

ARTICLES

Intramolecular Energy Hopping in Polyphenylene Dendrimers with an Increasing Number of Peryleneimide Chromophores

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The photophysical properties and the excitation energy transfer (EET) process of a series of first generation dendrimers having a rigid tetrahedral core and different numbers of peryleneimide chromophores at the rim were investigated by steady-state and time-resolved fluorescence spectroscopy in toluene. It was shown that the model dendrimer containing only one chromophore decays monoexponentially with a time constant of 4.2 ns. In the multichromophoric compounds an additional long decay time of 7.4 ns assigned to an “excimer-like” state was observed, and its contribution increases with the number of chromophores. The time-resolved polarization measurements resulted in a nanosecond relaxation component due to the rotation of the entire dendritic structure. However, for the compounds having more than one chromophore, an additional fast anisotropy decay time of 100–200 ps was observed which corresponds to Förster type EET among the chromophores.

Introduction

Recently much attention has been devoted to the controlled synthesis and characterization of three-dimensional molecular systems known as dendrimers since their highly branched regular structures can be used as building blocks for the fabrication of target designed molecules.^{1,2} The three-dimensional shape of the dendrimers is determined by the interior building block (core) and the peripheral groups. In contrast to the other macromolecules, the size, shape, and functionality of the dendrimers can be defined at the molecular level. Application of the dendrimers have been suggested in several fields such as guest–host chemistry,³ analytical chemistry,² optoelectronics,⁴ catalysis,⁵ biology,⁶ and medicine.⁷ Dendrimers are also used to mimic the photophysical processes of the basic photosynthetic light-harvesting antennae systems.^{8–12}

Synthetic procedures¹³ and physical organic studies^{14–17} on dendrimer molecules with a polyphenylene core decorated with different numbers of peryleneimide chromophores on the surface have been reported. The combination of picosecond single photon timing (SPT) and femtosecond transient absorption,¹⁴ femtosecond fluorescence upconversion,^{15,16} single-molecule spectroscopy,^{12,17} and atomic force microscopy¹⁸ lead to an understanding of the time- and space-resolved behavior of these dendrimers. Particularly for the time-resolved properties of the peryleneimide dendrimers with a biphenyl core,^{14–16} it was shown that although the excited state decays mainly with the fluorescence decay time of about 4 ns for the single chromophoric system along with a 7.5 ns component for the

“dimerlike” state resulting from chromophore–chromophore interaction for the dendrimers having multiple chromophores, a number of other kinetic components are present on the ultrashort time scale.

However, the studies mentioned above are for dendrimers involving peryleneimide chromophores and polyphenylene arms but with a biphenyl core leading to a three-dimensional structure with a biconical shape.¹² It was shown that due to the conformational dynamics the peryleneimide chromophores are distributed unevenly in different environments at the rim of the dendrimer.¹²

To get a better control over the spatial distribution of chromophores and a better insight into the interaction in the excited state, it is desirable to have them distributed on a more spherical surface. Recently, this goal has been achieved with the synthesis of a series of first generation peryleneimide dendrimers with a rigid tetrahedral central core.¹⁹ The number of peryleneimide chromophores (**PI**, Chart 1) is varied from 1 to 4 (**G1R_x**; $x = 1–4$, **I–IV**, Chart 1) along the branches of the dendrimer in the direction of the corners of a tetrahedron due to the central sp^3 carbon of the core. In this way, it is possible to get a more controlled orientation of the polyphenylene building blocks around the central core. As a result, the possibility of energy transfer in multichromophoric systems can be studied more systematically.

