

# Product Vibrational State-to-State Correlation in the $F + SiH_4 \rightarrow HF(v_{HF}) + SiH_3(0v_20)$ Reaction: A Crossed Molecular Beam Ion-Imaging Study<sup>†</sup>

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The dynamics of the  $F + SiH_4 \rightarrow HF + SiH_3$  reaction has been studied using the crossed molecular beam technique with slice imaging at collision energies from 1.25 to 8.17 kcal/mol. The product silyl radical,  $SiH_3(v_2 = 0-5)$ , was detected using the  $(2 + 1)$  resonance enhanced multiphoton ionization technique. The product velocity distributions and angular distributions of this reaction were obtained from the recorded images. Experimental results show that the silyl radical product is mainly forward scattered relative to the silane beam direction, and the majority of the available energy was partitioned into the vibration of the HF product. The state-to-state correlation between the two reaction products,  $SiH_3$  and HF, was also determined. In addition, we found that the reaction cross section goes down as the collision energy increases. From these results, we conclude that the  $F + SiH_4 \rightarrow HF + SiH_3$  reaction proceeds through a direct abstraction mechanism with little or no reaction barrier.

## 1. Introduction

The reactivity of the silane molecule has been the subject of many experimental studies because of its importance in the semiconductor manufacturing and electronic industry. In the process of chemical vapor deposition (CVD), the reactions involving silane are essential to the quality of the deposited films. Kinetics and dynamics of these reactions have been extensively investigated. For example, the reaction rate constant of silane and an oxygen atom was measured by Ding and co-workers in 1993,<sup>1</sup> and the dynamics was investigated using the crossed molecular beam method in 2000.<sup>2</sup> Fluorine atom reactions with different species have also been extensively studied. The F-atom reaction with  $H_2$  provided us an excellent example of reaction resonances,<sup>3-9</sup> while interesting dynamics of the F-atom reactions with methane,<sup>10-13</sup> ethylene,<sup>14,15</sup> propyne,<sup>16</sup> propene,<sup>17</sup> and other species<sup>18-20</sup> have been revealed in the crossed molecular beam scattering studies.

Similar to the F-atom reaction with methane, the  $F + SiH_4 \rightarrow HF + SiH_3$  reaction is exothermic with an exothermicity of 45.2 kcal/mol. In 1977, Smith and co-workers measured the rate constant of  $F + SiH_4$  using the HF infrared chemiluminescence method at room temperature.<sup>21</sup> From this work, the reaction rate of  $F + SiH_4$  was found to be about 7 times faster than that of  $F + CH_4$ , and the product HF from this reaction was clearly vibrationally excited. In 2006, Shen and co-workers carried out crossed molecular beam scattering studies of the F-atom reaction with  $SiH_4$  using the universal crossed molecular beam method.<sup>22</sup> Two different channels,  $HF + SiH_3$  and  $SiH_3F + H$ , have been observed in this reaction. Experimental results indicated that the HF product of the  $HF + SiH_3$  channel is predominantly forward scattered relative to the F-atom beam direction, whereas the  $SiH_3F$  product of the  $SiH_3F + H$  channel is backward scattered, relative to the F-atom beam direction,

suggesting that two reaction channels occur with very different reaction mechanisms. The  $SiH_3F + H$  channel takes place seemingly via a typical  $S_N2$ -type reaction mechanism, while the  $HF + SiH_3$  reaction channel proceeds via a direct harpoon-type abstraction mechanism that produces a forward scattering product distribution. The relative branching ratios of the two channels was estimated to be 61% to 39% for  $HF + SiH_3$  and  $SiH_3F + H$ , respectively. In the above studies, no vibrational state resolved information was obtained for both channels.

In the recent years, crossed molecular beam technique combined with the slicing ion velocity mapping method has been proved to be an extremely powerful tool for investigating the dynamics of polyatomic reactions, especially  $X (= F, Cl, O) + CH_4$  reactions.<sup>23-25</sup> Interesting dynamics in the F-atom reaction with methane has been observed, providing an excellent example of state-to-state correlation in a polyatomic chemical reaction. In this work, we would like to report our recent experimental results on the  $HF + SiH_3$  channel in the F-atom reaction with silane using a newly developed, crossed-beams slicing imaging apparatus. The  $SiH_3$  product at specific vibrational states was detected via a  $(2 + 1)$  resonance enhanced multiphoton ionization (REMPI) technique.<sup>26</sup> By measuring the vibrational state specific product ( $SiH_3$ ) ion velocity images, the correlated HF product vibrational state distributions were clearly determined. From these measurements, the vibrational state-to-state correlated dynamics in the  $F + SiH_4 \rightarrow HF + SiH_3$  reaction were investigated in details.

