above.<sup>4-15</sup> Recently Bae and Lee<sup>16</sup> and Cavenati et al.<sup>17,18</sup> have reported equilibrium and kinetic sorption measurements with N<sub>2</sub>, CH<sub>4</sub>, and CO<sub>2</sub> on carbon molecular sieves to moderate pressures at ambient temperatures and above. However, very little work has been conducted studying the efficacy of adsorption processes at cryogenic temperatures and high pressures, and few adsorption equilibria data exist at these conditions even for simple conventional sieves; some examples are the works of Habgood<sup>19,20</sup> in 1958 and Lederman and Williams<sup>21</sup> in 1964. One aim of our research program is to help address this deficiency and produce adsorption equilibria data at low temperatures and pressures up to saturation. Such data would help in evaluating the feasibility of using pressure swing adsorption to separate nitrogen and carbon dioxide from methane at cryogenic temperatures.

This paper details the design and commissioning of the system that we have developed to characterize the equilibrium adsorption behavior at cryogenic temperatures of pure gases at pressures up to 7 MPa. We also present experimental adsorption isotherms for each of the pure gases  $N_2$ ,  $CO_2$ , and  $CH_4$  on a carbon molecular sieve (CMS), from 115 K to 323 K and at pressures up to 5 MPa, provided that the operating pressure is below the gas saturation temperature at the corresponding temperature.

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Figure 1. Schematic diagram of volumetric-type sorption apparatus.

#### 2. Apparatus and Materials

A schematic of the volumetric-type apparatus used to measure the isothermal adsorption capacities is shown in Figure 1. It consisted primarily of two stainless steel pressure vessels, which we referred to as the adsorption cell and the loading cell. The adsorption cell had an internal volume of approximately 700 cm<sup>3</sup>. The CMS was held in a basket made from stainless steel mesh suspended from the lid of the adsorption cell. The temperature of the adsorption cell was monitored with two 100  $\Omega$  platinum resistance thermometers (PRT), which were calibrated to ITS-90 over the range 77 K to 303 K with an estimated uncertainty of 0.08 K. (All uncertainties reported here are one standard uncertainty.) Both thermometers were mounted on the interior walls of the adsorption cell, one immediately adjacent to the basket holding the sieve, and the other on the cell's lid.

The adsorption cell was suspended from the lid of a cryogenic dewar and cooled by the boil-off vapor of a small amount of liquid nitrogen in the bottom of the dewar. The liquid nitrogen level was maintained at a constant value using an autonomous liquid nitrogen pumping system, and the boil-off vapor was prevented from stratifying by using a cryogenically compatible fan suspended from the bottom of the adsorption cell. An electrical resistance heater and a third PRT-100 element were mounted on the external surface of the cell and used in a Proportional Integral (PI) loop to control the cell's temperature. At steady state, the temperature stability of the adsorption cell was approximately  $\pm$  0.05 K; however steady-state thermal gradients were also observed with the two calibrated thermometers. These gradients dominate the temperature uncertainty of our data, which is estimated to be 1 K. Gas pressure in the adsorption cell was monitored by a quartz crystal transducer located outside the dewar with a full scale of 7 MPa and an estimated uncertainty over this range of 0.4 kPa.

The loading cell, which had an internal volume of approximately 1800 cm<sup>3</sup>, was immersed in a stirred water bath

maintained at ambient temperature. Its temperature was monitored using a PRT-100 element, calibrated to ITS-90 with an estimated uncertainty of 0.08 K, mounted on one of the internal walls. The uncertainty of the gas temperature inside the loading cell was estimated to be 0.1 K. Pressure in the loading cell was monitored by a second quartz crystal transducer with a full scale of 14 MPa and an estimated uncertainty of 0.8 kPa. The adsorption and loading cells were connected to each other, to their respective pressure transducers and to the gas handling manifold by 1/8 in. stainless steel tubing and high pressure valves and fittings. The total volume of these unthermostatted lines was estimated to be 5.7 cm<sup>3</sup>, and thus their impact on the uncertainty of the adsorption measurements was less than 0.2 %. The temperatures and pressures measured by each of the aforementioned transducers were logged using commerical data acquisition software. The PI control loop for the adsorption cell temperature was also implemented in this software.

