

# Effective Rate Constants for the Surface Reaction between Solid Methanol and Deuterium Atoms at 10 K

Akihiro Nagaoka,<sup>†</sup> Naoki Watanabe,\* and Akira Kouchi

*Institute of Low Temperature Science, Hokkaido University, N19W8, Kita-ku, Sapporo, Hokkaido 060-0819, Japan*

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The surface reactions of CH<sub>3</sub>OH, CH<sub>2</sub>DOH, and CHD<sub>2</sub>OH with cold D atoms at 10 K were investigated using an atomic beam source and FTIR. Methyl-deuterated isotopologues CH<sub>2</sub>DOH, CHD<sub>2</sub>OH, and CD<sub>3</sub>OH were produced by exposure of amorphous solid CH<sub>3</sub>OH to D atoms at 10 K, and the pseudo-first-order rates for the reactions CH<sub>3</sub>OH + D → CH<sub>2</sub>OH + HD, CH<sub>2</sub>DOH + D → CHDOH + HD, and CHD<sub>2</sub>OH + D → CD<sub>2</sub>OH + HD were estimated. The ratios of the reaction rates of the second and third reactions to the first reaction were 0.69 ± 0.11 and 0.52 ± 0.14, respectively. The difference in reaction rates is thought to be due to a secondary kinetic isotope effect on the H-abstraction reaction from the methyl side by D atoms.

## 1. Introduction

The surface chemistry of solids at low temperatures is of physical and chemical interest as it allows observation of the quantum mechanical tunneling effect and also plays an important role in the astrophysical environment.<sup>1,2</sup> At low temperatures, the de Broglie wavelength of a light particle such as a hydrogen atom becomes comparable to typical reaction barrier widths, and thus the reactions can proceed via the atom-tunneling effect. Since the temperature in the cold regions of space is as low as 10 K, the tunneling reaction of hydrogen (H) atoms becomes very efficient in these areas. For example, in interstellar molecular clouds, which typically have a temperature of 10 K, tunneling reactions of H atoms with carbon monoxide (CO) on solid surfaces are a key process in the formation of interstellar formaldehyde (H<sub>2</sub>CO) and methanol (CH<sub>3</sub>OH) molecules, which are precursors of more complex organic molecules.<sup>3–7</sup>

Surface reactions between the deuterium (D) atom and other molecules at low temperatures are of interest because the kinetic isotope effect helps us to obtain information such as reaction barrier heights and widths. For reactions of CO with H and D atoms on solid water surfaces, the difference in the reaction rate constants due to the kinetic isotope effect has been determined experimentally as  $k_D/k_H \sim 0.1$  at 15 K, and the barrier height and width for the reaction are constrained.<sup>8</sup> From the astrochemical viewpoint, the surface reaction CH<sub>3</sub>OH + D is especially interesting because this reaction process on cold interstellar grain surfaces contributes greatly to deuterium fractionation in interstellar methanol.<sup>9,10</sup>

Although low-temperature reactions of solid CH<sub>3</sub>OH with radicals such as CH<sub>3</sub> have been studied extensively,<sup>11</sup> information about its reactions with H/D atoms is limited. Hiraoka and co-workers performed an experiment involving exposure of solid CH<sub>3</sub>OH to D atoms at 10 K and reported that CH<sub>2</sub>DOH was the only product obtained after 300 min of exposure.<sup>12</sup> Recently, we performed a similar experiment at 10 K and observed efficient formation of CH<sub>2</sub>DOH, CHD<sub>2</sub>OH, and CD<sub>3</sub>OH at

D-exposure times as short as a few minutes.<sup>9</sup> The difference between these two sets of results is thought to be due to a large difference in the flux of the D atoms. The flux in our experiment was measured as approximately  $1 \times 10^{14}$  atoms cm<sup>-2</sup> s<sup>-1</sup>, while the other group did not measure the flux of D atoms in their experiment but deduced it to be  $< 10^{13}$  atoms cm<sup>-2</sup> s<sup>-1</sup> based on the literature. In this paper, we report further experimental results for the surface reactions of methanol (CH<sub>3</sub>OH, CH<sub>2</sub>DOH, CHD<sub>2</sub>OH, CH<sub>3</sub>OD, and CD<sub>3</sub>OD) with D atoms on amorphous solid samples at 10 K.

## 2. Experimental Section

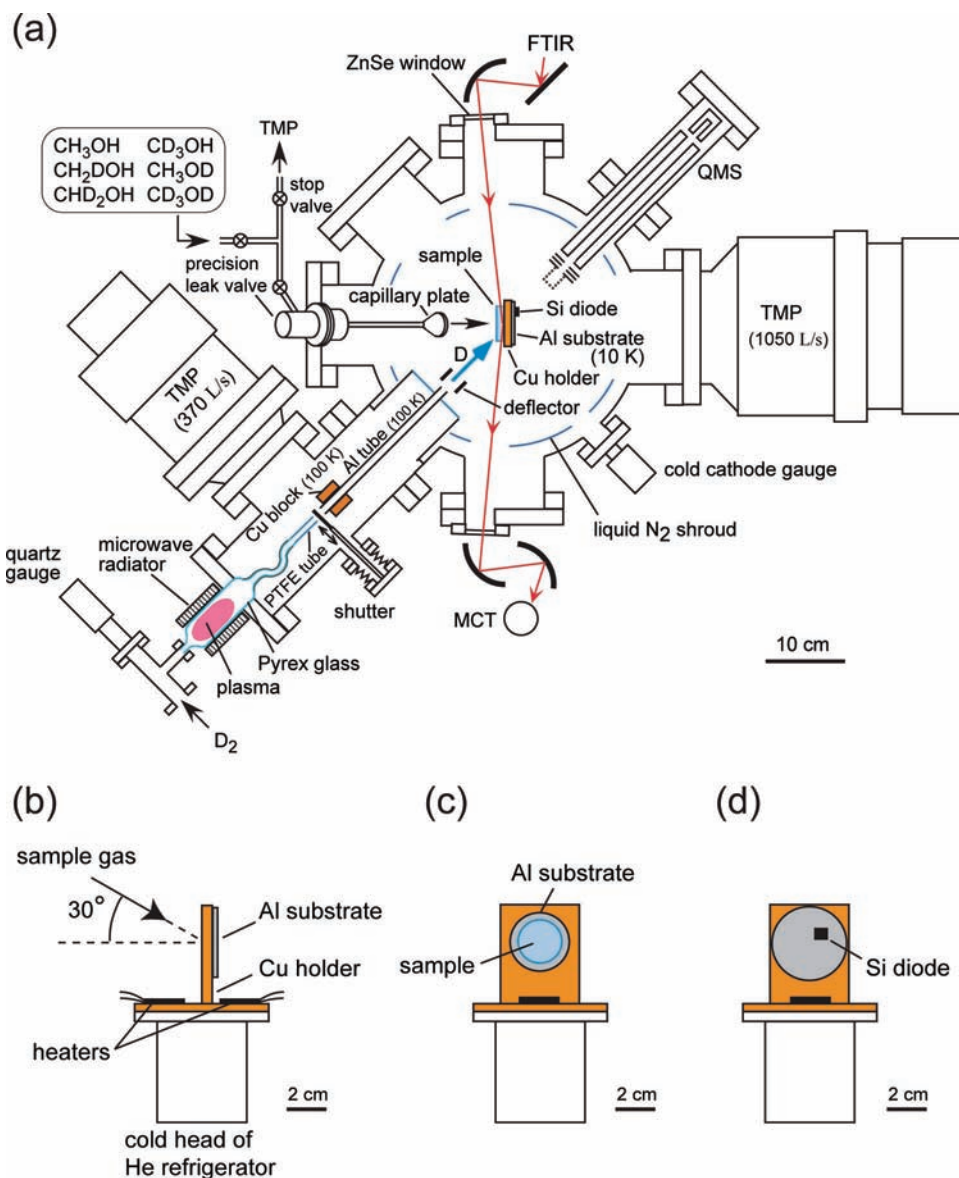
**2.1. Experimental Setup.** Experiments were performed using the ASURA (apparatus for surface reactions in astrophysics) system (Figure 1). The details of the ASURA system are described in a previous report.<sup>7</sup> Briefly, the system consists of a main chamber, an atomic source chamber, a Fourier transform infrared (FTIR) spectrometer (Spectrum One, Perkin-Elmer), and a quadrupole mass spectrometer (QMS) (M-400QA-M, Anelva).

At the center of the main chamber, a mirror-finished aluminum (Al) substrate (30 mm in diameter) is mounted on an oxygen-free copper (Cu) holder attached to the cold head of a He refrigerator (RDK-415, Sumitomo Heavy Industries). A solid sample was vapor-deposited on the substrate at 10 K through a capillary plate (J5022-09, Hamamatsu Photonics; deposition angle of 30° to the surface normal of the substrate). The sample liquids used were CH<sub>3</sub>OH (99.8% purity, Kishida Chemical Co., Ltd.), CH<sub>2</sub>DOH (98% purity, Isotec), CHD<sub>2</sub>OH (98% purity, Isotec), CH<sub>3</sub>OD (99.5% purity, Aldrich), and CD<sub>3</sub>-OD (99.95% purity, Acros). Before deposition, we conducted several freeze-thaw cycles to eliminate contaminating gases in the sample liquids. The thickness of all solid samples was 3 monolayer (ML), which was measured by FTIR (see section 2.3 for details). The deposition rate was approximately 0.18 monolayer min<sup>-1</sup>. In this paper, ML represents the amount of molecules deposited, on the assumption that 1 monolayer =  $1 \times 10^{15}$  molecule cm<sup>-2</sup>, because the solid samples were not crystalline but amorphous solids.<sup>13</sup>

After deposition, the sample (10 K) was exposed to cold D atoms. D atoms were generated by a microwave discharge of

\* Corresponding author. E-mail: watanabe@lowtem.hokudai.ac.jp. Tel.: +81-11-706-5501. Fax: +81-11-706-7142.

<sup>†</sup> JSPS Research Fellow.



