# First Principles Examination of the Acetylene-Water Clusters, $\mathbf{H C C H}-\left(\mathbf{H}_{2} \mathrm{O}\right)_{x}, \boldsymbol{x}=\mathbf{2}, \mathbf{3}$, and 4 

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#### Abstract

The acetylene-water (A-W) interactions have been investigated by examining the van der Waals clusters $\mathrm{AW}_{x}, x=2,3$, and 4, at the second order (MP2) perturbation theory using the correlation-consistent basis sets, aug-cc-pVnZ, $n=\mathrm{D}\left(\mathrm{AW}_{2}, \mathrm{AW}_{3}\right.$, and $\left.\mathrm{AW}_{4}\right)$, $\mathrm{T}\left(\mathrm{AW}_{2}\right)$. We located 4 minima (m) and 2 saddle points (sp), 10 m and 3 sp , and 30 m and 3 sp on the potential energy surfaces of the $\mathrm{AW}_{2}, \mathrm{AW}_{3}$, and $\mathrm{AW}_{4}$ clusters, respectively. We report the fully optimized geometries and interaction energies $\Delta E_{\mathrm{e}}$, including corrections for basis set superposition error, $\Delta E_{\mathrm{e}}(\mathrm{BSSE})$, as well as zero-point energies, $\Delta E_{0}(\mathrm{BSSE})$, for the various stationary points. The global minima of the $\mathrm{AW}_{2}$ and $\mathrm{AW}_{3}$ clusters are cyclic configurations in which the acetylene molecule inserts into the water hydrogen bonding network. The corresponding interaction energies $\Delta E_{\mathrm{e}}(\mathrm{BSSE})\left[\Delta E_{0}(\mathrm{BSSE})\right]$ are $\mathrm{AW}_{2},-10.37[-6.70] \mathrm{kcal} / \mathrm{mol}(\mathrm{MP} 2 / \mathrm{aug}-\mathrm{cc}-\mathrm{pVTZ})$ and $\mathrm{AW}_{3},-17.80[-11.46]$ $\mathrm{kcal} / \mathrm{mol}$ (MP2/aug-cc-pVDZ). The global minimum of $\mathrm{AW}_{4}$ corresponds to a van der Waals complex between a cyclic water tetramer $\mathrm{W}_{4}$ and A with an interaction energy of $-28.01[-18.67] \mathrm{kcal} / \mathrm{mol}$ (MP2/aug-cc$\mathrm{pVDZ})$. The 4 and 10 local minima for the $x=2$ and 3 clusters span an energy range of 4.3 and $6.1 \mathrm{kcal} / \mathrm{mol}$ above the respective global minima. For $\mathrm{AW}_{4}$, the energy range for the 30 minima is $14.1 \mathrm{kcal} / \mathrm{mol}$; however, the first 28 lie within $8.4 \mathrm{kcal} / \mathrm{mol}$ above the global minimum. The analysis of the many-body interaction energy terms suggests that the global and low-lying ring networks are stabilized by the maximization of the many-body (mainly the 3-body) terms, whereas the higher lying minima are mainly described by 2-body interactions.


## 1. Introduction

The present work is a continuation of our study on the aqueous microsolvation of acetylene. ${ }^{1,2}$ In the first paper, ${ }^{1}$ we presented an extensive study of the potential energy surface (PES) of the $\mathrm{C}_{2} \mathrm{H}_{2}-\mathrm{H}_{2} \mathrm{O}$ (acetylene-water, AW) dimer, where two minima, a global AW-Y (the water acting as proton acceptor), a local AW-T (the water acting as proton donor), and three transition states (shown in Figure 1) were located. However, because of basis set superposition errors (BSSE) and zero-point energy (ZPE) corrections, the AW-T local minimum is destabilized, and it can "slip" into the global minimum (AWY ), thus explaining the experimental observation of a single isomer. ${ }^{3-5}$ In a subsequent paper, ${ }^{2}$ we identified the spectroscopic signature for the predicted change in the structural pattern between $\mathrm{AW}_{3}$ and $\mathrm{AW}_{4}$ from a cyclic configuration that incorporates acetylene (A) into the water hydrogen bonding network $\left(\mathrm{AW}_{3}\right)$ to a van der Waals complex between A and a cyclic water tetramer $\left(\mathrm{AW}_{4}\right)$. That study also produced qualitative differences between the ab initio results and those with an empirical potential as regards the structures of the global minima of the first few $\mathrm{AW}_{x}$ clusters.

To the best of our knowledge, there currently exists limited experimental and/or theoretical work on acetylene-water

[^0]

Figure 1. Geometries of the $A W-Y$ and $A W-T$ minima. Bond distances in $\AA$ at the MP2/avdz(avtz)[avqz] level.
clusters $\left(\mathrm{AW}_{x}, x>1\right) .{ }^{6-8}$ Besides our previous work, ${ }^{2}$ we are aware of only three publications on the $\mathrm{AW}_{x}$ clusters. Choi et al. ${ }^{6}$ studied the ion/molecule reactions within the acetylene-water heterocluster ions $\left(\mathrm{C}_{2} \mathrm{H}_{2}\right)_{n} \cdot\left(\mathrm{H}_{2} \mathrm{O}\right)_{m}{ }^{+}$using electron impact time-of-flight mass spectrometry. Dykstra, ${ }^{7}$ employing the molecular mechanics for clusters (MMC) model potential, calculated an interaction energy of $\Delta E_{\mathrm{e}}=-12.65 \mathrm{kcal} / \mathrm{mol}\left(\Delta E_{\mathrm{o}}=-8.17\right.$ $\mathrm{kcal} / \mathrm{mol}$, including ZPE corrections) for $\mathrm{AW}_{2}$. van Voorhis and Dykstra, ${ }^{8}$ used the MMC model potential for $\mathrm{AW}_{3}$, comparing it with a number of other four-membered water containing clusters. They found four local minima within $0.8 \mathrm{kcal} / \mathrm{mol}$ of their global minimum, with one of those only $0.24 \mathrm{kcal} / \mathrm{mol}$ higher. Their characterization of the global minimum has it much like $\mathrm{W}_{3}$ with an adjacent acetylene. Our results show five local
minima within $0.8 \mathrm{kcal} / \mathrm{mol}$ of the global minimum, but the global minimum is a four-membered ring, and the secondary minimum is $0.26 \mathrm{kcal} / \mathrm{mol}(0.37 \mathrm{kcal} / \mathrm{mol}$ including ZPE) higher in energy. The same authors refer also to a $\mathrm{AW}_{4}$ minimum as a cyclic water tetramer interacting with acetylene but without reporting any structural or energetic results.

In view of the rich structural patterns found in the first few $\mathrm{AW}_{x}$ clusters and the disagreement between the earlier ab initio results and the ones obtained with empirical models, we have extended our previous work into a detail investigation of the global and local minima of the $\mathrm{AW}_{x}, x=2,3$, and 4 , clusters. The present investigation is organized as follows: in section 2, we outline the computational approach; in section 3, we report the structures and energetics of the various stationary points; and finally in section 4 , we summarize our findings and main conclusions.

## 2. Computational Approach

A preliminary sampling of the configuration space and the different hydrogen bonding networks was performed with the smaller 4-31G basis set. The resulting geometries were used as starting points and were subsequently fully optimized with the augmented correlation consistent basis sets, aug-cc-pVnZ (=avnz), $n=\mathrm{D}$ and T, of Dunning and co-workers. ${ }^{9}$ For $\mathrm{AW}_{2}$, both the avdz and avtz sets were used, whereas the avdz set was employed for the $\mathrm{AW}_{3}$ and $\mathrm{AW}_{4}$ clusters. In some instances, configurations of higher symmetry were probed, and this resulted in obtaining saddle points (sp) of higher order (i.e., configurations for which the Hessian matrix has more than one negative eigenvalues). All calculations were performed at the secondorder perturbation (MP2) level of theory with the Gaussian $98^{10}$ programs. The "very tight" or "tight" options were used in all geometry optimizations. Energies were converged to about 0.01 $\mu$ hartree, and the corresponding root-mean-square deviations of energy gradients with respect to nuclear coordinates were $\sim 11$ $\mu$ hartree/bohr. Harmonic vibrational frequencies were computed for all $\mathrm{AW}_{2}$ and $\mathrm{AW}_{3}$ minima and saddle points and for three minima for $\mathrm{AW}_{4}$ at the MP2/avdz level. Furthermore, for the four $\mathrm{AW}_{2}$ minima, the harmonic frequencies were also obtained with the avtz basis set.

Corrections due to basis set superposition error (BSSE), ${ }^{11}$ which are important for weakly bound van der Waals complexes, ${ }^{12}$ are taken into account following a procedure described earlier ${ }^{13}$ and briefly outlined below.

The interaction energy $\Delta E_{e}\left(\mathrm{AW}_{x}\right)$ of the $\mathrm{AW}_{x}$ cluster is defined as

$$
\begin{equation*}
\Delta E_{\mathrm{e}}\left(\mathrm{AW}_{x}\right)=E_{\mathrm{AW}_{x}}^{\mathrm{aW}_{x}}\left(\mathrm{AW}_{x}\right)-E_{\mathrm{A}}^{\mathrm{a}}(\mathrm{~A})-x E_{\mathrm{W}}^{\mathrm{w}}(\mathrm{~W}) \tag{1}
\end{equation*}
$$

where, $E_{\mathrm{G}}^{s}(\mathrm{M})$ refers to the total energy of the molecule M at the geometry G, computed with basis set $s$; the above relation is modified appropriately when the BSSE correction is taken into account. For instance, the (BSSE)-corrected interaction energy, $\Delta E_{\mathrm{e}}$ (BSSE), for the $\mathrm{AW}_{4}$ cluster can be written as

$$
\begin{array}{r}
\Delta E_{\mathrm{e}}(\mathrm{BSSE})=E_{\mathrm{AW}_{4}}^{\mathrm{aw}_{4}}\left(\mathrm{AW}_{4}\right)-E_{\mathrm{AW}_{4}}^{\mathrm{aw}_{4}}(\mathrm{~A})-\sum_{i=a}^{d} E_{\mathrm{AW}_{4}}^{\mathrm{aw}_{4}}\left(\mathrm{~W}_{i}\right)+ \\
R_{\mathrm{A}}+\sum_{i=a}^{d} R_{\mathrm{W}_{i}} \tag{2}
\end{array}
$$

where $R$ are relaxation or deformation terms defined by the
relations

$$
\begin{gather*}
R_{\mathrm{A}}=E_{\mathrm{AW}_{x}}^{\mathrm{a}}(\mathrm{~A})-E_{\mathrm{A}}^{\mathrm{a}}(A)  \tag{3a}\\
R_{\mathrm{W}_{i}}=E_{\mathrm{AW}_{x}}^{\mathrm{w}_{\mathrm{i}}}\left(\mathrm{~W}_{i}\right)-\mathrm{E}_{\mathrm{W}_{i}}^{\mathrm{w}_{i}}\left(\mathrm{~W}_{\mathrm{i}}\right), i=\mathrm{a}, \mathrm{~b}, \mathrm{c}, \mathrm{~d} \tag{3b}
\end{gather*}
$$

and $W_{i}(i=\mathrm{a}, \mathrm{b}, \mathrm{c}$, and d$)$ refers to the four water molecules, respectively.