The following list of pure gases and supplier-reported purities (by mole) were used over the course of the measurements: He, 99.9995 %; Ar, 99.999 %; CH<sub>4</sub>, 99.99 %; N<sub>2</sub>, 99.999 %; and CO<sub>2</sub>, 99.995 %. When changing the gas in the system, the cells and manifold were flushed several times by evacuating them to a pressure below 10 Pa (measured using a separate thermal convection vacuum gauge) and then pressurizing the system with the new gas to about 500 kPa (depending on the fluid's saturation pressure). The estimated uncertainty in the adsorption measurements due to gas purity was negligible.

The adsorbent used in this work was "Shirasagi MSC 3K-161" (Japan EnviroChemicals, Japan), a coconut-shell based carbon molecular sieve in the form of extruded pellets of cylindrical shape (length  $\approx 4$  mm, diameter  $\approx 1.8$  mm, lot number M683). This CMS has the same average micropore diameter as "Shirasagi MSC 3K-162",<sup>18</sup> which Bae and Lee<sup>16</sup> reported as being less than 4 Å with a BET surface area of 522.59 m<sup>2</sup> · g<sup>-1</sup>. Before introducing the adsorbent in the adsorption cell, the adsorbent was regenerated under vacuum (10 Pa) at 673 K for 24 h. The measured mass of regenerated CMS loaded into the cell was  $m_{ads} = (41.17 \pm 0.01)$  g.

## 3. Methodology and Analysis

**3.1.** Calibration of the Loading Cell. The volumetric technique for measuring equilibrium adsorption capacities requires an accurate knowledge of the volumes of the vessels between which gas is being transferred. We therefore determined first the volume of the loading cell at ambient temperature. Then, prior to measuring the CMS isothermal adsorption capacity for  $N_2$ , CH<sub>4</sub>, or CO<sub>2</sub>, we determined the volume of the adsorption cell at the measurement temperature.

The volume of the loading cell was determined by connecting it to a separate gas cylinder containing approximately 1 MPa of either helium or argon. Before connecting it to the loading cell, the mass of the cylinder plus the gas it contained was measured using an electronic balance with a full scale of 1100 g and a resolution of 0.001 g. Gas was then transferred from the cylinder to the evacuated loading cell, and once equilibrium was reached, the loading cell was isolated and its pressure and temperature recorded. The gas cylinder was reweighed and the mass of gas transferred to the loading cell was calculated; a small correction was made to account for gas lost between the valves isolating the cylinder and loading cell. Approximately 2 g of helium or 30 g of argon were transferred, depending on the initial pressure in the gas cylinder. The reference equations of state (EOS) for helium<sup>22</sup> and for argon,<sup>23</sup> as implemented in the software REFPROP 8.0,<sup>24</sup> were used to calculate the gas

density in the loading cell from the measured temperature and pressure. The loading cell volume was calculated from the mass transferred and the EOS densities with an estimated relative uncertainty of 0.2 %, based on the quadrature combination of the mass, temperature, pressure, and EOS uncertainties. The volume determination was repeated six times, thrice with helium and thrice with argon. The average of the six loading cell volume determinations was  $(1822 \pm 2) \text{ cm}^3$ .

*Volume and Sorption Measurements.* The volumetric method of measuring equilibrium adsorption capacity is based upon the transfer of a known amount of gas from one vessel (the loading cell) to another vessel containing the adsorbent (the adsorption cell). The two vessels constitute a closed system, and thus the material balance for the gas in the two vessels leads to

$$V_{\rm L}(\rho(p,T)_{\rm L}^{\rm init} - \rho(p,T)_{\rm L}^{\rm fin}) + V_{\rm A}(\rho(p,T)_{\rm A}^{\rm init} - \rho(p,T)_{\rm A}^{\rm fin}) + (n_{\rm ads}^{\rm init} - n_{\rm ads}^{\rm fin}) = 0 \quad (1)$$