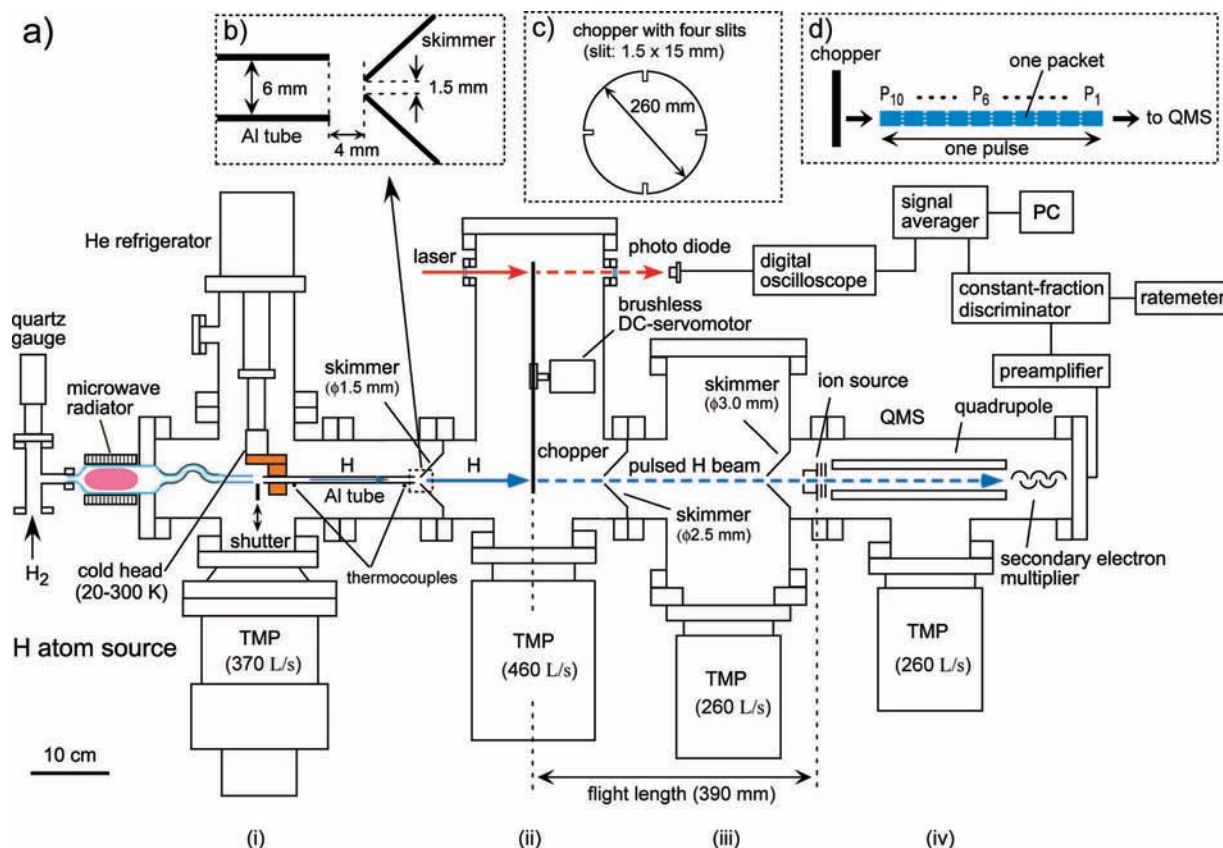
**Figure 1.** Schematic diagram of the ASURA system. (a) Top view. Side views of the substrate and the sample holder at the center of the main chamber from the (b) MCT side, (c) front side, and (d) back (TMP) side. The deposition angle of sample gases was  $30^\circ$  to the surface normal of the substrate.

$D_2$  (99.5% purity, Sumitomo Seika).  $D_2$  gas of 0.3 Torr, measured using a quartz gauge (GC-210, Vacuum Products), was introduced into a Pyrex glass discharge tube after passing through a cold trap cooled with liquid nitrogen to eliminate contaminating gases. Microwaves of 2.45 GHz (MP-201, Arios) with a power of  $\sim 100$  W was fed to a water-cooled copper radiator<sup>14</sup> surrounding the Pyrex glass discharge tube. Before exposure, D atoms were cooled to  $\sim 100$  K by multiple collisions with the inner wall of a cold aluminum tube (Al: 99.70%) connected to a He refrigerator (V204SC5L, Daikin). The temperature of the Al tube can be measured using two Au/Fe 0.07% Chromel thermocouples attached to the vicinity of the refrigerator and the exit of the Al tube (see Figure 2) and is controllable from 20 to 300 K. In the present experiment, the tube temperature was set at 100 K. The flux of D atoms was measured by QMS using the same method as Hidaka et al.<sup>8</sup> and was approximately  $1 \times 10^{14}$  atoms  $cm^{-2} s^{-1}$ . Charged particles and metastable atoms which may be formed in the discharge tube can be eliminated by a deflector of  $100 V cm^{-1}$  located at the exit of the Al tube. We used a photodiode (IRD AXUV-100G) to check whether UV photons from the  $D_2$  plasma

in the discharge tube reached the substrate, but no UV photons were detected.

The change in the chemical composition of the sample was measured using FTIR, in situ, with a resolution of  $4 cm^{-1}$ . The incident angle of the IR beam was  $83^\circ$  to the surface normal of the substrate. The main chamber and the atomic source chamber were evacuated by turbomolecular pumps, TG-900M (Osaka Vacuum, 900 L/s for  $N_2$  and 1050 L/s for  $H_2$ ) and TURBOVAC-340M (Leybold, 400 L/s for  $N_2$  and 370 L/s for  $H_2$ ), and the base pressures were  $3-4 \times 10^{-10}$  and  $5-6 \times 10^{-9}$  Torr, respectively. The pressures during D atom exposure were  $2-3 \times 10^{-7}$  and  $8-9 \times 10^{-6}$  Torr for the main chamber and the atomic source chamber, respectively. The temperature of the Al substrate was measured using a Si-diode sensor (DT-470, Lakeshore, accuracy  $\pm 0.2$  K) and was controlled with two ceramic heaters (100 W) attached at the cold head of the He refrigerator.

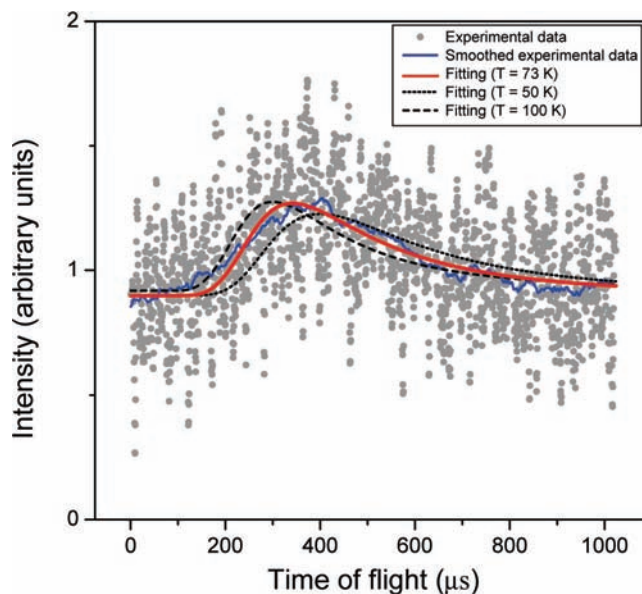
**2.2. Measurement of Translational Temperature of Hydrogen (Deuterium) Atoms.** We used cold D atoms of  $\sim 100$  K to simulate surface reactions on dust grains in cold interstellar clouds. To confirm the kinetic temperature of atoms and evaluate



**Figure 2.** (a) Schematic diagram of the experimental system for TOF measurement (side view). The apparatus consists of (i) atomic source chamber, (ii) chopping chamber, (iii) differential pumping chamber, (iv) QMS chamber, and signal processing system. (b) Close-up of the vicinity of the first skimmer. (c) Chopper disk. (d) Conceptual illustration of the H pulse beam. We regard one pulse as a series of 10 packets in analysis of the TOF spectrum (see section 2.2 for details).

the performance of our cooling system, we performed a time-of-flight (TOF) measurement for H atoms using the experimental system shown in Figure 2. The apparatus consists of a signal processing system and four vacuum chambers: (i) atomic source chamber, (ii) chopping chamber, (iii) differential pumping chamber, and (iv) QMS chamber. Three skimmers ( $\phi$  1.5, 2.5, and 3.0 mm) were mounted between the chambers. The chambers were differentially pumped using turbomolecular pumps: (i) TURBOVAC-340M, Leybold, 370 L/s for H<sub>2</sub>; (ii) STP-451, Seiko Seiki, 460 L/s; and (iii), (iv) PT-300, Mitsubishi Heavy Industries, 260 L/s. Base pressures were (i)  $7 \times 10^{-8}$ , (ii)  $5 \times 10^{-8}$ , (iii)  $5 \times 10^{-9}$ , and (iv)  $7 \times 10^{-9}$  Torr. The pressures in the chambers during TOF measurement were (i)  $5 \times 10^{-5}$ , (ii)  $7 \times 10^{-7}$ , (iii)  $6 \times 10^{-9}$ , and (iv)  $9 \times 10^{-9}$  Torr when the source pressure of H<sub>2</sub> was 0.3 Torr.

The atomic H beam was collimated by the first skimmer ( $\phi$  1.5 mm) and pulsed in the chopping chamber by a four-slit disk chopper mounted on a brushless DC-servomotor (Faulhaber, 3564K048B-K179-K1155). The timing of H pulse formation was measured using a He–Ne laser and a photodiode. The flight time was started at the moment of detection of the laser pulse. The frequency of the pulse formation and the pulse width were 66.66 Hz and 100  $\mu$ s, respectively. After ionization at the ion source of a QMS (AQA-360, Anelva), the H pulses reached a secondary electron multiplier (SEM). The pulse signals were integrated by a signal averager (ELK-5120AVE, Electronica). We defined the flight length as the distance between the disk chopper and the ion source. The TOF of H<sup>+</sup> from the ion source to the SEM was very short compared with that from the disk chopper and the ion source, because the H<sup>+</sup> formed in the ion source was accelerated with 3 V at the ion lens.



**Figure 3.** TOF spectrum of H atoms (Al tube: 50 K) and the fitting results by the least-square method using the modified Maxwell–Boltzmann transmission function (2). The blue continuous line represents the smoothed experimental data. The best-fitting result was at  $T = 73$  K (red continuous), which clearly reproduces the experiment better than the fitted curves at  $T = 50$  (black dotted) and 100 K (black dashed).

Figure 3 shows the TOF spectrum of H atoms. The velocity distribution function of gas molecules that have come through the skimmer is represented by the Maxwell–Boltzmann trans-

mission function:

$$f(v) = 4\pi \left( \frac{m}{2\pi k_B T} \right)^{3/2} v^3 \exp\left(-\frac{mv^2}{2k_B T}\right) \quad (1)$$

where  $v$  is the translational velocity of the gas molecule,  $m$  is the mass of the gas molecule,  $k_B$  is the Boltzmann constant, and  $T$  is the temperature. Under the present experimental conditions, eq 1 cannot be used directly for analysis of the TOF spectrum for H atoms, because the pulse width of 100  $\mu$ s is comparable to the TOF of H ( $\sim 270 \mu$ s at  $T = 100$  K); that is, the TOF spectrum is broadened to a longer time because of the wide time span in one pulse. We therefore used a modified function that takes account of the time difference:

$$f(v) = A \sum_{n=0}^9 4\pi \left( \frac{m}{2\pi k_B T} \right)^{3/2} \left( \frac{l}{t - 10n} \times 10^6 \right)^3 \times \exp\left(-m \left( \frac{l}{t - 10n} \times 10^6 \right)^2 \frac{1}{2k_B T}\right) \quad (2)$$

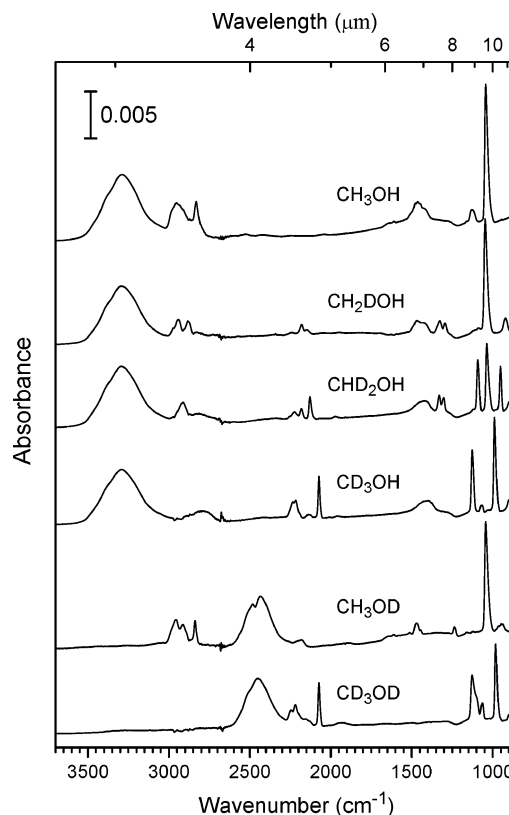
where  $A$  is the proportional constant,  $t$  is the time ( $\mu$ s), and  $l$  is the flight length of 0.39 m. In eq 2, we regard one pulse as a series of 10 packets  $P_{n+1}$  ( $n = 0-9$ , see Figure 2d). The packet  $P_{n+1}$  has a delay of  $10 \times n \mu$ s ( $n = 0-9$ ) in its start time.

