

## LETTERS

Anomalous Magnetic Behaviors of O<sub>2</sub>–CO<sub>2</sub> Mixed SolidA. Tohdoh and K. Kaneko<sup>\*,†</sup>*Physical Chemistry, Material Science, Graduate School of Natural Science and Technology, and Center for Frontier Electronics and Photonics, Chiba University, 1-33 Yayoi, Inage, Chiba 263, Japan**Received: April 19, 2001; In Final Form: June 26, 2001*

The temperature dependence of the magnetic susceptibility ( $\chi$ ) of O<sub>2</sub> and CO<sub>2</sub> mixture was measured as a function of the mole fraction of CO<sub>2</sub> over 1.9–100 K. The  $\chi$  value was remarkably depressed by the coexistence of CO<sub>2</sub>, although a change in  $\chi$  due to  $\alpha$ – $\beta$  and  $\beta$ – $\gamma$  phase transitions of solid O<sub>2</sub> was observed. The remarkable depression of the  $\chi$  of the CO<sub>2</sub>–O<sub>2</sub> mixture indicated the presence of new compound having the temperature independent-magnetic susceptibility. The relationship between the temperature-independent magnetism and the CO<sub>2</sub>–O<sub>2</sub> composition intensively suggests the formation of stoichiometric compounds of CO<sub>2</sub> and O<sub>2</sub>, such as clathrate compounds.

## Introduction

The ground state of O<sub>2</sub> is triplet, showing a typical paramagnetism in the gas phase. It is well-known that solid O<sub>2</sub> (*s*-O<sub>2</sub>) has three polymorphs of  $\alpha$ ,  $\beta$ , and  $\gamma$ . The  $\alpha$ – $\beta$  and  $\beta$ – $\gamma$  transitions show explicit anomalies in the magnetic susceptibility due to the corresponding phase transitions at 23.8 and 43.8 K.<sup>1,2</sup>

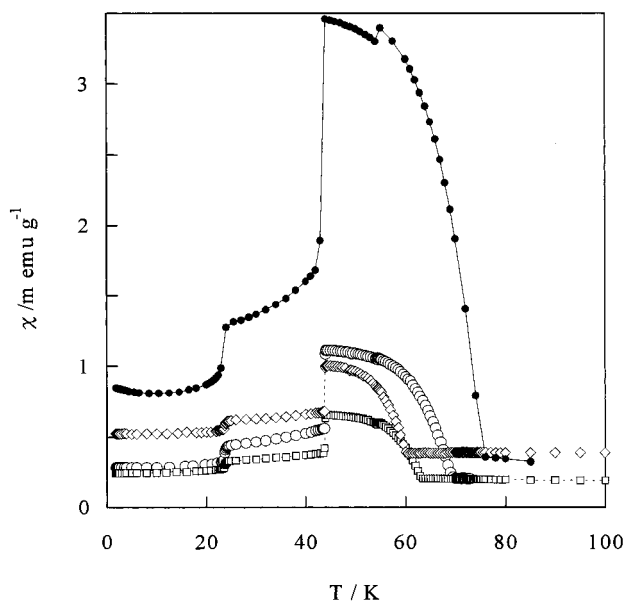
The binary mixture of simple gases such as O<sub>2</sub>, N<sub>2</sub>, CO<sub>2</sub>, and rare gases have been studied as model molecular systems. X-ray diffraction study gave the phase diagram of O<sub>2</sub> and N<sub>2</sub>, which indicates many phase transitions in the solid phase.<sup>3</sup> Raman study<sup>4</sup> suggested the presence of a new phase transition that was not observed by the X-ray diffraction approach. The fundamental understanding of the O<sub>2</sub> and N<sub>2</sub> mixture should be helpful to design a new air separation technology.