Here, the subscripts "L" and "A" refer to measurements of the loading and adsorption cells, respectively; the superscripts "init" and "fin" refer to measurements taken before and after the gas was transferred, respectively; V is the cell volume; p is the measured cell pressure; T is the measured cell temperature;  $\rho$  is the molar gas density determined using an EOS at the measured pressure and temperature; and  $n_{ads}$  is the number of moles of gas adsorbed on the CMS. Once the volume of the loading cell was determined, it was necessary to determine  $V_A$ via helium calibration before eq 1 could be used to analyze subsequent experiments with N<sub>2</sub>, CH<sub>4</sub>, and CO<sub>2</sub> to determine  $n_{ads}$ .

The volume of the adsorption cell at each temperature was determined by transferring helium from the loading cell into the adsorption cell with a estimated relative uncertainty of 0.3 %. Fundamental to the volumetric technique is the widely used assumption that for helium at low pressures,  $n_{ads} \equiv 0.25 - 29$  The validity of this assumption may be questionable, particularly at cryogenic conditions;<sup>30</sup> however there does not to seem be any better method for this purpose at present. More importantly, an adequate objective of these measurements is to determine  $n_{ads}$ for N<sub>2</sub>, CH<sub>4</sub>, and CO<sub>2</sub> relative to helium, which this methodology achieves. At each temperature, we made several such helium transfers into the adsorption cell with  $p_A^{fin}$  ranging from 50 kPa to 400 kPa and the calculated value of  $V_A$  on each occasion was consistent within the estimated uncertainty. The apparent volume of the adsorption cell decreased slightly with temperature with  $(1/V)(\Delta V/\Delta T) \approx 4 \cdot 10^{-5} \text{ K}^{-1}$  over the temperature range studied. Corrections due to the compressibility of the adsorbent and the pressure vessel were negligible for our working pressures.<sup>2</sup>

Between each measurement of an adsorption isotherm, the adsorption cell was evacuated to a pressure below 10 Pa and heated to 353 K for a period of 24 h. Then the adsorption cell was returned to the required temperature with several hours required for stabilization. The loading cell was filled to a predetermined pressure with the pure gas of interest and then both cells were isolated from the gas handling manifold. The sorption isotherm was measured in two phases; for adsorption measurements, gas was transferred from the loading cell to the adsorption cell. Once the upper pressure on the isotherm was reached in the adsorption cell, the desorption phase of the measurement began with gas being transferred from the adsorption cell to the loading cell. After each transfer, which was enumerated by the index *i*, eq 1 was used to calculate  $\Delta n_i \equiv$ 