Similarly, the analysis of the many-body interaction energy terms was performed using a procedure described before ${ }^{14}$ which is based on casting the total energy of the $n$-body cluster $\mathrm{X}_{1} \mathrm{X}_{2} \mathrm{X}_{3} \ldots \mathrm{X}_{n}$ as

$$
\begin{align*}
& E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{x}_{1} \mathrm{X}_{2} \ldots}\left(\mathrm{X}_{1} \mathrm{X}_{2} \ldots\right)=\sum_{i=1}^{n} E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{x}_{i}}\left(\mathrm{X}_{i}\right)+\sum_{i}^{n-1} \sum_{j>i}^{n} \Delta^{2} E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{x}_{1} \mathrm{X}_{j}}\left(\mathrm{X}_{i} \mathrm{X}_{j}\right)+ \\
& \sum_{i}^{n-2 n-1} \sum_{j>i} \sum_{k>j}^{n} \Delta^{3} E_{\mathrm{X}_{1} \mathrm{X}_{2} \mathrm{X}_{j} \ldots}^{\mathrm{x}_{j}}\left(\mathrm{X}_{i} \mathrm{X}_{j}\right)+\ldots \tag{4}
\end{align*}
$$

where $\Delta^{2} E, \Delta^{3} E, \ldots$ are two-, three-, etc. body terms, respectively, defined as

$$
\begin{align*}
& \Delta^{2} E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{x}_{\mathrm{i}} \mathrm{X}_{j}}\left(\mathrm{X}_{i} \mathrm{X}_{j}\right)=E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{x}_{\mathrm{i}} \mathrm{X}_{j}}\left(\mathrm{X}_{i} \mathrm{X}_{j}\right)-\left\{E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{x}_{i}}\left(\mathrm{X}_{i}\right)+\right. \\
& \left.E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{X}_{j}}\left(\mathrm{X}_{j}\right)\right\}  \tag{5}\\
& \Delta^{3} E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{x}_{\mathrm{i}} \mathrm{x}_{\mathrm{j}} \mathrm{X}_{k}}\left(\mathrm{X}_{i} \mathrm{X}_{j} \mathrm{X}_{k}\right)=E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{X}_{\mathrm{i}} \mathrm{x}_{j} \mathrm{x}_{k}}\left(\mathrm{X}_{i} \mathrm{X}_{j} \mathrm{X}_{k}\right)-\left\{E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{X}_{i}}\left(\mathrm{X}_{i}\right)+\right. \\
& \left.E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{X}_{j}}\left(\mathrm{X}_{j}\right)+E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{X}_{k}}\left(\mathrm{X}_{k}\right)\right\}-\left\{\Delta^{2} E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{x}_{\mathrm{i}} \mathrm{X}_{j}}\left(\mathrm{X}_{i} \mathrm{X}_{j}\right)+\right. \\
& \left.\Delta^{2} E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{x}_{\mathrm{i}} \mathrm{~K}_{k}}\left(\mathrm{X}_{i} \mathrm{X}_{k}\right)+\Delta^{2} E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{x}_{\mathrm{X}} \mathrm{X}_{k}}\left(\mathrm{X}_{j} \mathrm{X}_{k}\right)\right\} \tag{6}
\end{align*}
$$

In the preceding, the one-, two-, three-, etc. body term summations contain $\binom{n}{1},\binom{n}{2},\binom{n}{3}, \ldots$ terms, respectively, for a total of $\sum_{m=1}^{n}\binom{n}{m}=2^{n}-1$ terms.

In this notation, the BSSE-corrected, two- and three-body terms are ${ }^{14 \mathrm{~b}, \mathrm{c}}$

$$
\begin{align*}
& \Delta^{2} E_{X_{1} X_{2} \ldots}^{\mathrm{x}_{1} X_{2} \ldots}\left(\mathrm{X}_{i} \mathrm{X}_{j}, \mathrm{BSSE}\right)=E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{x}_{1} \mathrm{X}_{2} \ldots}\left(\mathrm{X}_{i} \mathrm{X}_{j}\right)-\{ E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{x}_{1} \mathrm{X}_{2} \ldots}\left(\mathrm{X}_{i}\right)+ \\
&\left.E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots \ldots}^{\mathrm{x}_{1} \mathrm{X}_{2} \ldots}\left(\mathrm{X}_{j}\right)\right\} \tag{7}
\end{align*}
$$

$$
\begin{align*}
& \Delta^{3} E_{X_{1} X_{2} \ldots}^{\mathrm{x}_{1} \mathrm{X}_{2} \ldots}\left(\mathrm{X}_{i} \mathrm{X}_{j} \mathrm{X}_{k}, \mathrm{BSSE}\right)=E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{x}_{1} \mathrm{X}_{2} \ldots}\left(\mathrm{X}_{i} \mathrm{X}_{j} \mathrm{X}_{k}\right)- \\
& \left\{E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{x}_{1} \mathrm{X}_{2} \ldots}\left(\mathrm{X}_{i}\right)+E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{x}_{1} \mathrm{X}_{2} \ldots}\left(\mathrm{X}_{j}\right)+E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{x}_{1} \mathrm{X}_{2} \ldots}\left(\mathrm{X}_{k}\right)\right\}- \\
& \left\{\Delta^{2} E_{X_{1} X_{2} \ldots}^{\mathrm{x}_{1} \mathrm{X}_{2} \ldots}\left(\mathrm{X}_{i} \mathrm{X}_{j}\right)+\Delta^{2} E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{x}_{1} \mathrm{X}_{2} \ldots}\left(\mathrm{X}_{i} \mathrm{X}_{k}\right)+\Delta^{2} E_{\mathrm{X}_{1} \mathrm{X}_{2} \ldots}^{\mathrm{x}_{1} \mathrm{X}_{2} \ldots}\left(\mathrm{X}_{j} \mathrm{X}_{k}\right)\right\} \tag{8}
\end{align*}
$$

## 3. Results and Discussion

a. $\mathrm{C}_{2} \mathbf{H}_{\mathbf{2}}, \mathrm{H}_{\mathbf{2}} \mathrm{O}$, and $\mathrm{C}_{2} \mathrm{H}_{\mathbf{2}}-\mathbf{H}_{\mathbf{2}} \mathrm{O}$. For the purpose of analyzing our current data, we list the results of our previous work ${ }^{1}$ on the $\mathrm{C}_{2} \mathrm{H}_{2}-\mathrm{H}_{2} \mathrm{O}$ cluster in Table 1. Specifically, the total and interaction energies of $\mathrm{A}, \mathrm{W},(\mathrm{AW}-\mathrm{Y}),(\mathrm{AW}-\mathrm{T}), \mathrm{W}_{x},{ }^{15} x=2$, 3 , and 4 , and selected optimal internal coordinates of (AW-Y) and (AW-T) at the MP2/avnz, $n=\mathrm{D}$ and T level of theory are listed (see Figure 1 for the corresponding structures). Table 1 contains a subset of the results of the exhaustive study of the PES of (AW) previously reported in ref 1.
b. $\mathbf{C}_{2} \mathbf{H}_{\mathbf{2}}\left(\mathbf{H}_{\mathbf{2}} \mathrm{O}\right)_{2}$. We have located six stationary points, four minima (m), and two saddle points (sp) on the PES of $\mathrm{AW}_{2}$; their structures are shown in Figure 2. The geometries of all six stationary points were fully optimized with the avdz basis set, whereas the four minima were also optimized with the larger avtz set. ${ }^{16 \mathrm{a}}$ Their total energies $\left(E_{\mathrm{e}}\right)$, interaction energies with respect to the isolated fragments ( $\Delta E_{\mathrm{e}}$ ), and BSSE-corrected

TABLE 1: Total Energies $\boldsymbol{E}_{\mathrm{e}}$ (hartree), Interaction Energies $\Delta E_{\mathrm{e}}(\mathrm{kcal} / \mathrm{mol})$, Corrected for BSSE, $\Delta E_{\mathrm{e}}(\mathrm{BSSE})(\mathrm{kcal} / \mathrm{mol})$, and Harmonic ZPE, (kcal/mol), of $\mathrm{C}_{2} \mathbf{H}_{2}, \mathrm{H}_{2} \mathrm{O}, \mathrm{C}_{2} \mathrm{H}_{2}-\mathrm{H}_{2} \mathrm{O}$ $(\mathrm{AW}-\mathrm{Y}, \mathrm{AW}-\mathrm{T})$, and $\left(\mathrm{H}_{2} \mathrm{O}\right)_{x}, x=2,3$, and $4^{a}$

${ }^{a}$ van der Waals Geometries, $R_{\mathrm{vdW}}(\AA)$ and $\phi_{\mathrm{vdw}}$ (degrees) of AW-Y and AW-T clusters at the MP2/avnz, $n=\mathrm{D}$ and T level. ${ }^{b}$ See Figure 1. ${ }^{c}$ Experimental value 2.229 A. ref $4 .{ }^{d} \mathrm{H}_{1} \ldots \mathrm{O}$ distance in $\mathrm{AW}-\mathrm{Y}$ (Figure 1a), $\mathrm{H}_{3}$...middle of the triple bond in AW-T (Figure 1b). ${ }^{e}$ Angle between the $\mathrm{C} \equiv \mathrm{C}$ and $\mathrm{C}_{2}\left(\mathrm{H}_{2} \mathrm{O}\right)$ axes in $\mathrm{AW}-\mathrm{Y}($ Figure 1a); angle between $\mathrm{O}-\mathrm{H}_{3}$-middle of the triple bond in AW-T (Figure 1b). ${ }^{f}$ References 13 and 15.
$\mathrm{C}_{2} \mathrm{H}_{2}\left(\mathrm{H}_{2} \mathrm{O}\right)_{2}$

a. $\mathrm{AW}_{2}$-m1 $\left(\mathrm{C}_{1}\right)$


c. $\mathrm{AW}_{2-\mathrm{sp}}\left(\mathrm{C}_{s}\right)$


Figure 2. Geometries of the four minima $\left(\mathrm{AW}_{2} \_\mathrm{m} n, n=1-4\right)$ and the two saddle points $\left(\mathrm{AW}_{2} \_\operatorname{spn}, n=1-2\right)$ of $\mathrm{AW}_{2}$. Bond distances in $\AA$ at the MP2/avdz(avtz) level.
interaction energies, $\Delta E_{\mathrm{e}}(\mathrm{BSSE})$ are listed in Table 2; van der Waals (vdW) intermolecular distances are also displayed in Figure 2. Harmonic vibrational frequencies and IR intensities of the global minimum are shown in Table 3. ${ }^{16 \mathrm{~b}}$ Finally, the decomposition of the interaction energy of the structures into many-body terms is reported in Table 4. Note that the minima

TABLE 2: Total Energies $\boldsymbol{E}_{\mathrm{e}}$ (hartree), Interaction Energies $\Delta E_{\mathrm{e}}(\mathrm{kcal} / \mathrm{mol})$, Corrected for BSSE [ $\Delta E_{\mathrm{e}}($ BSSE $\left.)\right]$, and Zero-Point Energy [ $\Delta E_{0}$ (BSSE)] for the AW $_{2}$ Cluster at the MP2 Level with the avdz and avtz Basis Sets

| $\mathrm{AW}_{2}{ }^{a}$ | $E_{\mathrm{e}}$ | $\Delta E_{\mathrm{e}}$ | $\Delta E_{\mathrm{e}}(\mathrm{BSSE})$ | $\Delta E_{0}(\mathrm{BSSE})$ |
| :--- | :---: | ---: | :---: | :---: |
|  |  | avdz |  |  |
| m 1 | -229.633966 | -12.02 | -9.60 | -5.76 |
| m 2 | -229.629037 | -8.92 | -6.88 | -3.71 |
| $\mathrm{sp1}^{b}$ | -229.627971 | -8.25 | -6.23 | -3.57 |
| m 3 | -229.625556 | -6.74 | -4.81 | -2.71 |
| $\mathrm{sp2}^{c}$ | -229.625506 | -6.71 | -4.47 | -2.47 |
| m 4 | -229.623800 | -5.64 | -3.90 | -1.75 |
|  |  | avtz |  |  |
| m 1 | -229.840473 | -11.57 | -10.37 | -6.70 |
| m 2 | -229.835290 | -8.31 | -7.34 | -4.37 |
| m 3 | -229.831277 | -5.79 | -4.91 | -3.03 |
| m 4 | -229.830802 | -5.50 | -4.61 | -2.45 |

${ }^{a} \mathrm{~m}$ (minimum) and sp (saddle point) according to Figure 2. ${ }^{b}$ Two imaginary frequencies. ${ }^{c}$ One imaginary frequency.

