# CASSCF and CAS $+1+2$ Studies on the Potential Energy Surface and the Rate Constants for the Reactions between $\mathrm{CH}_{\mathbf{2}}$ and $\mathrm{O}_{2}$ 

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Received: November 9, 2001; In Final Form: January 16, 2002


#### Abstract

CAS $(14,12) /$ cc-pvdz calculations are reported for the reaction of ${ }^{3} \mathrm{CH}_{2}+{ }^{3} \mathrm{O}_{2} \rightarrow$ products. On the singlet potential energy surface, a transition state has been located with an energy barrier of $1.65 \mathrm{kcal} / \mathrm{mol}$, which is in good agreement with the experimental estimation of $1.0-1.5 \mathrm{kcal} / \mathrm{mol}$. The rearrangement and metathesis of the singlet intermediates have been also investigated at the same level of theory. For the triplet case, the formation of $\mathrm{CH}_{2} \mathrm{OO}$ has an energy barrier of $5.79 \mathrm{kcal} / \mathrm{mol}$, and the formed triplet $\mathrm{CH}_{2} \mathrm{OO}$ could be further decomposed into $\mathrm{CH}_{2} \mathrm{O}+\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with an energy barrier of $2.92 \mathrm{kcal} / \mathrm{mol}$. The geometries of some key points have been relocated at the $\mathrm{CAS}(8,6)+1+2 / \mathrm{cc}-\mathrm{pvdz}$ level of theory for comparison. The present theoretical results for the total reaction rates, at the $\operatorname{CAS}(8,6)+1+2 / c c-p v d z$ level, can be expressed by the three-parameter expression: $k(T)=4.273 \times 10^{-18} T^{2.245} \exp (-185 / T)$ within $\pm 5 \%$ error at the temperature range $295-2600 \mathrm{~K}$.


## I. Introduction

Triplet methylene $\left({ }^{3} \mathrm{CH}_{2}\right)$ is a radical, found in chemical processes such as the combustion of acetylene. Direct measurements of the rate constants for the reactions of ${ }^{3} \mathrm{CH}_{2}$ with some common species as $\mathrm{O}_{2}, \mathrm{NO}$, and $\mathrm{C}_{2} \mathrm{H}_{2}$ have been made within the past 20 years. ${ }^{1-6}$ As for the titled reaction, a possible reaction mechanism with four intermediates was suggested by Alvarez et al. ${ }^{7}$ (see Figure 1). Four distinct isomeric species exist as potential minima on the $\mathrm{CH}_{2} \mathrm{O}_{2}$ hypersurface as shown in Figure 1. Ab initio calculations on the energies and geometries of the low-lying electronic states of the species I-IV have been reviewed. ${ }^{8}$ Moreover, energy barriers to some of the interconversions between species I-IV and barriers for their decomposition have also been calculated previously. ${ }^{9-18}$

Recently, SuH . et al. ${ }^{19}$ have studied the reactions of ${ }^{3} \mathrm{CH}_{2}$ and ${ }^{1} \mathrm{CH}_{2}$ with $\mathrm{O}_{2}$ by time-resolved FTIR spectroscopy, which is mainly focused on the vibrational populations of the reaction products. The branching ratio of the two main products, $\mathrm{CO} /$ $\mathrm{CO}_{2}$ was reported to be 4 by Dombrowsky et al. ${ }^{6}$ and near 1 by Alvarez et al. ${ }^{7}$ Alvarez et al. also reported $16 \%$ yield of $\mathrm{H}_{2} \mathrm{CO}$ and indirectly estimated $30 \%$ yield of OH .

The rate constant of the reaction of ${ }^{3} \mathrm{CH}_{2}+{ }^{3} \mathrm{O}_{2}$ has been measured to be $(3.2 \pm 0.3) \times 10^{-12} \mathrm{~cm}^{3}$ molecule ${ }^{-1} \mathrm{~s}^{-1} .{ }^{2-4} \mathrm{~A}$ slight temperature dependence of the rate constant was measured over the ranges $233-433^{4}$ and $295-600 \mathrm{k}$, ${ }^{5}$ suggesting the presence of a small barrier $\left(1.0-1.5^{4}\right.$ or $\left.1.5 \pm 0.3^{5} \mathrm{kcal} / \mathrm{mol}\right)$ to the formation of the initial $\mathrm{CH}_{2} \mathrm{OO}$ adduct (I) or to its subsequent isomerization or fragmentation. Although several theoretical study papers have been published on the interconversions between species $\mathbf{I}-\mathbf{I V}$; however, no theoretical data has been reported for the formation process of $\mathrm{CH}_{2} \mathrm{OO}(\mathbf{I})$. Therefore, in the present paper, ab initio calculations based on CASSCF wave functions have been performed to explore the potential energy surfaces for the reactions between ${ }^{3} \mathrm{CH}_{2}$ and ${ }^{3} \mathrm{O}_{2}$.