Table 1. Equilibrium Sorption Data for Nitrogen on Shirasagi MSC3K-161

р	$Q_{\rm ex}$	$Q_{ m abs}$	р	$Q_{\rm ex}$	$Q_{ m abs}$	
kPa	$mmol \cdot g^{-1}$	$mmol \cdot g^{-1}$	kPa	$mmol \cdot g^{-1}$	$mmol \cdot g^{-1}$	
T = 298  K						
150.5	0.42	0.42	5039.6	2 19	2 38	
222.2	0.71	0.71	1195.9	2.12	2.50	
522.2	0.71	0.71	4105.0	2.13	2.29	
505.4	0.93	0.94	1885.8	1.70	1.75	
666.2	1.08	1.09	1586.7	1.60	1.64	
883.6	1.25	1.26	1315.2	1.49	1.52	
1093.2	1.37	1.39	1089.8	1.38	1.40	
1318.1	1.48	1.51	881.9	1.25	1.27	
1601.2	1 59	1 64	670.1	1 10	1 11	
1050.3	1.59	1.01	401.6	0.03	0.04	
2216.0	1.71	1.70	228.0	0.93	0.94	
2316.9	1.81	1.88	328.0	0.73	0.74	
3289.0	2.02	2.14	154.8	0.44	0.44	
3967.8	2.10	2.25				
		T = 2	48 K			
229.0	1.28	1.28	2982.3	2.86	3.04	
684.8	2.01	2.04	2630.2	2.81	2.96	
1096 7	2 30	2 35	2300.1	2 71	2 84	
1533.0	2.50	2.55	2031.7	2.66	2.01	
1996.0	2.50	2.57	1757.2	2.00	2.77	
1000.9	2.01	2.71	1/3/.2	2.00	2.09	
2313.1	2.71	2.84	1509.2	2.53	2.61	
2609.1	2.76	2.91	1249.7	2.45	2.51	
3035.2	2.88	3.06	990.2	2.31	2.36	
3570.7	2.90	3.12	822.5	2.20	2.24	
4246.7	2.93	3.20	608.3	2.01	2.03	
4611.8	2.95	3 25	397.8	1 74	1 75	
1011.0	2.93	3 21	303.1	1.56	1.73	
3554.8	2.94	3.13	150.3	1.50	1.57	
5554.0	2.90	5.15	157.5	1.10	1.10	
05.1	1.90	T = 1	98 K	2.46	2.07	
95.1	1.89	1.90	3975.7	3.40	3.87	
220.3	2.42	2.44	1862.5	3.47	3.64	
409.9	2.80	2.83	729.2	3.12	3.18	
623.1	3.03	3.08	298.2	2.62	2.64	
1024.6	3.27	3.36	135.1	2.12	2.13	
2979.0	3.54	3.84	66.4	1.67	1.68	
5023.7	3.36	3.89				
		T = 1	48 K			
14.0	2.51	2 5 1	3785 1	1 17	5.00	
14.0	2.51	2.31	2052.2	4.17	5.00	
122.7	3.08	3.70	2052.2	4.44	4.80	
232.9	3.97	4.00	987.2	4.43	4.58	
402.0	4.18	4.23	456.1	4.23	4.30	
604.7	4.31	4.40	217.9	3.96	3.99	
999.9	4.42	4.57	108.8	3.65	3.67	
2912.9	4.35	4.93	57.7	3.33	3.34	
4542.9	3.98	5.09	0,11	0.00	0.0	
-5-2.7	5.70	5.05	1.5.17			
0.01	2.54	T = 1	15 K	5 41	5 (0	
0.01	2.54	2.54	/98.2	5.41	5.62	
8.9	4.26	4.26	13/3.4	5.32	5.74	
64.1	4.92	4.93	1115.8	5.27	5.58	
152.6	5.17	5.20	631.4	5.40	5.56	
246.6	5.32	5.38	336.7	5.40	5.48	
347.8	5.40	5.48	179.6	5.31	5.35	
440.8	5.10	5 53	1, 7.0	0.01	2.55	

 $(n_{ads}^{fin} - n_{ads}^{init})_i$ . For i = 1,  $n_{ads}^{init}$  was assumed to be zero (based on the lengthy period of evacuation and heating of the adsorption cell prior to the start of each isotherm).

The amount of gas transferred between each step along the isotherm ranged between about 50 mmol at low pressures and 500 mmol at high pressures. The uncertainty of the reference EOS for N<sub>2</sub>,<sup>31</sup> CH<sub>4</sub>,<sup>32</sup> and CO<sub>2</sub><sup>33</sup> was negligible compared to the uncertainties of  $V_A$  and  $V_L$  and those of our pressure and temperature measurements. The cumulative excess surface adsorption  $Q_{ex}^{25}$  on the CMS was then calculated using