TABLE 3: Harmonic Vibrational Frequencies $\omega\left(\mathrm{cm}^{-1}\right)$, IR Intensities IR-I ( $\mathbf{k m} / \mathbf{m o l}$ ), and Zero-Point Energies ZPE ( $\mathrm{kcal} / \mathrm{mol}$ ) of the Global Minima of $\mathrm{AW}_{x}, x=2,3$, and 4 at the MP2 Level of Theory with the avdz and avtz (for $\mathrm{AW}_{2}$ ) Basis Sets

|  | $\mathrm{AW}_{2}$ _ml |  |  | $\mathrm{AW}_{3} \mathrm{~m}$ m1 |  | $\mathrm{AW}_{4 \_ \text {ml }}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | avdz |  | avtz | avdz |  | avdz |  |
|  | $\omega$ | IR-I | $\omega$ | $\omega$ | IR-I | $\omega$ | IR-I |
| $\omega_{1}$ | 88 | 3.75 | 91 | 32 | 0.591 | 23 | 0.228 |
| $\omega_{2}$ | 94 | 18.1 | 95 | 57 | 3.06 | 39 | 0.211 |
| $\omega_{3}$ | 118 | 46.0 | 116 | 115 | 18.0 | 66 | 0.132 |
| $\omega_{4}$ | 148 | 30.2 | 151 | 122 | 2.79 | 80 | 1.83 |
| $\omega_{5}$ | 170 | 13.4 | 173 | 151 | 36.8 | 99 | 1.54 |
| $\omega_{6}$ | 183 | 88.2 | 192 | 163 | 8.00 | 122 | 4.24 |
| $\omega_{7}$ | 192 | 18.4 | 194 | 170 | 35.6 | 146 | 0.398 |
| $\omega_{8}$ | 240 | 69.7 | 251 | 191 | 95.9 | 198 | 2.50 |
| $\omega_{9}$ | 348 | 144 | 351 | 207 | 24.8 | 221 | 75.9 |
| $\omega_{10}$ | 467 | 38.6 | 495 | 236 | 21.2 | 237 | 15.5 |
| $\omega_{11}$ | 476 | 6.09 | 622 | 252 | 129 | 254 | 9.43 |
| $\omega_{12}$ | 530 | 42.3 | 636 | 283 | 65.4 | 264 | 86.5 |
| $\omega_{13}$ | 687 | 132 | 691 | 402 | 52.1 | 269 | 96.9 |
| $\omega_{14}$ | 775 | 105 | 804 | 418 | 22.3 | 285 | 144 |
| $\omega_{15}$ | 794 | 77.9 | 827 | 476 | 13.4 | 356 | 70.9 |
| $\omega_{16}$ | 1629 | 69.5 | 1634 | 490 | 29.0 | 413 | 3.50 |
| $\omega_{17}$ | 1644 | 28.7 | 1651 | 554 | 54.5 | 440 | 18.6 |
| $\omega_{18}$ | 1931 | 5.21 | 1952 | 729 | 119 | 452 | 37.7 |
| $\omega_{19}$ | 3363 | 228 | 3365 | 809 | 158 | 463 | 7.92 |
| $\omega_{20}$ | 3488 | 3.55 | 3497 | 821 | 95.2 | 479 | 6.28 |
| $\omega_{21}$ | 3647 | 285 | 3657 | 862 | 20.5 | 482 | 18.5 |
| $\omega_{22}$ | 3734 | 147 | 3741 | 1632 | 55.3 | 713 | 123 |
| $\omega_{23}$ | 3891 | 138 | 3901 | 1644 | 60.1 | 767 | 92.4 |
| $\omega_{24}$ | 3899 | 104 | 3910 | 1664 | 11.6 | 775 | 165 |
| $\omega_{25}$ |  |  |  | 1922 | 17.6 | 807 | 173 |
| $\omega_{26}$ |  |  |  | 3297 | 379 | 850 | 134 |
| $\omega_{27}$ |  |  |  | 3478 | 12.6 | 994 | 4.92 |
| $\omega_{28}$ |  |  |  | 3532 | 570 | 1639 | 84.6 |
| $\omega_{29}$ |  |  |  | 3581 | 534 | 1652 | 48.7 |
| $\omega_{30}$ |  |  |  | 3694 | 277 | 1661 | 80.6 |
| $\omega_{31}$ |  |  |  | 3883 | 142 | 1685 | 10.6 |
| $\omega_{32}$ |  |  |  | 3886 | 99.9 | 1934 | 3.97 |
| $\omega_{33}$ |  |  |  | 3894 | 88.7 | 3367 | 142 |
| $\omega_{34}$ |  |  |  |  |  | 3381 | 260 |
| $\omega_{35}$ |  |  |  |  |  | 3445 | 1378 |
| $\omega_{36}$ |  |  |  |  |  | 3484 | 633 |
| $\omega_{37}$ |  |  |  |  |  | 3493 | 7.41 |
| $\omega_{38}$ |  |  |  |  |  | 3591 | 393 |
| $\omega_{39}$ |  |  |  |  |  | 3832 | 129 |
| $\omega_{40}$ |  |  |  |  |  | 3881 | 92.5 |
| $\omega_{41}$ |  |  |  |  |  | 3882 | 84.9 |
| $\omega_{42}$ |  |  |  |  |  | 3885 | 84.0 |
| ZPE | 46.51 |  | 47.18 | 62.39 |  | 78.78 |  |

and saddle points are ordered according to their uncorrected (for BSSE and ZPE) interaction energies ( $\Delta E_{\mathrm{e}}$ ).

TABLE 4: Many-Body Decomposition of the Interaction Energies (in $\mathrm{kcal} / \mathrm{mol}$ ) of the $\mathrm{AW}_{2}$ at the MP2/avdz and avtz Levels of Theory ${ }^{a}$

| $\mathrm{AW}_{2}{ }^{\text {b }}$ | m1 | m2 | sp1 | m3 | sp2 | m4 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | avdz |  |  |  |
| $\mathrm{A}-\mathrm{W}_{\mathrm{a}}$ | -2.94(-2.09) | -3.16(-2.40) | $-2.48(-1.92)$ | $-3.58(-2.64)$ | -3.59(-2.64) | -3.04(-2.19) |
| A $-\mathrm{W}_{\mathrm{b}}$ | -3.00(-2.18) | $-1.05(-0.46)$ | $-2.56(-1.78)$ | -3.58(-2.64) | -2.98(-2.65) | -3.04(-2.19) |
| $\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}$ | $-5.20(-4.33)$ | $-5.11(-4.39)$ | $-3.36(-2.69)$ | 0.00 (0.12) | 0.17 (0.17) | 0.15 (0.17) |
| total two body | $-11.13(-8.60)$ | -9.32(-7.25) | -8.41(-6.38) | -7.17(-5.15) | $-6.40(-5.13)$ | -5.93(-4.22) |
| A-W-W | -1.06(-1.18) | 0.36(0.34) | 0.07(0.07) | 0.41(0.32) | -0.33 0.63) | 0.29(0.31) |
| relaxation | 0.18 | 0.03 | 0.08 | 0.02 | 0.02 | 0.01 |
| $\Delta \mathrm{E}_{\mathrm{e}}\left[\Delta \mathrm{E}_{\mathrm{e}}(\mathrm{BSSE})\right]$ | -12.02(-9.60) | -8.92(-6.88) | $-8.25(-6.23)$ | $-6.74(-4.81)$ | -6.71(-4.47) | -5.64(-3.90) |
|  |  |  | avtz |  |  |  |
| $\mathrm{A}-\mathrm{W}_{\mathrm{a}}$ | $-2.68(-2.33)$ | -2.73(-2.39) |  | $-3.16(-2.71)$ |  | -3.01(-2.57) |
| $\mathrm{A}-\mathrm{W}_{\mathrm{b}}$ | $-2.89(-2.46)$ | $-1.05(-0.80)$ |  | $-3.16(-2.71)$ |  | -3.01(-2.57) |
| $\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}$ | -5.10(-4.63) | $-5.02(-4.61)$ |  | 0.12(0.12) |  | 0.16(0.17) |
| total two body | $-10.67(-9.43)$ | -8.80(-7.80) |  | $-6.20(-5.31)$ |  | -5.86(-4.97) |
| A-W-W | $-1.09(-1.14)$ | 0.44 (0.41) |  | $0.37(0.36)$ |  | 0.35(0.34) |
| relaxation | 0.20 | 0.04 |  | 0.03 |  | 0.02 |
| $\Delta \mathrm{E}_{\mathrm{e}}\left[\Delta \mathrm{E}_{\mathrm{e}}(\mathrm{BSSE})\right]$ | -11.57(-10.37) | $-8.31(-7.34)$ |  | $-5.79(-4.91)$ |  | -5.50(-4.61) |

${ }^{a}$ BSSE-corrected values are shown in parentheses. ${ }^{a}$ m(minimum) and sp(saddle point), Figure 2.

The two lowest minima $\mathrm{AW}_{2} \mathrm{~m} 1$ and $\mathrm{AW}_{2}$ _m2 (Figures 2 a , 2 b ) are cyclic trimers of $\mathrm{C}_{1}$ and $\mathrm{C}_{s}$ symmetry, respectively, in which the acetylene molecule acts simultaneously as a proton donor and a proton acceptor to neighboring water molecules. In essence, both structures can be viewed as resulting from the interaction of a water dimer (albeit with different orientation) with A . The difference between m 1 and m 2 lies in the orientation of the T- and Y-water molecules (cf. Figure 2). In the former, the water molecules are simultaneously proton donors and proton acceptors to A, whereas in the latter, the T-water is a double H donor and the Y -water a double H acceptor.

The van der Waals (vdW) $\mathrm{H}_{1} \ldots \mathrm{O}_{\mathrm{a}}$ distance in $\mathrm{AW}_{2 \_} \mathrm{m} 1$ is increased by $\sim 0.02 \AA$ and the $\mathrm{C}_{2} \mathrm{C}_{1} \mathrm{O}_{1}$ angle $\left(151.9^{\circ}\right)$ is decreased by $27.6^{\circ}$ with respect to the $\mathrm{AW}-\mathrm{Y}$ minimum (Figure $1)^{1}$ at the MP2/avtz level. The intermolecular distance between the two water molecules, $\mathrm{O}_{\mathrm{b}} \ldots \mathrm{H}_{1 \mathrm{a}}$, is $1.916 \AA, 0.017 \AA$ shorter than the corresponding value in the water dimer $\left(\mathrm{W}_{2}\right){ }^{15 b}$ and the angle $\mathrm{H}_{1 \mathrm{a}}-\mathrm{O}_{\mathrm{a}} \ldots \mathrm{O}_{\mathrm{b}}$ is $13.6^{\circ}$ as contrasted to $5.7^{\circ}$ in $\mathrm{W}_{2} .{ }^{15 \mathrm{c}}$ In general, the structure of the water dimer fragment within the $\mathrm{AW}_{2 \_} \mathrm{m} 1$ minimum is very similar to that of the free $\mathrm{W}_{2} .{ }^{15 \mathrm{~b}, \mathrm{c}}$ In contrast, the vdW distances in m 2 are $0.1-0.3 \AA$ longer than the corresponding distances in m 1 , mainly because the T-water in ml acts as double H donor in that configuration.

The MP2 interaction energies $\Delta E_{\mathrm{e}}\left[\Delta E_{\mathrm{e}}(\mathrm{BSSE})\right]$ for the global minimum ( $\mathrm{AW}_{2} \_\mathrm{m} 1$ ) with the avdz and avtz sets are $-12.02[-$ $9.60]$ and $-11.57[-10.37] \mathrm{kcal} / \mathrm{mol}$, respectively. Upon corrections for ZPE , these become $\Delta E_{0}(\mathrm{BSSE})=-5.76(\mathrm{avdz})$ and -6.70 (avtz) $\mathrm{kcal} / \mathrm{mol}$. The corresponding $\Delta E_{\mathrm{e}}\left[\Delta E_{\mathrm{e}}(\mathrm{BSSE})\right]$ $\left\{\Delta E_{0}(\mathrm{BSSE})\right\}$ energies for $\left(\mathrm{AW}_{2} \_\mathrm{m} 2\right)$ are $-8.92[-6.88]$ $\{-3.71\} \mathrm{kcal} / \mathrm{mol}$ with the avdz and $-8.31[-7.34]\{-4.37\}$ $\mathrm{kcal} / \mathrm{mol}$ with the avtz sets, respectively. Dykstra ${ }^{7}$ previously reported $\Delta E_{\mathrm{e}}\left(\Delta E_{0}\right)$ values of $-12.65(-8.17) \mathrm{kcal} / \mathrm{mol}$ for the $\mathrm{AW}_{2}$ (cyclic) cluster using the MMC approach but without referring to any specific geometry. The energetic stabilization of m 1 with respect to m 2 by $\sim 2.3 \mathrm{kcal} / \mathrm{mol}$ (MP2/avtz) is also reflected in the more "open" structure of the latter when compared to the former. From Table 4, it is seen that a large portion ( $\sim 74 \%$ ) of this difference arises from the $\mathrm{A}-\mathrm{W}_{\mathrm{b}}$ interaction, which is weaker by $1.7 \mathrm{kcal} / \mathrm{mol}$ (MP2/avtz, see Table 4) in m 2 . Furthermore, there is a $1.55 \mathrm{kcal} / \mathrm{mol}$ difference between the two minima in the three-body term $\mathrm{A}-\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}$ which is attractive ( $-1.14 \mathrm{kcal} / \mathrm{mol}$ ) for m 1 but repulsive $(+0.41 \mathrm{kcal} /$ mol ) for m 2 . This is consistent with the "homodromic" topology of the ring in ml and the fact that these networks have been
previously reported ${ }^{14 \mathrm{c}}$ to exhibit larger nonadditivities than other hydrogen bonding arrangements. The two-body $\left(\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}\right)$ term is nearly identical in the two isomers ( $-4.63 \mathrm{vs}-4.61 \mathrm{kcal} /$ mol, MP2/avtz, BSSE-corrected) and to the free water dimer interaction ( $-4.71 \mathrm{kcal} / \mathrm{mol}$, Table 1), indicating that an almost unperturbed water dimer exists within the cluster, a fact that is also evident by the intermolecular $\mathrm{W}-\mathrm{W}$ separations discussed previously.

