Kinetics and Thermodynamics of Redox Processes in Room Temperature Ionic Liquids: The Use of Voltammetry and the Disproportionation of Radical Cations of *N*,*N*-Dimethyl-*p*-phenylenediamine in 1-Butyl-3-methylimidazolium Tetrafluoroborate

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The voltammetry of *N*,*N*-dimethyl-*p*-phenylenediamine, DMPD, in the room temperature ionic liquid (RTIL) 1-butyl-3-methylimidazolium tetrafluoroborate, $[C_4mim][BF_4]$ is studied by means of microdisc voltammetry over a wide range of voltage scan rates incorporating the transition from linear to convergent diffusion. Two voltammetric waves were recorded corresponding to the formation of the radical cation DMPD⁺ and dication DMPD²⁺; both are stable on the voltammetric time scale. Double potential step chronoamperometry was used to measure the diffusion coefficients of DMPD and DMPD⁺. These values, in conjunction with simulations of the cyclic voltammetry, permit the inference of the equilibrium constant and kinetics for the comproportionation of DMPD and DMPD⁺ to be studied.

1. Introduction

Room temperature ionic liquids (RTILs) have led to substantial synthetic innovations in both academic^{1–3} and industrial⁴ contexts. Advantages such as solvent tunability,⁵ greenness,⁶ recyclability, and negligible volatility are seen over conventional nonaqueous solvents. However, unlike reactions in the latter media, the measurement of reaction parameters, both kinetic and thermodynamic, has lagged, and comparatively less data is available to guide the planning and interpretation of chemistry in the RTIL environment.

Voltammetry, at least in principle, offers an experimentally simple yet accurate route into the quantitative measurement of both kinetic and thermodynamic data. For example, in the case of an arbitrary species, A, which can undergo oxidation to species B and C:

$$A - e^{-} \rightleftharpoons B(E_{f,A/B}^{\Theta})$$
(1)

$$\mathbf{B} - \mathbf{e}^{-} \rightleftharpoons \mathbf{C}(E_{\mathbf{f},\mathbf{B}/\mathbf{C}}^{\Theta}) \tag{2}$$

If two separate voltammetric features are attributable to these processes, then their analysis can give information about the comproportionation reaction:⁷

$$A + C \rightleftharpoons 2B$$
 (3)

where, specifically, the formal potentials $E_{f, B/C}^{\Theta}$ and $E_{f, B/C}^{\Theta}$ will provide information about the equilibrium constant:

$$K_{\rm eqm} = \frac{[B]^2}{[A][C]} \tag{4}$$

and ΔG^{Θ} for reaction 3:

$$\Delta G^{\ominus} = -RT \ln(K_{\text{eqm}}) = -F(E_{\text{f,B/C}}^{\ominus} - E_{\text{f,A/B}}^{\ominus}) \qquad (5)$$

The kinetics of reaction 3 may sometimes be discernible from the precise voltammetric waveshapes and their separation.⁸ However, in the context of voltammetry using large planar electrodes ("macroelectrodes"), Savéant and Andrieux have shown in a paper, long regarded as a classic,⁹ that if species A, B, and C have identical diffusion coefficients, D_A , D_B , and D_C , then the voltammetry is completely insensitive to the rate of comproportionation. The same is true of voltammetry recorded under diffusional modes other than linear.⁷ Only if there is an appreciable disparity in the diffusion coefficients can the kinetics of reaction 3 become apparent in the voltammetry. Accordingly, voltammetry has been relatively little used in the context of measuring comproportionation kinetics except insofar that it has been coupled to spectroscopic methods, such as electron spin resonance (ESR)⁹ and UV-visible spectroscopy.¹⁰ At the heart of this neglect is the fact that, in most aqueous and nonaqueous systems, species A, B, and C often have closely similar diffusion coefficients unless they are structurally very dissimilar.