The fitting of the results by the least-square method using eq 2 for H atoms is shown in Figure 3. The best fit for the measured TOF spectrum was obtained at 73 K; hence, the H atoms are effectively cooled to near the tube temperature. We also measured the TOF of D atoms and H<sub>2</sub> and D<sub>2</sub> molecules at 50 K and obtained the same result as that for H. When the Al tube was set at 20 K, the intensity of the TOF signal for H was 1 order of magnitude smaller than that at 50 K. This indicates that approximately 90% of the H atoms stuck and recombined at the inner wall of the Al tube.

We also evaluated the performance of the previously used atom cooling system (see Figure 1 in Hidaka et al.<sup>6</sup>) using the TOF measurement system. In our past experiments, a polytetrafluoroethylene (PTFE) tube was used for atom cooling. The use of PTFE is known to suppress hydrogen atom recombination ( $H + H \rightarrow H_2$ ) at the surface.<sup>15</sup> The PTFE tube was tightly covered with a copper tube which was connected to the cold head. The translational temperature of H through the PTFE tube was found to be approximately 200 K when the outer copper tube was at 30 K, indicating that this type of cooling tube does not work well. The reason for this may be the difference in the rates of thermal expansion of copper and PTFE. Since the rate of the thermal expansion of PTFE is approximately 1 order of magnitude greater than that of copper at 30 K,<sup>16</sup> thermal contact between the PTFE and the copper deteriorates at low temperatures, which means that the PTFE is not cooled to the temperature of the copper tube.

**2.3. Infrared Spectra and Relative Integrated Band Strengths of Solid CH<sub>3</sub>OH and Deuterated Isotopologues (CH<sub>2</sub>DOH, CHD<sub>2</sub>OH, CD<sub>3</sub>OH, CH<sub>3</sub>OD, and CD<sub>3</sub>OD).** To calculate the column densities of solid CH<sub>3</sub>OH and its deuterated isotopologues, the integrated band strengths of these species are required. As these values have not yet been reported for the solid deuterated isotopologues, we measured the integrated band strengths of pure solid CH<sub>2</sub>DOH, CH<sub>3</sub>OD, CHD<sub>2</sub>OH, CD<sub>3</sub>OH, and CD<sub>3</sub>OD relative to CH<sub>3</sub>OH at 10 K using the ASURA system.

The solid samples were produced by vapor deposition of CH<sub>3</sub>-OH, CH<sub>2</sub>DOH, CH<sub>3</sub>OD, CHD<sub>2</sub>OH, CD<sub>3</sub>OH (99.5% purity,

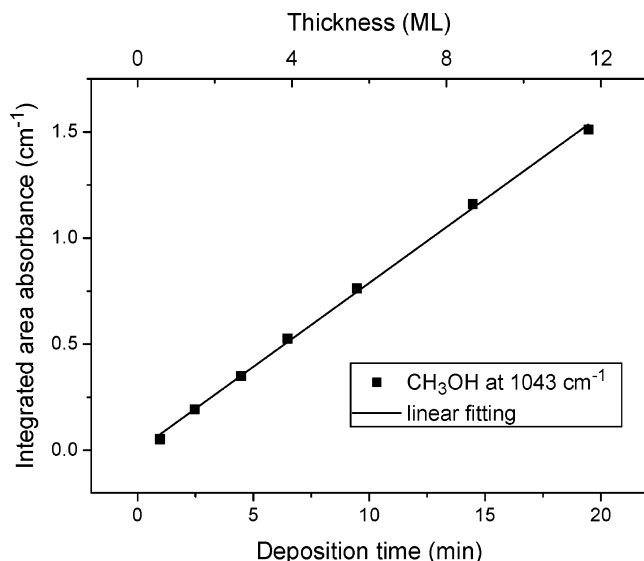


**Figure 4.** Infrared absorption spectra of amorphous solid CH<sub>3</sub>OH, CH<sub>2</sub>DOH, CH<sub>3</sub>OD, CHD<sub>2</sub>OH, CD<sub>3</sub>OH, and CD<sub>3</sub>OD deposited on the Al substrate at 10 K. Absorbance was obtained as common logarithms. The thickness of the samples is 3 monolayer. Spikes near 2675  $\text{cm}^{-1}$  are noise caused by the vibration of the He refrigerator.

Acros), and CD<sub>3</sub>OD through a capillary plate. To maintain a constant deposition rate during the experiments, the sample gas was turned on and off using a stop valve (SS-BNVCR4, Swagelok), and no adjustments were made to the precision leak valve (951-7172, Anelva). After the gas was turned off, the gas line was immediately evacuated. To avoid contamination, we measured one isotopologue per day.

Figure 4 shows infrared absorption spectra of solid CH<sub>3</sub>OH and isotopologues of 3 monolayer at 10 K. The shape of the -OH (-OD) stretching mode in the vicinity of 3300 (2450)  $\text{cm}^{-1}$  for the isotopologues indicates an amorphous structure in the samples.<sup>13</sup> The thickness of solid CH<sub>3</sub>OH was estimated from the integrated area absorbance of the C-O stretching mode at 1043.6  $\text{cm}^{-1}$  and the integrated band strength of  $1.8 \times 10^{-17} \text{ cm molecule}^{-1}$ .<sup>17,18</sup> The deposition rate was approximately 0.6 monolayer  $\text{min}^{-1}$ .

To quantify the amount of molecules in the solid samples (thickness of the samples) before and after D exposure using integrated band strengths, the linearity between the thickness and the integrated area absorbance had to be checked. Figure 5 shows the relationship between deposition time/thickness and integrated area absorbance for solid CH<sub>3</sub>OH at 10 K. The linearity between the deposition time and the integrated area absorbance was fairly good within the present experimental duration, which suggests that the sticking coefficients of CH<sub>3</sub>-OH molecules to the substrate and the solid CH<sub>3</sub>OH surfaces are constant at least up to a deposition time of 20 min. We also estimated the thickness of the samples using the temperature-programmed desorption (TPD) method;<sup>19</sup> the results are consistent with those obtained using integrated band strengths. We thus conclude that the linearity between the integrated area absorbance and the sample thickness is good, and the relative



**Figure 5.** Increase in integrated area absorbance at 1043  $\text{cm}^{-1}$  in the infrared spectra of amorphous solid  $\text{CH}_3\text{OH}$  at 10 K with deposition time. The upper abscissa represents the thickness of the solid  $\text{CH}_3\text{OH}$  samples estimated by the TPD method.<sup>19</sup> The solid line is the result of linear fitting.

integrated band strengths measured in the present work can be used for quantification. Similar measurements were also performed for all of the deuterated isotopologues listed above, and good linearity was confirmed. Hence, the relative integrated band strengths listed in Table 1 were obtained from the gradients of the linear fitted lines for all bands in the solid samples (e.g., see Figure 5).

### 3. Results and Discussion

#### 3.1. Exposure of Amorphous Solid $\text{CH}_3\text{OH}$ to D Atoms.

Figure 6 shows the infrared absorption spectrum of amorphous solid  $\text{CH}_3\text{OH}$  and the change in the spectrum after exposure to cold D atoms at 10 K. As shown in Figure 6b,  $\text{CH}_2\text{DOH}$ ,  $\text{CHD}_2\text{OH}$ , and  $\text{CD}_3\text{OH}$  appear with the consumption of  $\text{CH}_3\text{OH}$ . Isotopologues containing an -OD group, such as  $\text{CH}_3\text{OD}$  and  $\text{CD}_3\text{OD}$ , were not detected. Other molecules (e.g.,  $\text{H}_2\text{CO}$ ,  $\text{HDCO}$ ,  $\text{D}_2\text{CO}$ , and  $\text{CO}$ ) were also undetected. The detection limit of the ASURA system is  $\sim 1 \times 10^{12}$  molecule  $\text{cm}^{-2}$  ( $\sim 0.001$  ML) for  $\text{CH}_3\text{OH}$ . After D atom exposure, we maintained the solid samples at 10 K for several hours, but no change was observed in the spectrum, suggesting that no additional slow reaction proceeds at 10 K after termination of exposure. Exposure of solid  $\text{CH}_3\text{OH}$  to  $\text{D}_2$ , HD, and  $\text{H}_2$  molecules at 10 K did not induce any change in the spectrum. We performed a blank test<sup>20</sup> to check for contamination in the discharge tube and the Al tube, but no  $\text{CH}_3\text{OH}$ , deuterated isotopologues, or other molecules were detected; this is in agreement with results reported elsewhere.<sup>3–10</sup> This shows that both the discharge tube and the Al tube were uncontaminated by sample molecules; therefore, the obtained methanol isotopologues were produced via surface reactions with D atoms.

Figure 7 demonstrates the procedure for assignment of products and determination of molecular column densities. As can be seen in spectra a, b, and c in Figure 7, the bands at (ii) 920.8, (iv) 951.9, and (v) 988.4  $\text{cm}^{-1}$  appear only for  $\text{CH}_2\text{DOH}$ ,  $\text{CHD}_2\text{OH}$ , and  $\text{CD}_3\text{OH}$ , respectively. Therefore, these products are easily assigned and quantified in the exposed  $\text{CH}_3\text{OH}$  sample (spectrum d). In contrast, the decrease in  $\text{CH}_3\text{OH}$  should be carefully estimated because all of the  $\text{CH}_3\text{OH}$  bands overlap with those of the isotopologues. For example, the main

band of  $\text{CH}_3\text{OH}$  is located at 1043.6  $\text{cm}^{-1}$ , with the width of the base lying between 1000 and 1070  $\text{cm}^{-1}$ , while  $\text{CH}_2\text{DOH}$  and  $\text{CHD}_2\text{OH}$  have bands at (i) 1046.9  $\text{cm}^{-1}$  (base width 1005–1070  $\text{cm}^{-1}$ ) and (iii) 1036.7  $\text{cm}^{-1}$  (base width 1000–1060  $\text{cm}^{-1}$ ), respectively; the dip (vi) at 1050  $\text{cm}^{-1}$  in the spectrum of the exposed sample d consists of a decrease in the main band ( $\text{CH}_3\text{OH}$ ) and an increase in product bands (i) and (iii) ( $\text{CH}_2\text{DOH}$  and  $\text{CHD}_2\text{OH}$ , respectively). Fortunately, we had obtained the integrated band strengths of bands (i) and (iii) relative to those of (ii) and (iv), respectively, as listed in Table 1. Therefore, in the spectrum of the exposed sample d, the contribution of bands (i) and (iii) in the dip (vi) can be estimated from the increase in bands (ii) and (iv).