## 2. Experimental Methods

The reaction of  $F + SiH_4$  has been studied in our laboratory using the crossed molecular beam scattering technique with slicing ion imaging, in combination with REMPI detection. The experiments were carried on a newly developed apparatus that has been described in detail elsewhere.<sup>27</sup> The new crossed beams ion imaging apparatus has three vacuum regions: two beam source chambers and one reaction/detection chamber. One of

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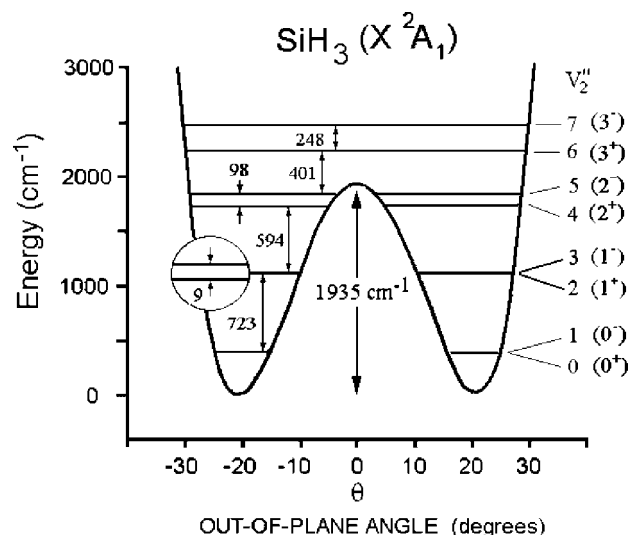
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the source chambers is fixed, and the other one can be rotated from  $60^\circ$  to  $180^\circ$  relative to the fixed beam source so that collision energy can well controlled during the experiment. For the current experiment, the F-atom beam was generated by a double-stage pulsed discharge before the expansion of a premixed sample of  $F_2$  (5%) and He (95%) (Spectra Gases Inc.)<sup>28</sup> or a premixed sample of  $NF_3$  (50%),  $F_2$  (2.5%), and He (47.5%) in the fixed source chamber. No contaminant that could affect the results has been found in the F-atom beam. The two samples were used to cover the difference collision energy range. The F-atom beam was led to the reaction chamber through a 1.5 mm diameter skimmer and was then crossed with a skimmed  $SiH_4$  beam in the reaction chamber from the rotating source chamber. The  $SiH_4$  molecular beam was generated by expanding a sample of 50%  $SiH_4$  in Ar at a stagnation pressure of about 3 atm through a carefully adjusted pulsed valve (General Valve) with a risetime of about 50  $\mu s$ . The silane beam was then skimmed once by a 1.5 mm orifice skimmer before entering the reaction chamber. The speed of the  $SiH_4$  beam was about 680 m/s. The speed of the F-atom beam from the discharge of the premixed sample of  $F_2$  (5%) and He (95%) was about 1900 m/s, whereas that from the discharge of the premixed sample of  $NF_3$  (50%),  $F_2$  (2.5%), and He (47.5%) was about 1000 m/s. The collision energy was varied by changing the crossing angle between the F-atom beam and the silane beam, as well as exchanging the two premixed samples for the F-atom beam. Using the premixed sample of  $NF_3$  (50%),  $F_2$  (2.5%), and He (47.5%), collision energy as low as 1.25 kcal/mol can be achieved with a beam crossing angle of  $60^\circ$ . Using the premixed sample of  $F_2$  (5%) and He (95%), a collision energy of 8.17 kcal/mol can be reached with a crossing angle of  $125^\circ$  between the two beams.