In the case of O<sub>2</sub>–Ar binary system, X-ray diffraction study showed the absence of solid-phase transition in the mole fraction range of Ar from 0.5 to 0.8. The mixture of Ar and O<sub>2</sub> of the above composition has a hexagonal close packed structure, which is completely different from the structure of pure solid

O<sub>2</sub>.<sup>5</sup> On the other hand, the heat capacity measurement of the Ar–O<sub>2</sub> system indicated the presence of a “hump”-type phase transition at 0.6 of Ar mole fraction and 14–25 K, not evidenced with X-ray diffraction.<sup>6</sup> Thus, the fundamental study on the O<sub>2</sub> based-molecular solid should be carried out not only by X-ray diffraction, but also by other techniques such as Raman spectroscopy.<sup>7–11</sup> Furthermore, magnetic susceptibility measurement is a potential method for O<sub>2</sub> based molecular solids because the magnetic interaction is quite sensitive to the local structure around an oxygen molecule. The structural study on the O<sub>2</sub> based-molecular solids should provide a new aspect for the intermolecular interaction.

As a CO<sub>2</sub> molecule has a great quadrupole moment,<sup>12</sup> the presence of CO<sub>2</sub> molecules should intensively affect the solid O<sub>2</sub> structure. The magnetic susceptibility of the solid O<sub>2</sub> can quite sensitively detect the local structural change of an oxygen molecule. The low-temperature magnetic susceptibility of the O<sub>2</sub>–CO<sub>2</sub> mixture was measured over the temperature range from 1.9 to 100 K in this study. The present work describes a remarkable effect of the coexistent CO<sub>2</sub> on the magnetism of solid O<sub>2</sub>.

<sup>†</sup> Center for Frontier Electronics and Photonics.



**Figure 1.** Temperature dependencies of the magnetic susceptibility of the solid O<sub>2</sub>-CO<sub>2</sub> system as a function of the CO<sub>2</sub> mole fraction  $x$ . ●:  $x = 0$ , ○:  $x = 0.2$ , □:  $x = 0.3$ , and ◇:  $x = 0.5$ .

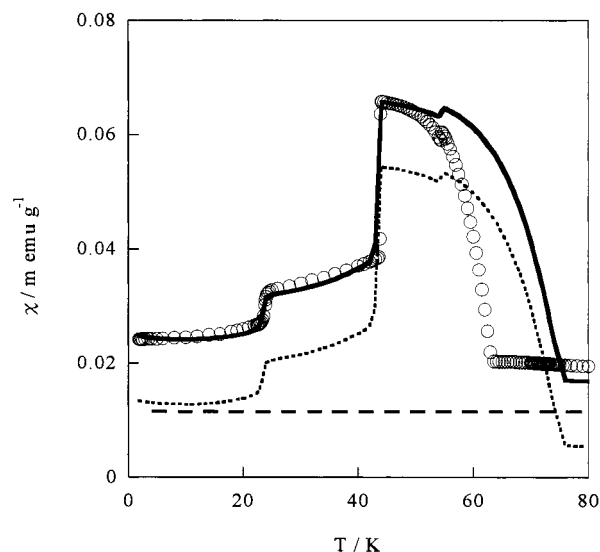
### Experimental Section

CO<sub>2</sub> and O<sub>2</sub> gases of high purity (99.99%) were used after repeated vacuum distillations. The mixed gas of CO<sub>2</sub> and O<sub>2</sub> were kept for 90 min at 303 K after their mixing. The gas mixture was introduced into an ESR quartz tube (length: 90 mm, diameter: 5 mm), and then the quartz tube was sealed. The composition of the gas mixture was determined through measurement of the O<sub>2</sub> and CO<sub>2</sub> partial pressures at 303 K. The mole fraction of CO<sub>2</sub> is designated as  $x$ . The quantity of O<sub>2</sub> was determined from the O<sub>2</sub> pressure and the volume of the sealed quartz tube. The magnetic susceptibility,  $\chi$ , was measured with a SQUID magnetometer system MPMSR2 (Quantum Design) over the temperature range from 1.9 to 100 K at a magnetic field of 1T. The magnetic susceptibility,  $\chi$ , of the quartz tube was subtracted.