Variations in the column densities of molecules in the sample ( $\Delta N_t$ , molecule  $\text{cm}^{-2}$ ) due to D atom exposure were calculated using the integrated band strengths shown in Table 1; these are plotted in Figure 8.  $\Delta N_t$  is defined as  $\Delta N(X)_t = N(X)_t - N(X)_{t=0}$  for molecule X, where  $N(X)_t$  is the column density at an exposure time of  $t$  min. Since the spectra in Figure 6b were obtained by subtracting the initial ( $t = 0$ ) spectrum (Figure 6a) from the spectra after  $t$  min exposure,  $\Delta N_t$  values were calculated from the integrated area absorbance in the spectra shown in Figure 6b and these integrated band strengths using the following equations:

$$\Delta N(\text{CH}_3\text{OH})_t = \alpha \left[ \Delta A(1000-1070 \text{ cm}^{-1})_t - \Delta A(920.8 \text{ cm}^{-1})_t \times \frac{0.863}{0.105} - \Delta A(951.9 \text{ cm}^{-1})_t \times \frac{0.432}{0.181} \right] \quad (3a)$$

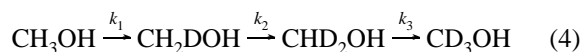
$$\Delta N(\text{CH}_2\text{DOH})_t = \frac{\alpha}{0.105} \times \Delta A(920.8 \text{ cm}^{-1})_t \quad (3b)$$

$$\Delta N(\text{CHD}_2\text{OH})_t = \frac{\alpha}{0.181} \times \Delta A(951.9 \text{ cm}^{-1})_t \quad (3c)$$

$$\Delta N(\text{CD}_3\text{OH})_t = \frac{\alpha}{0.484} \times \Delta A(988.4 \text{ cm}^{-1})_t \quad (3d)$$

$$\alpha = \frac{\cos 83^\circ}{2\epsilon} \quad (3e)$$

where  $\Delta A(X \text{ cm}^{-1})_t$  ( $\text{cm}^{-1}$ ) is the integrated area absorbance at  $X \text{ cm}^{-1}$  and  $\epsilon$  ( $\text{cm molecule}^{-1}$ ) is the integrated band strength of the C–O stretching mode in  $\text{CH}_3\text{OH}$ . Variations in the column density at the beginning of exposure ( $t < 5$  min), as shown in Figure 8, indicate that  $\text{CH}_2\text{DOH}$  forms first, and subsequently  $\text{CHD}_2\text{OH}$  and  $\text{CD}_3\text{OH}$  appear, with the concurrent consumption of  $\text{CH}_3\text{OH}$ . In addition,  $\text{CH}_2\text{DOH}$  and  $\text{CHD}_2\text{OH}$  have maxima at approximately 5 and 20 min, respectively. These features strongly suggest that  $\text{CH}_2\text{DOH}$ ,  $\text{CHD}_2\text{OH}$ , and  $\text{CD}_3\text{OH}$  are produced by the following successive H–D substitution reactions:



where  $k_1$ ,  $k_2$ , and  $k_3$  are the reaction rate constants. If deuterated isotopologues were produced in parallel by reactions such as  $\text{CH}_3\text{OH} \rightarrow \text{CH}_2\text{DOH}$ ,  $\text{CH}_3\text{OH} \rightarrow \text{CHD}_2\text{OH}$ , and  $\text{CH}_3\text{OH} \rightarrow \text{CD}_3\text{OH}$ , then  $\text{CH}_2\text{DOH}$  and  $\text{CHD}_2\text{OH}$  would not have maxima. The sum of the products ( $\text{CH}_2\text{DOH} + \text{CHD}_2\text{OH} + \text{CD}_3\text{OH}$ ; diamonds in Figure 8) is very close to the amount by which  $\text{CH}_3\text{OH}$  decreased in the same exposure time, showing that desorption of  $\text{CH}_3\text{OH}$  and products during D exposure is

TABLE 1: Infrared Band Positions and Relative Integrated Band Strengths

molecule	band position <sup>a</sup> (cm <sup>-1</sup> )	assignment <sup>b</sup>	relative integrated band strength <sup>c</sup>	no. for Figure 7	
CH <sub>3</sub> OH	3397.4 ± 10.9 <sup>d</sup>	OH stretching	1.068		
	3274.3 ± 2.9 <sup>d</sup>	OH stretching	3.348		
	2987.6 ± 1.0 <sup>d</sup>	CH <sub>3</sub> asym. stretching	0.053		
	2956.5 ± 4.0 <sup>d</sup>	CH <sub>3</sub> asym. stretching or CH <sub>3</sub> sym. bending	0.476		
	2907.0 ± 4.8 <sup>d</sup>	CH <sub>3</sub> asym. stretching or CH <sub>3</sub> asym. bending + CH <sub>3</sub> sym. bending	0.404		
	2862.6 ± 1.9 <sup>d</sup>	CH <sub>3</sub> sym. stretching or sym. bending	0.073		
	2831.2 ± 0.2 <sup>d</sup>	CH <sub>3</sub> sym. stretching	0.298		
	2821.3 ± 4.5 <sup>d</sup>	CH <sub>3</sub> sym. stretching	0.139		
	2606.9 ± 1.7 <sup>d</sup>	OH bending + CH <sub>3</sub> rocking	0.049		
	2530.2 ± 1.0 <sup>d</sup>	OH bending + CH <sub>3</sub> rocking	0.062		
	2422.2 ± 1.6 <sup>d</sup>	CH <sub>3</sub> rocking	0.088		
	2247.2 ± 2.9 <sup>d</sup>	CH <sub>3</sub> rocking	0.016		
	2046.2 ± 1.4 <sup>d</sup>	CO stretching	0.014		
	1491.4 ± 1.0 <sup>d</sup>	CH <sub>3</sub> asym. bending	0.147		
	1477.4 ± 0.4 <sup>d</sup>	CH <sub>3</sub> asym. bending or CH <sub>3</sub> sym. bending	0.051		
	1463.5 ± 0.3 <sup>d</sup>	CH <sub>3</sub> asym. bending	0.035		
	1449.6 ± 0.5 <sup>d</sup>	CH <sub>3</sub> sym. bending	0.070		
	1421.4 ± 0.3 <sup>d</sup>	OH bending	0.170		
	1163.5 ± 1.4 <sup>d</sup>	CH <sub>3</sub> rocking	0.007		
	1125.8 ± 0.2 <sup>d</sup>	CH <sub>3</sub> rocking	0.142		
	CH <sub>2</sub> DOH	1043.6 ± 0.1	CO stretching	1.000	
		3418.8 ± 3.1 <sup>d</sup>	OH stretching	0.263	
		3291.2 ± 3.2 <sup>d</sup>	OH stretching	3.300	
2977.3 ± 0.6 <sup>d</sup>		CH <sub>2</sub> asym. stretching	0.046		
2941.2 ± 0.3 <sup>d</sup>		CH <sub>2</sub> asym. stretching	0.235		
2905.1 ± 2.1 <sup>d</sup>		CH <sub>2</sub> sym. stretching	0.020		
2882.6 ± 0.2 <sup>d</sup>		CH <sub>2</sub> sym. stretching	0.191		
2818.9 ± 1.0 <sup>d</sup>		CH <sub>2</sub> sym. stretching?	0.139		
2725.2 ± 1.6 <sup>d</sup>		?	0.058		
2241.5 ± 0.2 <sup>d</sup>		?	0.027		
2180.2 ± 0.1 <sup>d</sup>		CH <sub>2</sub> rocking	0.077		
2147.2 ± 0.3 <sup>d</sup>		CD stretching	0.041		
1467.7 ± 0.4 <sup>d</sup>		CH <sub>2</sub> bending	0.175		
1419.1 ± 0.5 <sup>d</sup>		OH bending?	0.129		
1328.3 ± 0.1 <sup>d</sup>		CH <sub>2</sub> twisting	0.119		
1292.7 ± 0.1 <sup>d</sup>		OH bending	0.058		
1110.9 ± 6.0 <sup>d</sup>		CH <sub>2</sub> rocking	0.043		
1082.4 ± 1.9 <sup>d</sup>		CH <sub>2</sub> rocking	0.033		
1046.9 ± 0.1 <sup>d</sup>		CO stretching	0.863	i	
920.8 ± 0.2		CD bending	0.105	ii	
CHD <sub>2</sub> OH		3420.9 ± 2.5 <sup>d</sup>	OH stretching	0.352	
		3292.6 ± 3.1 <sup>d</sup>	OH stretching	3.335	
		2977.3 ± 1.4 <sup>d</sup>	CH stretching	0.012	
	2951.3 ± 0.8 <sup>d</sup>	CO stretching	0.032		
	2915.5 ± 0.4 <sup>d</sup>	CH stretching	0.255		
	2867.5 ± 1.0 <sup>d</sup>	?	0.004		
	2815.9 ± 1.4 <sup>d</sup>	?	0.403		
	2713.7 ± 1.3 <sup>d</sup>	?	0.038		
	2256.0 ± 1.5 <sup>d</sup>	CD <sub>2</sub> rocking	0.003		
	2224.8 ± 0.5 <sup>d</sup>	CD <sub>2</sub> asym. stretching	0.073		
	2181.4 ± 0.2 <sup>d</sup>	CD <sub>2</sub> bending	0.057		
	2128.4 ± 0.1	CD <sub>2</sub> sym. stretching	0.124		
	1973.0 ± 2.0	CO stretching + CD <sub>2</sub> wagging?	0.012		
	1448.4 ± 1.1 <sup>d</sup>	OH bending?	0.222		
	1404.9 ± 0.3 <sup>d</sup>	OH bending?	0.072		
	1330.4 ± 0.1 <sup>d</sup>	CH bending	0.090		
	1303.4 ± 0.1 <sup>d</sup>	CH bending	0.080		
	1120.6 ± 1.4 <sup>d</sup>	CD <sub>2</sub> bending	0.016		
	1090.2 ± 0.1 <sup>d</sup>	CD <sub>2</sub> bending	0.273		
	1036.7 ± 0.1	CO stretching	0.432	iii	
	951.9 ± 0.1 <sup>d</sup>	CD <sub>2</sub> wagging	0.181	iv	
	898.3 ± 4.1 <sup>d</sup>	CD <sub>2</sub> twisting	0.059		
	887.2 ± 2.3 <sup>d</sup>	CD <sub>2</sub> twisting?	0.008		
CD <sub>3</sub> OH	3421.3 ± 6.2 <sup>d</sup>	OH stretching	0.341		
	3289.3 ± 2.1 <sup>d</sup>	OH stretching	3.403		
	2802.1 ± 0.9 <sup>d</sup>	OH bending	0.504		
	2246.1 ± 0.4 <sup>d</sup>	CD <sub>3</sub> asym. stretching	0.072		
	2235.1 ± 0.1 <sup>d</sup>	CD <sub>3</sub> asym. stretching	0.036		
	2215.7 ± 0.1 <sup>d</sup>	CD <sub>3</sub> asym. stretching	0.130		
	2191.7 ± 0.1 <sup>d</sup>	?	0.016		
	2146.1 ± 4.5 <sup>d</sup>	CD <sub>3</sub> asym. bending	0.009		
	2129.7 ± 7.0 <sup>d</sup>	CD <sub>3</sub> asym. bending	0.024		
	2093.5 ± 1.4 <sup>d</sup>	CD <sub>3</sub> sym. stretching	0.003		