The geometries of the third and the fourth minima $\mathrm{AW}_{2}$ m3 and $\mathrm{AW}_{2}$ m 4 (both of $\mathrm{C}_{2 h}$ symmetry) are shown in parts d and f of Figure 2. In m3, two equivalent (AW-Y) bonds are formed, whereas in m 4 , two equivalent ( $\mathrm{AW}-\mathrm{T}$ ) bonds are formed. When compared to the AW-Y and AW-T structures, ${ }^{1}$ both the Y - and T -vdW bond distances in m 3 and m 4 minima increase by approximately $0.04 \AA$. Practically, the geometries of A and W molecules within the m 3 and m 4 clusters are identical with those of AW-Y and AW-T structures, respectively.

With regard to the MP2/avtz interaction energies $\Delta E_{\mathrm{e}}\left[\Delta E_{\mathrm{e}}\right.$ (BSSE)] of m 3 and m 4 , these are $-5.79[-4.91]$ and -5.50 $[-4.61] \mathrm{kcal} / \mathrm{mol}$, respectively (cf. Table 4). As expected, these are almost twice as large as the corresponding $\Delta E_{\mathrm{e}}(\mathrm{BSSE})$ values of the AW -Y and AW -T isomers, viz. $-2.72 \times 2$ and -2.56 $\times 2 \mathrm{kcal} / \mathrm{mol},{ }^{1}$ a result consistent with the fact that the rest of the terms in the many-body expansion (two-body $\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}<$ $0.2 \mathrm{kcal} / \mathrm{mol}$ and three-body $\left.\mathrm{A}-\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}<0.4 \mathrm{kcal} / \mathrm{mol}\right)$ are quite small and repulsive. These terms are responsible for the destabilization of m 3 and m 4 with respect to the global minimum (m1), although the former two have slightly larger two-body $\mathrm{A}-\mathrm{W}_{\mathrm{a}}$ and $\mathrm{A}-\mathrm{W}_{\mathrm{b}}$ terms with respect to m 1 because of the more optimal orientation of the water molecules on either side of (A) when compared to the ring ml structure. Inclusion of ZPE corrections produces $\Delta E_{0}(\mathrm{BSSE})=-3.03 \mathrm{kcal} / \mathrm{mol}$ (m3) and -2.46 (m4) kcal/mol at the MP2/avtz level.
Finally, two saddle point structures, $\mathrm{AW}_{2}$ sp1 and $\mathrm{AW}_{2}$ sp2, both of $\mathrm{C}_{s}$ symmetry (see Figure $2 \mathrm{c}, \mathrm{e}$ ) were located. The sp1 is a second-order saddle point (two imaginary frequencies) constituting a three-member ring, whereras the sp2 is a transition state (one imaginary frequency) resulting from the m 3 by a $90^{\circ}$ rotation of the $\sigma$ plane of one of the water molecules around the acetylene axis. The MP2/avdz interaction energies of sp 1 and $\mathrm{sp} 2, \Delta E_{\mathrm{e}}\left[\Delta E_{\mathrm{e}}(\mathrm{BSSE})\right]\left\{\Delta E_{0}(\mathrm{BSSE})\right\}$, are $-8.25[-6.23]$ $\{-3.57\}$ and $-6.71[-4.47]\{-2.47\} \mathrm{kcal} / \mathrm{mol}$, respectively. Judging from the variation of the energetics of the minima with
$\mathrm{C}_{2} \mathrm{H}_{2}\left(\mathrm{H}_{2} \mathrm{O}\right)_{3}$

a. $\mathrm{AW}_{3}$ _m1



d. $\mathrm{AW}_{3}$-m4



m. $\mathbf{A W}_{3}$ m10 $\left(\mathrm{C}_{\mathrm{s}}\right)$


k. $\mathrm{AW}_{3} \_\mathrm{m} 9\left(\mathrm{C}_{s}\right)$

l. $\mathbf{A W}_{3-} \operatorname{sp} \mathbf{3}\left(\mathrm{C}_{s}\right)$

Figure 3. Geometries of the 10 minima $\left(\mathrm{AW}_{3-} \mathrm{m} n, n=1-10\right)$ and the three saddle points $\left(\mathrm{AW}_{3_{-}} \mathrm{sp} n, n=1-3\right)$ of $\mathrm{AW}_{3}$. Bond distances in $\AA$ at the MP2/avdz level.
basis set, we estimate that these energies are $\sim 0.5 \mathrm{kcal} / \mathrm{mol}$ weaker than the analogous values with the avtz set.
c. $\mathbf{C}_{\mathbf{2}} \mathbf{H}_{\mathbf{2}}\left(\mathbf{H}_{\mathbf{2}} \mathbf{O}\right)_{3}$. We located 10 minima and 3 saddle points on the PES of $\mathrm{AW}_{3}$. The optimal structures, together with representative intermolecular vdW distances, are shown in Figure 3. The depicted configurations were derived by combining the following sets of building units: $\mathrm{W}_{3}+\mathrm{A}, \mathrm{W}_{2}+\mathrm{A}$ $+\mathrm{W}, \mathrm{W}_{2}+\mathrm{AW}, \mathrm{AW}_{2}+\mathrm{W}, \mathrm{AW}+2 \mathrm{~W}$, and $\mathrm{A}+3 \mathrm{~W}$. All structures were fully optimized at the MP2/avdz level. ${ }^{16 c}$ Total energies ( $E_{\mathrm{e}}$ ) and interaction energies $\Delta E_{\mathrm{e}}, \Delta E_{\mathrm{e}}(\mathrm{BSSE})$, and $\Delta E_{0}$ (BSSE) are listed in Table 5; the vibrational frequencies and the IR intensities of the global minimum ( m 1 ) are given in Table 3. ${ }^{16 \mathrm{~d}}$ Finally, the decomposition of interaction energies into two-, three- and four-body terms for 6 selected minima are presented in Table 6. The structures of the stationary points fall into groups according to the various hydrogen bonding networks that are formed and are discussed as such below.

The $\mathrm{AW}_{3}$ m1 and $\mathrm{AW}_{3} \_\mathrm{m} 4$ minima and the $\mathrm{AW}_{3}$ _sp2 saddle point (two imaginary frequencies) form a four member heavy atom ring (see Figure 3 parts a, d, and j). The global (m1) minimum resembles the $\mathrm{W}_{4}$ cyclic tetramer ${ }^{15}$ arrangement but with one W replaced by the $\mathrm{C}-\mathrm{H}$ bond of A . The m 1 and m 4 isomers are quite similar, with the main difference being the relative positions of the hydrogens $\mathrm{H}_{1 \mathrm{a}}$ and $\mathrm{H}_{1 \mathrm{~b}}$ with respect to

TABLE 5: Total Energies $\boldsymbol{E}_{\mathrm{e}}$ (hartree), Interaction Energies $\Delta E_{\mathrm{e}}(\mathrm{kcal} / \mathrm{mol})$, Corrected for BSSE [ $\left.\Delta E_{\mathrm{e}}(\mathrm{BSSE})\right]$, and Zero-Point Energy [ $\Delta E_{0}($ BSSE $)$ ] for AW $_{3}$ at the MP2/avdz Level of Theory

| $\mathrm{AW}_{3}{ }^{a}$ | $E_{\mathrm{e}}$ | $\Delta E_{\mathrm{e}}$ | $\Delta E_{\mathrm{e}}(\mathrm{BSSE})$ | $\Delta E_{0}(\mathrm{BSSE})$ |
| :---: | :---: | :---: | :---: | :---: |
| m 1 | -305.910717 | -21.96 | -17.80 | -11.46 |
| m 2 | -305.909968 | -21.49 | -17.54 | -11.09 |
| m 3 | -305.909879 | -21.43 | -17.47 | -11.05 |
| m 4 | -305.909764 | -21.36 | -17.25 | -11.07 |
| m 5 | -305.909479 | -21.18 | -17.12 | -10.63 |
| m 6 | -305.903302 | -17.30 | -13.81 | -8.47 |
| m 7 | -305.902655 | -16.90 | -13.55 | -8.10 |
| m 8 | -305.901287 | -16.04 | -12.60 | -7.65 |
| $\mathrm{sp1}^{b}$ | -305.901011 | -15.87 | -12.43 | -7.79 |
| sp2 $^{c}$ | -305.899517 | -14.93 | -11.55 | -6.66 |
| m 9 | -305.898487 | -14.28 | -10.92 | -5.91 |
| sp3 $^{b}$ | -305.896882 | -13.28 | -10.11 | -5.80 |
| $\mathrm{~m}^{2} 0$ | -305.895900 | -12.66 | -9.62 | -5.35 |

${ }^{a} \mathrm{~m}$ (minimum) and sp (saddle point), Figure 3. ${ }^{b}$ Three imaginary frequencies. ${ }^{c}$ Two imaginary frequencies.
the almost planar $\mathrm{C}_{1} \mathrm{O}_{a} \mathrm{O}_{b} \mathrm{O}_{c}$ ring: $\mathrm{H}_{1 \mathrm{a}}$ is up and $\mathrm{H}_{1 \mathrm{~b}}$ is down (ud) in m1, whereas $H_{1 a}$ is down and $H_{l b}$ is up (du) in m4. In this notation, the $\mathrm{W}_{3}$ moiety in m 1 and m 4 can be characterized as udu and duu, respectively. A noted difference between (m1, $\mathrm{m} 4)$ and sp 2 is that, in the former, every water acts as a donor-

TABLE 6: Many-Body Decomposition of the Interaction Energies ( $\mathrm{kcal} / \mathrm{mol}$ ) for the $\mathrm{m} 1, \mathrm{~m} 2, \mathrm{~m} 5, \mathrm{~m} 6, \mathrm{~m} 7$, and m 9 Minima of $\mathrm{AW}_{3}$ at the MP2/avdz Level of Theory ${ }^{a}$

| $\mathrm{AW}_{3}{ }^{b}$ | m 1 | m 2 | m 5 | m 6 | m 7 | m 9 |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{~A}-\mathrm{W}_{\mathrm{a}}$ | $-3.29(-2.19)$ | $-1.27(-0.59)$ | $-2.54(-1.82)$ | $-2.49(-1.99)$ | $-3.02(-2.08)$ | $-2.99(-2.21)$ |
| $\mathrm{A}-\mathrm{W}_{\mathrm{b}}$ | $-1.18(-1.03)$ | $-3.07(-2.25)$ | $-2.92(-2.22)$ | $-2.62(-1.54)$ | $-2.91(-2.09)$ | $0.25(0.30)$ |
| $\mathrm{A}-\mathrm{W}_{\mathrm{c}}$ | $-2.94(-1.95)$ | $-1.03(-0.87)$ | $0.34(0.52)$ | $-2.49(-1.99)$ | $-0.29(-0.24)$ | $-2.99(-2.21)$ |
| $\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}$ | $-5.14(-4.14)$ | $-4.60(-3.91)$ | $-4.87(-3.82)$ | $-4.82(-4.02)$ | $-5.12(-4.33)$ | $-5.18(-4.45)$ |
| $\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{c}}$ | $-1.25(-1.15)$ | $-4.43(-3.50)$ | $-4.67(-3.87)$ | $0.47(0.49)$ | $-5.29(-4.48)$ | $1.13(1.31)$ |
| $\mathrm{W}_{\mathrm{b}}-\mathrm{W}_{\mathrm{c}}$ | $-5.14(-4.13)$ | $-4.92(-3.96)$ | $-4.60(-3.79)$ | $-4.82(-4.02)$ | $0.48(0.49)$ | $-5.18(-4.44)$ |
| total two body | $-18.94(-14.60)$ | $-19.32(-15.08)$ | $-19.25(-15.00)$ | $-16.78(-13.07)$ | $-16.52(-12.73)$ | $-14.96(-11.70)$ |
| $\mathrm{A}-\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}$ | $-0.72(-0.81)$ | $0.26(0.20)$ | $-0.90(-0.94)$ | $-0.74(-0.83)$ | $-1.02(-1.15)$ | $0.31(0.29)$ |
| $\mathrm{A}-\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{c}}$ | $-0.55(-0.60)$ | $-0.17(-0.22)$ | $0.33(0.30)$ | $0.23(0.21)$ | $-0.45(-0.45)$ | $-0.14(-0.04)$ |
| $\mathrm{A}-\mathrm{W}_{\mathrm{b}}-\mathrm{W}_{\mathrm{c}}$ | $-0.64(-0.70)$ | $-0.49(-0.55)$ | $0.26(0.27)$ | $-0.74(-0.83)$ | $0.05(0.04)$ | $0.31(0.29)$ |
| $\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}-\mathrm{W}_{\mathrm{c}}$ | $-1.20(-1.25)$ | $-2.17(-2.33)$ | $-2.13(-2.28)$ | $0.37(0.31)$ | $0.47(0.50)$ | $0.23(0.28)$ |
| total three body | $-3.11(-3.36)$ | $-2.57(-2.91)$ | $-2.44(-2.66)$ | $-0.87(-1.14)$ | $-0.95(-1.05)$ | $0.71(0.81)$ |
| $\mathrm{A}-\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}-\mathrm{W}_{\mathrm{c}}$ | $-0.33(-0.27)$ | $-0.04(0.01)$ | $0.06(0.08)$ | $0.01(0.07)$ | $-0.03(0)$ | $-0.07(-0.06)$ |
| relaxation | 0.43 | 0.44 | 0.46 | 0.34 | 0.23 | -164 |
| $\Delta \mathrm{E}_{\mathrm{e}}\left[\Delta \mathrm{E}_{\mathrm{e}}(\mathrm{BSSE})\right]$ | $-21.96(-17.80)$ | $-21.49(-17.54)$ | $-21.18(-17.12)$ | $-17.30(-13.81)$ | $-16.90(-13.55)$ | $-14.28(-10.92)$ |

${ }^{a}$ BSSE-corrected values are shown in parentheses. ${ }^{b} \mathrm{~m} 1$ to m 9 minima as shown in Figure 3.
acceptor to nearest neighbors, whereas in the latter, the $\mathrm{W}_{\mathrm{c}}$ and $\mathrm{W}_{\mathrm{a}}$ fragments are double H -donors and double acceptors, respectively.