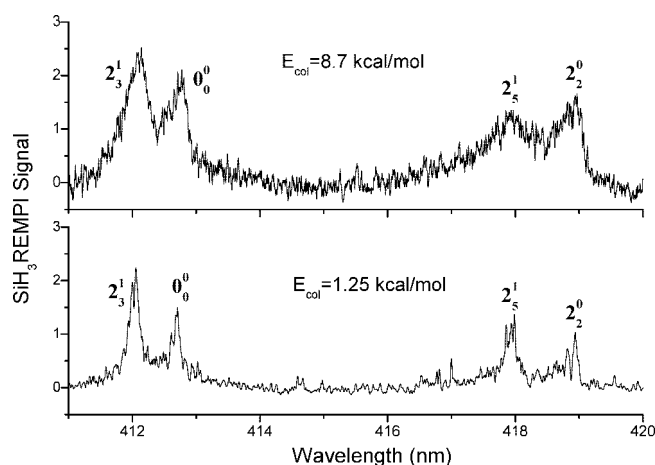
The crossing region of two beams was located at the center of the ion-imaging optics. The reaction product,  $SiH_3$ , was detected by a (2 + 1) REMPI scheme via the  $\tilde{E}^2A''_2$  (4p) Rydberg state.<sup>26</sup> The UV photon of 410–420 nm was generated by a YAG-pumped dye laser output. The typical UV pulse energy was 5–6 mJ, and a lens with a focal length of 43 cm was used to focus the laser beam into the crossing region. The  $SiH_3$  product ion was velocity mapped to an imaging quality microchannel plate (MCP) detector. The ion imaging signal was coupled out via a phosphor screen and detected by a CCD camera. The image signal was then transferred to a computer and accumulated over a significant amount of laser shots. The total voltage of ion optics was held at 700 V, at which the whole image could be recorded. The turnaround time of the products ion packet was more than 200 ns at this voltage setting, and a slicing window of 30 ns was used for the  $SiH_3$  product ions for proper slicing imaging. The 30 ns slicing window was sufficiently narrow for high resolution and 3D velocity distribution reconstruction.<sup>29</sup> The (2 + 1) REMPI spectrum of the  $SiH_3$  product was recorded by a photomultiplier tube looking on the phosphor screen. During the experiments, the vacuum was maintained at about  $1 \times 10^{-5}$  torr in the source chambers and at about  $4 \times 10^{-7}$  torr in the detection chamber.

### 3. Results and Discussion

**3.1.  $SiH_3$  REMPI Spectrum.** The (2 + 1) REMPI spectroscopy of the  $SiH_3$  radical has been carefully studied by Hudgens and co-workers.<sup>26</sup> From the detailed REMPI spectra, it was found that the  $SiH_3$  radical in the ground electronic state is pyramidal with its potential minimum at the angle of  $\theta_m = 20.7^\circ$ , where  $\theta$  is defined as the angle between the line through the Si–H axis and the plane formed from the three hydrogen



**Figure 1.** One-dimension, double minimum potential function of  $SiH_3$  along the  $C_3$  axis. This plot also shows  $\nu_2$  umbrella vibrational levels of  $SiH_3$  as a function of symmetric out-of-plane bending “umbrella” angle.



**Figure 2.** (2 + 1) REMPI spectra of the  $SiH_3$  product via the  $\tilde{E}^2A''_2$  (4p) Rydberg state at the collision energy of 1.25 and 8.17 kcal/mol.