As electrochemical research using RTILs as solvents^{11–17} has matured, it has become apparent that, unlike conventional solvents such as DMF, DMSO, and acetonitrile, the difference in diffusion coefficients between structurally similar species in a redox couple is not always insignificant. An extreme example, discovered by Buzzeo et al.,¹⁸ is that of the oxygen/superoxide (O_2/O_2^{-}) couple in which the anion has a diffusion coefficient about 30 times smaller than the neutral parent in the RTIL *n*-hexyltriethylammonium bis(trifluoromethanesulfonyl)imide $([N_{6222}][N(Tf)_2])$. Work by Evans et al.¹² and by Hapiot et al.¹⁹ have shown that even for aromatic molecules, such as N, N, N', N'tetramethyl-para-phenylenediamine, the diffusion coefficients differ significantly between the neutral parent and the corresponding radical cation and dication despite close structural similarity. Analyte ions of higher charge typically diffuse less rapidly because of electrostatic interaction with the ions in the solvent. As a consequence, it can be anticipated that, in ionic liquids, voltammetric analysis can be used as a means of obtaining information about the rate of comproportionation in solution.

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Figure 1. Structure of N,N-dimethyl-p-phenylenediamine (DMPD).



Figure 2. Structure of the ionic liquid 1-butyl-3-methylimidazolium tetrafluoroborate ($[C_4mim][BF_4]$).



Figure 3. Circular disk electrode inlaid on an insulating support.

In this paper we consider the oxidation of *N*,*N*-dimethyl-*p*-phenylenediamine, DMPD (Figure 1), in the ionic liquid 1-butyl-3-methylimidazolium tetrafluoroborate, $[C_4mim][BF_4]$ (Figure 2), using a microdisc electrode. It is shown that comparison between experimental and simulated data can be used to extract both kinetic and thermodynamic data, including information about comproportionation. Also noted is the value of double potential step chronoamperometry in elucidating the diffusion coefficients of DMPD and DMPD⁺.¹⁴

2. Theory

We consider a stepwise two electron oxidation reaction at a microdisc electrode (Figure 3), as defined by eqs 1 and 2, where we additionally allow for comproportionation (eq 3). The heterogeneous rate constants for the first two steps are $k_{A/B}^0$ and $k_{B/C}^0$, respectively. The homogeneous rate constant for comproportionation is k_{comp} and for disproportionation is k_{disp} . The equilibrium constant for eq 3 is therefore $K_{eqm} = (k_{comp})/(k_{disp})$. Specifically we consider the case of cyclic voltammetry.

$$\frac{\partial c_{\rm A}}{\partial t} = D_{\rm A} \left(\frac{\partial^2 c_{\rm A}}{\partial r^2} + \frac{1}{r} \frac{\partial c_{\rm A}}{\partial r} + \frac{\partial^2 c_{\rm A}}{\partial z^2} \right) + k_{\rm disp} c_{\rm B}^2 - k_{\rm comp} c_{\rm A} c_{\rm C}$$
(6)

$$\frac{\partial c_{\rm B}}{\partial t} = D_{\rm B} \left(\frac{\partial^2 c_{\rm B}}{\partial r^2} + \frac{1}{r} \frac{\partial c_{\rm B}}{\partial r} + \frac{\partial^2 c_{\rm B}}{\partial z^2} \right) - 2(k_{\rm disp} c_{\rm B}^2 - k_{\rm comp} c_{\rm A} c_{\rm C})$$
(7)

$$\frac{\partial c_{\rm C}}{\partial t} = D_{\rm C} \left(\frac{\partial^2 c_{\rm C}}{\partial r^2} + \frac{1}{r} \frac{\partial c_{\rm C}}{\partial r} + \frac{\partial^2 c_{\rm C}}{\partial z^2} \right) + k_{\rm disp} c_{\rm B}^2 - k_{\rm comp} c_{\rm A} c_{\rm C}$$
(8)

where all symbols are defined in Table 1.