TABLE 1: (Continued)

molecule	band position <sup>a</sup> (cm <sup>-1</sup> )	assignment <sup>b</sup>	relative integrated band strength <sup>c</sup>	no. for Figure 7	
CD <sub>3</sub> OH (contd)	2073.0 ± 0.1 <sup>d</sup>	CD <sub>3</sub> sym. stretching	0.172		
	2009.3 ± 3.0 <sup>d</sup>	?	0.002		
	1957.1 ± 2.0 <sup>d</sup>	CO stretching	0.011		
	1427.2 ± 0.4 <sup>d</sup>	OH bending	0.274		
	1384.1 ± 0.2 <sup>d</sup>	OH bending	0.056		
	1126.4 ± 0.1	CD <sub>3</sub> sym. bending	0.328		
	1068.0 ± 0.5	CD <sub>3</sub> asym. bending	0.036		
	1030.1 ± 1.0	CD <sub>3</sub> asym. bending	0.003		
	988.4 ± 0.1	CO stretching	0.484	v	
	916.6 ± 5.4 <sup>d</sup>	CD <sub>3</sub> rocking	0.011		
	899.9 ± 0.6 <sup>d</sup>	CD <sub>3</sub> rocking	0.052		
	883.0 ± 1.0 <sup>d</sup>	CD <sub>3</sub> rocking	0.028		
	CH <sub>3</sub> OD	2985.6 ± 1.6 <sup>d</sup>	CH <sub>3</sub> asym. stretching	0.095	
		2955.8 ± 0.6 <sup>d</sup>	CH <sub>3</sub> asym. stretching	0.251	
		2912.7 ± 0.4 <sup>d</sup>	CH <sub>3</sub> asym. stretching	0.208	
		2864.1 ± 0.6 <sup>d</sup>	CH <sub>3</sub> sym. stretching	0.023	
		2837.2 ± 0.1 <sup>d</sup>	CH <sub>3</sub> sym. stretching	0.136	
2536.0 ± 9.6 <sup>d</sup>		OD stretching	0.263		
2500.6 ± 1.5 <sup>d</sup>		OD stretching	0.179		
2430.6 ± 0.6 <sup>d</sup>		OD stretching	1.260		
2405.9 ± 5.0 <sup>d</sup>		OD stretching	1.042		
2199.6 ± 14.4 <sup>d</sup>		CH <sub>3</sub> rocking + OD bending?	0.060		
2174.4 ± 3.0 <sup>d</sup>		CH <sub>3</sub> rocking + OD bending	0.034		
1513.4 ± 0.1 <sup>d</sup>		?	0.007		
1477.8 ± 0.3 <sup>d</sup>		CH <sub>3</sub> asym. bending	0.041		
1463.5 ± 0.3 <sup>d</sup>		CH <sub>3</sub> asym. bending	0.057		
1235.8 ± 0.5		CH <sub>3</sub> rocking	0.038		
1162.0 ± 2.0		CH <sub>3</sub> rocking	0.008		
1124.0 ± 3.0		CH <sub>3</sub> rocking	0.005		
1043.5 ± 0.1	CO stretching	0.797			
966.3 ± 1.3 <sup>d</sup>	OD bending	0.025			
941.5 ± 0.9 <sup>d</sup>	OD bending	0.067			
CD <sub>3</sub> OD	2527.3 ± 3.3 <sup>d</sup>	OD stretching	0.228		
	2451.6 ± 2.2 <sup>d</sup>	OD stretching	0.918		
	2421.3 ± 2.5 <sup>d</sup>	OD stretching	1.502		
	2244.7 ± 0.8 <sup>d</sup>	CD <sub>3</sub> asym. stretching	0.134		
	2218.7 ± 0.4 <sup>d</sup>	CD <sub>3</sub> asym. stretching	0.073		
	2202.8 ± 4.8 <sup>d</sup>	CD <sub>3</sub> asym. stretching	0.105		
	2155.0 ± 2.5 <sup>d</sup>	CD <sub>3</sub> asym. stretching	0.136		
	2073.0 ± 0.1 <sup>d</sup>	CD <sub>3</sub> sym. stretching	0.190		
	1934.0 ± 5.0 <sup>d</sup>	CO stretching	0.055		
	1127.5 ± 0.1 <sup>d</sup>	CD <sub>3</sub> sym. stretching	0.180		
	1107.2 ± 0.8 <sup>d</sup>	CD <sub>3</sub> asym. stretching	0.268		
	1064.4 ± 0.2 <sup>d</sup>	CD <sub>3</sub> asym. stretching	0.077		
	981.5 ± 0.1	CO stretching	0.400		
	900.0 ± 0.5	CD <sub>3</sub> rocking	0.021		

<sup>a</sup> Values are in 3 monolayer. Bands at 1700–1600 cm<sup>-1</sup> in CH<sub>3</sub>OH, CH<sub>2</sub>DOH, and CH<sub>3</sub>OD are noise. <sup>b</sup> Bands were assigned on the basis of literature values.<sup>26–29</sup> <sup>c</sup> Values are normalized to the CO stretching mode in CH<sub>3</sub>OH at 1043.6 cm<sup>-1</sup>. <sup>d</sup> Band positions and integrated area absorbance (relative integrated band strength) were obtained by the least-square method (peak deconvolution).

inefficient. Therefore, the maxima mentioned above are not due to desorption from the solid sample but are due to the successive H–D substitution reaction 4.

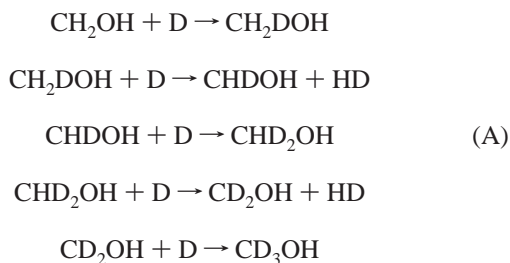
It can be seen that the decrease in CH<sub>3</sub>OH was saturated at approximately –0.1 after the long exposure time. This saturation probably resulted from a chemical equilibrium and/or a slow diffusion rate of D atoms to the inside of the solid sample. To examine the former possibility, we performed an experiment in which solid CH<sub>2</sub>DOH was exposed to cold D atoms at 10 K (see section 3.2 for details). No CH<sub>3</sub>OH was observed, indicating that the unknown backward process CH<sub>2</sub>DOH → CH<sub>3</sub>OH does not proceed. Thus, the origin of saturation must be the latter; it seems that CH<sub>3</sub>OH molecules at the surface of the solid sample easily react with D atoms, but buried CH<sub>3</sub>OH molecules hardly react because of the very slow diffusion rate of D atoms into the bulk.

To investigate the influence of the translational temperature of the D atoms on reactivity, we performed the same experiment using D atoms at 50, 200, and 300 K. The results of these

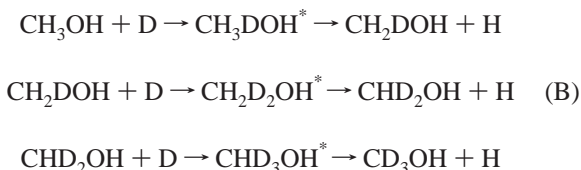
experiments are consistent with each other, suggesting that reaction 4 proceeds not via the Eley–Rideal (ER) process but via the Langmuir–Hinshelwood (LH) process: the surface reaction between adsorbed atoms and molecules on a surface. If the ER process is dominant in reaction 4, the reaction rate should become large with an increase in the translational temperature of D, because the ER process is a direct reaction between atoms impinging from the gas phase and adsorbed molecules on the surface. To confirm this scenario, we exposed solid CH<sub>3</sub>OH to cold (100 K) D atoms at 30 K. Reaction 4 was not observed in this experiment within a duration of more than 180 min. This is due to a drop in the sticking coefficient of the D atom to the sample surface above ~20 K,<sup>5,7</sup> and supports the dominance of the LH process in reaction 4.

The possible mechanisms of the H–D substitution reaction 4 are as follows:

