

LETTERS

Effects of Na₂SO₄ and NaClO₄ on the Molecular Organization of H₂O

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We study thermodynamic behaviors of 1-propanol (abbreviated as 1P) in ternary 1P–salt–H₂O in order to elucidate the effects of salt on the molecular organization of H₂O. We determine the excess partial molar enthalpy of 1P, H_{1P}^E , as a function of the mole fraction of 1P, x_{1P} , and the salt. We examine the x_{1P} dependence of H_{1P}^E , and changes induced by the presence of salt are used to elucidate the effect of the salt on the molecular organization of H₂O. We previously found that an NaCl ion pair binds 7 to 8 molecules of H₂O, but has no effect on H₂O outside of the hydration shell. Na₂SO₄ and NaClO₄, on the other hand, modify the molecular organization of H₂O each in different ways. In particular, ClO₄[−] seems to participate in the hydrogen bond network of H₂O and to retard the degree of fluctuation inherent in liquid H₂O.

Omta et al. in a recent article¹ concluded that the addition of Mg(ClO₄)₂, NaClO₄, or Na₂SO₄ has no effect on the rotational dynamics of H₂O molecules outside the first hydration shell, by means of femto second pump–probe spectroscopy. This observation is potentially of great importance, since SO₄^{2−} and ClO₄[−] are known to have strong effects on respectively kosmotropic and chaotropic properties.² Therefore, direct dynamic information on aqueous solutions of these ions could be useful for molecular level understandings of the ubiquitous, but still poorly understood, Hofmeister effects.^{2–4} In this letter, we report thermodynamic data suggesting that, although NaCl, which exhibits weak chaotropic/kosmotropic effects, has indeed no effect, NaClO₄ and Na₂SO₄ do have qualitatively different effects on the molecular organization of H₂O. Our strategy^{5–7} is to determine the excess partial molar enthalpy of 1-propanol (1P) in a 1P–salt–H₂O system. We titrate a small enough portion

of 1P successively into a mixed solvent of salt + H₂O. For each titration, we determine the thermal response of the system and obtain the excess partial molar enthalpy of 1P, H_{1P}^E , a change in enthalpy of the entire system per mole of added 1P. Thus, H_{1P}^E provides information about the actual enthalpic situation of 1P in the mixture. The composition dependence of H_{1P}^E thus determined, or the slope of the tangent at a given mole fraction of 1P, x_{1P} , signifies the effect of additional 1P on the actual enthalpic situation of existing 1P. We name this quantity as the enthalpic interaction, H_{1P-1P}^E , which provides a measure of the 1P–1P interaction in terms of enthalpy. This information is obtained purely experimentally and is completely free from presumption inherent in any model system.

We then inspect the x_{1P} dependence of H_{1P-1P}^E on addition of a salt. The resulting changes on H_{1P-1P}^E and on its x_{1P} dependence are used as a probe to elucidate the effect of an added salt on the molecular organization of H₂O. As explained briefly below, this methodology was powerful in elucidating the effect of NaCl on the molecular organization of H₂O. Namely, we showed earlier⁸ that an Na⁺ and a Cl[−] ions bind the total of 7–8 molecules of H₂O but leave the bulk H₂O away

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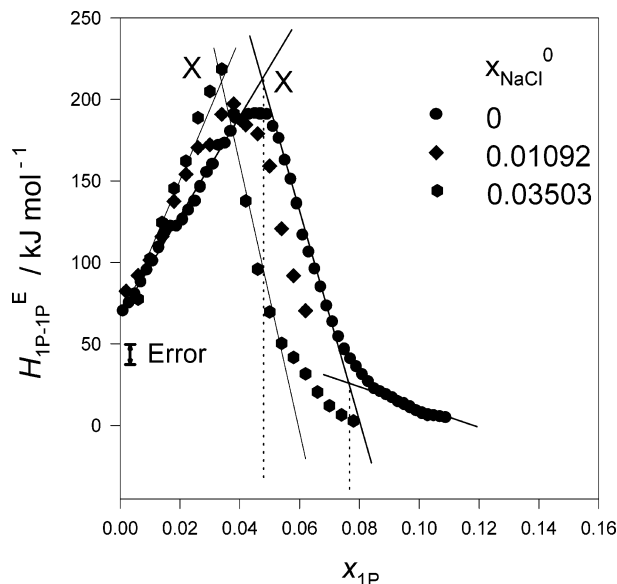