$$Q_{\rm ex} = \frac{\sum_{i} \Delta n_i}{m_{\rm ads}} \tag{2}$$

Table 2. Equilibrium Sorption Data for Methane on Shirasagi MSC3K-161

 Table 3. Equilibrium Sorption Data for Carbon Dioxide on

 Shirasagi MSC 3K-161

р

 $Q_{\rm ex}$ 

 $Q_{\rm abs}$ 

 $Q_{\rm abs}$ 

p

 $Q_{\rm ex}$ 

p	$Q_{\rm ex}$	$Q_{ m abs}$	р	Qex	$Q_{ m abs}$		
kPa	$mmol \cdot g^{-1}$	$mmol \cdot g^{-1}$	kPa	$mmol \cdot g^{-1}$	$mmol \cdot g^{-1}$		
T = 298  K							
241.5	1.45	1.46	3716.2	3.01	3.24		
550.3	1.96	1.98	3023.9	2.94	3.12		
848.5	2.25	2.29	2484.0	2.87	3.01		
1124.5	2.43	2.48	1981.9	2.76	2.87		
1377.5	2.54	2.61	1480.6	2.62	2.70		
1610.3	2.63	2.71	1053.9	2.44	2.49		
1883.6	2.74	2.84	831.8	2.29	2.33		
2267.9	2.82	2.95	609.5	2.10	2.13		
2659.5	2.88	3.04	503.1	1.98	2.00		
3063.9	2.93	3.12	399.0	1.83	1.84		
3740.8	2.98	3.22	306.7	1.66	1.67		
4267.8	3.02	3.30	206.0	1.41	1.42		
5155.6	3.03	3.37	105.5	1.02	1.02		
4377.1	3.00	3.28					
T = 248  K							
43.5	1.45	1.45	3462.8	3.56	3.89		
144.8	2.20	2.20	2293.2	3.55	3.76		
382.8	2.79	2.81	1079.0	3.33	3.42		
843.9	3.20	3.26	599.6	3.07	3.11		
1588.9	3.44	3.57	505.7	2.98	3.02		
2063.6	3.51	3.69	302.2	2.69	2.71		
3131.6	3.59	3.88	147.0	2.25	2.26		
4054.3	3.55	3.96	71.9	1.80	1.80		
5046.2	3.48	4.02	40.3	1.46	1.46		
4433.3	3.51	3.97	25.2	1.20	1.20		
		T =	198 K				
25.5	2.48	2.48	4499.1	3.70	4.67		
112.7	3.31	3.32	3753.3	3.89	4.58		
227.2	3.62	3.64	1992.5	4.13	4.41		
427.7	3.86	3.91	827.4	4.05	4.15		
624.9	3.98	4.05	324.4	3.76	3.79		
1033.4	4.11	4.24	139.6	3.41	3.43		
3004.7	3.98	4.45	63.8	3.04	3.05		
T = 148  K							
38.4	3.56	3.56	701.0	4.35	4.48		
115.6	4.05	4.07	813.7	4.32	4.48		
201.2	4.23	4.26	5/4./	4.37	4.48		
293.2	4.51	4.30	275.9	4.28	4.33		
403.5	4.30	4.43	128.9	4.11	4.13		
505.5	4.37	4.40	04.9	3.91	3.92		
003.8	4.30	4.48					

After each pressure change, equilibrium was considered to be reached when  $\Delta n_i$  was stable within 0.2 %. At pressures above 1 MPa, this criterion was usually achieved within 1 h, whereas at pressures below 100 kPa approximately 10 h was required. The repeatability of our  $Q_{ex}$  data along each isotherm (including adsorption and desorption steps) was better than 1 %. On the basis of an error propagation analysis, we estimate the average relative uncertainty of the  $Q_{ex}$  data to be 4 %.

To ensure that the values of  $Q_{ex}$  determined at each pressure by the cumulative application of eq 1 were not adversely affected by error propagation, a check measurement was made at the conclusion of the adsorption isotherm. The adsorption cell was re-evacuated and then gas was transferred to it from the loading cell, causing a pressure increase of between 1 and 2 MPa. The value of  $Q_{ex}$  at the equilibrium pressure obtained from this single transfer was always found to be within 2 % of the value determined from the multiple small transfers.