2.1. Cyclic Voltammetry. In the cyclic voltammetry experiment, the applied potential, *E*, is swept from an initial value,

Table 1. Dimensional Parameters

dimensional			
parameter	definition		
$\alpha_{X/Y}$	transfer coefficient for $X - e^- \rightleftharpoons Y/unitless$		
$c_{\rm X}$	concentration of species X/mol·m ⁻³		
$c_{\rm X,0}$	concentration of species X at the electrode surface/mol·m ⁻³		
$c_{\rm X}^*$	concentration of species X in bulk solution/mol·m ⁻³		
$D_{\rm X}$	diffusion coefficient of species $X/m^2 \cdot s^{-1}$		
Ε	applied potential/V		
$E_{\rm f, X/Y}^{\Theta}$	formal reduction potential for $X - e^- \rightleftharpoons Y/V$		
$k_{\rm X/Y}^0$	electrochemical rate constant for $X - e^- \rightleftharpoons Y/m \cdot s^{-1}$		
$k_{\rm comp}$	rate constant for comproportionation/dm ³ ·mol ⁻¹ ·s ⁻¹		
$k_{\rm disp}$	rate constant for comproportionation/dm ³ ·mol ⁻¹ ·s ⁻¹		
i	current/A		
j	flux across the electrode/mol \cdot m ² \cdot s ⁻¹		
r	radial coordinate/m		
t	time/s		
υ	scan rate/V \cdot s ⁻¹		
z	axial coordinate/m		

 $E_{\rm i}$, to a more oxidizing value, $E_{\rm f}$, and then back to the initial value. The value of *E* is therefore calculated at any time on the forward sweep using eq 9 and on the reverse sweep using eq 10, where *v* is the voltage scan rate and *t* is the time:

$$E_{\rm forward} = E_{\rm i} + vt \tag{9}$$

$$E_{\text{reverse}} = 2E_{\text{f}} + E_{\text{i}} - vt \tag{10}$$

When t > 0, z = 0, and $r \le r_e$, the boundary conditions are the Butler–Volmer conditions describing species A and C and conservation of mass for species B:

$$\begin{split} D_{\rm A} & \left(\frac{\partial c_{\rm A}}{\partial z}\right)_0 = k_{\rm A/B}^0 c_{\rm B,0} \exp\left(-\alpha_{\rm A/B} \frac{F(E - E_{\rm f,A/B}^{\ominus})}{RT}\right) - \\ & k_{\rm A/B}^0 c_{\rm A,0} \exp\left((1 - \alpha_{\rm A/B}) \frac{F(E - E_{\rm f,A/B}^{\ominus})}{RT}\right) \\ D_{\rm B} & \left(\frac{\partial c_{\rm B}}{\partial z}\right)_0 = -D_{\rm A} & \left(\frac{\partial c_{\rm A}}{\partial z}\right)_0 - D_{\rm C} & \left(\frac{\partial c_{\rm C}}{\partial z}\right)_0 \\ D_{\rm C} & \left(\frac{\partial c_{\rm C}}{\partial z}\right)_0 = -k_{\rm B/C}^0 c_{\rm C,0} \exp\left(-\alpha_{\rm B/C} \frac{F(E - E_{\rm f,B/C}^{\ominus})}{RT}\right) + \\ & k_{\rm B/C}^0 c_{\rm B,0} \exp\left((1 - \alpha_{\rm B/C}) \frac{F(E - E_{\rm f,B/C}^{\ominus})}{RT}\right) \end{split}$$

The remaining boundary conditions are:

$$t = 0, \text{ all } r, \text{ all } z$$

$$c_{A} = c_{A}^{*} \qquad c_{B} = 0 \qquad c_{C} = 0$$

$$t > 0, r > r_{e}, z = 0$$

$$\frac{\partial c_{A}}{\partial z} = 0 \qquad \frac{\partial c_{B}}{\partial z} = 0 \qquad \frac{\partial c_{C}}{\partial z} = 0$$

$$t > 0, \text{ all } r, z \to \infty$$

$$\frac{\partial c_{A}}{\partial z} = 0 \qquad \frac{\partial c_{B}}{\partial z} = 0 \qquad \frac{\partial c_{C}}{\partial z} = 0$$