### Results and Discussion

Temperature dependencies of the magnetic susceptibility  $\chi$  of the O<sub>2</sub>-CO<sub>2</sub> system are shown in Figure 1 as a function of mole fraction  $x$  of CO<sub>2</sub>. Pure O<sub>2</sub> shows explicit anomalies at the phase transition temperatures of 23.9 K ( $\alpha$ - $\beta$ ), 43.8 K ( $\beta$ - $\gamma$ ), and 54.4 K ( $\gamma$ -liquid), which agree with the literature values.<sup>1,2</sup> The anomalies due to  $\alpha$ - $\beta$  and  $\beta$ - $\gamma$  transitions are preserved at the same temperatures, even in the O<sub>2</sub>-CO<sub>2</sub> system, although the  $\chi$  is remarkably depressed by the coexistence of CO<sub>2</sub>. As the CO<sub>2</sub> molecule has no spin, the observed changes in the magnetism stem from the change in the O<sub>2</sub>-O<sub>2</sub> and/or O<sub>2</sub>-CO<sub>2</sub> intermolecular structure. In this work, the magnetic susceptibility data of liquid and gaseous O<sub>2</sub> above 54.4 K are not sufficiently discussed.

The magnetic susceptibility depression is not proportional to the CO<sub>2</sub> mole fraction  $x$ . Therefore, we assumed that the coexistent CO<sub>2</sub> changes the magnetic interaction between O<sub>2</sub> molecules through the formation of a weakly interacted compound such as clathrate compound. We introduce the temperature-independent magnetic susceptibility ( $\chi_{ti}$ ) due to the above compound formation. As the O<sub>2</sub>-CO<sub>2</sub> system should consist of pure O<sub>2</sub> and the O<sub>2</sub>-CO<sub>2</sub> compound, the observed magnetic



**Figure 2.** Fitting procedure of the magnetic susceptibility of the O<sub>2</sub>-CO<sub>2</sub> system of  $x = 0.3$ . ○: Experimental data, — — —:  $\chi_{ti}$ , - - -:  $b_0\chi_O$ , and - · - : ( $\chi_{ti} + b_0\chi_O$ ).

susceptibility  $\chi_{ob}$  is expressed by  $\chi_{ti}$  and the magnetic susceptibility of pure O<sub>2</sub>,  $\chi_O$ , as given by

$$\chi_{ob}(T) = \chi_{ti} + b_0 \cdot \chi_O(T) \quad (1)$$

Here  $\chi_{ti}$  and the constant  $b_0$  are determined by the fitting procedure, as shown in Figure 2. The fitting over the temperature range from 1.9 to 50 K is well done. Nevertheless, the calculated  $\chi_{ob}(T)$  using eq 1 deviates above 50 K, and the deviation becomes significant with increasing temperature. This is because the O<sub>2</sub> partial pressure of the O<sub>2</sub>-CO<sub>2</sub> system is different from that of pure O<sub>2</sub>, shifting the boiling temperature of liquid O<sub>2</sub>. Figure 3 shows  $\chi_{ti}$  and  $b_0$  with  $x$ . The  $\chi_{ti}$  increases with the increase of  $x$ . Hence, O<sub>2</sub> molecules must form a compound with CO<sub>2</sub> molecules, as suggested above. On the other hand, the  $b_0$  decreases until  $x = 0.2$  and oscillates with the increase of  $x$ .

The above interesting data must be explained by the constant  $b_0$ , which indicates the contribution of pure O<sub>2</sub> in the O<sub>2</sub>-CO<sub>2</sub> system. We can presume that the pure O<sub>2</sub>, which is not associated with the compound formation with CO<sub>2</sub>, shows the same magnetism as the pure solid O<sub>2</sub>. Then, the contribution of the pure O<sub>2</sub> and the component O<sub>2</sub> in the O<sub>2</sub>-CO<sub>2</sub> compound to the observed magnetism can be separately determined as follows. The total weight of O<sub>2</sub> is described by the sum of the weight of pure O<sub>2</sub> and that of the component O<sub>2</sub> in the O<sub>2</sub>-CO<sub>2</sub> compound, as given by eqs 2 and 3