To a first approximation, the term  $V_A$  in eq 1 was assumed to be constant for all pressures on the adsorption isotherm. However, this assumption ignores the volume occupied by the adsorbed phase, which increases with gas pressure, and causes the actual value of  $V_A$  available to the gas phase to decrease. The absolute adsorption capacity of the CMS,  $Q_{abs}$ , is obtained

kPa	$mmol \cdot g^{-1}$	$mmol \cdot g^{-1}$	kPa	$mmol \cdot g^{-1}$	$mmol \cdot g^{-1}$
		T = 3	323 K		
10.9	0.34	0.34	4117.2	3.99	4.35
24.2	0.62	0.62	2004.7	3.77	3.91
94.7	1.42	1.42	1388.6	3.61	3.70
183.6	1.94	1.95	627.9	3.09	3.12
256.5	2.24	2.25	292.7	2.46	2.47
465.2	2.76	2.78	125.7	1.77	1.77
648.1	3.04	3.07	66.2	1 30	1 30
1668.4	3.68	3.80	38.8	0.98	0.98
3526.4	3.00	4 28	50.0	0.90	0.70
5520.4	5.77	T _ (	000 IZ		
12.1	0.72	1 = 1	4022 5	4 20	171
12.1	0.72	0.72	4052.5	4.29	4.74
20.2	1.12	1.12	2313.3	4.29	4.50
96.3	2.08	2.09	1/65.4	4.21	4.37
186.9	2.65	2.65	829.8	3.85	3.92
2/0.1	2.96	2.98	327.3	3.20	3.22
480.4	3.41	3.44	142.5	2.49	2.49
6/1.5	3.64	3.69	93.7	2.13	2.13
1786.4	4.13	4.29	54.5	1.69	1.69
3469.2	4.33	4.69			
		T = 2	273 K		
16.6	1.47	1.47	2169.9	4.83	5.09
32.2	1.96	1.96	1376.4	4.69	4.84
110.2	3.05	3.05	740.6	4.45	4.52
259.4	3.81	3.83	340.1	4.00	4.03
535.1	4.26	4.31	153.0	3.36	3.37
719.5	4.42	4.48	81.9	2.84	2.84
1181.7	4.62	4.74	53.2	2.40	2.40
		T = 2	748 K		
8.0	1.90	1.90	859.0	5.04	5.15
21.2	2.59	2.59	322.5	4.66	4.70
86.5	3.85	3.85	142.1	4.19	4.20
195.2	4 42	4 44	74.8	3 74	3 74
467.1	4.86	4 92	57.4	3 51	3 51
685.3	5.01	5.09	33.2	3.07	3.08
1107.5	5.01	5 30	25.2	2.84	2.85
1107.0	5.10	T = 2	223 K	2.01	2.05
1.8	1.73	1.73	440.9	5.49	5.55
4.9	2.61	2.61	315.2	5.38	5.42
11.7	3.43	3.44	191.6	5.20	5.23
81.6	4.78	4.79	91.9	4.88	4.89
303.7	5.35	5.39	52.3	4.60	4.60
515.1	5.54	5.61	- =		

from  $Q_{ex}$  by correcting for the volume occupied by the adsorbed phase.<sup>34</sup>

$$Q_{\rm abs} = \frac{Q_{\rm ex}}{1 - \rho_{\rm M,gas}/\rho_{\rm M,ads}} \tag{3}$$

Here  $\rho_{M,gas}$  and  $\rho_{M,ads}$  are the mass densities of the gas phase and adsorbed phase, respectively. The density of the adsorbed phase cannot be measured directly but can be inferred using one of several different methods described in the literature.<sup>28,34–36</sup> We used the Ono–Kondo model<sup>34</sup> to obtain  $\rho_{M,ads} = 0.354$ , 0.701, and 0.996 g·cm<sup>-3</sup> for CH<sub>4</sub>, N<sub>2</sub>, and CO<sub>2</sub>, respectively. These densities were assumed to be independent of pressure and temperature. The values of  $\rho_{M,ads}$  obtained with the various methods<sup>34</sup> (Ono–Kondo model, ZGR EOS, VDW covolume) vary by about 10 %, increasing the relative uncertainty of  $Q_{abs}$ to 5 %.