The interaction energies $\Delta E_{\mathrm{e}}\left[\Delta E_{\mathrm{e}}(\mathrm{BSSE})\right]\left\{\Delta E_{0}(\mathrm{BSSE})\right\}$ are $-21.96[-17.80]\{-11.46\},-21.36[-17.25]\{-11.07\}$, and $-14.93[-11.55]\{-6.66\} \mathrm{kcal} / \mathrm{mol}$ for $\mathrm{m} 1, \mathrm{~m} 4$, and sp 2 , respectively. It is interesting to point out that the $\Delta E_{0}(\mathrm{BSSE})$ interaction energies for the $\mathrm{m} 2, \mathrm{~m} 3$, and m 4 minima are almost identical ( $-11.09,-11.05$, and $-11.07 \mathrm{kcal} / \mathrm{mol}$, cf. Table 5). The many-body analysis of the global minimum (m1) configuration, listed in Table 6, shows that (BSSE-corrected) total two-, three-, and four-body terms are $-14.60,-3.36$, and -0.27 $\mathrm{kcal} / \mathrm{mol}$, respectively, summing up to $-18.23 \mathrm{kcal} / \mathrm{mol}$, yielding a $\Delta E_{\mathrm{e}}(\mathrm{BSSE})$ of $-17.80 \mathrm{kcal} / \mathrm{mol}$ when the total relaxation energy of $+0.43 \mathrm{kcal} / \mathrm{mol}$ is taken into account. Furthermore, we obtain an interaction of A with the three water molecules of $-9.56(-7.46$ including ZPE) $\mathrm{kcal} / \mathrm{mol}$ by summing together the two-body terms $\mathrm{A}-\mathrm{W}_{i}$, the three-body terms $\mathrm{A}-\mathrm{W}_{i}-\mathrm{W}_{j}$, the 4-body term $\mathrm{A}-\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}-\mathrm{W}_{\mathrm{c}}$, and the relaxation energy.

The $\mathrm{AW}_{3 \_} \mathrm{m} 2, \mathrm{AW}_{3} \_\mathrm{m} 3$, and $\mathrm{AW}_{3 \_} \mathrm{m} 5$ minima all have a cyclic water trimer structure that interacts with A via two vdW bonds (Figure 3 parts b, c, and e). The dihedral angles between the planes $\mathrm{C}_{1} \mathrm{O}_{\mathrm{b}} \mathrm{O}_{\mathrm{a}}$ and $\mathrm{O}_{\mathrm{b}} \mathrm{O}_{\mathrm{a}} \mathrm{O}_{\mathrm{c}}$ are $108.5^{\circ}, 107.4^{\circ}$, and $96.3^{\circ}$ for $\mathrm{m} 2, \mathrm{~m} 3$, and m 5 , respectively. The directions of the hydrogen-bonded H atoms within the $\mathrm{W}_{3}$ fragment are, for a fixed orientation of A with respect to $\mathrm{W}_{3}$, clockwise (m2 and m 3 ) and counterclockwise (m5). The main difference between the m 2 and m 3 isomers lies in the direction of the free H -atom of the $\mathrm{W}_{\mathrm{c}}$ fragment which is "down" for m 2 and "up" for m 3 , with respect to the $\mathrm{O}_{a} \mathrm{O}_{b} \mathrm{O}_{c}$ plane; the direction of the rest of the "free" H atoms for $\mathrm{W}_{\mathrm{a}}$ and $\mathrm{W}_{\mathrm{b}}$ being the same.

The interaction energies $\Delta E_{\mathrm{e}}\left[\Delta E_{\mathrm{e}}(\mathrm{BSSE})\right]\left\{\Delta E_{0}(\mathrm{BSSE})\right\}$ are $-21.49[-17.54]\{-11.09\},-21.43[-17.47]\{-11.05\}$, and $-21.18[-17.12]\{-10.63\} \mathrm{kcal} / \mathrm{mol}$ for $\mathrm{m} 2, \mathrm{~m} 3$, and m 5 , respectively. We note that the BSSE-corrected interaction energies of the five minima $(\mathrm{m} 1-\mathrm{m} 5)$ discussed so far are within an energy range of $0.68 \mathrm{kcal} / \mathrm{mol}(0.83 \mathrm{kcal} / \mathrm{mol}$ with ZPE corrections). The many-body analysis of the configurations of the m 2 and m 5 minima (the m 3 being very similar to m 2 ), shown in Table 6, indicates that the BSSE-corrected total contributions of the two-, three- and four-body terms are -15.08 , $-2.91, \sim 0.0 \mathrm{kcal} / \mathrm{mol}$ (for m 2 ) and $-15.0,-2.66,+0.1 \mathrm{kcal} /$ mol (for m 5 ). The corresponding relaxation terms are similar $(+0.44$ and $+0.46 \mathrm{kcal} / \mathrm{mol}$ respectively for m 2 and m 5$)$. At the m 2 and m 5 minimum configurations, the interaction of A with the three water molecules is $-5.78(-4.24$ including ZPE $)$
$\mathrm{kcal} / \mathrm{mol}(\mathrm{m} 2)$ and $-5.34(-3.78) \mathrm{kcal} / \mathrm{mol}(\mathrm{m} 5)$, values that are about half of the analogous interaction in m 1 . As expected, the water-only $\left(\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}-\mathrm{W}_{\mathrm{c}}\right)$ three-body term is larger for m 2 and $m 5$ than for $m 1$ because of the formation of the water trimer ring in the first two. However, the total four-body term for m 1 is larger than for m 2 because of the formation of the homodromic four heavy-atom ring in the former, again in accordance with previous conclusions ${ }^{14 \mathrm{c}}$ suggesting the maximization of the nonadditivities for homodromic hydrogen-bonding networks.

The $\mathrm{AW}_{3}$ _m6 minimum (Figure 3f) is the first minimum of higher symmetry $\left(C_{s}\right)$ found. The five heavy atoms form a pentagonal structure of trapezoidal topology divided in two equal parts by the $\mathrm{H}_{2 \mathrm{~b}} \ldots \equiv$ (middle of the A-triple bond) line. Its interaction energy $\Delta E_{\mathrm{e}}\left[\Delta E_{\mathrm{e}}(\mathrm{BSSE})\right]\left\{\Delta E_{0}(\mathrm{BSSE})\right\}$ is $-17.30[-$ $13.81]\{-8.47\} \mathrm{kcal} / \mathrm{mol}$, placing it $4.0 \mathrm{kcal} / \mathrm{mol}$ ( 3.0 when ZPE corrections are included) above the global minimum. The total (BSSE-corrected) two-, three- and four-body interactions, listed in Table 6, are $-13.07,-1.14$, and $+0.1 \mathrm{kcal} / \mathrm{mol}$, respectively, showing the dominance of the two-body contributions. The interaction of A with the three water molecules is $-8.67(-6.73$ including ZPE) $\mathrm{kcal} / \mathrm{mol}$.

The $\mathrm{AW}_{3}$ _m7 minimum originates from the interaction of $\mathrm{AW}_{2}$ m1 with a water molecule via a $\mathrm{H}_{1 \mathrm{a}} \ldots \mathrm{O}_{\mathrm{c}}$ vdW bond (Figure 2 g ), whereas the $\mathrm{AW}_{3 \_} \mathrm{m} 8$ and $\mathrm{AW}_{3 \_} \mathrm{m} 10$ minima (the second of $C_{s}$ symmetry) are reminiscent of the $\mathrm{AW}_{2 \_} \mathrm{m} 1$ and $\mathrm{AW}_{2 \_} \mathrm{m} 2$ configurations interacting with a water molecule via a $\mathrm{H}_{2} \ldots \mathrm{O}_{c}$ (Y-like) vdW bond (Figure 3h,m). The interaction energies $\Delta E_{\mathrm{e}}$ $\left[\Delta E_{\mathrm{e}}(\mathrm{BSSE})\right]\left\{\Delta E_{0}(\mathrm{BSSE})\right\}$ are $-16.90[-13.55]\{-8.10\},-16.04$ [ -12.60$]\{-7.65\}$, and $-12.66[-9.62]\{-5.35\} \mathrm{kcal} / \mathrm{mol}$ for m 7 , m 8 , and m 10 , respectively. By summing up the interaction energies of $\mathrm{AW}_{2}$ m1 and $\mathrm{AW}-\mathrm{Y},{ }^{1}$ we obtain $\Delta E_{\mathrm{e}}(\mathrm{BSSE})=$ $-9.60-2.62=-12.22 \mathrm{kcal} / \mathrm{mol}$, just $0.4 \mathrm{kcal} / \mathrm{mol}(0.3 \mathrm{kcal} /$ mol including ZPE) weaker than $\Delta E_{\text {e }}(\mathrm{BSSE})$ of $\mathrm{AW}_{3}$ m8 (Table 5), rationalizing our previous characterization of the m 8 minimum as the combination of the interaction between the $\mathrm{AW}_{2} \mathrm{~m} 1$ moiety and a water molecule in the Y arrangement. The same holds for $\mathrm{AW}_{3} \mathrm{~m} 10$, for which the corresponding difference is just $0.12 \mathrm{kcal} / \mathrm{mol}(0.05 \mathrm{kcal} / \mathrm{mol}$ including ZPE). The analysis of the many-body energy terms for selected $\mathrm{AW}_{3}$ minima suggests that the nonadditive (three-body and higher) terms are larger for the global minimum (m1) for which they amount to $18 \%$, again in accordance with the formation of the homodromic ring incorporating all fragments. The percentage contribution of the many-body terms decreases with increasing separation from the global minimum, becoming $6 \%$ for m 7 .

The last group of the $\mathrm{AW}_{3}$ isomers is composed of the $\mathrm{AW}_{3}$ m m 9 minimum, and the third order (3 imaginary frequencies) saddle structures, $\mathrm{AW}_{3}$ sp1 and $\mathrm{AW}_{3}$ sp3 (Figure 3 parts $\mathrm{k}, \mathrm{i}$, and l ). All three have $C_{s}$ symmetry and four member rings $\mathrm{O}_{\mathrm{a}} \mathrm{O}_{\mathrm{b}} \mathrm{O}_{\mathrm{c}} \equiv$, where $\equiv$ represents the center of the acetylene triple bond. Their energetics, $\Delta E_{\mathrm{e}}\left[\Delta E_{\mathrm{e}}(\mathrm{BSSE})\right]\left\{\Delta E_{0}(\mathrm{BSSE})\right\}$, are $-14.28[-10.92]\{-5.91\},-15.87[-12.43]\{-7.79\}$,and $-13.28[-10.11]$ $\{-5.80\} \mathrm{kcal} / \mathrm{mol}$ for m 9 , sp1, and sp3, respectively. The total three-body interaction for m 9 is positive $(+0.8 \mathrm{kcal} / \mathrm{mol}$, Table 6 ), notably the only destabilizing three-body interaction in all $10 \mathrm{AW}_{3}$ m configurations studied.
d. $\mathbf{C}_{2} \mathbf{H}_{\mathbf{2}}\left(\mathbf{H}_{\mathbf{2}} \mathbf{O}\right)_{\mathbf{4}}$. For the $\mathrm{AW}_{4}$ cluster, we located 30 minima $\left(\mathrm{AW}_{4 \_} \mathrm{m} 1\right.$ to $\left.\mathrm{AW}_{4 \_\mathrm{m} 30}\right)$ and three saddle point structures $\left(\mathrm{AW}_{4}\right.$ _sp1, 2, 3) which are depicted in Figure 4. With the exception of the $\mathrm{m} 19, \mathrm{~m} 26$, and m 27 minima and sp 1 and sp 2 saddle points, which are obtained at the MP2/4-31G level, ${ }^{16 e}$ all other structures were fully optimized at the MP2/avdz level of theory. ${ }^{16 e}$ Out of the 30 minima studied, only two (m28 and m 30 ) belong to the $S_{2}$ and $C_{2 h}$ point groups, whereas the rest lack any symmetry elements $\left(C_{1}\right)$. The minima were obtained by considering the following building units: $\mathrm{W}_{4}+\mathrm{A}, \mathrm{W}_{3}+$ $\mathrm{W}+\mathrm{A}, \mathrm{W}_{3}+\mathrm{AW}, \mathrm{W}_{2}+\mathrm{AW}_{2}, 2 \mathrm{~W}_{2}+\mathrm{A}, \mathrm{W}_{2}+2 \mathrm{~W}+\mathrm{A}$, $4 \mathrm{~W}+\mathrm{A}, \mathrm{AW}_{3}+\mathrm{W}, \mathrm{AW}_{2}+2 \mathrm{~W}$, and $\mathrm{AW}+3 \mathrm{~W}$. The total and interaction energies are listed in Table 7, whereas the manybody analysis for eight selected minima is presented in Table 8. The harmonic vibrational frequencies and IR intensities for m 1 are listed in Table 3, whereas corresponding values for m 8 and $m 9$ are available in the Supporting Information. ${ }^{16 f}$ Because ZPE corrections have been computed only for the $\mathrm{m} 1, \mathrm{~m} 8$, and m 9 minima, the main body of the interaction energy analysis is based on $\Delta E_{\mathrm{e}}$ and $\Delta E_{\mathrm{e}}(\mathrm{BSSE})$ values. Below, we discuss the different structures according to their grouping into similar morphologies.