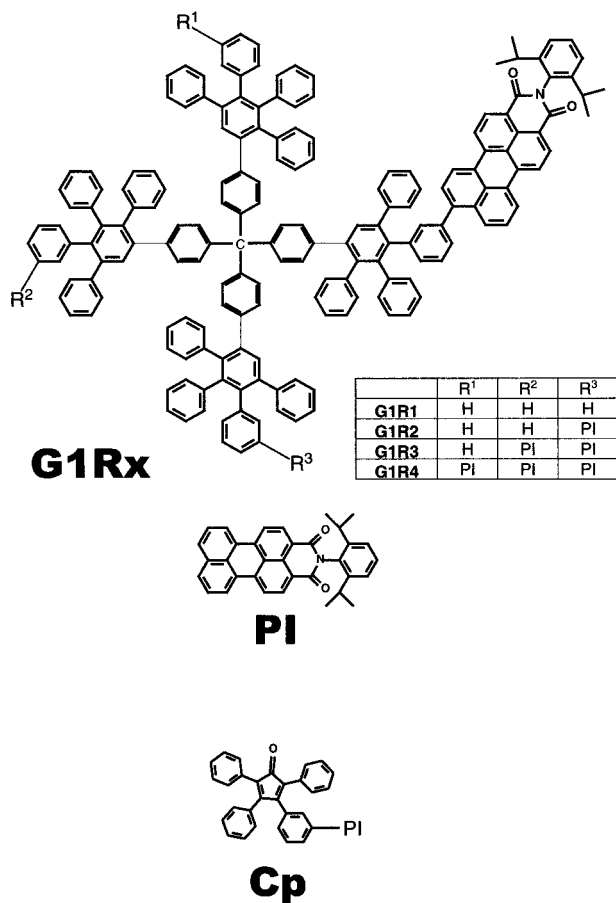
In this contribution, we report the solution phase photophysics of these dendrimers studied by steady-state and picosecond time-resolved spectroscopic techniques.

Experimental Section

Materials. The synthesis of these rigid dendrimers having a tetrahedral core decorated with a different number of perylene-

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CHART 1: Structures of the Molecules Used in This Study

imide chromophores on the polyphenylene rim utilizes the tetraphenylcyclopentadienone (**CP**, Chart 1) which carries the peryleneimide chromophore. While the symmetric **G1R4** is made via Diels–Alder cycloaddition of **CP** to a core molecule with four ethynyl functions, the “desymmetrized” species **G1R1**, **G1R2**, **G1R3** require a novel approach which is being published separately.¹⁹

The samples were dissolved in toluene (Aldrich) to have an optical density just below 0.1 in a 1 cm cell at the absorption maximum (490 nm) which corresponds to a concentration of $\sim 10^{-7}$ M.

Steady-State Measurements. Steady-state absorption and corrected fluorescence spectra were recorded with Lambda 40 (Perkin-Elmer) and SPEX spectrophotometers, respectively. The fluorescence quantum yields have been determined using Rhodamine 101 and Cresyl violet perchlorate as a reference.²⁰

Picosecond Time-Resolved Measurements. The fluorescence decay times have been determined by the single photon timing method using a laser and detection setup described previously.^{14,21} In brief, the second harmonic of a Ti:sapphire laser (Tsunami, Spectra Physics) has been used to excite the samples at 488 nm with a repetition rate of 4.09 MHz.¹⁴ The detection system²¹ consists of a subtractive double monochromator (9030DS, Scientech) and a microchannel plate photomultiplier (R3809U, Hamamatsu). A time-correlated single photon counting PC module (SPC 630, Picoquant GmbH) which has the two constant fraction discriminators (CFD), a time-to-amplitude converter (TAC), and an analog-to-digital converter (ADC) on board was used to obtain the fluorescence decay histograms in 4096 channels with time increments of 5 or 10

ps. The fluorescence decays have been recorded at three different orientations of the emission polarizer relative to the polarization plane of the excitation light, i.e., 54.7° (magic angle), 0° (I_{par}), and 90° (I_{perp}). The magic angle decays were analyzed globally with a time-resolved fluorescence analysis (TRFA) software.²² The quality of the fits has been judged by the fitting parameters such as χ^2 (< 1.2) and $Z\chi^2$ (< 3) as well as by the visual inspection of the residuals and autocorrelation function.²³ The anisotropy decay as calculated from eq 1 were fitted by a linear combination of exponentially decaying functions using the Levenberg–Marquard χ^2 minimization on the basis of eq 2, where N is the number of data points and P is the number of free parameters in the fit function. The anisotropy decay analysis was performed with the commercial ORIGIN 6.0 (Microcal Software Inc.) program, and the results were double-checked with the TRFA program²² which also takes pulse deconvolution into account.