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Figure 1. The possible reaction paths suggested by Alvarez et al. ${ }^{7}$

## II. Computational Details

All calculations employ the Dunning, ${ }^{20}$ cc-pvdz basis set. Two levels of complete active space self-consistent field ${ }^{21-23}$ (CASSCF) calculations were done. The larger CASSCF reference wave function is a nearly full valence CAS(except 2 s orbitals of O atoms), which consists of 14 active electrons and 12 active orbitals (denoted as $\operatorname{CAS}(14,12)$ ). The geometric locations and vibrational analysis were mainly using CAS(14,12) method. It was not feasible to perform multireference configuration interaction (MR-CI) calculations using this reference wave function. For the MR-CI calculations, the reference wave function was reduced to a six-orbital, eight-electron active space, denoted as $\operatorname{CAS}(8,6)$. The internally contracted MR-CI calculations ${ }^{24,25}$ include all single and double excitations relative to the $\operatorname{CAS}(8,6)$ reference wave function. These calculations will be denoted as $\mathrm{CAS}+1+2$. For the key points, $\mathrm{CAS}+1+2$ calculations have been employed to locate the geometric parameters for comparison. The CAS and CAS $+1+2$ calculations were done using the MOLPRO ${ }^{26}$ program.

The rate constants were calculated with conventional transition state theory using the POLYRATE $8.2^{27}$ program. The effect of replacing the harmonic oscillator partition function for the out-of-plane mode with a hindered rotation partition function ${ }^{28,29}$ was also investigated.

## III. Results and Discussions

1. Singlet Surface. On this surface, the extensive search revealed that there are two different approaching modes: (1)


Figure 2. The optimized geometries for the stationary points on the singlet energy surface (the values in parentheses are obtained from CAS$(8,6)+1+2 /$ cc-pvdz calculations $)$.
all atoms on the same plane with Cs symmetry and (2) twisted $\mathrm{CH}_{2}$ group, but also with $C_{s}$ symmetry. The geometries of all the possible stationary points obtained on this surface are depicted in Figure 2. Comparing the geometries optimized at different methods, the largest difference $(0.04 \AA$ ) could be found for the distance of $\mathrm{C}-\mathrm{O}$ in TS1 and all others are less than $0.02 \AA$. The differences of angles are less than $4^{\circ}$. Therefore, CAS $(14,12) /$ cc-pvdz is good enough to locate the geometries. The distance of $\mathrm{C}-\mathrm{O}$ is quite large ( $2.253 \AA$ ) in TS1, but it is $1.281 \AA$ in ${ }^{1} \mathrm{CH}_{2} \mathrm{OO}(\mathbf{I})$, which indicates that $\mathrm{C}-\mathrm{O}$ bond in ${ }^{1}$ $\mathrm{CH}_{2} \mathrm{OO}$ exhibits some double-bond character. For twisted case, a stationary point has been found along the reaction path, but vibrational analysis shows that such species is a second saddle point $(\mathrm{C}-\mathrm{O}=2.082 \AA)$, i.e., there are two imaginary frequencies $\left(-333.6 \mathrm{i},-212.8 \mathrm{i} \mathrm{cm}^{-1}\right)$. We changed the geometric parameters slightly along the first imaginary mode, a saddle point $(\mathrm{C}-\mathrm{O}=1.419 \AA)$ has also been found, which corresponds to the transition state of the rotation of $\mathrm{CH}_{2}$ group. The rotation barrier is quite high because the $\mathrm{C}-\mathrm{O}$ bond in ${ }^{1} \mathrm{CH}_{2} \mathrm{OO}$ is of some double-bond character. However, if it is shifted slightly along the second imaginary mode, it leads to a true transition state (TS1). Therefore, TS1 is the only true transition state on the singlet potential energy surface. The energy barrier is only $1.67 \mathrm{kcal} / \mathrm{mol}$ at $\operatorname{CAS}(14,12) / \mathrm{cc}-\mathrm{pvdz}$ level as listed in Table 1. Our best estimate one with zero-point energy correction is $0.7 \mathrm{kcal} / \mathrm{mol}$ at $\operatorname{CAS}(8,6)+1+2 / \mathrm{cc}-\mathrm{pvdz}$ level, which is quite close to the experimental estimation $(1.0-1.5 \mathrm{kcal} / \mathrm{mol})$. $^{4,5}$