### 4. Results and Discussion

Measured values of the  $Q_{ex}$  and  $Q_{abs}$  for the CMS as a function of temperature and pressure are presented in Table



**Figure 2.** (a) Isothermal absolute adsorption data for N<sub>2</sub> on Shirasagi MSC 3K-161 measured in this work:  $\blacklozenge$ , 298 K;  $\blacksquare$ , 248 K;  $\blacktriangle$ , 198 K;  $\bigcirc$ , 148 K; ×, 115 K; and –, absolute adsorption capacity  $Q_{calc}$  determined using eq 4 with the N<sub>2</sub> parameters listed in Table 4. (b) Deviations ( $Q_{meas} - Q_{calc}$ ) for the  $Q_{meas}$  data shown in (a) and from two literature sources. ---, estimated uncertainty, this work; –, ref 16 (293 K); – – –, ref 17 (298 K).

1 for  $N_2$ , in Table 2 for  $CH_4$ , and in Table 3 for  $CO_2$ . The measurement temperatures were 115 K, 148 K, 198 K, 248 K, and 298 K for  $N_2$ ; 148 K, 198 K, 248 K, and 298 K for  $CH_4$ ; and 223 K, 248 K, 273 K, 298 K, and 323 K for  $CO_2$ . The maximum pressure on each isotherm was either between 4000 kPa and 5200 kPa, or slightly below the gas saturation pressure at the corresponding temperature.

At low temperatures, some of the excess adsorption isotherms we measured exhibited a maximum; such maxima have been reported by many workers (for example Menon,<sup>26</sup> Zhou et al.,<sup>36</sup> and Sakurovs et al.<sup>37</sup>) and are a consequence of the increasing volume occupied by the adsorbed phase. As pressure increases, the excess adsorption typically increases to a maximum and then decreases linearly with increasing gas density, reaching zero again at the pressure where the gas density becomes the same as the adsorbed phase density.

In the top panels of Figures 2 to 4, the absolute adsorption isotherms are shown as a function of the pressure for each pure gas. Absolute adsorption increases with increasing pressure and decreasing temperature. The shape of these isotherms corresponds to type I in the IUPAC classification,<sup>38</sup> characterized by a saturation limit at high pressures when the micropores of the CMS are completely filled. For each



**Figure 3.** (a) Isothermal absolute adsorption data for CH<sub>4</sub> on Shirasagi MSC 3K-161 measured in this work:  $\blacklozenge$ , 298 K;  $\blacksquare$ , 248 K;  $\blacktriangle$ , 198 K;  $\bigcirc$ , 148 K; and -, absolute adsorption capacity  $Q_{calc}$  determined using eq 4 with the CH<sub>4</sub> parameters listed in Table 4. (b) Deviations ( $Q_{meas} - Q_{calc}$ ) for the  $Q_{meas}$  data shown in (a) and from two literature sources. ---, estimated uncertainty, this work; -, ref 16 (293 K); - -, ref 17 (298 K).

gas, the repeatability of the adsorption data (measured as the pressure was increased) and the desorption data (measured as the pressure was decreased) was significantly better than the estimated uncertainty of the data. Thus the equilibrium sorption behavior of all three gases was completely reversible.

For each pure gas, the entire set of  $Q_{ads}$  data were regressed to the Toth equation<sup>39</sup>

$$Q_{\rm abs} = Q_{\rm max} \frac{Kp}{\left(1 + (Kp)^n\right)^{1/n}}$$
 (4a)

$$K = K_0 \exp\left(\frac{-\Delta H}{RT}\right) \tag{4b}$$

Here *R* is the molar gas constant and  $\Delta H$  is the isosteric enthalpy of adsorption. In the regression,  $\Delta H$  was treated as an adjustable parameter along with the three empirical parameters  $Q_{\text{max}}$ ,  $K_0$ , and *n*. The Toth equation is recommended for microporous adsorbents<sup>40</sup> because it has several advantages. It reduces to Henry's law in the limit of low pressure, it contains only four parameters, and it provides an analytical expression for the  $Q_{\text{ads}}$  that is an explicit function of temperature and pressure. The Toth equation



**Figure 4.** (a) Isothermal absolute adsorption data for CO<sub>2</sub> on Shirasagi MSC 3K-161 measured in this work:  $\triangle$ , 323 K;  $\blacklozenge$ , 298 K;  $\blacklozenge$ , 273 K;  $\blacksquare$ , 248 K; \*, 223 K; and —, absolute adsorption capacity  $Q_{calc}$  determined using eq 4 with the CO<sub>2</sub> parameters listed in Table 4. (b) Deviations ( $Q_{meas} - Q_{calc}$ ) for the  $Q_{meas}$  data shown in (a) and from two literature sources. ---, estimated uncertainty, this work; –, ref 16 (293 K); – – –, ref 18 (298 K).