$$t > 0, r \to \infty, \text{ all } z$$

$$\frac{\partial c_{A}}{\partial r} = 0 \qquad \frac{\partial c_{B}}{\partial r} = 0 \qquad \frac{\partial c_{C}}{\partial r} = 0$$

$$t > 0, r = 0, \text{ all } z$$

$$\frac{\partial c_{A}}{\partial r} = 0 \qquad \frac{\partial c_{B}}{\partial r} = 0 \qquad \frac{\partial c_{C}}{\partial r} = 0$$

2.2. Computational Details. The problem is generalized by means of a transformation into nondimensional form using a conventional set of normalized parameters (defined in Table 2).

The resulting set of normalized diffusion equations were published by Belding et al.⁷ The boundary conditions given in the previous sections are readily formulated in terms of normalized variables using the definitions given in Table 2. The problem is then discretized using the alternating direct implicit method^{20,21} (ADI) and solved numerically using the iterative Newton–Raphson scheme.²² The discretized spatial mesh is analogous to that reported by Gavaghan²³ (Figure 4) and converged to within 0.1 %. Each temporal increment is optimized during run time such that convergence to within 0.01 % is obtained. All programs were written in C++ and compiled using a Borland compiler. The simulations were run on a desktop personal computer with a processor speed of \approx 3 GHz. Approximately 20 min of central processing unit time were required to simulate a single voltammogram.

3. Experimental Section

3.1. Chemical Reagents. 1-Butyl-3-methylimidazolium tetrafluoroborate ($[C_4mim][BF_4]$) was kindly donated by Merck KGaA. DMPD (Aldrich, > 97 %), ferrocene (Aldrich, 98 %), tetrabutylammonium perchlorate (TBAP, Fluka, Puriss electrochemical grade, > 99.99 %), and acetonitrile (Fischer Scientific, dried and distilled, > 99.99 %) were used as received without further purification.

3.2. *Instrumental Section.* Electrochemical experiments were performed using a computer controlled μ -Autolab potentiostat (Eco-Chemie, Netherlands). A conventional two-electrode system was used, typically with a platinum electrode (10 μ m diameter) as the working electrode and a 0.3 mm diameter platinum wire as a quasi-reference electrode. The platinum microdisc working electrode was polished on soft lapping pads (Kemet Ltd., U.K.) using alumina powder (Buehler, IL) of (5.0, 1.0, and 0.3) μ m sizes. The electrode diameter was calibrated electrochemically by analyzing the steady-state voltammetry of a 2 mM solution of ferrocene in acetonitrile containing 0.1 M TBAP, with a diffusion coefficient for ferrocene of 2.3 $\cdot 10^{-5}$ cm² \cdot s⁻¹ at 293 K.²⁴

The electrodes were housed in a glass cell "T-cell" designed for investigating microsamples of ionic liquids under a controlled atmosphere²⁵ (Figure 5). RTILs are sensitive to water,^{25,26} and the presence of water can alter the viscosity of the ionic liquid and reduce the electrochemical window; therefore, the samples are purged under vacuum before voltammetry is carried out. In addition, this procedure removes electroactive oxygen from solution.²⁷ The working electrode was modified with a section of disposable micropipet tip to create a small cavity above the disk into which a drop (20 μ L) of ionic liquid was placed. DMPD was directly dissolved in [C₄mim][BF₄] at concentration of 20 mM. Prior to voltammetric scanning, the RTIL solution was purged under vacuum (Edwards high vacuum pump, Model ES 50) for about 90 min, which served to remove trace