The IRC curve at $\operatorname{CAS}(14,12) /$ cc-pvdz level in Figure 3a shows that before TS1 the energy potential curve is very flat, and afterward it becomes very steep after TS1. Figure 4a plots
the population changes of the orbitals $(11,12,13$, and 14$)$ along reaction coordinate $s$ on singlet energy surface, from which one can realize that when $s<-0.75$, the occupation of orbitals is the main reason keeping the curve flat. However near peroxymethylene ( $s \approx 1.6$ ), the occupation of numbers 11 and 12 are close to 2.0 and those of numbers 13 and 14 are approaching 0.0 , which means singlet peroxymethylene is in fact not a diradical, but some sort of zwitterions. In our MP2/cc-pvdz calculations both geometries have been obtained, the diradical one has higher energy than that of zwitterion. On the other hand, only a zwitterion one was located at $\operatorname{CAS}(14,12) /$ cc-pvdz level of theory. From Figure 4b, one can see that triplet peroxymethylene is a true diradical.

Two-dimensional plot of the potential surface near the TS1 at CAS(14,12)/cc-pvdz level of theory is shown in Figure 5, from which one can realize that there is a maximum near $\mathrm{C}-\mathrm{O}$ $=2.25 \AA$ and another at $\mathrm{O}-\mathrm{O}=1.22 \AA$, and a minimum is located near $\mathrm{C}-\mathrm{O}=1.3 \AA$ and another at $\mathrm{O}-\mathrm{O}=1.35 \AA$. In fact, the global optimization leads to a true minimum at $\mathrm{C}-\mathrm{O}$ $=1.28$ and $\mathrm{O}-\mathrm{O}=1.36 \AA$. When $\mathrm{C}-\mathrm{O}$ bond is kept near 1.3 $\AA$ and $\mathrm{O}-\mathrm{O}$ bond is increased, the curve shows that the energy is always increased, which indicates that this minimum ${ }^{1} \mathrm{CH}_{2-}$ OO is not related to $\mathrm{CH}_{2} \mathrm{O}+\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ but $\mathrm{CH}_{2} \mathrm{O}+\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ since the energy of the first one is lower than ${ }^{1} \mathrm{CH}_{2} \mathrm{OO}$. Actually, this minimum can also lead to a three-membered ring intermediate $\mathrm{CH}_{2} \mathrm{OO}(\mathrm{II})$ via a transition state TS2 (see Figure 2). The energy barrier for this process is $19.7 \mathrm{kcal} / \mathrm{mol}$ at $\mathrm{CAS}(14,12) / \mathrm{cc}-\mathrm{pvdz}$ level(close to $21.5 \mathrm{kcal} / \mathrm{mol}$ in ref 11). Intermediate $\mathrm{CH}_{2} \mathrm{OO}-$ (III) seems less stable than $\mathrm{CH}_{2} \mathrm{OO}(\mathbf{I I})$ as shown in Table 2,

TABLE 1: Total Energies (au) and the Relative Energies $(\mathbf{k c a l} / \mathrm{mol})^{a}$