The $\mathrm{AW}_{4 \_} \mathrm{m} 1, \mathrm{AW}_{4 \_} \mathrm{m} 4, \mathrm{AW}_{4 \_m} \mathrm{~m}$, and $\mathrm{AW}_{4 \_} \mathrm{m} 7$ minima span an energy difference range of $1 \mathrm{kcal} / \mathrm{mol}$ and are composed of a cyclic water tetramer $\left(\mathrm{W}_{4}\right)$ interacting with the A moiety via two adjacent W molecules. In all four previous structures, the four-member oxygen rings $\left(\mathrm{O}_{\mathrm{a}} \mathrm{O}_{\mathrm{b}} \mathrm{O}_{\mathrm{c}} \mathrm{O}_{\mathrm{d}}\right)$ are almost planar with the A molecule located on the $\mathrm{O}_{a} \mathrm{O}_{\mathrm{b}} \mathrm{C}_{1}$ plane, with the dihedral angles between these two planes ranging from 111 (m1) to $91^{\circ}(\mathrm{m} 7)$. Within ( m 1 and m 4 ) and ( m 6 and m 7 ), the $\mathrm{W}_{4}$ fragment assumes the structure of the free $\mathrm{W}_{4}$ global (udud) and first local (uudd) minima, respectively, according to the MP2/avdz//CCSD(T) and MP2-R12 calculations of Schultz et al. ${ }^{17}$ These four minima result from all possible combinations of the direction of the "hydrogen-bonded" H atoms in the $\mathrm{W}_{4}$ ring with respect to a fixed A orientation (clockwise, counterclockwise) and the position of the "free" H-atoms with respect to the $\mathrm{O}_{\mathrm{a}} \mathrm{O}_{\mathrm{b}} \mathrm{O}_{\mathrm{c}} \mathrm{O}_{\mathrm{d}}$ plane, viz., (udud) and (uudd).

The interaction energies $\Delta E_{\mathrm{e}}\left[\Delta E_{\mathrm{e}}(\mathrm{BSSE})\right]$ are $-33.89[-$ 28.01], $-33.51[-27.57],-33.08[-27.25]$, and $-32.96[-26.96]$ $\mathrm{kcal} / \mathrm{mol}$ for $\mathrm{m} 1, \mathrm{~m} 4, \mathrm{~m} 6$, and m 7 , respectively (cf. Table 7). When ZPE corrections are included, the interaction energy $\Delta E_{0^{-}}$ (BSSE) of ml becomes $-18.67 \mathrm{kcal} / \mathrm{mol}$. At the same level of theory (MP2/avdz), the free water tetramer $\mathrm{W}_{4}$ has an interaction energy $\Delta E_{0}($ BSSE $)=16.03 \mathrm{kcal} / \mathrm{mol} .{ }^{15}$ Therefore, the difference of $-2.64 \mathrm{kcal} / \mathrm{mol}$ closely resembles the binding of $\mathrm{AW}_{4 \_\mathrm{m}} \mathrm{m}$ with respect to $\mathrm{W}_{4}+\mathrm{A}$. This value is practically equal to the sum of $\Delta E_{0}(\mathrm{BSSE})$ interactions for $\mathrm{AW}-\mathrm{T}+\mathrm{AW}-\mathrm{Y}(-2.59$ $\mathrm{kcal} / \mathrm{mol}) .{ }^{1}$ The analysis of many-body energy terms for m 1 (cf. Table 8), reveals that the sum of two-, three-, four-, and five-body interactions are $-21.90,-6.65,-0.51$, and +0.01 $\mathrm{kcal} / \mathrm{mol}$, respectively, with the total relaxation R being +1.03 $\mathrm{kcal} / \mathrm{mol}$.

The $\mathrm{AW}_{4 \_} \mathrm{m} 2, \mathrm{AW}_{4 \_} \mathrm{m} 3$, and $\mathrm{AW}_{4 \_} \mathrm{m} 5$ minima are also composed of water tetramers, but with two diagonal water molecules of the $\mathrm{W}_{4}$ ring interacting with A through two vdW bonds, forming puckered four-member rings with dihedral angles between the $\mathrm{O}_{\mathrm{a}} \mathrm{O}_{\mathrm{b}} \mathrm{O}_{\mathrm{d}}$ and $\mathrm{O}_{\mathrm{b}} \mathrm{O}_{\mathrm{c}} \mathrm{O}_{\mathrm{d}}$ planes of $150.4^{\circ}, 150.6^{\circ}$, and $155.0^{\circ}$, respectively (see Figure 4 parts b, c, and e). As with the previous group, these minima arise from the different choices of the direction of the "hydrogen-bonded" (clockwise, counterclockwise) and "free" (up, down, and planar (p)) H atoms. Although the m 2 and m 3 minima are practically degenerate, their energy difference being $0.11 \mathrm{kcal} / \mathrm{mol}(0.12 \mathrm{kcal} / \mathrm{mol}$ including BSSE), they correspond to structures that have different directions as regards the "hydrogen-bonded" H atoms within the $\mathrm{W}_{4}$ fragment (clockwise for m 2 and counterclockwise for m3). It should be mentioned that the corresponding structures (dpud) and (ddup) of the gas-phase water tetramer $\mathrm{W}_{4}$, lying $\sim 1.2 \mathrm{kcal} / \mathrm{mol}$ higher than the global minimum, are saddle points of second and first order on the water tetramer PES, ${ }^{17}$ but are stabilized as minima m 2 and m 3 in $\mathrm{AW}_{4}$. The corresponding interaction energies $\Delta E_{\mathrm{e}}\left[\Delta E_{\mathrm{e}}(\mathrm{BSSE})\right]$ are $-33.68[-27.47]$, $-33.57[-27.35]$, and $-33.46[-27.25] \mathrm{kcal} / \mathrm{mol}$ for m 2 , m 3 , and m5, see Table 7. Note that BSSE corrections alter the order of the minima to $\mathrm{m} 4, \mathrm{~m} 2, \mathrm{~m} 3$, and $\mathrm{m} 5=\mathrm{m} 6$.

The $\mathrm{AW}_{4}$ m8 and $\mathrm{AW}_{4 \_} \mathrm{m} 18$ minima are formed by the interaction of a perturbed "cage" water tetramer with A (Figure 4 parts h and t ). The cage $\mathrm{W}_{4}$ structure was found to be a minimum on the PES of the ASP-P (anisotropic site potential neglecting all nonpairwise additive effects) potential by Gregory and Clary; ${ }^{18}$ however, at the MP2/dzp level, the cage morphology was found to collapse to the global ring minimum. ${ }^{18}$ In the $\mathrm{W}_{4}$ cage, the $\mathrm{H}_{\text {la }}$ hydrogen forms a vdW bond with the $\mathrm{O}_{\mathrm{c}}$ atom, but in the m 8 and m 18 minima, this interaction is altered because of the intervention of the A molecule and the formation of two vdW bonds between $\mathrm{W}_{4}$ and A . The difference between m 8 and m18 lies again in the topology of the hydrogen bonding network. Furthermore, the $W_{4}$ ring is more "puckered" than the one in $\mathrm{m} 2, \mathrm{~m} 3$, and m 5 minima: the dihedral angles between the two oxygen planes $\mathrm{O}_{\mathrm{a}} \mathrm{O}_{\mathrm{b}} \mathrm{O}_{\mathrm{d}}$ and $\mathrm{O}_{\mathrm{b}} \mathrm{O}_{\mathrm{d}} \mathrm{O}_{\mathrm{c}}$ are $95.6^{\circ}(\mathrm{m} 8)$ and $110.4^{\circ}$ (m18).

The interaction energies $\Delta E_{\mathrm{e}}\left[\Delta E_{\mathrm{e}}(\mathrm{BSSE})\right]$ are $-30.81[-$ $25.01] \mathrm{kcal} / \mathrm{mol}$ for m 8 and $-28.63[-23.01] \mathrm{kcal} / \mathrm{mol}$ for m 18 . Including ZPE corrections, the m8 interaction energy becomes $-15.89 \mathrm{kcal} / \mathrm{mol}$. The energy difference of $2.18 \mathrm{kcal} / \mathrm{mol}(2.0$ $\mathrm{kcal} / \mathrm{mol}$ including the BSSE) between the two isomers is rather due to the extra H bond in m 8 .

The $\mathrm{AW}_{4 \_} \mathrm{m} 9$ minimum (Figure 4i) forms a cyclic pentamer ring that incorporates $A$ and resembles the free $W_{5}$ duduu configuration. ${ }^{15}$ The replacement of any other water molecule by A collapses to the previous minimum. All six heavy atoms are approximately on the same plane and the interaction energy $\Delta E_{\mathrm{e}}\left[\Delta E_{\mathrm{e}}(\mathrm{BSSE})\right]\left\{\Delta E_{0}(\mathrm{BSSE})\right\}$ is $-30.64[-25.14]\{-16.57\}$ $\mathrm{kcal} / \mathrm{mol}$. Correcting for BSSE and ZPE was found to invert the ordering between m 8 and m 9 . As expected from the homodromic ring configuration of this structure, all nonadditive (three-body and higher) energy terms are negative; that is, they all stabilize this particular network.

The common feature among the $\mathrm{AW}_{4 \_} \mathrm{m} 10, \mathrm{AW}_{4 \_} \mathrm{m} 16$, and $\mathrm{AW}_{4 \_} \mathrm{m} 17$ minima (Figure 4 parts j, p, and q) and the two firstorder saddle point structures, $\mathrm{AW}_{4 \_} \mathrm{sp} 1$ and $\mathrm{AW}_{4}$ sp2 (Figure 4 parts r and s ), is a water tetramer with identical topology, whose gas-phase configuration has symmetry $C_{s}$ This structure was previously identified by Clementi et al. ${ }^{19}$ and later discussed by Gregory and Clary. ${ }^{18}$ The intervention of the A molecule destroys the $C_{s}$ symmetry of the $\mathrm{W}_{4}$ fragment in the $\mathrm{m} 10, \mathrm{~m} 16$,

## $\mathbf{C}_{2} \mathbf{H}_{\mathbf{2}}\left(\mathbf{H}_{\mathbf{2}} \mathrm{O}\right)_{\mathbf{4}}$








m. $\mathrm{AW}_{4}$ m13

n. $\mathrm{AW}_{4}$ m14

o. $\mathrm{AW}_{4}$ m15


u. $\mathrm{AW}_{4}$ m19





bb. $\mathrm{AW}_{4}$ m26







Figure 4. Geometries of the 30 minima $\left(\mathrm{AW}_{4} \mathrm{~m} n, n=1-30\right)$ and three saddle points $\left(\mathrm{AW}_{4} \mathrm{sp} n, n=1-3\right)$ of $\mathrm{AW}_{4}$ cluster. Bond distances in $\AA$ at the MP2/avdz level.
and m17 minima, but this symmetry is maintained in the sp1 and sp2 conformations. The difference between these five structures lies in which water molecule(s) of the $\mathrm{W}_{4}$ fragment interacts with the acetylene. The dihedral angles between the oxygen planes $\mathrm{O}_{\mathrm{a}} \mathrm{O}_{\mathrm{b}} \mathrm{O}_{\mathrm{d}}$ and $\mathrm{O}_{\mathrm{b}} \mathrm{O}_{\mathrm{d}} \mathrm{O}_{\mathrm{c}}$ are $97.3^{\circ}, 99.3^{\circ}$, and $115.3^{\circ}$ in the $\mathrm{m} 10, \mathrm{~m} 16$, and m 17 minima, respectively.