Table 2	2.	Normalized	Parameters

normalized parameter	definition
J	$i/(Fc_{\rm A}^*D_{\rm A}r_{\rm e})$
K^0	$(k^0 r_{\rm e})/D_{\rm A}$
$K_{\rm comp}$	$(k_{\rm comp}c_{\rm A}^*r_{\rm e}^2)/D_{\rm A}$
$K_{\rm disp}$	$(k_{\rm disp}c_{\rm A}^*r_{\rm e}^2)/D_{\rm A}$
θ	$(F(E - E_{f,A/B}^{\Theta}))/(RT)$
R	$r/r_{\rm e}$
σ	$(r_e^2 F v)/(D_A RT)$
τ	$(D_{\rm A}t)/r_{\rm e}^2$
x	$c_{\rm X}/c_{\rm X}^*$
x_0	$c_{\rm X, 0}/c_{\rm X}^{*}$
Ζ	$z/r_{\rm e}$

atmospheric moisture naturally present in the RTIL. All experiments were performed inside a fume cupboard, in a thermostatted box (previously described by Evans et al.)¹² which also functioned as a Faraday cage. The temperature was maintained at (298 \pm 1.0) K.

3.3. Results and Discussion.

3.3.1. Chronoamperometry. Double potential step chronoamperometric transients at the platinum electrode (10 μ m diameter) were achieved using a sample time of 0.01 s. A solution of approximately 20 mM DMPD in [C₄mim][BF₄] was pretreated by holding the potential at a point of zero current for 20 s, after which the potential was stepped to a position after the oxidative peak for DMPD corresponding to transportcontrolled diffusion, and the current was measured for 5 s. The potential was then stepped to a position after the reductive peak for DMPD⁺ corresponding to transport-controlled diffusion, and the current response was measured for a further 5 s. To extract diffusion coefficients and solubility data from these transients, the first potential step was fitted using the nonlinear curve-fitting function in the software package Origin 7.0 (Microcal Software, Inc.) according to the following equations proposed by Shoup and Szabo:²⁸

$$i = -4nFD_{\rm DMPD}c_{\rm DMPD}r_{\rm e}f(\tau) \tag{11}$$

$$f(\tau) = 0.7854 + 0.4432\tau^{-1/2} + 0.2146 \exp(-0.3912\tau^{-1/2})$$
(12)

$$\tau = \frac{D_{\text{DMPD}}t}{r_{\text{e}}^2} \tag{13}$$

where *n* is the number of electrons transferred, *F* is the Faraday constant, D_{DMPD} is the diffusion coefficient, c_{DMPD} is the initial concentration of parent species, r_{e} is the radius of the disk electrode, *t* is the time, and $f(\tau)$ is a function describing the slope of the transient. The equations used in this approximation are sufficient to give D_{DMPD} and c_{DMPD} within an uncertainty of ± 0.6 %. The value of the parameter D_{DMPD^+} was elucidated by fitting the transient for the second potential step using simulation software developed by Klymenko et al.¹⁴

The fitted transient is shown in Figure 6 and the corresponding values of c_{DMPD} , D_{DMPD} , and D_{DMPD^+} recorded in Table 3. It is important to note the value of $c_{\text{DMPD}} = 19.5$ mM obtained from fitting the chronoamperometric transient to the Shoup and Szabo equation is in good agreement with the value of $c_{\text{DMPD}} \approx 20$ mM used in the experiment. As noted in other situations,¹⁸ despite structural similarity between DMPD and DMPD⁺, the discrepancy in diffusion coefficients is not insignificant. It is important to note that the potentials applied over the course of the double-step experiment are not sufficient to bring about the transfer of a second electron (DMPD⁺ – e⁻ = DMPD²⁺). The concentration of $D_{\text{DMPD}^{2+}}$ in solution is negligible, and consequently, comproportionation cannot occur: Figure 6 is independent of the effects of comproportionation regardless of the values of D_{DMPD^+} , and $D_{\text{DMPD}^{2+}}$.