$$m_O = m_{O,p} + m_{O,c} \quad (2)$$

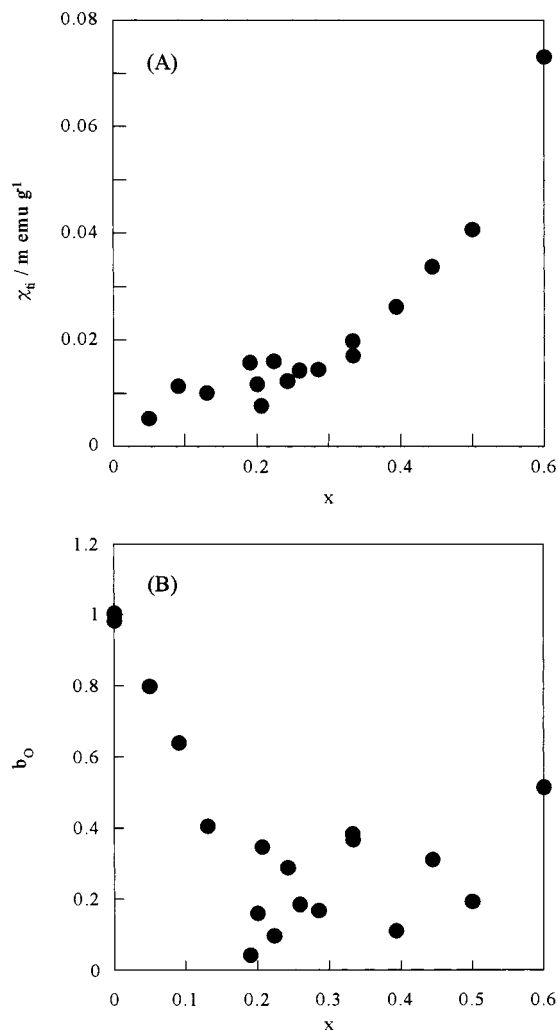
$$m_{O,p} = b_0 \cdot m_O \quad (3)$$

Here,  $m_O$  is the total weight of O<sub>2</sub> in the sample cell,  $m_{O,p}$  is the weight of the pure O<sub>2</sub>, and  $m_{O,c}$  is the weight of the component O<sub>2</sub>. Then,  $b_0$  and  $m_O$  can express  $m_{O,c}$

$$m_{O,c} = (1 - b_0) \cdot m_O \quad (4)$$

The mole fraction  $x_{comp}$  of CO<sub>2</sub> in the O<sub>2</sub>-CO<sub>2</sub> compound is described by  $m_{O,c}$  and  $m_{CO}$ , as shown by eq 5

$$x_{comp} = \frac{m_{CO}}{m_{O,c} + m_{CO}} = \frac{m_{CO}}{(1 - b_0) \cdot m_O + m_{CO}} = \frac{x}{b_0(x - 1) + 1} \quad (5)$$