|  | CAS $(14,12)$ | CAS $+1+2 \mathrm{OPT}^{\text {a }}$ |
| :---: | :---: | :---: |
| ${ }^{3} \mathrm{CH}_{2}+{ }^{3} \mathrm{O}_{2}$ | $\begin{aligned} & -188.669414 \\ & (0.0) \end{aligned}$ | $\begin{aligned} & -188.974803 \\ & (0.0) \end{aligned}$ |
| TS1 | $\begin{aligned} & -188.666790 \\ & (1.64) \end{aligned}$ | $\begin{aligned} & -188.974634 \\ & (0.11) \end{aligned}$ |
| ${ }^{1} \mathrm{CH}_{2} \mathrm{OO}\left({ }^{\mathbf{1}} \mathbf{I}\right)$ | $\begin{aligned} & -188.756231 \\ & (-54.38) \end{aligned}$ | $\begin{aligned} & -188.073362 \\ & (-60.85) \end{aligned}$ |
| TS2 | $\begin{aligned} & -188.724873 \\ & (-34.80) \end{aligned}$ |  |
| ${ }^{1} \mathrm{CH}_{2} \mathrm{OO}($ III $)$ | $\begin{aligned} & -188.798993 \\ & (-81.31) \end{aligned}$ |  |
| TS3 | $\begin{aligned} & -188.765728 \\ & (-60.50) \end{aligned}$ |  |
| ${ }^{1} \mathrm{CH}_{2} \mathrm{OO}$ (III) | $\begin{aligned} & -188.781949 \\ & (-70.62) \end{aligned}$ |  |
| TS4 | $\begin{aligned} & -188.778748 \\ & (-68.61) \end{aligned}$ |  |
| ${ }^{1} \mathrm{CH}_{2} \mathrm{OO}(\mathbf{I V})$ <br> trans | $\begin{aligned} & -188.953631 \\ & (-178.35) \end{aligned}$ |  |
| TS5 | $\begin{aligned} & -188.927954 \\ & (-162.24) \end{aligned}$ |  |
| $\begin{aligned} & { }^{1} \mathrm{CH}_{2} \mathrm{OO}(\mathbf{I V}) \\ & \text { cis } \\ & \text { TS6 } \end{aligned}$ | $\begin{aligned} & -188.944812 \\ & (-172.81) \\ & -188.836950 \\ & (-105.13) \end{aligned}$ |  |
| $\mathrm{CO}+\mathrm{H}_{2} \mathrm{O}$ | $\begin{aligned} & -188.956243 \\ & (-179.99) \end{aligned}$ |  |
| TS7 | $\begin{aligned} & -188.834879 \\ & (-103.83) \end{aligned}$ |  |
| $\mathrm{CO}_{2}+\mathrm{H}_{2}$ | $\begin{aligned} & -188.971955 \\ & (-189.85) \end{aligned}$ |  |
| TS8 | $\begin{aligned} & -188.660189 \\ & (5.79) \end{aligned}$ | $\begin{aligned} & -188.969347 \\ & (3.42) \end{aligned}$ |
| ${ }^{3} \mathrm{CH}_{2} \mathrm{OO}\left({ }^{3} \mathbf{I}\right)$ | $\begin{aligned} & -188.715355 \\ & (-28.83) \end{aligned}$ | $\begin{aligned} & -189.029157 \\ & (-34.11) \end{aligned}$ |
| TS9 | $\begin{aligned} & -188.710702 \\ & (-25.91) \end{aligned}$ |  |
| $\mathrm{H}_{2} \mathrm{CO}+\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ | $\begin{aligned} & -188.794576 \\ & (-78.54) \end{aligned}$ |  |

${ }^{a}$ The geometries were optimized at $\operatorname{CAS}(8,6)+1+2 / \mathrm{cc}-\mathrm{pvdz}$ level.
and a transition state TS3 could be found between them with energy barrier of $20.8 \mathrm{kcal} / \mathrm{mol}$ from II to III.