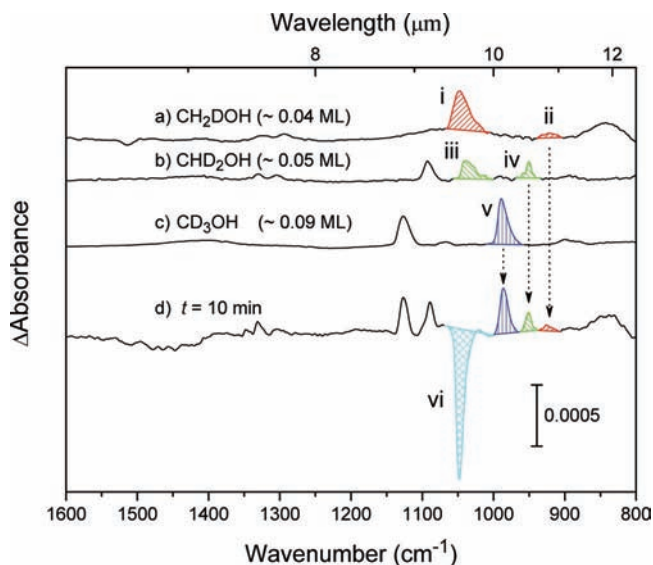




and

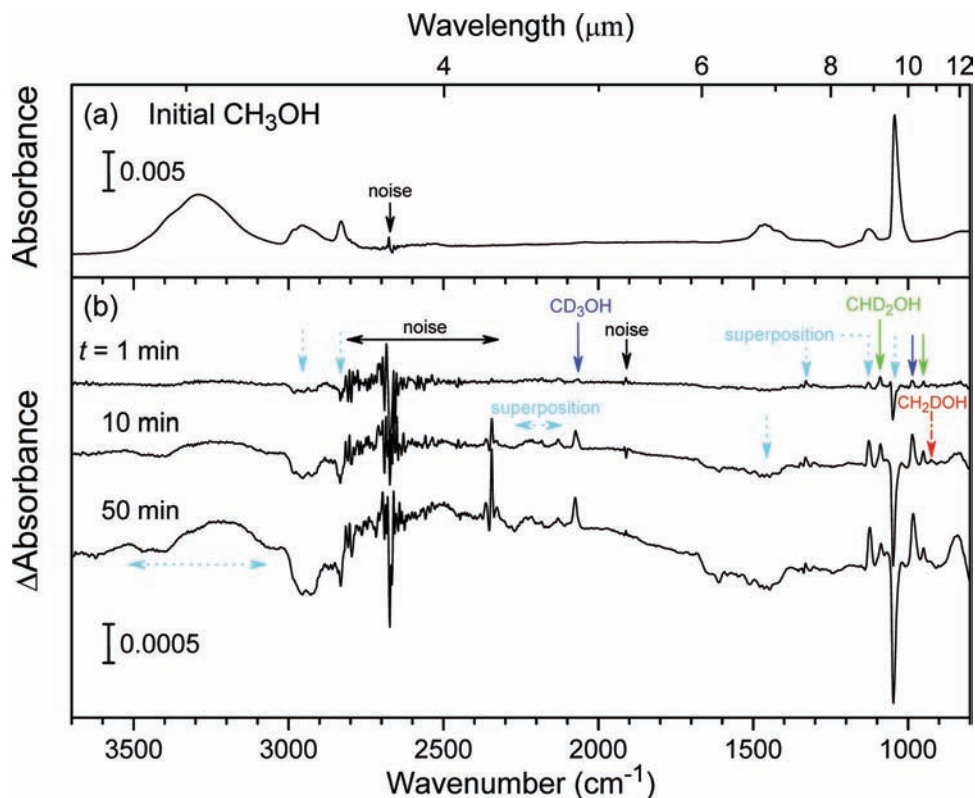


Process A consists of the repetition of the hydroxymethyl radical formation by H abstraction from the methyl side in methanol and the subsequent D addition to the radical to form methanol isotopologues; process B consists of the formation of excited intermediates such as  $\text{CH}_3\text{DOH}^*$  by D addition and the subsequent H–D exchange in the methyl group. Since the difference between A and B is in the formation of hydroxymethyl radicals, the detection of such radicals enables us to determine which process dominates the H–D substitution reaction. However, although hydroxymethyl radicals were not detected in this experiment, we cannot conclude that process B is the dominant process in the H–D substitution reaction. Even when these radicals are produced in process A, they will be



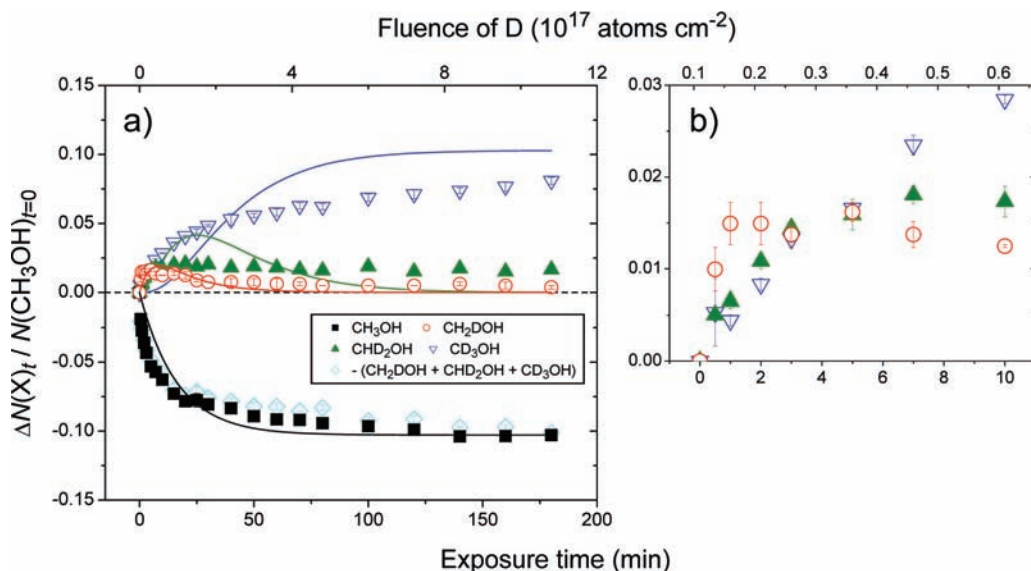
**Figure 7.** Example of band assignment in the spectrum obtained after D exposure. Spectra a, b, and c represent pure amorphous solid samples of  $\text{CH}_2\text{DOH}$  (0.04 ML),  $\text{CHD}_2\text{OH}$  (0.05 ML), and  $\text{CD}_3\text{OH}$  (0.09 ML) at 10 K, respectively, corresponding to the increase of these isotopologues in the solid  $\text{CH}_3\text{OH}$  sample exposed to D atoms for 10 min (d).

consumed immediately by the reaction with D atoms because of the very fast rate of the radical atom reaction. Therefore, the average surface density of the radicals will be too low for detection, similar to the nondetection of  $\text{HCO}$  and  $\text{CH}_3\text{O}$  radicals in the experiment on successive hydrogenation of  $\text{CO}$ .<sup>5</sup> We detected particles of mass 3 (HD) and 1 (H) by QMS during the exposure. However, the atomic beam used in the present

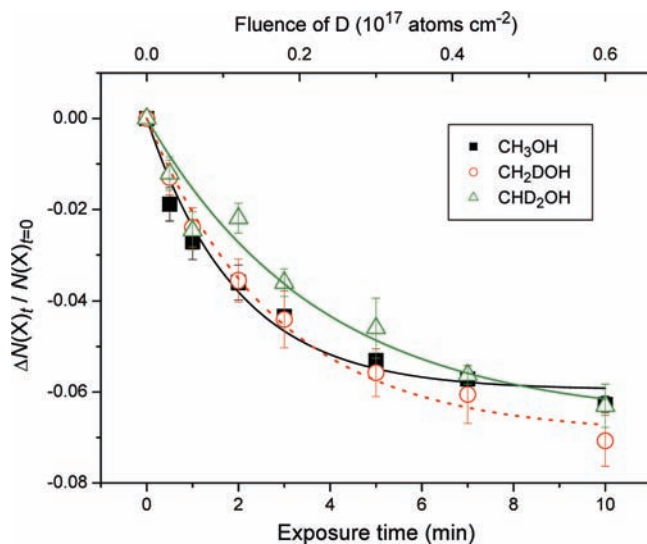


**Figure 6.** (a) Infrared absorption spectrum of initial (pre-exposure) amorphous solid  $\text{CH}_3\text{OH}$  (3 ML) deposited onto the Al substrate at 10 K. Absorbance is obtained as common logarithms. (b) Change in the spectrum after D exposure of 1, 10, and 50 min at 10 K. Spectra were obtained by subtracting the initial spectrum from D-exposed spectra. Peaks below and above the baseline represent decreases and increases, respectively, compared to the initial spectrum. Red (dashed-dotted), green (continuous), and blue (continuous) arrows represent increases in  $\text{CH}_2\text{DOH}$ ,  $\text{CHD}_2\text{OH}$ , and  $\text{CD}_3\text{OH}$ , respectively. Sky-blue (dotted) arrows are a superposition of  $\text{CH}_3\text{OH}$  decreases and increases in isotopologues. Noise, mainly caused by the vibration of the He refrigerator, is denoted by black (continuous) arrows.





**Figure 8.** (a) Variation in column densities normalized to initial CH<sub>3</sub>OH in the experiment in which amorphous solid CH<sub>3</sub>OH (3 ML) was exposed to D atoms at 10 K. The upper abscissa represents the fluence of D atoms. The lines are the results of fitting using eq 8a–d. Error bars represent statistical error. The diamonds represent the sum of the products multiplied by  $-1$ . (b) Close-up of the products for exposure time of  $t = 0$ –10 min.



**Figure 9.** Decay of CH<sub>3</sub>OH, CH<sub>2</sub>DOH, and CHD<sub>2</sub>OH in experiments in which amorphous solid CH<sub>3</sub>OH, CH<sub>2</sub>DOH, and CHD<sub>2</sub>OH were exposed to D atoms at 10 K. The upper abscissa represents the fluence of D atoms. The lines are the results of fitting using single-exponential decay function (8a). Error bars represent statistical error.

experiment is a continuous beam, and thus the pressure during D exposure increases up to  $\sim 2 \times 10^{-7}$  Torr. Therefore, it is unclear whether the origin of those masses is processes A and B or the background. As a result, it is difficult to evaluate the contribution of processes A and B to reaction 4 from the present experiment.

Isotopologues containing D in a hydroxyl group (e.g., CH<sub>3</sub>-OD) were not detected in the present experiment. This is due in part to the larger dissociation energy of the O–H bond (CH<sub>3</sub>-OH  $\rightarrow$  CH<sub>3</sub>O + H) compared with the C–H bond (CH<sub>3</sub>OH  $\rightarrow$  CH<sub>2</sub>OH + H) in CH<sub>3</sub>OH (by 4–11 kcal mol<sup>-1</sup>).<sup>21</sup> In addition, the vibrationally adiabatic barrier heights of the following two channels for the reactions of CH<sub>3</sub>OH with D atoms in the gas phase were calculated using an ab initio method:<sup>22</sup>



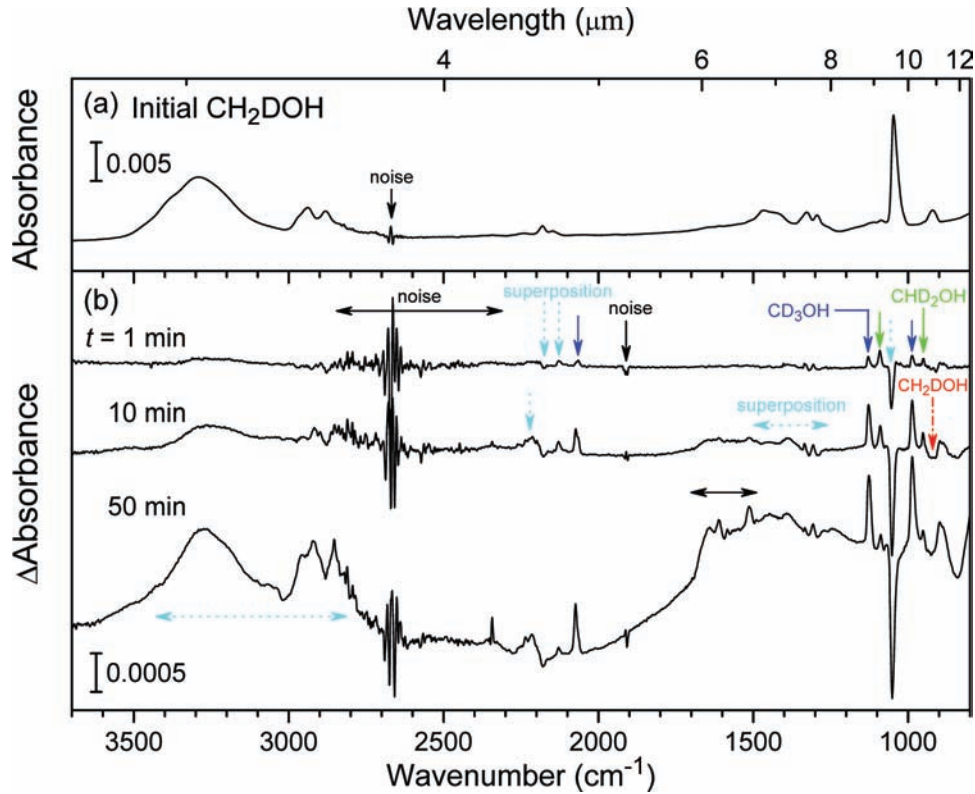
The calculated barrier heights for reactions 5a,b are 7.72 and 12.2 kcal mol<sup>-1</sup>,<sup>22</sup> respectively, suggesting that H abstraction from the methyl group (5a) generally proceeds faster than from the hydroxyl group (5b).