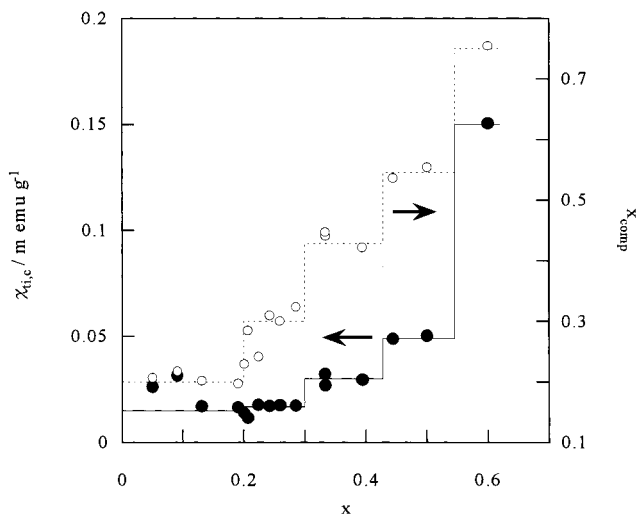


**Figure 3.** Changes of  $\chi_{ti}$  and  $b_O$  with  $x$ . (A):  $\chi_{ti}$  vs  $x$  and (B):  $b_O$  vs  $x$ .

Here,  $m_{CO}$  is the weight of  $CO_2$  in the sample cell. The temperature independent magnetism should stem from the  $O_2$ – $CO_2$  compound. Therefore, the observed  $\chi_{ti}$  must be recalculated using the weight of  $O_2$  in the  $O_2$ – $CO_2$  compound to get the correct magnetic susceptibility  $\chi_{ti,c}$  per unit mass of the component  $O_2$  in the  $O_2$ – $CO_2$  compound with eq 6

$$\chi_{ti,c} = \chi_{ti} \cdot \frac{m_{O_2}}{m_{O_2,c}} = \chi_{ti} \cdot \frac{1}{1 - b_O} \quad (6)$$

Thus, the variations of  $\chi_{ti,c}$  and  $x_{comp}$  with  $x$  are obtained, as shown in Figure 4. There are clear stepwise relationships for both of  $\chi_{ti,c}$  and  $x_{comp}$ . Although the relationship between  $b_O$



**Figure 4.** Changes of the temperature independent magnetism  $\chi_{ti,c}$  of  $O_2$  and the  $CO_2$  mole fraction,  $x_{comp}$ , of the  $O_2$ – $CO_2$  compound with the mole fraction  $x$  of  $CO_2$ .

and  $x$  above  $x = 0.2$  is complex, simple  $\chi_{ti,c}$  versus  $x$  and  $x_{comp}$  versus  $x$  relations can be obtained by the separation of  $O_2$  into pure  $O_2$  and the component  $O_2$ . So far, we cannot determine precisely the step position in Figure 4. However, these stepwise changes strongly suggest the formation of the stoichiometric compound of  $O_2$  and  $CO_2$  whose composition depends on the  $CO_2$  content. Unfortunately, only magnetic susceptibility data cannot determine directly the structure of the  $O_2$ – $CO_2$  compound showing the temperature-independent magnetism, although the temperature-independent magnetism suggests the antiferromagnetic interaction between  $O_2$  spins.

**Acknowledgment.** This work was funded by the Grant in-Aid for Scientific Research, Grant B from the Japanese Government.

#### References and Notes

- (1) DeFotis, G. C. *Phys. Rev.* **1981**, *B23*, 4717.
- (2) Kanda, E.; Haseda, T.; Otsubo, A. *Physica* **1954**, *20*, 131.
- (3) Barrett, C. S.; Meyer, L.; Greer, S. C.; Wasserman, J. *J. Chem. Phys.* **1968**, *48*, 2670.
- (4) Baer, B. J.; Nicol, M. *J. Phys. Chem.* **1990**, *94*, 1073.
- (5) Barrett, C. S.; Jordan, T. H.; Meyer, L. *J. Chem. Phys.* **1969**, *51*, 2941.
- (6) Xie, J.; Enderle, M.; Knorr, K. *Phys. Rev.* **1997**, *B55*, 8194.
- (7) Damde, K.; Jodl, H.-J. *J. Low Temp. Phys.* **1998**, *111*, 327.
- (8) Pace, E. L.; Bivens, R. L. *J. Chem. Phys.* **1970**, *53*, 748.
- (9) Barrett, C. S.; Meyer, L. *J. Chem. Phys.* **1965**, *42*, 107.
- (10) Klee, H.; Knorr, K. *Phys. Rev.* **1990**, *B42*, 3152.
- (11) Jin, L.; Knorr, K. *Phys. Rev.* **1993**, *B47*, 14 142.
- (12) Stone, A. J. *The Theory of Intermolecular Forces*; Clarendon Press: Oxford, 1996, p 14.