The reaction of intermediate III might proceed in several ways directly. Keeping the $C_{2 v}$ symmetry, we did locate a stationary point for the formation of $\mathrm{H}_{2}$ and $\mathrm{CO}_{2}$, but the vibrational analysis showed that this stationary point has two imaginary frequencies ( $-1272 \mathrm{i},-597 \mathrm{i} \mathrm{cm}^{-1}$ ). To eliminate the second imaginary mode, another true transition state TS4 has been located, which corresponds to that for the rearrangement from intermediate III to trans-IV. Another possible metathesis process is forming H and $\mathrm{HC}(\mathrm{O}) \mathrm{O}$, however, no transition state could be found between them. In fact, if $\mathrm{HC}(\mathrm{O}) \mathrm{O}$ were involved, the reaction would proceed to $\mathrm{CO}_{2}$ more easily than to $\mathrm{CO},{ }^{30,31}$ which is inconsistent with the branching ratio of CO to $\mathrm{CO}_{2} \cdot{ }^{6}$ Therefore, the possible reaction path is the rearrangement from III to IV. One transition state TS4 has been located along the reaction path, which is only $2 \mathrm{kcal} / \mathrm{mol}$ higher than the intermediate III. The potential energy surfaces for the reactions from HCOOH to $\mathrm{CO}+\mathrm{H}_{2} \mathrm{O}$ or to $\mathrm{CO}_{2}+\mathrm{H}_{2}$ have been reported in some literatures. ${ }^{17,18}$ For comparison our $\mathrm{CAS}(14,12) / \mathrm{cc}-$ pvdz results are also given. The schematic description of the potential energy surfaces for this reaction is shown in Figure 6, from which one can realize that the reaction to CO or $\mathrm{CO}_{2}$ are all exothermic. The calculated heats of reactions $\mathrm{CH}_{2}+\mathrm{O}_{2} \rightarrow \mathrm{CO}+\mathrm{H}_{2} \mathrm{O}$ and $\mathrm{CH}_{2}+\mathrm{O}_{2} \rightarrow \mathrm{CO}_{2}+\mathrm{H}_{2}$ are -176.4 and $-188.9 \mathrm{kcal} / \mathrm{mol}$ at CAS $(14,12) /$ cc-pvdz level (with ZPE correction), respectively, which are quite close to the experimental values. ${ }^{7,32}$ The activation energies for the formation of $\mathrm{CO}+\mathrm{H}_{2} \mathrm{O}$ and $\mathrm{CO}_{2}+\mathrm{H}_{2}$


Figure 3. (a) The energy change along $s$ (the reaction coordinate) for ${ }^{3} \mathrm{CH}_{2}+{ }^{3} \mathrm{O}_{2} \rightarrow \mathbf{I}$ on the singlet surface. (b) The energy change along $s$ on the triplet surface.
from trans-IV are 67.9 and $71.3 \mathrm{kcal} / \mathrm{mol}$, respectively, which is almost the same as the results from Goddard et al. ${ }^{17}$.
2. Triplet Surface. On the triplet surface, we tried to locate all the possible stationary points. Vibrational analysis shows that the true transition state and minimum are with twisted $\mathrm{CH}_{2}$ as shown in Figure 7. The activation barrier is $5.8 \mathrm{kcal} / \mathrm{mol}-$ (see Table 1), about 2 times more than that for the singlet surface. MRCI calculation gives about $3.4 \mathrm{kcal} / \mathrm{mol}$ of energy barrier. IRC calculation shows that apart from TS8, one direction could lead to reactants, and another direction leads to the twisted ${ }^{3} \mathrm{CH}_{2} \mathrm{OO}$ (see Figure 3b).

Two-dimensional plot of the potential surface near the TS8 at $\operatorname{CAS}(14,12) / \mathrm{cc}-\mathrm{pvdz}$ level is shown in Figure 8. One can realize that there is a maximum near $\mathrm{C}-\mathrm{O}=2.1$ and another $\mathrm{O}-\mathrm{O}=1.22 \AA$ and a minimum near $\mathrm{C}-\mathrm{O}=1.39$ and another $\mathrm{O}-\mathrm{O}=1.42 \AA$, which is very close to the optimized one ( $\mathrm{C}-\mathrm{O}$ $=1.38$ and $\mathrm{C}-\mathrm{O}=1.42 \AA)$. Another maximum is near $\mathrm{C}-\mathrm{O}$ $=1.39$ and $\mathrm{O}-\mathrm{O}=1.60 \AA$, and the optimized one is at $\mathrm{C}-\mathrm{O}$ $=1.35$ and $\mathrm{O}-\mathrm{O}=1.51 \AA$. The metathesis of the triplet $\mathrm{CH}_{2}-$ $\mathrm{OO}\left({ }^{3} \mathbf{I}\right)$ could take place via a transition state TS9 with an activation energy of $1.44 \mathrm{kcal} / \mathrm{mol}$ (with zero-point energy correction), which is easier to proceed.