The interaction energies $\Delta E_{\mathrm{e}}\left[\Delta E_{\mathrm{e}}(\mathrm{BSSE})\right]$ are $-30.63[-$ 24.70], $-28.78[-23.28]$, and $-28.73[-23.27] \mathrm{kcal} / \mathrm{mol}$ for $\mathrm{m} 10, \mathrm{~m} 16$, and m 17 , respectively; note that m 16 and m 17 are practically degenerate at this level of theory. The many-body analysis for the m 10 isomer is given in Table 8. We remind the reader that the sp 1 and sp 2 structures were calculated at the MP2/4-31G level, and they might be lying higher than m17 and near to m18.

In the $\mathrm{AW}_{4 \_} \mathrm{m} 11$ to $\mathrm{AW}_{4 \_\mathrm{m}} 15$ series (Figure 4 parts $\mathrm{k}-\mathrm{o}$ ), a water trimer ring interacts with A through two vdW bonds forming an almost planar four member ring. The dihedral angles between the oxygen "triangles" and the "squares" range from 95 to $113^{\circ}$. It is interesting to note that all five minima lie within an energy range of $0.3 \mathrm{kcal} / \mathrm{mol}$ (with or without BSSE corrections) at the MP2/avdz level (Table 7). The five structures can be divided in three groups (m11), (m12 and m14), and (m13 and m 15 ) according to the vdW bonds formed between the five molecules. The difference between m 12 and $\mathrm{m} 14, \mathrm{~m} 13$ and m 15 is the $\mathrm{H}_{1 \mathrm{~d}}$ hydrogen direction (up or down) with respect to the $\mathrm{O}_{\mathrm{b}} \mathrm{O}_{\mathrm{c}} \mathrm{O}_{\mathrm{d}}$ plane.

The $\mathrm{AW}_{4 \_} \mathrm{m} 19, \mathrm{AW}_{4 \_\mathrm{m}} 20, \mathrm{AW}_{4 \_\mathrm{m}} 21$, and $\mathrm{AW}_{4 \_} \mathrm{m} 23$ minima (Figure 4 parts $u$, $v, w$, and $y$ ) can be considered as $\mathrm{AW}_{3} \mathrm{~m} 1+\mathrm{W}_{\mathrm{d}}$ and $\mathrm{AW}_{3 \_} \mathrm{m} 4+\mathrm{W}_{\mathrm{d}}$. The four-member rings $\mathrm{O}_{\mathrm{a}} \mathrm{O}_{\mathrm{b}} \mathrm{O}_{\mathrm{c}} \mathrm{C}_{1}$ in $\mathrm{m} 19, \mathrm{~m} 21$, and m 23 are almost planar, but in m20, the dihedral angle between the $\mathrm{O}_{\mathrm{a}} \mathrm{O}_{\mathrm{b}} \mathrm{O}_{\mathrm{c}}$ and $\mathrm{O}_{\mathrm{b}} \mathrm{O}_{\mathrm{c}} \mathrm{C}_{1}$ planes is $-28.0^{\circ}$. The similarity between the m 21 and m 23 structures is

TABLE 7: Total Energies $\boldsymbol{E}_{\mathrm{e}}$ (hartree), Interaction Energies $\Delta E_{\mathrm{e}}(\mathrm{kcal} / \mathrm{mol})$, Corrected for BSSE [ $\Delta E_{\mathrm{e}}($ BSSE $\left.)\right]$, and Zero-Point Energy [ $\Delta_{0}($ BSSE $)$ ] for $\mathbf{A W}_{4}$ at the MP2/avdz Level of Theory

| $\mathrm{AW}_{4}{ }^{a}$ | $E_{\mathrm{e}}$ | $\Delta E_{\mathrm{e}}$ | $\Delta E_{\mathrm{e}}(\mathrm{BSSE})$ | $\Delta E_{0}(\mathrm{BSSE})$ |
| :--- | :---: | :---: | :---: | :---: |
| m 1 | -382.190641 | -33.89 | -28.01 | -18.67 |
| m 2 | -382.190305 | -33.68 | -27.47 |  |
| m 3 | -382.190136 | -33.57 | -27.35 |  |
| m 4 | -382.190037 | -33.51 | -27.57 |  |
| m 5 | -382.189952 | -33.46 | -27.25 |  |
| m 6 | -382.189348 | -33.08 | -27.25 |  |
| m 7 | -382.189163 | -32.96 | -26.96 |  |
| m 8 | -382.185732 | -30.81 | -25.01 | -15.89 |
| m 9 | -382.185459 | -30.64 | -25.14 | -16.57 |
| m 10 | -382.185443 | -30.63 | -24.70 |  |
| m 11 | -382.185102 | -30.41 | -24.65 |  |
| m 12 | -382.184782 | -30.21 | -24.52 |  |
| m 13 | -382.184743 | -30.19 | -24.49 |  |
| m 14 | -382.184703 | -30.16 | -24.41 |  |
| m 15 | -382.184659 | -30.13 | -24.36 |  |
| m 16 | -382.182504 | -28.78 | -23.28 |  |
| m 17 | -382.182414 | -28.73 | -23.27 |  |
| m 18 | -382.182267 | -28.63 | -23.01 |  |
| m 20 | -382.181013 | -27.85 | -22.52 |  |
| m 21 | -382.178107 | -26.02 | -20.82 |  |
| m 22 | -382.177154 | -25.43 | -20.41 |  |
| m 23 | -382.177139 | -25.42 | -20.28 |  |
| m 24 | -382.176850 | -25.23 | -20.26 |  |
| m 25 | -382.176520 | -25.03 | -19.95 |  |
| m 28 | -382.175652 | -24.48 | -19.60 |  |
| m 29 | -382.170289 | -21.12 | -16.58 |  |
| m 30 | -382.165349 | -18.02 | -13.87 |  |
| $\mathrm{sp} 33^{b}$ | -382.161304 | -15.48 | -11.65 |  |

${ }^{a} \mathrm{~m}$ (minimum) and sp (saddle point), Figure $4 .{ }^{b}$ Four imaginary frequencies.

TABLE 8: Many-Body Decomposition of Interaction Energies $\Delta E_{\mathrm{e}}(\mathrm{kcal} / \mathrm{mol})$ of the $\mathrm{m} 1, \mathrm{~m} 2, \mathrm{~m} 3, \mathrm{~m} 4, \mathrm{~m} 8, \mathrm{~m} 9, \mathrm{~m} 10$, and m 11 Minima of $\mathrm{AW}_{4}$ at the MP2/avdz Level of Theory ${ }^{a}$

| $\mathrm{AW}_{4}{ }^{\text {b }}$ | m1 | m2 | m3 | m4 | m8 | m9 | m10 | m11 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| A- | -1.05(-0.37) | 0.29(0.56) | $-1.36(-1.14)$ | -2.71(-2.02) | $-2.57(-1.66)$ | -3.34(-2.19) | -3.44(-2.45) | $-2.88(-1.95)$ |
| $\mathrm{A}-\mathrm{W}_{\mathrm{b}}$ | -2.95(-2.20) | -3.22(-2.32) | -3.29(-2.40) | -2.88(-2.17) | -0.26(-0.18) | -0.72(-0.64) | -0.33(-0.17) | -1.10(-0.96) |
| $\mathrm{A}-\mathrm{W}_{\mathrm{c}}$ | $-0.67(-0.58)$ | $-1.17(-0.97)$ | 0.44(0.70) | 0.23(0.38) | -3.16(-2.12) | $-0.70(-0.60)$ | -2.52(-1.65) | -3.37(-2.32) |
| $\mathrm{A}-\mathrm{W}_{\mathrm{d}}$ | $-0.30(-0.24)$ | $-2.34(-1.61)$ | $-2.28(-1.51)$ | 0.17(0.24) | $-0.96(-0.84)$ | $-2.88(-1.85)$ | $-0.40(-0.30)$ | $0.10(0.17)$ |
| $\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}$ | -4.99(-4.07) | -4.88(-3.97) | -4.60(-3.48) | -4.74(-3.59) | -4.65(-4.00) | -5.00(-3.95) | -3.76(-3.26) | -4.66(-3.53) |
| $\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{c}}$ | -1.75(-1.59) | -1.85(-1.69) | -1.90(-1.73) | $-1.69(-1.52)$ | -0.82(-0.61) | -1.03(-0.96) | 0.10(0.27) | -1.07(-0.95) |
| $\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{d}}$ | $-4.87(-3.78)$ | $-4.23(-3.25)$ | $-4.37(-3.25)$ | -4.84(-3.81) | $-4.82(-3.73)$ | $-0.80(-0.74)$ | $-4.50(-3.63)$ | -4.75(-4.00) |
| $\mathrm{W}_{\mathrm{b}}-\mathrm{W}_{\mathrm{c}}$ | -4.84(-3.73) | -4.44(-3.32) | -4.53(-3.59) | -5.06(-4.05) | -1.92(-1.38) | $-4.90(-3.81)$ | -3.91(-3.02) | $-5.25(-4.36)$ |
| $\mathrm{W}_{\mathrm{b}}-\mathrm{W}_{\mathrm{d}}$ | -1.71(-1.57) | -2.09(-1.88) | -2.08(-1.87) | -1.68(-1.54) | -4.41(-3.57) | -0.94(-0.86) | -4.51(-3.34) | -4.57(-3.81) |
| $\mathrm{W}_{\mathrm{c}}-\mathrm{W}_{\mathrm{d}}$ | $-4.85(-3.76)$ | $-4.68(-3.55)$ | $-4.63(-3.67)$ | -5.04(-4.00) | $-4.40(-3.71)$ | -4.98(-3.94) | -4.49(-3.92) | 0.54(0.69) |
| total two body | -27.99(-21.90) | -28.63(-22.00) | -28.59(-21.94) | -28.25(-22.08) | -27.98(-21.81) | $-25.28(-19.56)$ | -27.77(-21.47) | 27.01(-21.01) |
| $\mathrm{A}-\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}$ | 0.50(0.46) | 0.51(0.47) | -0.68(-0.75) | $-1.09(-1.16)$ | 0.27(0.26) | -0.73(-0.80) | 0.19(0.14) | -0.61(-0.62) |
| $\mathrm{A}-\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{c}}$ | 0.03(0.02) | 0.02(0.06) | 0.02(0.06) | 0.13(0.11) | -0.39(-0.49) | $-0.18(-0.20)$ | $-0.24(-0.32)$ | $-0.46(-0.52)$ |
| $\mathrm{A}-\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{d}}$ | -0.32(-0.29) | 0.50(0.43) | -0.46(-0.52) | 0.28(0.27) | $-0.51(-0.53)$ | -0.43(-0.46) | -0.38(-0.40) | 0.34(0.32) |
| $\mathrm{A}-\mathrm{W}_{\mathrm{b}}-\mathrm{W}_{\mathrm{c}}$ | $-0.59(-0.61)$ | -0.53(-0.60) | 0.51(0.47) | 0.34(0.34) | $-0.19(-0.22)$ | -0.15(-0.18) | $-0.25(-0.30)$ | $-0.61(-0.70)$ |
| $\mathrm{A}-\mathrm{W}_{\mathrm{b}}-\mathrm{W}_{\mathrm{d}}$ | 0.00(0.00) | -0.21(-0.27) | -0.18(-0.24) | 0.05(0.07) | -0.02(0.00) | -0.14(-0.15) | 0.02(-0.01) | 0.02(0.05) |
| $\mathrm{A}-\mathrm{W}_{\mathrm{c}}-\mathrm{W}_{\mathrm{d}}$ | -0.10(-0.12) | $-0.53(-0.59)$ | 0.55(0.49) | 0.11(0.09) | $-0.51(-0.58)$ | $-0.59(-0.63)$ | 0.19(0.17) | 0.10(0.08) |
| $\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}-\mathrm{W}_{\mathrm{c}}$ | -1.38(-1.46) | -1.42(-1.49) | -1.19(-1.28) | -1.50(-1.55) | 0.34(0.32) | -1.30(-1.34) | 0.42(0.35) | -1.17(-0.18) |
| $\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}-\mathrm{W}_{\mathrm{d}}$ | -1.35(-1.43) | $-1.04(-1.14)$ | $-1.66(-1.74)$ | $-1.47(-1.53)$ | -1.97(-2.13) | -0.23(-0.25) | $-1.76(-1.85)$ | $-2.20(-2.35)$ |
| $\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{c}}-\mathrm{W}_{\mathrm{d}}$ | $-1.53(-1.60)$ | $-1.20(-1.28)$ | -1.44(-1.50) | -1.38(-1.46) | $-0.85(-0.94)$ | -0.27(-0.28) | 0.30(0.31) | 0.16(0.15) |
| $\mathrm{W}_{\mathrm{b}}-\mathrm{W}_{\mathrm{c}}-\mathrm{W}_{\mathrm{d}}$ | -1.54(-1.62) | -1.66(-1.74) | -1.03(-1.12) | $-1.36(-1.44)$ | 0.47(0.43) | -1.26(-1.29) | -1.96(-2.07) | 0.44(0.44) |
| Total three body | $-6.27(-6.65)$ | -5.57(-6.14) | $-5.55(-6.14)$ | -5.89(-6.26) | -3.36(-3.88) | -5.27(-5.57) | -3.47(-3.96) | $-4.00(-4.34)$ |
| $\mathrm{A}-\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}-\mathrm{W}_{\mathrm{c}}$ | 0.00(0.06) | 0.05(0.06) | 0.03(0.04) | -0.01(0.04) | -0.01(0.02) | $-0.18(-0.15)$ | $-0.07(-0.04)$ | $-0.28(-0.23)$ |
| $\mathrm{A}-\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}-\mathrm{W}_{\mathrm{d}}$ | 0.02(0.04) | 0.08(0.12) | -0.18(-0.13) | 0.00(0.05) | -0.01(0.03) | -0.12(-0.10) | -0.04(0.00) | 0.00(0.05) |
| $\mathrm{A}-\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{c}}-\mathrm{W}_{\mathrm{d}}$ | $-0.05(-0.04)$ | 0.02(0.04) | 0.04(0.05) | 0.02(0.04) | $-0.22(-0.15)$ | $-0.11(-0.10)$ | $-0.04(-0.03)$ | 0.02(0.04) |
| $\mathrm{A}-\mathrm{W}_{\mathrm{b}}-\mathrm{W}_{\mathrm{c}}-\mathrm{W}_{\mathrm{d}}$ | $-0.07(-0.04)$ | $-0.20(-0.16)$ | 0.11(0.12) | 0.05(0.03) | -0.2(0.00) | $-0.13(-0.12)$ | -0.03(0.02) | 0.04(0.03) |
| $\begin{aligned} & \mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}- \\ & \mathrm{W}_{\mathrm{c}}-\mathrm{W}_{\mathrm{d}} \end{aligned}$ | $-0.61(-0.53)$ | $-0.51(-0.44)$ | $-0.51(-0.44)$ | $-0.59(-0.50)$ | 0.04(0.07) | -0.19(-0.18) | 0.07(0.09) | 0.00(0.04) |
| total four body | $-0.71(-0.51)$ | $-0.57(-0.38)$ | $-0.52(-0.36)$ | $-0.53(-0.34)$ | $-0.22(-0.03)$ | $-0.72(-0.65)$ | -0.11(0.04) | $-0.21(-0.06)$ |
| $\begin{gathered} \mathrm{A}-\mathrm{W}_{\mathrm{a}}-\mathrm{W}_{\mathrm{b}}- \\ \mathrm{W}_{\mathrm{c}}-\mathrm{W}_{\mathrm{d}} \end{gathered}$ | 0.06(0.01) | 0.05(0.03) | 0.04(0.03) | 0.05(0.01) | 0.04(0.00) | $-0.04(-0.04)$ | 0.01(-0.01) | 0.06(0.01) |
| relaxation | 1.03 | 1.03 | 1.05 | 1.11 | 0.72 | 0.68 | 0.71 | 0.75 |
| $\begin{aligned} & \Delta E_{\mathrm{e}}\left[\Delta E_{\mathrm{e}}\right. \\ & \quad(\mathrm{BSSE})] \end{aligned}$ | -33.89(-28.01) | -33.68(-27.47) | -33.57(-27.35) | -33.51(-27.57) | -30.81(-25.01) | -30.64(-25.14) | -30.63(-24.70) | -30.41(-24.65) |