Figure 1. Enthalpic interaction H_{1P-1P}^E in 1-propanol–NaCl–H₂O at 25 °C, as a function of mole fraction of 1-propanol, x_{1P} , given the initial mole fraction of NaCl, x_{NaCl}^0 . In the region from $x_{1P} = 0$ to point X, H₂O keeps the hydrogen bond percolation nature.

from NaCl ions unaltered. This is in complete accordance with the conclusion by a first principle simulation study.⁹ However, our recent studies using the same methodology on Na₂SO₄ and NaClO₄ indicate that their effects are different from each other and in turn different from that by NaCl.

Figure 1 shows the enthalpic interaction, H_{1P-1P}^E , in 1P–NaCl–H₂O.⁸ The region from the start at $x_{1P} = 0$ to point X in the figure describes the process of 1P modifying the molecular organization of H₂O. In this region, 1P enhances the hydrogen bond network of H₂O in its immediate vicinity. At the same time, the hydrogen bond probability of bulk H₂O away from 1P is progressively reduced.^{5–7} However, one of the characteristics of liquid H₂O that the hydrogen bond network is bond-percolated¹⁰ is still retained. This proceeds to point X whereupon the hydrogen bond probability of bulk H₂O away from 1P is reduced to the percolation threshold of ice Ih bond connectivity, and the percolation nature of the network is lost. A new mixing scheme pertinent in the intermediate region sets in.^{5–7} As is evident in Figure 1, addition of NaCl shifts point X to a smaller value of x_{1P} without changing the value of H_{1P-1P}^E , in comparison with the uncertainty. Its value at $x_{1P} = 0$ also remains unchanged.

Earlier we observed¹¹ that the x_{1P} -dependence of H_{1P-1P}^E and that of the partial molar entropy-volume cross fluctuation calculated using thermal expansivity data are identical in a relative scale. This suggests that the underlying mechanism for the x_{1P} -dependence of H_{1P-1P}^E and that of the partial molar entropy-volume cross fluctuation share the same cause. One of the characteristics of liquid H₂O manifests itself in a negative contribution in the entropy-volume cross fluctuation due to putative formation/destruction of ice-like patches in addition to the positive part as normal liquids. Hence, the increase in H_{1P-1P}^E is associated with the decrease of the negative contribution or the increase in the positive part in the cross fluctuation. Returning to Figure 1, the shift of point X to a smaller value of x_{1P} suggests that the presence of NaCl retards the availability of H₂O toward 1P for the latter to drive the system to point X. From the proportionality factor, about 7.5 molecules of H₂O are bound to one ion pair of NaCl and are protected from interacting with 1P. From the fact that the values of H_{1P-1P}^E

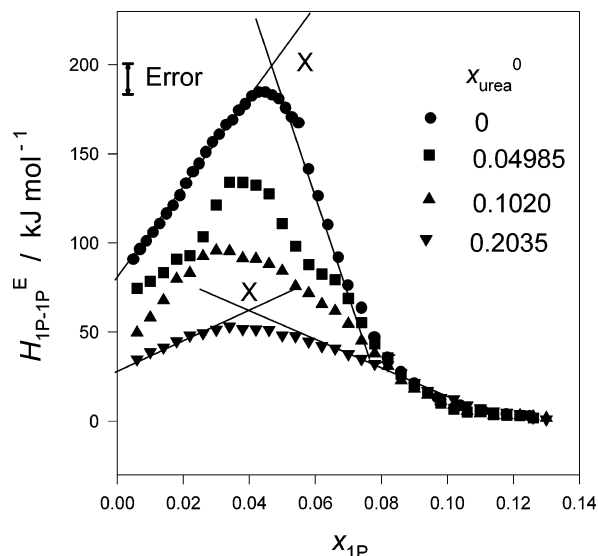


Figure 2. Enthalpic interaction H_{1P-1P}^E in 1-propanol–urea–H₂O at 25 °C, as a function of mole fraction of 1-propanol, x_{1P} , given the initial mole fraction of urea, x_{urea}^0 . In the region from $x_{1P} = 0$ to point X, H₂O keeps the hydrogen bond percolation nature.