TABLE 2: Harmonic Vibrational Frequencies $\left(\mathbf{c m}^{-1}\right)$ for All the Possible Stationary Points at CAS(14,12)/cc-pvdz Level ${ }^{a}$

|  | $\omega_{1}$ | $\omega_{2}$ | $\omega_{3}$ | $\omega_{4}$ | $\omega_{5}$ | $\omega_{6}$ | $\omega_{7}$ | $\omega_{8}$ | $\omega_{9}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ${ }^{3} \mathrm{CH}_{2}$ | 3260 | 3050 | 1123 |  |  |  |  |  |  |
|  | (3446) | (3220) | (1192) |  |  |  |  |  |  |
| ${ }^{3} \mathrm{O}_{2}$ | 1534 |  |  |  |  |  |  |  |  |
|  | (1688) |  |  |  |  |  |  |  |  |
| TS1 | 3237 | 3033 | 1478 | 1140 | 384 | 198 | 152 | 141 | -234i |
|  | (3430) | (3209) | (1657) | (1196) | (314) | (262) | (122) | (90) | (-162i) |
| ${ }^{1}$ I | 3215 | 3065 | 1465 | 1269 | 1233 | 849 | 793 | 618 | 537 |
|  | (3434) | (3258) | (1513) | (1319) | (1253) | (955) | (741) | (612) | (548) |
| TS2 | 3143 | 2929 | 1530 | 1347 | 1194 | 951 | 721 | 542 | -770i |
| ${ }^{1}$ II | 3249 | 2895 | 1572 | 1305 | 1288 | 1222 | 1021 | 923 | 694 |
| TS3 | 3189 | 3130 | 1568 | 1440 | 1117 | 1083 | 1063 | 947 | -1764i |
| ${ }^{1}$ III | 2873 | 2814 | 1274 | 1192 | 1120 | 912 | 739 | 575 | 516 |
| TS4 | 3100 | 2625 | 1319 | 1296 | 1134 | 1011 | 628 | 624 | -923i |
| ${ }^{1}$ IV | 3667 | 3265 | 1850 | 1478 | 1359 | 1152 | 1105 | 683 | 642 |
| trans |  |  |  |  |  |  |  |  |  |
| TS5 | 4132 | 3215 | 1861 | 1479 | 1288 | 1121 | 966 | 684 | -602i |
| ${ }^{1}$ IV | 3772 | 3167 | 1886 | 1496 | 1335 | 1121 | 1084 | 663 | 519 |
| cis |  |  |  |  |  |  |  |  |  |
| TS6 | 3665 | 2717 | 1918 | 1120 | 795 | 687 | 365 | 314 | -1764i |
| CO | 2169 |  |  |  |  |  |  |  |  |
| $\mathrm{H}_{2} \mathrm{O}$ | 3851 | 3735 | 1709 |  |  |  |  |  |  |
| TS7 | 3756 | 3214 | 1767 | 1297 | 1142 | 845 | 720 | 608 | -2465i |
| $\mathrm{CO}_{2}$ | 2425 | 1339 | 654 | 654 |  |  |  |  |  |
| $\mathrm{H}_{2}$ | 4246 |  |  |  |  |  |  |  |  |
| TS8 | 3227 | 3026 | 1384 | 1149 | 515 | 320 | 218 | 88 | -374i |
|  | (3424) | (3205) | (1578) | (1197) | (467) | (294) | (186) | (82) | (-318i) |
| ${ }^{3} \mathbf{I}$ | 3362 | 3121 | 1461 | 1186 | 1057 | 928 | 787 | 488 | 195 |
|  | (3482) | (3308) | (1523) | (1225) | (1072) | (932) | (780) | (475) | (191) |
| TS9 | 3361 | 3043 | 1500 | 1190 | 1104 | 694 | 426 | 226 | -745i |
| $\mathrm{H}_{2} \mathrm{CO}$ | 2806 | 2522 | 1814 | 1557 | 1291 | 1201 |  |  |  |

${ }^{a}$ The values in parentheses are obtained by CAS $+1+2 / \mathrm{cc}-\mathrm{pvdz}$ method.


Figure 4. The population changes of the orbitals along $s$ : (a) singlet surface and (b) triplet surface.