${ }^{a}$ BSSE-corrected values are shown in parentheses. ${ }^{b}$ All minimum structures are shown in Figure 4.
striking, and it is remarkable that these constitute separate minima (at least at the MP2/avdz level of theory), their difference being in the directions of the $\mathrm{H}_{1 \mathrm{a}}$ and $\mathrm{H}_{1 \mathrm{~b}}$ hydrogen atoms with respect to the $\mathrm{O}_{\mathrm{a}} \mathrm{O}_{b} \mathrm{O}_{\mathrm{c}}$ plane. Finally, the interaction energies $\Delta E_{\mathrm{e}}\left[\Delta E_{\mathrm{e}}(\mathrm{BSSE})\right]$ are $-27.85[-22.52],-26.02[-$ 20.82], and -25.42 [-20.28] kcal/mol for $\mathrm{m} 20, \mathrm{~m} 21$, and m 23 , respectively. As expected, for $\mathrm{m} 21(\mathrm{~m} 23)$, these interaction energies are equal to the energies of $\mathrm{AW}_{3} \_\mathrm{m} 1(\mathrm{~m} 4)+\mathrm{AW}-\mathrm{Y}$. Again note that the m19 isomer has been computed at the MP2/ 4-31G level of theory.

The $\mathrm{AW}_{4 \_} \mathrm{m} 22, \mathrm{AW}_{4 \_} \mathrm{m} 24$, and $\mathrm{AW}_{4 \_m} 25$ minima (Figure 4 parts $x, z$, and aa) can be viewed as the $\mathrm{AW}_{3} \_\mathrm{m} 2, \mathrm{AW}_{3} \_\mathrm{m} 3$, and $\mathrm{AW}_{3} \mathrm{~m} 5$ structures, respectively, in which the A molecule is interacting with an additional water molecule in an Y arrangement (as an acceptor to acetylene). Correspondingly, the $\mathrm{AW}_{4 \_\mathrm{m}} 26$ and $\mathrm{AW}_{4 \_\mathrm{m}} 27$ isomers (obtained at the MP2/4-31G level) are seen as $\mathrm{AW}_{3}$ m 3 and $\mathrm{AW}_{3} \_\mathrm{m} 2$ isomers, respectively, interacting with a water molecule in a T fashion (Figures 4bb and 4 cc$)$. The first three minima lie within an energy range of $0.40 \mathrm{kcal} / \mathrm{mol}(0.46 \mathrm{kcal} / \mathrm{mol}$ including BSSE corrections), having interaction energies in the range -25.43 to $-25.03 \mathrm{kcal} /$ mol ( -20.41 to $-19.95 \mathrm{kcal} / \mathrm{mol}$ including BSSE corrections, cf. Table 7).

The $\mathrm{AW}_{4 \_} \mathrm{m} 28, \mathrm{AW}_{4 \_\mathrm{m}} 29$, and $\mathrm{AW}_{4 \_} \mathrm{m} 30$ (Figure 4 parts dd, ee, and ff) are the energetically highest and most well separated minima studied in the present work, their energy differences being approximately $3 \mathrm{kcal} / \mathrm{mol}$ (Table 7). The m 28 $\left(S_{2}\right)$ and m30 ( $C_{2 h}$ ) minima can be thought of as "double" $\mathrm{AW}_{2} \mathrm{~m} 1$ and $\mathrm{AW}_{2} \mathrm{~m} 2$ configurations, whereas the m 29 minimum can be thought of as a combination of the $\mathrm{AW}_{2} \_\mathrm{m} 1$ and $\mathrm{AW}_{2 \_} \mathrm{m} 2$ structures. Although all of the heavy atoms in m 30 and the four oxygen atoms in m 28 are coplanar by
symmetry, in the latter, all six heavy atoms are also, practically, coplanar. The interaction energies $\Delta E_{\mathrm{e}}\left[\Delta E_{\mathrm{e}}(\mathrm{BSSE})\right]$ are $-24.48[-$ 19.60], $-21.12[-16.58]$, and $-18.02[-13.87] \mathrm{kcal} / \mathrm{mol}$ for $\mathrm{m} 28, \mathrm{~m} 29$, and m 30 , respectively. By summing up twice the interaction energies of $\mathrm{AW}_{2}$ m1 and $\mathrm{AW}_{2}$ m2, we obtain $-24.04[-19.20],-20.94[-16.48]$, and $-17.84[-13.76] \mathrm{kcal} /$ mol , values that are very close indeed to the interaction energies of $\mathrm{m} 28, \mathrm{~m} 29$, and m 30 minima.

Finally, the $\mathrm{AW}_{4}$ sp3 (Figure 4 gg ) is a fourth order saddle point (four imaginary frequencies) of $C_{2 h}$ symmetry, and is composed of two independent Y and two independent T arrangements (Figure 1). Alternatively, it can be seen as a "superposition" of the $\mathrm{AW}_{2}$ m3 and $\mathrm{AW}_{2}$ m4 arrangements (Figure 2d,f). Its interaction energy $\Delta E_{\mathrm{e}}\left[\Delta E_{\mathrm{e}}(\mathrm{BSSE})\right]$ is $-15.48[-$ $11.65] \mathrm{kcal} / \mathrm{mol}$, as contrasted to the sum of interaction energies of two AW-Y and two AW-T isomers, which is $-13.30[-$ $9.56] \mathrm{kcal} / \mathrm{mol} .{ }^{1}$

## 4. Summary

Extended parts of the multidimensional PESs of a number of acetylene-water, $\mathrm{AW}_{x}, x=2,3$, and 4 clusters have been probed by ab initio calculations. Using chemical intuition and extended systematic searches on the multidimensional PESs with smaller basis sets to obtain candidates of stationary points, we have located 4,10 , and 30 minima for $\mathrm{AW}_{2}, \mathrm{AW}_{3}$, and $\mathrm{AW}_{4}$ clusters, respectively. Although we cannot claim that every possible isomer has been located within the energy range considered here, we can ascertain that the global minima for these clusters have been identified. The energy separation between the located stationary points (minima and saddle points) is shown in Figure 5. As expected, ${ }^{20}$ the density of the local


Figure 5. Interaction energies $\Delta E_{\mathrm{e}}(\mathrm{BSSE})$ and $\Delta E_{0}(\mathrm{BSSE})$ of all stationary points of the $\mathrm{AW}_{2}, \mathrm{AW}_{3}$, and $\mathrm{AW}_{4}$ clusters relative to the corresponding global minina (m1).


Figure 6. Variation of the "average" dissociation energies, $\Delta E_{\mathrm{e}} /(x+$ $1), \Delta E_{\mathrm{e}}(\mathrm{BSSE}) /(x+1)$, and $\Delta E_{0}(\mathrm{BSSE}) /(x+1)$ of the global minima (m1) with cluster size $x$ at the MP2/avdz level.
minima is increasing dramatically with cluster size, a manifestation of the energetic competition between different hydrogen bonding networks. For instance, for $\mathrm{AW}_{4}$, the first 28 minima lie within a range of $\sim 8 \mathrm{kcal} / \mathrm{mol}$. To this end, corrections for BSSE and ZPE can alter the order of these closely spaced minima. Such is the case for the range of m10-m15 of $\mathrm{AW}_{4}$ where six minima are packed within an energy range of just $0.34 \mathrm{kcal} / \mathrm{mol}$ and the change in order for m 8 and m 9 upon including ZPE corrections. The "average" interaction energies, $\Delta E_{\mathrm{e}} /(x+1), \Delta E_{\mathrm{e}}(\mathrm{BSSE}) /(x+1)$, and $\Delta E_{0}(\mathrm{BSSE}) /(x+1)$ for the $\mathrm{AW}, \mathrm{AW}_{2}, \mathrm{AW}_{3}$, and $\mathrm{AW}_{4}$ global mimima are plotted in Figure 6. The analysis of the many-body interactions for several hydrogen bonded networks indicates that there are different requirements as regards an empirical interaction potential needed to reproduce the relative cluster energetics: usually the low(er) lying energy structures are stabilized because of the maximization of the nonadditive (mainly the three-body) components of the interaction energies, whereas higher lying structures are composed from mainly two-body interactions. Therefore, empirical interaction potentials that aim to reproduce the relative cluster energetics will need to include many-body effects because they need to describe both the ring configurations (large nonadditivities, mainly three-body) as well as other more compact configurations (mainly two-body) with the same accuracy.

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Supporting Information Available: Optimal geometries of all $\mathrm{AW}_{2}$ structures (Table 1). Harmonic vibrational frequencies and IR intensities of $\mathrm{AW}_{2}$ stationary points (Table 2). Cartesian coordinates of all $\mathrm{AW}_{3}$ structures (Table 3). Harmonic vibrational frequencies and IR intensities of $\mathrm{AW}_{3}$ structures (Table 4). Cartesian coordinates of all $\mathrm{AW}_{4}$ structures (Table 5). Harmonic vibrational frequencies and IR intensities of two AW 4 structures (Table 6). This material is available free of charge via the Internet at http://pubs.acs.org.

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