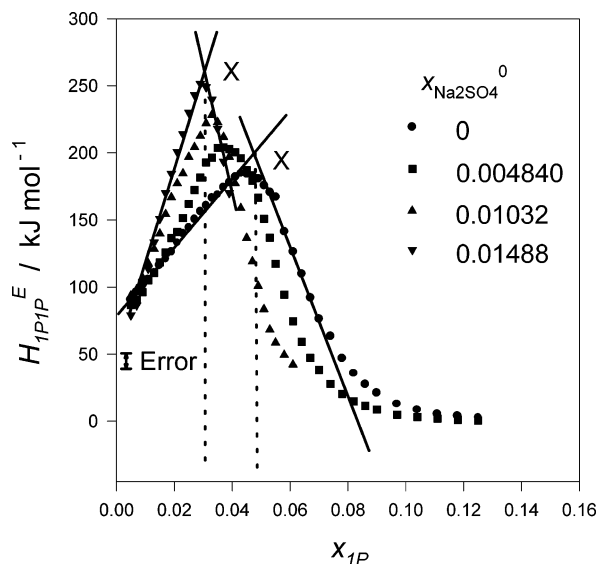


Figure 3. Enthalpic interaction H_{1P-1P}^E in 1-propanol–Na₂SO₄–H₂O at 25 °C, as a function of mole fraction of 1-propanol, x_{1P} , given the initial mole fraction of Na₂SO₄, $x_{Na_2SO_4}^0$. In the region from $x_{1P} = 0$ to point X, H₂O keeps the hydrogen bond percolation nature.

and hence of the partial molar entropy-volume cross fluctuation were independent of the NaCl concentration both at $x_{1P} = 0$, and the end point, point X, we concluded that NaCl has very little effect on the molecular organization of H₂O beyond the hydration shell.⁸

Figure 2 shows the same type of graph for 1P–urea–H₂O.¹² The fact that the locus of point X remains almost unchanged suggests that the percolated network of hydrogen bonding remains intact in the presence of urea. However, the value of H_{1P-1P}^E and hence the degree of fluctuation inherent in liquid H₂O is retarded. Thus, urea seems to participate in forming the hydrogen bond network of H₂O and to reduce the degree of fluctuation by breaking the proton donor/acceptor symmetry in pure H₂O. Another simulation study¹³ concluded that urea leads to stiffening of the short-time dynamics of both urea and H₂O.

Figures 3 and 4 show our recent analogous results for Na₂SO₄ and NaClO₄. The behaviors of the x_{1P} dependence of H_{1P-1P}^E for either case are quite different from that for NaCl. Thus,