Figure 5. Contour plot of $\operatorname{CAS}(14,12) / c c-p v d z$ energies in the vicinity of TS1 on the singlet energy surface for the reaction of ${ }^{3} \mathrm{CH}_{2}+{ }^{3} \mathrm{O}_{2} \rightarrow{ }^{1} \mathbf{I}$. The contour increment is $1.0 \mathrm{kcal} / \mathrm{mol}$. The solid contours denote energies below ${ }^{3} \mathrm{O}_{2}+{ }^{3} \mathrm{CH}_{2}$; dashed contours denote energies above ${ }^{3} \mathrm{O}_{2}+{ }^{3} \mathrm{CH}_{2}$.
3. The Rate Constants for the Reaction ${ }^{3} \mathbf{O}_{2}+{ }^{3} \mathrm{CH}_{2} \rightarrow$ Products. To get more accurate rate constants, CAS $+1+2 /$ cc-pvdz has been employed to locate the some key species, such as $\mathrm{CH}_{2}, \mathrm{O}_{2}, \mathrm{TS} 1, \mathrm{TS} 8$, and intermediates ${ }^{1} \mathbf{I}$ and ${ }^{3} \mathbf{I}$. The vibrational analysis was done at the same level. The geometries, energies, and vibrational frequencies obtained at $\mathrm{CAS}(8,6)+1+2 /$ cc-pvdz level were used throughout this section.

Figure 9 shows the calculated rate constants, along with the recent experimental values. The calculated ones are in good


Figure 6. The schematic expression of the possible reaction path for the reaction $\mathrm{CH}_{2}+\mathrm{O}_{2} \rightarrow$ products (with ZPE correction).


Figure 7. The optimized geometries for the stationary points on the triplet energy surface(the values in parentheses are obtained by CAS$(8,6)+1+2 /$ cc-pvdz calculations).
agreement with the experimental values. To test the hindered rotor contribution to the rate constants, a rotation barrier of 0.99 $\mathrm{kcal} / \mathrm{mol}$ was calculated by changing the torsion angle HCOO from $0^{\circ}$ to $180^{\circ}$ for the TS1 geometry, which was located at torsion angle $\mathrm{HCOO}=90^{\circ}$. The hindered rotor rate constants are closer to the experimental ones. At higher temperatures ( $1000-1800 \mathrm{~K}$ ), a negative temperature dependence has been observed, ${ }^{6}$ which is inconsistent with those at lower temperatures. It should be kept in mind, however, that those rate constants at higher temperatures are not determined directly. Our calculations confirm that there is indeed a small barrier between the reactants ${ }^{3} \mathrm{CH}_{2}+{ }^{3} \mathrm{O}_{2}$ and intermediates ${ }^{1} \mathbf{I}$ or ${ }^{3} \mathbf{I}$. Therefore, our results are closer to the lower temperature case.


Figure 8. Contour plot of $\mathrm{CAS}(14,12) / \mathrm{cc}-\mathrm{pvdz}$ energies in the vicinity of TS2 on the triplet energy surface for the reaction of ${ }^{3} \mathrm{CH}_{2}+{ }^{3} \mathrm{O}_{2} \rightarrow 3$. Plotting conventions are the same as in Figure 5.


Figure 9. The comparison of the Arrhenius plots of the calculated hindered rotor (solid line) and harmonic(dot line) rate constants vs $1000 / T$ to the available experimental data for the titled reaction.

With the data obtained from hindered rotor rate constants, we can fit the rate constant expression to the following equation:

$$
k(T)=4.273 \times 10^{-18} T^{2.245} \exp (-185 / T)
$$

Acknowledgment. This work was supported by National Natural Science Foundation of China (Grant No. 20073006). The authors also thank Prof. L. B. Harding for providing the contour plotting program and Prof. T. -H. Tang for language polishing.