$$r(t) = \frac{I_{\text{par}}(t) - I_{\text{perp}}(t)}{I_{\text{par}}(t) + 2I_{\text{perp}}(t)} \quad (1)$$

$$\chi^2 = \frac{\sum_{i=1}^N [y_i - f(x_i)]^2}{N - P} \quad (2)$$

Theoretical Model for Energy Hopping

Our theoretical modeling of the observed dynamics involves excitation energy hopping among $i = 2, 3$, or 4 identical chromophores, which, in the simplest approach possible, are perfectly arranged with in average equal distances and orientations toward each other. This allows to consider a single rate constant k_{cet} of energy transfer. From this assumption it follows that dynamics of incoherent energy hopping are described by the master equation

$$\frac{d\mathbf{P}}{dt} = \mathbf{K}\mathbf{P} \quad (3)$$

where \mathbf{P} is the vector of the time-dependent individual excitation probabilities

$$\mathbf{P}(t) = \begin{pmatrix} P_1(t) \\ \vdots \\ P_i(t) \end{pmatrix} \quad (4)$$

and \mathbf{K} is the matrix of rate constants

$$\mathbf{K} = \begin{pmatrix} -(i-1)k_{\text{cet}} - \tau_f^{-1} & \cdots & k_{\text{cet}} \\ \cdots & \cdots & \cdots \\ k_{\text{cet}} & \cdots & -(i-1)k_{\text{cet}} - \tau_f^{-1} \end{pmatrix} \quad (5)$$

where all off-diagonal elements are equal to k_{cet} and the diagonal elements are dependent on k_{cet} , the number of chromophores i , and the fluorescence lifetime τ_f of the chromophores. The fluorescence intensity decay $I_f(t)$ can hence be described by a single-exponential function with arbitrary amplitude α according to

$$I_f(t) = \sum_{i=1}^i \frac{\alpha}{i} P_i(t) = \alpha e^{-t/\tau_f} \quad (6)$$

which does therefore not allow one to determine k_{cet} .

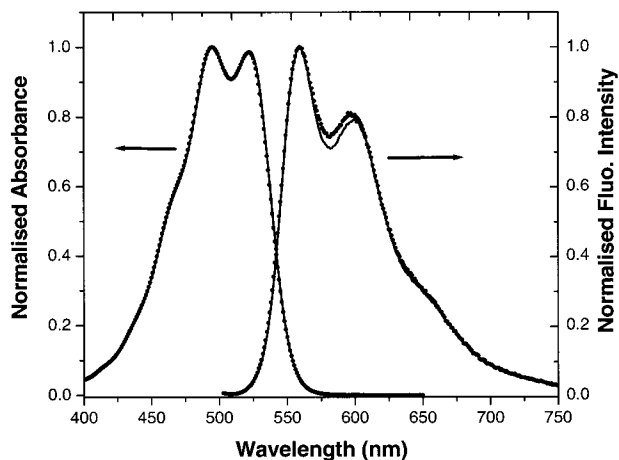


Figure 1. Steady-state absorption and emission spectra of the model compound **G1R1** (solid line) and **G1R4** (dots).