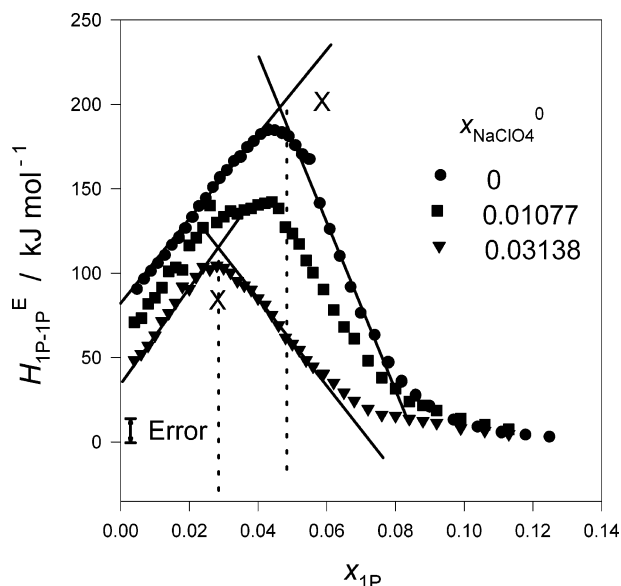


Figure 4. Enthalpic interaction H_{1P-1P}^E in 1-propanol– NaClO_4 – H_2O at 25 °C, as a function of mole fraction of 1-propanol, x_{1P} , given the initial mole fraction of NaClO_4 , $x_{\text{NaClO}_4}^0$. In the region from $x_{1P} = 0$ to point X, H_2O keeps the hydrogen bond percolation nature.

contrary to the claim by Omta et al.,¹ Na_2SO_4 and NaClO_4 seem to modify the molecular organization of H_2O . The shifts of point X to smaller values of x_{1P} may partly be due to Na^+ ions in both salts. However, for NaClO_4 , its effect is closer to that of urea, whereas Na_2SO_4 seems to either retard the negative contribution of the cross fluctuation strongly or to increase the normal positive part of the entropy-volume cross fluctuation. SO_4^{2-} and ClO_4^- ions are in almost at the opposite end of the Hofmeister series.^{2–4} Thus, it is natural that their effects on H_2O could be qualitatively different. A fuller account including other salts in the Hofmeister series is forthcoming. Attempts at correlating these types of plots (H_{1P-1P}^E vs x_{1P}) to the mechanism underlying Hofmeister effects await the accumulation of more data for comparative analysis.

Materials and Methods

Na_2SO_4 (AnalaR, Analytical Reagent, BDH) and NaClO_4 (Fluka, >99.0%) were used as supplied. Stock solutions were made by weighing in a dry atmosphere and successively diluted by weight as required. 1-Propanol (ACROS Organics, 99.5%) and (Fluka, 99.8%) as supplied were treated in a dry nitrogen atmosphere.

The excess partial molar enthalpies of 1P, H_{1P}^E , for Na_2SO_4 were determined by a homemade titration calorimeter of a similar design to an LKB Bromma 8700.¹⁴ The uncertainty is estimated as ± 0.03 kJ mol⁻¹. For NaClO_4 , a TAM titration

calorimeter was used. The uncertainty is ± 0.005 kJ mol⁻¹. The value of $H_{1P}^E = (\partial H^E / \partial n_{1P})$ is approximated as $(\delta H^E / \delta n_{1P})$, where δH_{1P}^E is the determined thermal response of the system on titrating a small amount δn_{1P} . The goodness of this approximation is tested by changing the size of δn_{1P} and by checking that the quotient $(\delta H^E / \delta n_{1P})$ converges. About 1 mol % of δn_{1P} to the total mole in the cell seems satisfactory.¹⁵ The uncertainties quoted above are obtained considering this procedure. The resulting data of H_{1P}^E are plotted on a graph. The data points are typically in the increment in mole fraction of 1P, $\delta x_{1P} = 0.003$. We then draw the best smooth curve through all of the data points by aid of a flexible ruler, and read the value off the curve drawn at a chosen interval δx_{1P} . The enthalpic interaction, H_{1P-1P}^E , is defined as, $H_{1P-1P}^E = N(\partial H_{1P}^E / \partial n_{1P}) = (1 - x_{1P})(\partial H_{1P}^E / \partial x_{1P})$. The last derivative is then approximated by $(\delta H_{1P}^E / \delta x_{1P})$, where δH_{1P}^E is the difference in H_{1P}^E read off at a given interval δx_{1P} . The most appropriate size of δx_{1P} for the best approximation was determined to be 0.01 earlier by trial and error.¹⁶ The caution against the method of curve fitting a function to H_{1P}^E vs x_{1P} data and arithmetically differentiating the fitted function is given earlier.⁵ Thus, we now obtained H_{1P-1P}^E data purely experimentally and are completely free from presumptions inherent in a fitting function.

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