## References and Notes

(1) Laufer, A. H. Rev. Chem. Intermed. 1981, 4, 225 and references therein.
(2) Darwin, D. C.; Young, A. T., Johnston, H. S.; Moore, C. B. J. Phys. Chem. 1989, 93, 1074
(3) Böhland, T.; Temps, F.; Wagner, H. Gg. Ber. Bunsen-Ges. Phys. Chem. 1984, 88, 455.
(4) Bley, U.; Temps, F.; Wagner, H. Gg.; Wolf, M. Ber. Bunsen-Ges. Phys. Chem. 1992, 96, 1043.
(5) Vinckier, C.; Debruyn, W. J. Phys. Chem. 1979, 83, 2057.
(6) Dombrowsky, Ch.; Wagner, H. Gg. Ber. Bunsen-Ges. Phys. Chem. 1992, 96, 1048.
(7) Alvarez, R. A.; Moore, C. B. J. Phys. Chem. 1994, 98, 174.
(8) Kafafi, S. A.; Martinez, R. I.; Herron, J. T. In Modern Models of Bonding and Delocalization; Leibman, J. F., Greenberg, A., Eds.; VCH: New York, 1988; p 283.
(9) Ha, T.-K.; Kühne, H.; Vaccane, H.; Günthard, Hs. H. Chem. Phys. Lett. 1974, 24, 172.
(10) Harding, L. B.; Goddard, W. A. J. Am. Chem. Soc. 1978, 100, 7180.
(11) Karlstrôm, G.; Emgstrôm, S.; Jônsson, B. Chem. Phys. Lett. 1979, 67, 343.
(12) Cremer, D.; Schmidt, T.; Gauss, J.; Radhakrishnan, T. P. Angew. Chem., Int. Ed. Engl. 1988, 27, 427.
(13) Bach, R. D.; Andrés, J. L.; Owensby, A. L.; Schlegel, H. B.; McDouall, J. J. W. J. Am. Chem. Soc. 1992, 114, 7207.
(14) Karlstrôm, G.; Roos, B. O. Chem. Phys. Lett. 1981, 79, 416.
(15) Cimiraqglia, R.; Ha, T.-K.; Günthard, Hs. H. Chem. Phys. Lett. 1982, 85, 262.
(16) Dupuis, M.; Lester, W. A. J. Chem. Phys. 1984, 80, 4193.
(17) Goddard, J. D.; Yamaguchi, Y.; Schaefer, H. F. J. Chem. Phys. 1992, 96, 1158.
(18) Francisco, J. S. J. Chem. Phys. 1992, 96, 1167.
(19) Su, H.; Mao, W.; Kong F. Chem. Phys. Lett. 2000, 322, 21.
(20) Woon, D. E.; Dunning, T. H., Jr. J. Phys. Chem. 1993, 98, 1358.
(21) Roos, B. O.; Taylor, P. R.; Siegbahn, P. E. M. Chem. Phys. 1980, 48, 157.
(22) Werner H. J.; Knowles, P. J. J. Chem. Phys. 1985, 82, 5053.
(23) Knowles, P. J.; Werner, H. J. Chem. Phys. Lett. 1985, 115, 259.
(24) Werner H. J.; Knowles, P. J. J. Chem. Phys. 1988, 89, 5803.
(25) Knowles, P. J.; Werner, H. J. Chem. Phys. Lett. 1988, 145, 514
(26) MOLPRO is a package of ab initio programs written by Werner, H.-J., and Knowles, P. J., with contributions from Almlof, J.; Amos, R. D.; Berning, A.; Cooper, D. L.; Deegan, M. J. O.; Dobbyn, A. J.; Eckert, F.; Elbert, S. T.; Hampel, C.; Lindh, R.; Lloyd, A. W.; Meyer, W. Nicklass, A.; Peterson, K.; Pitzer, R.; Stone, A. J.; Taylor, P. R.; Mura, M. E.; Pulay, P.; Schutz, M.; Stoll, H.; Thorsteinsson, T.
(27) Chuang, Y. Y.; Corchado, J. C.; Fast, P. L.; Villa, J.; Hu, W. P.; Liu, Y. P.; Lynch, G. C.; Jackels, C. F.; Nguyen, K. A.; Gu, M..Z.; Rossi, I.; Coitino, E. L.; Clayton, S.; Melissas, V. S.; Steckler, R.; Garret, B. C.; Isaacson, A. D.; Truhlar, D. G. POLYRATE, version 8.2; University of Minnesota: Minneapolis, 1999.
(28) Pitzer, K. S. J. Chem. Phys. 1946, 14, 239.
(29) Chuang, Y. Y.; Truhlar, D. G. J. Chem. Phys. 2000, 112, 1221.
(30) Schatz, G. C.; Fitzcharles, M. S.; Harding, L. B. Faraday Discuss., Chem. Soc. 1987, 74, 359.
(31) Aoyagi, M.; Kato, S. J. Chem. Phys. 1988, 88, 6409.
(32) Baulch, D.L.; Cox, R. A.; Hampson, R. F.; Kerr, J. A.; Troe, J.; Watson, R.T. J. Phys. Chem. Ref. Data, 1980, 9, 295.


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