The observable quantity that reflects the energy hopping dynamics is the time-resolved fluorescence anisotropy $r(t)$.^{8,24–29} Since the anisotropy contribution from excited chromophores populated by energy transfer for an energy-hopping system can be assumed to be zero with an error of less than 2% for random chromophore orientation,^{24,25} the ordinary differential equation system (eq 7) for the time-dependent anisotropy residing on individual chromophores, where the rate constants for depolarization are k_{eet} and the inverse of the anisotropy relaxation time Θ_{rot}^{-1} for the rotation of the whole molecule, needs to be solved only for the anisotropy r_1 of the initially excited chromophore^{8,26–28} with the initial conditions that $r_1(0)$ equals the limiting anisotropy r_0 and that $r_{j \neq 1} = 0$ yielding eq 8.

$$\frac{d}{dt} \begin{pmatrix} r_1(t) \\ \dots \\ r_i(t) \end{pmatrix} = \begin{pmatrix} -(i-1)k_{\text{eet}} - \Theta_{\text{rot}}^{-1} & \dots & k_{\text{eet}} \\ \dots & \dots & \dots \\ k_{\text{eet}} & \dots & -(i-1)k_{\text{eet}} - \Theta_{\text{rot}}^{-1} \end{pmatrix} \quad (7)$$

This simple model thus predicts a double-exponential anisotropy

$$r(t) = r_1(t) = \frac{r_0}{i} [e^{-t/\Theta_{\text{rot}}} + (i-1)e^{-(ik_{\text{eet}} + \Theta_{\text{rot}}^{-1})t}] \quad (8)$$

decay with an amplitude ratio of $(i-1)$. In case of absent rotational movement ($\Theta_{\text{rot}}^{-1} = 0$) a constant term r_0/i gives the well-known leveling-off at long times.^{8,26–29}

Results

Steady-State Spectra. The steady-state absorption and fluorescence spectra of all the dendrimers **I–IV** which contain only the peryleneimide chromophore are very similar. Representative examples for the model compound with one chromophore, **G1R1**, and with maximum number of four chromophores, **G1R4**, in toluene are shown in Figure 1. The absorption spectra of all the dendritic compounds in toluene give two vibronic maxima at 495 and 520 nm. The fluorescence spectra shows maxima at 560 and 595 nm. The spectral properties and the high fluorescence quantum yields ($\phi_f = 0.95 \pm 0.05$) of all dendrimers **I–IV** are practically the same within error. Only the fluorescence full width at half-maxima (fwhm) increases slightly from 2680 cm^{-1} for **G1R1** to 2750 cm^{-1} for **G1R4** with the number of chromophores.

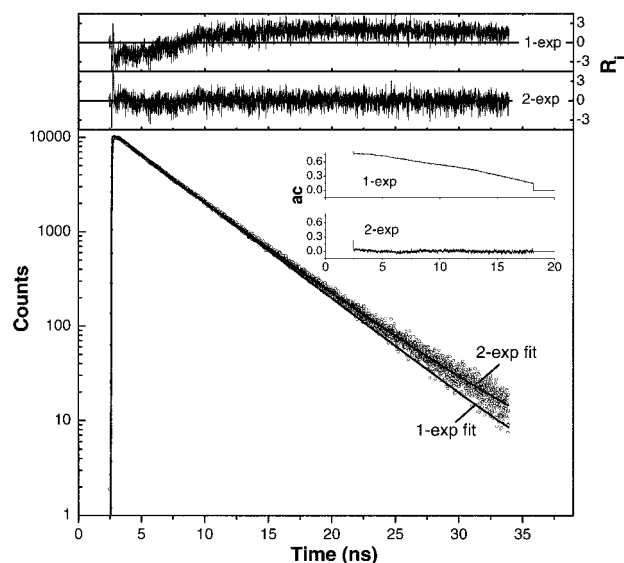


Figure 2. Time-resolved fluorescence decay of **G1R4** with mono- and biexponential fitting. The upper panel shows the distribution of weighted residuals (R_i), and in inset the autocorrelation (ac) functions are given in both cases.