Although the H–D exchange process (B) has never been reported on a solid surface, the gas-phase H–H exchange reaction was investigated by Osamura et al.<sup>23</sup> Using the ab initio method, they obtained the barrier height for H addition to the C side in CH<sub>3</sub>OH and subsequent exchange for another H atom bonded with C (CH<sub>3</sub>OH + H'  $\rightarrow$  CH<sub>3</sub>H'OH\*  $\rightarrow$  CH<sub>2</sub>H'OH + H) to be 40.4 kcal mol<sup>-1</sup>, and thus they concluded that this H–H exchange reaction does not occur at low temperatures. Therefore, it is reasonable to consider that the H–D substitution reaction 4 proceeds via process A rather than process B. Considering the theoretically calculated barrier height of 7.72 kcal mol<sup>-1</sup>,<sup>22</sup> we find that the H-abstraction reaction is likely to proceed via a quantum mechanical tunneling reaction at a low-temperature surface.

We determined the effective rate constants for the H–D substitution reaction 4 from the experimental results on the assumption that the reaction is dominated by process A:



where  $k_n'$  and  $k_n''$  (cm<sup>2</sup> molecule<sup>-1</sup> s<sup>-1</sup>,  $n = 1$ –3) are the reaction rate constants. Since the rate constants  $k_n'$  for hydroxymethyl radicals CH<sub>2</sub>OH, CHDOH, and CD<sub>2</sub>OH must be much smaller than those for D-addition  $k_n''$  ( $n = 1$ –3), the rate-



**Figure 10.** (a) Infrared absorption spectrum of initial (pre-exposure) amorphous solid  $\text{CH}_2\text{DOH}$  deposited at 10 K. (b) Spectral change after D exposure for 1, 10, and 50 min at 10 K. Absorbance was obtained as common logarithms. The spectra were obtained by subtracting the initial spectrum from D-exposed spectra. Red (dashed-dotted), green (continuous), and blue (continuous) arrows show a decrease in  $\text{CH}_2\text{DOH}$ , increases in  $\text{CHD}_2\text{OH}$  and  $\text{CD}_3\text{OH}$ . Sky-blue (dotted) arrows show a superposition of  $\text{CH}_2\text{DOH}$  decreases and increases in the other isotopologues. Noise, mainly caused by the vibration of the He refrigerator, is denoted by black (continuous) arrows.

**TABLE 2: Effective Rate Constants for the H-Abstraction Reaction by D Atoms at 10 K**

$n$	pseudo-first order: $k_n'N_D^a$ ( $10^{-3} \text{ s}^{-1}$ )	second order: $k_n'P^{1/2b}$ ( $10^{-11} \text{ cm}^2 \text{ molecule}^{-1} \text{ s}^{-1}$ )
1	$8.5 \pm 1.1$	$1.6 \pm 0.21$
2	$5.9 \pm 0.57$	$1.1 \pm 0.11$
3	$4.4 \pm 1.0$	$0.82 \pm 0.19$

<sup>a</sup> Fitting results for decay of sample molecule in each experiment.

<sup>b</sup> Effective rate constants calculated using  $N_D = 53.9(1 \times 10^{14}P)^{1/2}$ .

determining steps for the H–D substitution reaction 4 are 6a,c,e. Hence, the rate constants  $k_1'$ ,  $k_2'$ , and  $k_3'$  can be approximated as  $k_1$ ,  $k_2$ , and  $k_3$ , respectively. Finally, the rate equations for reaction 4 are expressed as follows:

$$\frac{dN(\text{CH}_3\text{OH})_t}{dt} = -k_1'N_D N(\text{CH}_3\text{OH})_t \quad (7a)$$

$$\frac{dN(\text{CH}_2\text{DOH})_t}{dt} = k_1'N_D N(\text{CH}_3\text{OH})_t - k_2'N_D N(\text{CH}_2\text{DOH})_t \quad (7b)$$

$$\frac{dN(\text{CHD}_2\text{OH})_t}{dt} = k_2'N_D N(\text{CH}_2\text{DOH})_t - k_3'N_D N(\text{CHD}_2\text{OH})_t \quad (7c)$$

$$\frac{dN(\text{CD}_3\text{OH})_t}{dt} = k_3'N_D N(\text{CHD}_2\text{OH})_t \quad (7d)$$

where  $N(X)_t$  (molecule  $\text{cm}^{-2}$ ) is the column density at  $t$ , and  $N_D$  is the surface density of D atoms (molecule  $\text{cm}^{-2}$ ). The

solution of 7a–d under the initial condition ( $t = 0$ ) of  $N(\text{CH}_2\text{DOH})_0 = N(\text{CHD}_2\text{OH})_0 = N(\text{CD}_3\text{OH})_0 = 0$  is given by

$$\frac{N(\text{CH}_3\text{OH})_t}{N(\text{CH}_3\text{OH})_0} = e^{-k_1'N_D t} \quad (8a)$$

$$\frac{N(\text{CH}_2\text{DOH})_t}{N(\text{CH}_3\text{OH})_0} = \frac{k_1'}{k_2' - k_1'} (e^{-k_1'N_D t} - e^{-k_2'N_D t}) \quad (8b)$$

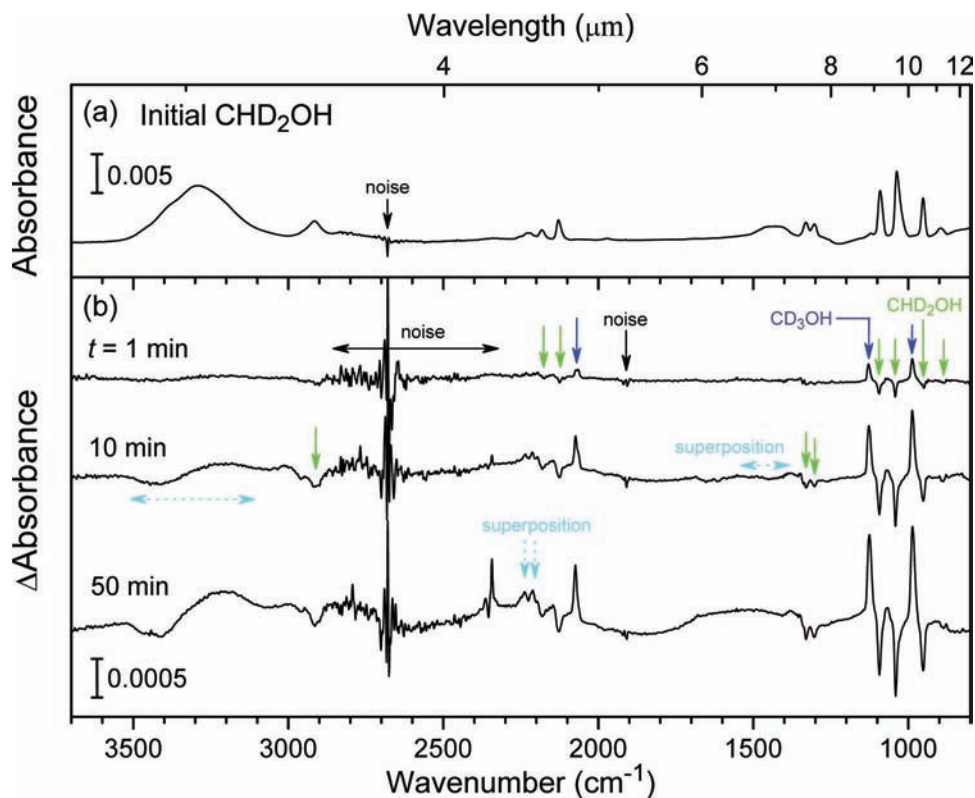
$$\frac{N(\text{CHD}_2\text{OH})_t}{N(\text{CH}_3\text{OH})_0} = \frac{k_1'k_2'}{k_2' - k_1'} \left( \frac{e^{-k_1'N_D t} - e^{-k_3'N_D t}}{k_3' - k_1'} - \frac{e^{-k_2'N_D t} - e^{-k_3'N_D t}}{k_3' - k_2'} \right) \quad (8c)$$

$$\frac{N(\text{CD}_3\text{OH})_t}{N(\text{CH}_3\text{OH})_0} = \frac{k_1'k_2'k_3'}{k_2' - k_1'} \left[ \frac{e^{-k_1'N_D t} - 1}{-k_1'(k_3' - k_1')} - \frac{e^{-k_3'N_D t} - 1}{-k_3'(k_3' - k_1')} - \frac{e^{-k_2'N_D t} - 1}{-k_2'(k_3' - k_2')} + \frac{e^{-k_3'N_D t} - 1}{-k_3'(k_3' - k_2')} \right] \quad (8d)$$

where  $N(\text{CH}_3\text{OH})_0$  is the column density of  $\text{CH}_3\text{OH}$  at  $t = 0$  (molecule  $\text{cm}^{-2}$ ). The time variation of  $N_D$  is expressed as follows:<sup>24</sup>

$$\frac{dN_D}{dt} = Pf_D - k_{D-D}N_D^2 - R_D N_D \quad (9)$$

where  $P$  is the sticking coefficient of the D atom,  $f_D$  is the flux of D atoms (molecule  $\text{cm}^{-2} \text{ s}^{-1}$ ),  $k_{D-D}$  is the rate constant for recombination of D atoms ( $\text{D} + \text{D} \rightarrow \text{D}_2$ ) ( $\text{cm}^2 \text{ molecule}^{-1} \text{ s}^{-1}$ ),

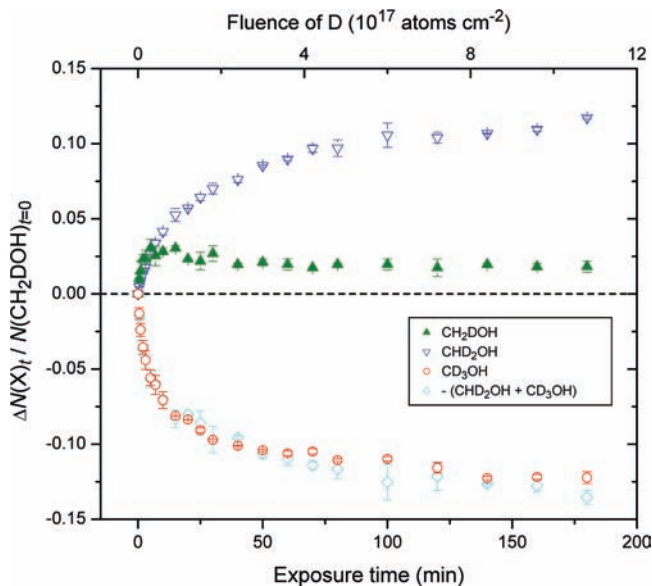


**Figure 11.** (a) Infrared absorption spectrum of initial (pre-exposure) amorphous solid  $\text{CHD}_2\text{OH}$  deposited at 10 K. (b) Spectral change after D exposure for 1, 10, and 50 min at 10 K. Spectra were obtained by subtracting the initial spectrum from the D-exposed spectra. Absorbance was obtained as common logarithms. Green (continuous) and blue (continuous) arrows show a decrease in  $\text{CHD}_2\text{OH}$  and an increase in  $\text{CD}_3\text{OH}$ . Sky-blue (dotted) arrows show a superposition of  $\text{CHD}_2\text{OH}$  decreases and increases in  $\text{CD}_3\text{OH}$ . Noise, mainly caused by the vibration of the He refrigerator, is denoted by black (continuous) arrows.

and  $R_D$  is the rate constant for desorption of D atoms (molecule  $\text{s}^{-1}$ ). In the right-hand side of eq 9, the terms for D loss by reaction with  $\text{CH}_3\text{OH}$ ,  $\text{CH}_2\text{DOH}$ , and  $\text{CHD}_2\text{OH}$  are neglected, because these reactions are very slow compared with recombination and desorption. The calculation by Chigai et al. demonstrates that  $N_D$  immediately becomes  $N_D = 53.9 \times (f_D P)^{1/2}$  at the beginning ( $t < 10^{-9}$  s) of D exposure and stays constant during exposure under the present experimental conditions.