3.3.2. Cyclic Voltammetry. The kinetic and thermodynamic parameters of the two redox couples were elucidated by means of cyclic voltammetry. The oxidation of DMPD was examined in [C₄mim][BF₄]. Cyclic voltammetry for the oxidation of 20 mM DMPD was carried out on a Pt microelectrode (diameter of 10 μ m) at scan rates ranging from 10 mV·s⁻¹ to 1 V·s⁻¹. The voltammograms are typically scanned from (0 to 1.5) V versus a Pt quasi-reference electrode. Two oxidation peaks were observed for scan rates of 10 mV·s⁻¹ to 1 V·s⁻¹, shown in Figure 7 at potentials of (0.4 and 0.9) V. These correspond,



Figure 4. Discretized grid used in the simulations (some points have been removed for clarity).



Figure 5. Cross section of the T-cell used to perform chronoamperometry.

respectively, to the oxidation of DMPD to DMPD⁺ and DMPD²⁺. Both electron transfers are chemically reversible as indicated by the presence of two back peaks. The reliability of the fitting procedure was ensured by fitting the voltammograms over a wide range of scan rates (two orders of magnitude). The fitted voltammograms are shown in Figure 7 and the corresponding data quoted in Table 3.

The onset of electrochemical reversibility occurs in RTILs at much lower values of the electrochemical rate constant compared to other nonaqueous solvents. For example, in a classic paper by Amatore et al.,²⁹ electrochemical rate constants as high as 4 cm·s⁻¹ could be distinguished voltammetrically in acetonitrile. The reason for this discrepancy lies in the high viscosity of ionic liquid solvents (e.g., $\mu_{[C_{4mim}][BF_4]} = 112$ mPa·s at 25 °C³⁰) compared to other nonaqueous solvents (e.g., $\mu_{acetonitrile} = 34$ mPa·s at 25 °C³¹). Mass transport is more likely to be rate-limiting in solvents of greater viscosity. The rate constant for the first electron transfer (DMPD – e⁻⁼ DMPD⁺) is electrochemically reversible on the voltammetric time scale



Figure 6. Chronoamperometric fit: -, experiment; O, simulation.

Table 3. Fitted Data for DMPD in [C₄mim][BF₄]

parameter	experimental value
$\alpha_{DMPD/DMPD^+}$	(0.5 ± 0.1)
$\alpha_{DMPD^+/DMPD^{2+}}$	(0.5 ± 0.1)
CDMPD	$(19.5 \pm 1.0) \text{ mM}$
$E_{\mathrm{f, DMPD/DMPD^+}}^{\Theta}$	$(0.205 \pm 0.01 \text{ V vs Pt})$
$E_{\rm f, DMPD^+/DMPD^{2+}}^{\Theta}$	(0.745 ± 0.01) V vs Pt
$k_{\rm DMPD/DMPD^+}^0$	$\geq 0.01 \text{ cm} \cdot \text{s}^{-1}$
$k^0_{\mathrm{DMPD}^+/\mathrm{DMPD}^{2+}}$	$(0.001 \pm 0.002) \text{ cm} \cdot \text{s}^{-1}$
D_{DMPD}	$(2.0 \pm 0.1) \cdot 10^{-7} \text{ cm}^2 \cdot \text{s}^{-1}$
$D_{ m DMPD^+}$	$(1.1 \pm 1) \cdot 10^{-7} \text{ cm}^2 \cdot \text{s}^{-1}$
$D_{ m DMPD^{2+}}$	$(7 \pm 0.1) \cdot 10^{-8} \text{ cm}^2 \cdot \text{s}^{-1}$
k _{comp}	$(250 \pm 150) \text{ dm}^3 \cdot \text{mol}^{-1} \cdot \text{s}^{-1}$
$\Delta G^{\dot{\Theta}}_{ m comp}$	$-52.1 \text{ kJ} \cdot \text{mol}^{-1}$

 $(k_{\text{DMPD/DMPD}^+}^0 \ge 0.01 \text{ cm} \cdot \text{s}^{-1})$, while the second (DMPD⁺ - e⁻ \Rightarrow DMPD²⁺) is quasi-reversible $(k_{\text{DMPD}^+/\text{DMPD}^{2+}}^0 = (0.001 \pm 0.002) \text{ cm} \cdot \text{s}^{-1})$.