TABLE 1: Fluorescence Decay Times (τ_i) and Associated Relative Amplitudes (α_i) for **G1R1–G1R4 Measured in Toluene at Room Temperature Using $\lambda_{\text{exc}} = 488 \text{ nm}^a$**

	τ_1 (ns)	τ_2 (ns)	$\lambda_{\text{flu}} = 600 \text{ nm}$		$\lambda_{\text{flu}} = 725 \text{ nm}$	
			α_1 (%)	α_2 (%)	α_1 (%)	α_2 (%)
I	4.25		100	0		
II	4.25	7.4 ^b	99.2	0.8	98.7	1.3
III	4.25	7.4 ^b	98.1	1.9	96.2	3.8
IV	4.25	7.4	96.0	4.0	93.7	6.3

^a The error of τ_1 is ± 0.05 ns and that of τ_2 is ± 0.6 ns. ^b Kept fixed for **II** and **III** as obtained from **IV** (see text for details).

Time-Resolved Fluorescence Measurements. To examine the properties of the fluorescent states for the dendrimers more closely, fluorescence decay times for all the dendrimers were determined in toluene by single-photon-timing detecting the emission under magic angle condition. A representative decay trace for **G1R4** is shown in Figure 2, and the respective decay parameters are collected in Table 1. The fluorescence intensity of the dendrimer having only one chromophore (**G1R1**) decays single exponentially with a decay time of 4.25 ± 0.05 ns. However, as the number of chromophores is increased in the dendrimer, a small contribution of an additional long decay component of 7.4 ± 0.6 ns is found essential to fit the experimental data as demonstrated by the distribution of weighted residuals and autocorrelation function for a mono- and biexponential fit in Figure 2. It has to be noted that the amplitude of this long decay component is however very small in **G1R2** and **G1R3**. Thus, to minimize the error in the fit procedure, the decay time of 7.4 ns obtained from **G1R4** was kept fixed in the analysis of **G1R2** and **G1R3** to compare better the corresponding amplitudes. In addition, the relative amplitude of the longer decay time is larger at the red edge of the fluorescence spectrum for all multichromophoric dendrimers as reported in Table 1 by the comparison of results obtained at 600 and 725 nm emission.

Time-Resolved Fluorescence Polarization Measurements. From time-resolved fluorescence polarization measurements, the anisotropy decay times (Θ_j) and the associated anisotropy (β_j) have been determined for all the compounds using eq 9.

For the dendrimer having only one chromophore (**G1R1**), a

TABLE 2: Fitting Parameters of the Fluorescence Anisotropy Decays Measured for G1R1–G1R4 in Toluene with $\lambda_{\text{exc}} = 488$ and $\lambda_{\text{flu}} = 600$ nm and Calculated Rotor Diameters (d_{rot}) and Average Chromophore–Chromophore Distances (d_{FRET})

compd	r_0	Θ_1 (ns)	Θ_2 (ns)	β_1	β_2	β_2/r_0 (%)	d_{rot} (Å) ^a	d_{FRET} (Å) ^b
I	0.38	0.95		0.38		0	23.3	
II	0.31	1.1	0.20	0.16	0.15	48	24.3	26.5
III	0.28	1.2	0.13	0.10	0.18	63	25.3	26.1
IV	0.24	1.3	0.11	0.08	0.16	66	25.9	26.5

^a Calculated from eqs 10 and 14. ^b Calculated from eqs 14 and 15 using i equal to 2, 3, and 4 for II–IV, respectively. See text for details.

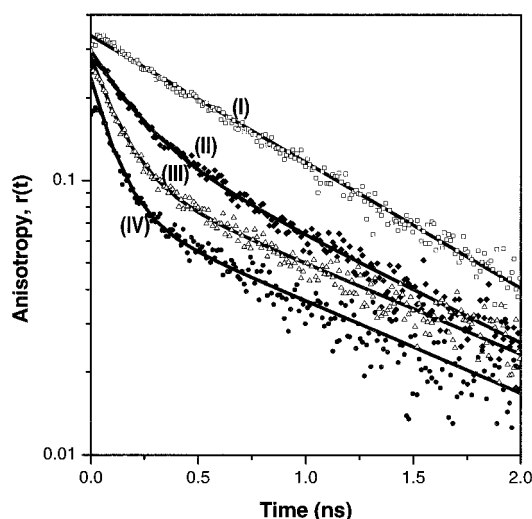


Figure 3. Time-resolved fluorescence anisotropy decays (scattered symbol) and fits (solid line) for the dendrimers **G1R1 (I)**, **G1R2 (II)**, **G1R3 (III)**, and **G1R4 (IV)** in toluene. The excitation and emission wavelengths are 488 and 600 nm, respectively.