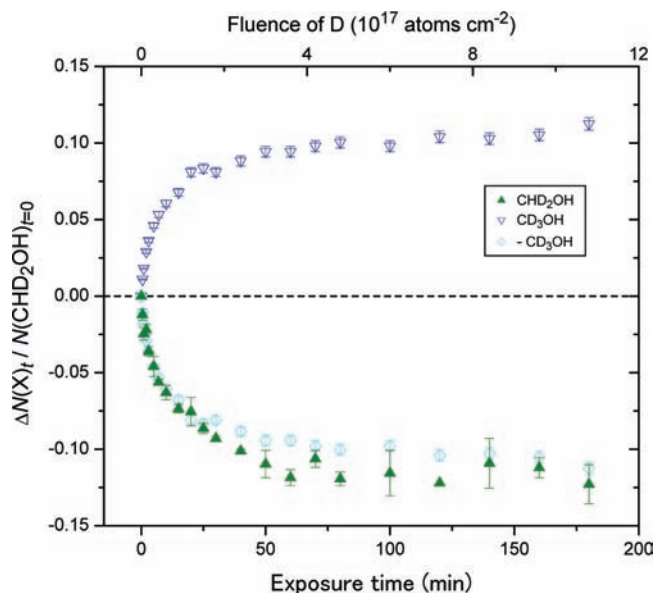
The solid lines in Figure 8 show the results of fitting by the least-square method using 8a–d. None of the fitted curves reproduce the experimental data well. The crucial difference between the fitted curves and the experimental data can be seen in  $\text{CH}_2\text{DOH}$  and  $\text{CHD}_2\text{OH}$ ; these curves approach zero beyond their maxima, while the experimental data for  $\text{CH}_2\text{DOH}$  and  $\text{CHD}_2\text{OH}$  are almost constant after  $\sim 50$  min and do not approach zero within the duration of the experiment. This difference may be attributable to the slow diffusion rate of D atoms to the inside of the solid samples, as mentioned above; molecules at the surface of the solid sample easily react with D atoms, but buried molecules react with more difficulty and tend to remain intact. Therefore, we fit the decay of  $\text{CH}_3\text{OH}$  at a duration of  $t = 0$ –10 min using eq 8a alone (black line in Figure 9). The fitted curve reproduces the decay of  $\text{CH}_3\text{OH}$  well. The obtained pseudo-first-order reaction rate for reaction 6a is  $(8.5 \pm 1.1) \times 10^{-3} \text{ s}^{-1}$ . Assuming that  $N_D = 53.9(1 \times 10^{14} P)^{1/2}$ , the effective second-order rate constant can be obtained as a function of the sticking coefficient of D atoms:  $k_1' P^{1/2} = (1.6 \pm 0.21) \times 10^{-11} \text{ cm}^2 \text{ molecule}^{-1} \text{ s}^{-1}$ .

**3.2. Exposure of Amorphous Solid  $\text{CH}_2\text{DOH}$  and  $\text{CHD}_2\text{OH}$  to D Atoms.** To obtain  $k_2'$  and  $k_3'$ , respectively, we exposed solid  $\text{CH}_2\text{DOH}$  and  $\text{CHD}_2\text{OH}$  at 10 K to cold D atoms. The experimental procedure and conditions were the same as those



**Figure 12.** (a) Variation in column densities normalized to initial  $\text{CH}_2\text{DOH}$  in the experiment in which amorphous solid  $\text{CH}_2\text{DOH}$  (3 ML) was exposed to D atoms at 10 K. The upper abscissa represents the fluence of D atoms. The error bars represent statistical error. Diamonds represent the sum of the products multiplied by  $-1$ .

for solid  $\text{CH}_3\text{OH}$ . The thickness of the samples was also the same as in the  $\text{CH}_3\text{OH} + \text{D}$  experiment. Figures 10 and 11 show the infrared absorption spectra of solid  $\text{CH}_2\text{DOH}$  and  $\text{CHD}_2\text{OH}$  before exposure and the change in the spectra after exposure to cold D atoms at 100 K, respectively. In the experiment using  $\text{CH}_2\text{DOH}$ , the formation of  $\text{CHD}_2\text{OH}$  and  $\text{CD}_3\text{OH}$  is observed with the consumption of  $\text{CH}_2\text{DOH}$  (Figure



**Figure 13.** (a) Variation in column densities normalized to initial  $\text{CHD}_2\text{OH}$  in the experiment in which amorphous solid  $\text{CHD}_2\text{OH}$  (3 ML) was exposed to D atoms at 10 K. The upper abscissa represents the fluence of D atoms. The error bars represent statistical error. Diamonds represent  $\text{CD}_3\text{OH}$  multiplied by  $-1$ .

**TABLE 3: Relative Rates for H-Abstraction Reaction by D Atoms at 10 K**

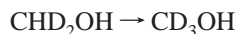
	relative rates <sup>a</sup>	calculation <sup>b</sup>
$k_2'/k_1'$	$0.69 \pm 0.11$	0.576
$k_3'/k_1'$	$0.52 \pm 0.14$	
$k_3'/k_2'$	$0.75 \pm 0.18$	

<sup>a</sup> Obtained from the pseudo-first-order rates in Table 2. <sup>b</sup> Ab initio calculation by Lendvay et al.<sup>25</sup>

10b); in the experiment using  $\text{CHD}_2\text{OH}$ , the only product is  $\text{CD}_3\text{OH}$  (Figure 11b). The backward process (e.g.,  $\text{CH}_2\text{DOH} \rightarrow \text{CH}_3\text{OH}$ ) and other products ( $\text{HDCO}$ ,  $\text{D}_2\text{CO}$ , etc.) are not observed at all. Variations in the column densities of the samples due to D atom exposure are shown in Figures 12 and 13. H–D substitution reactions were found to proceed as follows in each experiment:



and



For the reasons described in section 3.1, we fit the data of  $\text{CH}_2\text{DOH}$  and  $\text{CHD}_2\text{OH}$  decay to single-exponential decay (8a, Figure 9). The obtained pseudo-first-order rate and effective second-order rate constants are summarized in Table 2.

Assuming that  $P$ ,  $k_{\text{D-D}}$ , and  $R_{\text{D}}$  are independent of the initial solid sample types, the relationships between the rate constants  $k_n'$  ( $n = 1-3$ ) were found to be  $k_1' > k_2' > k_3'$ ; the ratios  $k_2'/k_1' = 0.69 \pm 0.11$  and  $k_3'/k_1' = 0.52 \pm 0.14$  were obtained from the results of fitting for  $\text{CH}_3\text{OH}$ ,  $\text{CH}_2\text{DOH}$ , and  $\text{CHD}_2\text{OH}$  samples (Table 3). This relationship is mainly due to the difference in activation energy for the H-abstraction reaction 6a,c,e (i.e., the secondary kinetic isotope effect). Using ab initio methods, Lendvay et al. calculated the secondary kinetic isotope effect for the reactions  $\text{CH}_3\text{OH} + \text{H} \rightarrow \text{CH}_2\text{OH} + \text{H}_2$  and  $\text{CH}_2\text{DOH} + \text{H} \rightarrow \text{CHDOH} + \text{H}_2$  at 300–2000 K (ref 25) and reported a  $k_2'/k_1'$  ratio of 0.576 at 300 K. This value is smaller than the experimentally obtained ratio of  $0.69 \pm 0.11$ . Although

the origin of the difference between the theory and our experiments is unclear, it may have arisen from differences in conditions such as reactants (H atom vs D atom), phase (gas vs solid), and temperature (300 K vs 10 K).

### 3.3. Exposure of Amorphous Solid $\text{CH}_3\text{OD}$ to D Atoms.

An experiment in which amorphous solid  $\text{CH}_3\text{OD}$  with a thickness of 3 monolayer was exposed to D atoms at 10 K was also carried out. Although the resulting changes in the IR spectrum are not shown here, we observed a decrease in  $\text{CH}_3\text{OD}$  and the appearance of deuterated molecules such as  $\text{CD}_3\text{OD}$ . However, kinetic data could not be obtained for this experiment because the integrated band strengths of  $\text{CH}_2\text{DOD}$  and  $\text{CHD}_2\text{OD}$  are unknown. The bands of  $\text{CH}_2\text{DOD}$  and  $\text{CHD}_2\text{OD}$  partially overlap with those of  $\text{CH}_3\text{OD}$ , and thus we could not quantify either the decrease in  $\text{CH}_3\text{OD}$  or the yields of  $\text{CH}_2\text{DOD}$ ,  $\text{CHD}_2\text{OD}$ , and  $\text{CD}_3\text{OD}$ .

### 3.4. Exposure of Amorphous Solid $\text{CD}_3\text{OD}$ to D Atoms.

We also exposed amorphous solid  $\text{CD}_3\text{OD}$  (3 ML) to D atoms at 10 K over a period of 180 min, but no change was observed in the IR spectra (data not shown). This result indicates that the D-abstraction reaction by D atoms ( $\text{CD}_3\text{OD} + \text{D} \rightarrow \text{CD}_2\text{OD} + \text{D}_2$ ) is inhibited or alternatively that the D-abstraction reaction does proceed within the experimental duration, but the resulting  $\text{CD}_2\text{OD}$  radicals immediately react with D atoms to form  $\text{CD}_3\text{OD}$ . It is not currently clear whether the D-abstraction reaction actually proceeds on a solid surface at 10 K.

## 4. Conclusion

We performed experiments in which solid  $\text{CH}_3\text{OH}$ ,  $\text{CH}_2\text{DOH}$ , and  $\text{CHD}_2\text{OH}$  at 10 K were exposed to cold D atoms at 100 K. The H–D substitution reactions  $\text{CH}_3\text{OH} \rightarrow \text{CH}_2\text{DOH} \rightarrow \text{CHD}_2\text{OH} \rightarrow \text{CD}_3\text{OH}$  were observed; these reactions are likely to proceed via quantum mechanical tunneling. The effective reaction rate constants for the H–D substitution reaction were estimated, and the secondary kinetic isotope effects for the reaction steps were derived: values of  $0.69 \pm 0.11$  and  $0.52 \pm 0.14$  were obtained for the second and the third step, respectively, compared with the first step.

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