The comproportionation of DMPD and DMPD²⁺ to form DMPD⁺ is thermodynamically downhill. Using eq 5 and the data in Table 3, $\Delta G^{\Theta} = -52.1 \text{ kJ} \cdot \text{mol}^{-1}$. However, from the Table 3, comproportionation is slow on the voltammetric time



Figure 7. Voltammetric fits: -, experiment; O, simulation ($k_{comp} = 250 \text{ dm}^3 \cdot \text{mol}^{-1} \cdot \text{s}^{-1}$). (a) 100 mV $\cdot \text{s}^{-1}$, (b) 100 mV $\cdot \text{s}^{-1}$, and (c) 10 mV $\cdot \text{s}^{-1}$.



Figure 8. Voltammetric fits: -, experiment; O, simulation ($k_{comp} = 0 \text{ dm}^3 \cdot \text{mol}^{-1} \cdot \text{s}^{-1}$). (a) 1000 mV $\cdot \text{s}^{-1}$, (b) 100 mV $\cdot \text{s}^{-1}$, and (c) 10 mV $\cdot \text{s}^{-1}$.

scale ($k_{\text{comp}} = 250 \text{ dm}^3 \cdot \text{mol}^{-1} \cdot \text{s}^{-1}$). When DMPD is oxidized, the lone electron pair on each nitrogen atom makes an increased contribution to the resonance within the aromatic ring: the double bond character of each of the C–N bonds increases. The rearrangement energy for the electron transfer is substantial and leads to a small value of the rate constant as laid out in Marcus theory.^{32,33}

In a recent paper⁷ it was shown that, in the limit of electrochemical reversibility, the presence of diffusionally controlled comproportionation can be discerned only at high scan rates when the diffusion coefficients of the species (D_{DMPD} , D_{DMPD^+} , and $D_{\text{DMPD}^{2+}}$) are significantly different such that $D_{\text{DMPD}^+}/D_{\text{DMPD}} > 1.5$ and $D_{\text{DMPD}^{2+}}/D_{\text{DMPD}} > 1.5$ or $D_{\text{DMPD}^+}/D_{\text{DMPD}} < 0.75$ and $D_{\text{DMPD}^{2+}}/D_{\text{DMPD}} < 0.75$. It is important to note that this conclusion refers to a comparison between $k_{\text{comp}} = 0$ dm³·mol⁻¹·s⁻¹ and *diffusion control*. For an arbitrary value of

 $k_{\rm comp}$ the aforementioned conditions for the observation of comproportionation are neither necessary nor sufficient. Figure 8 shows a comparison between the experimental voltammograms at a range of scan rates and the simulated transients corresponding to the data in Table 3 with $k_{\text{comp}} = 0$ $dm^3 \cdot mol^{-1} \cdot s^{-1}$. It is clear that, for the present system, the occurrence of comproportionation can be identified voltammetrically, and the discrepancy between $k_{\text{comp}} = 0$ $dm^3 \cdot mol^{-1} \cdot s^{-1}$ and $k_{comp} = 250 dm^3 \cdot mol^{-1} \cdot s^{-1}$ increases with decreasing scan rate. At low scan rates a steady-state voltammetric response is obtained, while at high scan rates comproportionation is outrun on the voltammetric time scale. Therefore, in each case, the presence of comproportionation cannot be discerned from the voltammetry. For DMPD in [C₄mim][BF₄], the effect of comproportionation is not expected to be apparent in the limit of very low scan rates.

4. Conclusion

In this paper, simulation studies have been used to show that comproportionation, in which DMPD reacts homogeneously with DMPD⁺ (DMPD + DMPD²⁺ \rightleftharpoons 2DMPD⁺), can be discerned for the two-electron oxidation of DMPD in the ionic liquid [C₄mim][BF₄]. In addition, the kinetic and thermodynamic parameters of the reaction were elucidated by means of double potential step chronoamperometry and cyclic voltammetry.

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