monoexponential fit of the anisotropy decay function is sufficient which gives a relaxation time of $\Theta_1 = 950 \pm 30$ ps with $\beta_1 = r_0 = 0.38$ (Table 2). However, the anisotropy decay functions for the dendrimers having more than one chromophore (**G1R2**–**G1R4**) can only be fitted with two exponential decay functions (Figure 3). The long depolarization time constant is similar to that obtained for **G1R1**. The magnitude of this long time component increases with the number of chromophores from 1.1 ± 0.04 ns for **G1R2** to 1.3 ± 0.07 ns for **G1R4** while the value of the fast component (Θ_2) changes from 200 ± 30 ps for **G1R2** to 110 ± 20 ps in **G1R4** (Table 2).

$$r(t) = \sum_j \beta_j \exp(-t/\Theta_j) \quad r_0 = \sum_j \beta_j \quad (9)$$

Discussion

Fluorescence from a Localized and Excimer-like Species.

The steady-state properties of all the dendrimers (**I**–**IV**) are practically the same (Figure 1) with the exception of some broadening of the fluorescence spectra with increasing number of chromophores. This result is different from the observations made for the higher generation peryleneimide dendrimers having a biphenyl flexible core and multiple chromophores as reported elsewhere,¹⁴ where a number of properties, such as spectral shape, fluorescence maxima, and fluorescence quantum yield are influenced by the formation of an excimer-like species. From the small difference in the spectral width, the assumption of an excited-state excimer-like chromophore–chromophore interaction is possible but not conclusive.

Better insight into the extent of excimer-like emission is obtained from the fluorescence decays where only for the multichromophoric dendrimers a long decay component of 7.4 ns is observed along with the typical peryleneimide fluorescence decay time of 4.25 ns as obtained for the monochromophoric model compound **G1R1**. The attribution of this long time constant can be made to an “excimer-like” species as the decay time is similar to that reported for the higher generation dendrimers having a flexible biphenyl core.¹⁴ Further evidence for this assignment can be derived from the dependence of the amplitude connected with the 7.4 ns component on the number of chromophores and the dependence on the emission wavelength (λ_{flu}), respectively. As reported in Table 1, this amplitude is 0.8% for **G1R2** and increases to 4% for **G1R4** at $\lambda_{\text{flu}} = 600$ nm. This is reasonable as the probability of formation of the “excimer-like” entity increases as the number of chromophores increases in the dendrimer. By detection at $\lambda_{\text{flu}} = 725$ nm, the values of the amplitude for the 7.4 ns component increase to 1.3% for **G1R2** and to 6.3% for **G1R4**. The larger contribution of that component at longer emission wavelengths is also consistent with a red-shifted fluorescence from excimer-like entities.

Size of the Dendrimers. Fluorescence depolarization measurements of the monochromophoric **G1R1** give a single-exponential anisotropy decay time (Θ_1) of about 1 ns (Table 2) and can therefore be associated with the time constant of rotation (Θ_{rot}) of the whole molecule. The corresponding long polarization time constant (Θ_1) increases to 1.3 ns in **G1R4**. Thus, we can calculate the diameter of the spherical rotor (d_{rot})