Table 4. Toth Equation Parameters (Equation 4) for the Absolute Adsorption Capacity of Shirasagi MSC 3K-161 for  $N_2$ ,  $CH_4$ , and  $CO_2$ 

n dioxide
$\pm 0.071$
$\pm 0.065$
$\pm 0.21$
$\pm 0.010$
1

reduces to the Langmuir equation when the parameter n = 1; this parameter characterizes whether the gas adsorption on the adsorbent's surface is heterogeneous  $(n \neq 1)$  or homogeneous (n = 1).

The optimized parameter values resulting from the regression of eq 4 to each pure fluid data set are shown, together with their statistical uncertainties in Table 4. The relative standard errors of the fits ranged from 1 % for CO<sub>2</sub> to 3.5 % for N<sub>2</sub>; all are smaller than the data's estimated relative standard uncertainty of 5 %. The deviation of our data from the regressed model are shown for all gases in the bottom panels of Figures 2 to 4. Systematic trends with decreasing temperature are apparent in the deviation plots, indicating

that the model's functional form is not able to exactly describe the data within the precision of the data. Nevertheless, the model is more than adequate for the purpose of engineering design calculations.

Also shown in the bottom panels of Figures 2 to 4 are comparisons of our results with those of other workers who have investigated the adsorption of N<sub>2</sub>, CH<sub>4</sub>, and CO<sub>2</sub> on carbon molecular sieves from the same manufacturer. Bae and Lee<sup>16</sup> studied Shirasagi MSC 3A-162 using a volumetrictype apparatus at 293 K and pressures to 1.6 MPa. According to the manusfacturer, the equilibrium sorption properties of that CMS are very similar to those of the CMS studied in this work, although the two adsorbents differ in their kinetic properties.<sup>41</sup> Cavenati et al.<sup>17,18</sup> used a gravimetric-type apparatus to investigate the same adsorbent studied here for N<sub>2</sub> and CH<sub>4</sub> at 298 K at pressures to 2 MPa, and CO<sub>2</sub> at 298 K and 323 K at pressures to 0.5 MPa. (Cavenati et al. referred to their CMS as Takeda CMS3K, and Bae and Lee referred to their CMS as Takeda CMS-T3A. These names are now obsolete.<sup>41</sup>) The salient properties of the two adsorbents are given in refs 16 and 17, respectively; further information is also available from the manufacturer.<sup>42</sup>

In general, the data of Bae and Lee<sup>16</sup> and Cavenati et al.<sup>17,18</sup> deviate from the model fit to our data by between one and two times the estimated standard uncertainty of our results. However, Bae and Lee<sup>16</sup> and Cavenati et al.<sup>17,18</sup> only reported the model parameters which describe their data, rather than the data themselves or their uncertainties. It is difficult therefore to assess whether our data are consistent with theirs within the combined experimental uncertainties. Furthermore, in several cases the deviations between our data and those of Bae and Lee<sup>16</sup> or those of Cavenati et al.<sup>17,18</sup> are comparable in magnitude to the deviations between the two literature data sets. Since the deviations are in most instances not more than about twice our estimated uncertainty, it is plausible that the data of of Bae and Lee<sup>16</sup> and Cavenati et al.<sup>17,18</sup> might agree with ours within the combined uncertainty.

Fundamental adsorption data are critical to the evaluation of potential Pressure Swing Absorption separation processes based on particular adsorbents. This work represents the first stage in an ongoing effort to develop new experimental systems to measure the equilibrium and kinetic sorption phenomena of various adsorbents at cryogenic temperatures and high pressures. The capabilities of the apparatus described here are currently being extended to allow us to investigate such phenomena in multicomponent gas mixtures.

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