$$d_{\text{rot}} = \sqrt[3]{\frac{6V}{\pi}} \quad (10)$$

using the Stokes–Einstein–Debye relation

$$\Theta_{\text{rot}} = \frac{V\eta}{\kappa T} \quad (11)$$

where Θ_{rot} is the rotational relaxation time obtained from Θ_1 according to eqs 8 and 9, V the hydrodynamic volume, η the solvent viscosity, κ the Boltzmann constant, and T the temperature.²³ On the basis of the usually applied assumption of a spherical volume, d_{rot} can be calculated from eq 10. Inserting the determined Θ_{rot} values of 1 to 1.3 ns from Table 2 in the above equation yields rotor diameters in the range of ~ 24 – 26 Å which are close to the values (22–24 Å) deduced from different conformations of **G1R4** by Molecular Modeling³⁰ assuming that the distance between the center of the chromophores and the sp^3 core carbon determines the length of a dendrimer branch. The increase in Θ_{rot} or d_{rot} reflects well the expected increase of the hydrodynamic volume with the introduction of more chromophores in the dendritic structure.

Intramolecular Excitation Energy Hopping. Energy transfer processes can be revealed by time-resolved anisotropy data. The large value for the limiting anisotropy (r_0) of **G1R1** confirms the parallel orientation of the absorption and emission transition dipole moment for a single chromophore.

In contrast to the fluorescence anisotropy decay of **G1R1**, which contains only one peryleneimide chromophore, a 100–200 ps anisotropy decay component is observed in the multichromophoric dendrimers. Therefore, this fast depolarization process can unambiguously be related to excitation energy hopping among the identical chromophores. The time scale of 100 ps for these processes suggests that the observed energy hopping occurs in terms of fluorescence resonance energy

transfer (FRET). Within the framework of the Förster formulation^{29,31–33} an effective interaction radius (R_0) can be calculated from the steady-state spectra and the fluorescence quantum yield of the donor chromophore (ϕ_D) with the following equations:

$$R_0^6 = 8.875 \times 10^{-5} \frac{\Gamma^2 \phi_D J}{n^4} \quad (12)$$

Here, Γ^2 equals 2/3 for the assumed random orientation of the chromophores, ϕ_D is the donor fluorescence quantum yield, n is the refractive index of the solvent (1.496 for toluene), and J is the spectral overlap integral defined by

$$J = \frac{\int F_D(\lambda) [\epsilon_A(\lambda)] \lambda^4 d\lambda}{\int F_D(\lambda) d\lambda} \quad (13)$$

where $\epsilon_A(\lambda)$ represents the molar extinction coefficient of the acceptor and $F_D(\lambda)$ denotes the donor fluorescence spectrum on a wavelength (λ) scale. The calculated values of $R_0 = 38 \text{ \AA}$ and $J = 2.5 \times 10^{14} \text{ M}^{-1} \text{ cm}^{-1} \text{ nm}^4$ using the spectral data for the monochromophoric **G1R1** model compound are on the typical order of magnitude where FRET has been reported in numerous publications.³⁴

Information about the rate constant of FRET (k_{FRET}) can be derived from the fast anisotropy decay time (Θ_2). However, to quantify k_{FRET} we have to take into account the possibility of multiple energy transfer channels in the case of a multichromophoric system containing equal chromophores among which efficient dipole–dipole interaction occurs. According to the energy hopping model described by eqs 3–8, the measured decay time Θ_2 can be related to $k_{\text{cet}} = k_{\text{FRET}}$ by combining eqs 8 and 9 to $\Theta_{\text{rot}} = \Theta_1$ and eq 14, where the value of i represents the number of chromophores fully interacting in both forward and backward direction.

$$k_{\text{FRET}} = \frac{1}{i\Theta_2} - \frac{1}{i\Theta_1} \quad (14)$$

If we take **G1R2** as the model system for energy hopping between two peryleneimide chromophores and further assume that the energy transfer occurs in both directions, then the rate constant (k_{FRET}) calculated from eq 14 with $i = 2$ results in $k_{\text{FRET}} = 2 \text{ ns}^{-1}$. Using this value for k_{FRET} we can calculate the expected anisotropy decay time (Θ_2) for the case of equally distributed and interacting chromophores in **G1R3** ($i = 3$) and **G1R4** ($i = 4$), which gives 143 and 112 ps, respectively. These results are in good agreement with the experimentally observed $\Theta_2 = 130 \text{ ps}$ of **G1R3** and $\Theta_2 = 110 \text{ ps}$ of **G1R4**, which indicates the validity of the proposed model of energy hopping among all chromophores. The theoretical energy hopping model can further be tested by the FRET anisotropy contribution β_2/r_0 (Table 2), which according to eq 8 should be equal to $(i - 1)/i$. While the values of the parameter β_2/r_0 for **G1R2** (48%) and **G1R3** (63%) agree quite well with the theoretically expected values 50% and 66%, respectively, β_2/r_0 for **G1R4** (66%) is somewhat smaller than the theoretical value of 75%. This might indicate that all the four chromophores in **G1R4** are not interacting equally; i.e., the fourth chromophore does not contribute in the energy-hopping process with the same efficiency as the other three. On the other hand, the comparatively small value of β_2/r_0 for **G1R4** can be due to the presence of the preformed dimer decaying with 7.4 ns (Table 1) and acting as a energy trap thus affecting the anisotropy decay.

On the basis of the same approximation of energy hopping in terms of FRET in the different dendrimers, i.e., full cross-talk among all chromophores in **G1R2** and **G1R3** but only cross-talk among mainly three chromophores in **G1R4**, the donor–acceptor distances (d_{FRET}) as calculated with the Förster eq 15^{29,31–33} using eqs 12–14 and the fluorescence decay time of the donor chromophore (τ_D) are listed in Table 2. All calculated values of d_{FRET} for the peryleneimide dendrimers are in the order of 26 Å which is in good agreement with the average distance between two chromophores in different conformations obtained from molecular modeling results. This agreement further substantiates the validity of the above proposed model for energy-hopping mechanism in the present dendrimers.

$$d_{\text{FRET}}^6 = \frac{R_0^6}{k_{\text{FRET}} \tau_D} \quad (15)$$

However, with an increasing number of chromophores r_0 decreases from 0.38 in **G1R1** to 0.31 in **G1R2** and to 0.24 in **G1R4**. This means that an additional fast depolarization processes on a time scale below the time resolution of 10 ps for the time-correlated single photon counting anisotropy experiments takes place. Most likely this is energy transfer between neighboring chromophores, which, due to their conformational freedom, can become very close lying in the order of 10 Å. In view of the determined overlap integral (eq 18) and the corresponding Förster radius (eq 12) the rate constant of FRET at such short distances is in the range of 1 ps⁻¹. Alternatively, the loss of limiting anisotropy may be related to direct excitation into a locally excited state of a preformed dimer which leads rapidly into a delocalized excimer-like state as has been discussed previously.¹⁵

Conclusions

The comparative steady-state and time-resolved study of a series of dendrimers with a rigid tetrahedral core containing a different number of chromophores at the rim reveals different intramolecular interactions between the chromophores. The fluorescence decays show emission from an “excimer-like” species with a time constant of 7.4 ns, beside the emission of the individual phenyl-substituted peryleneimide chromophores with 4.25 ns. A decrease in the limiting anisotropy for the dendrimers containing multiple peryleneimide chromophores indicates the presence of ultrafast energy transfer processes in these systems. The obtained rotational relaxation times varying between 0.95 and 1.3 ns are consistent with the size of the different molecules. Furthermore, an additional fast depolarization component with time constants in the range of 100 to 200 ps can be attributed to Förster resonance energy transfer among the attached chromophores, since it is observed only for the multichromophoric systems. This excitation energy hopping between the peryleneimide chromophores in **II–IV** can be understood by the calculated large value for the spectral overlap integral ($J = 2.5 \times 10^{14} \text{ M}^{-1} \text{ cm}^{-1} \text{ nm}^4$) associated with a sizable Förster radius of 3.8 nm. From the anisotropy analysis, it is also shown that while there is forward and backward energy transfer among all chromophores in the bichromophoric **G1R2** and the trichromophoric **G1R3**, the additional fourth chromophore in **G1R4** seems to interact to a little lesser extent with the others.

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