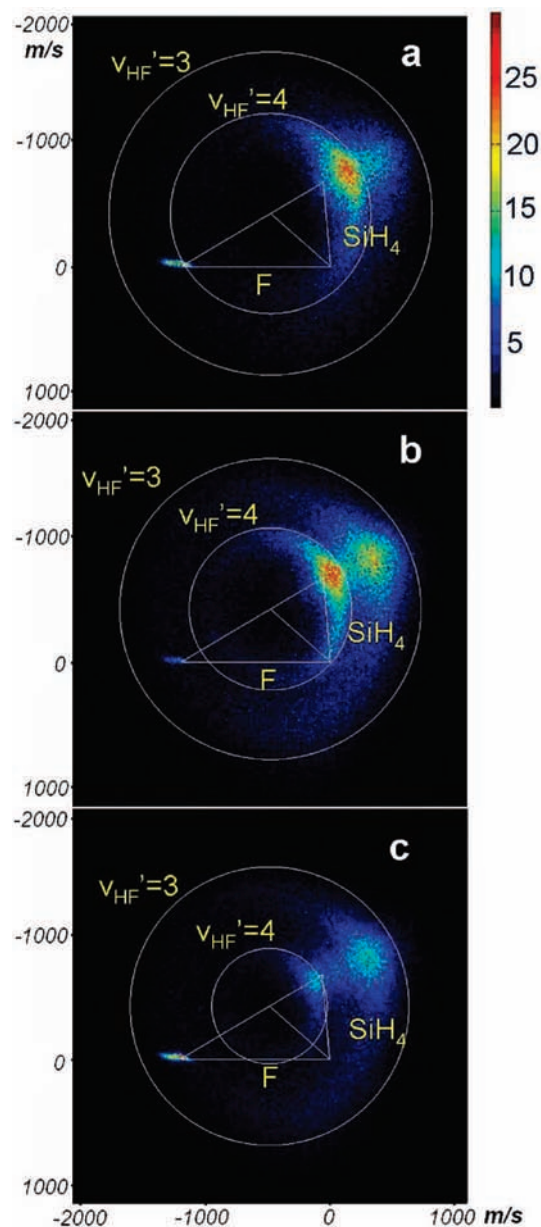
atoms. The  $SiH_3$  radical is similar to ammonia and has two potential energy minima along the  $C_3$  axis, the umbrella mode coordinate, shown in Figure 1.<sup>26</sup> For a typical double-minimum potential, the vibrational levels below the potential barrier have collapsed into closely spaced doublet levels. The splitting of these levels are related the potential barrier height. Those levels above the barrier are well separated. The (2 + 1) REMPI spectrum for the  $SiH_3$  radical has been measured and assigned by Hudgens and co-workers.<sup>26</sup> This is helpful for the current dynamical studies of the  $F + SiH_4 \rightarrow HF + SiH_3$  reaction. In this work, we have measured the REMPI spectrum of the  $SiH_3$  radical product from the  $F + SiH_4$  reaction. Figure 2 shows the (2 + 1) REMPI spectrum of the  $SiH_3$  product from the  $F + SiH_4$  reaction via the  $\tilde{E}^2A''_2$  (4p) Rydberg state. The assignment of the REMPI spectrum of the  $SiH_3$  radical was made according to Hudgens and co-workers.<sup>26</sup> From the REMPI spectrum, the  $SiH_3$  product is only excited in the umbrella mode. Vibrational excited states up to  $\nu_2 = 5$  in the umbrella mode are populated in the  $SiH_3$  product from the  $F + SiH_4$  reaction.

Using the (2 + 1) REMPI detection scheme, the  $SiH_3$  product at specific vibrational states from the  $F + SiH_4$  reaction can be efficiently detected. However, detection efficiency of  $SiH_3$  at

different vibrational states are different, and therefore relative product populations of SiH<sub>3</sub> in different vibrational states cannot be quantitatively measured. From the REMPI spectra shown in Figure 2, the shapes of the different vibrational bands are different at different collision energies. At 1.25 kcal/mol, the vibrational bands are quite narrow, and the different bands are well separated from each other in the spectrum. At a higher collision energy of 8.17 kcal/mol, each vibrational band becomes broader. Consequently, the vibrational features become somewhat overlapped. These results indicate that the SiH<sub>3</sub> product is less rotationally excited at lower collision energy and significantly more rotationally excited at higher collision energy. From the relative integrated intensities of the different vibrational bands, we found that the higher vibrationally excited state of SiH<sub>3</sub> becomes relatively more important at higher collision energy, suggesting that the SiH<sub>3</sub> product vibrational excitation in the umbrella mode increases as the collision energy.

**3.2. Reaction Product Images.** Sliced velocity map ion images for the SiH<sub>3</sub> radical at different vibrational states were measured at different collision energies using the REMPI detection scheme since the first six vibrational levels in the umbrella mode in SiH<sub>3</sub> appear in pairs,  $v_2 = 0, 1, 2$  and  $3, 4, 5$  with small energy difference. We have measured velocity ion images for each pair at the collision energy of 2.5 kcal/mol, and the images for the SiH<sub>3</sub> product at each pair vibrational levels are very similar. This suggests that the reaction dynamics producing the vibrationally excited SiH<sub>3</sub> products at these pair states are similar. Therefore, the data presented in this work are all based on the velocity map ion images of the SiH<sub>3</sub> product at  $v_2 = 0, 3, 5$  states, which were detected through (2 + 1) REMPI at the wavelengths of 412.7, 412.0, 417.9 nm respectively. Figure 3 shows the typical raw images for the SiH<sub>3</sub> product at the  $v_2 = 0, 3, 5$  levels at the collision energy of 2.5 kcal/mol. Two ringlike structures are well resolved in these images. From the well-known energetics of this reaction, these structures are assigned to the correlated HF product at  $v' = 4$  (inner ring) and  $v' = 3$  (outer ring). Clearly, the correlated HF product is produced exclusively in  $v' = 3$  and 4 states. From the raw images, the SiH<sub>3</sub> product is also mainly forward scattered with respect to the SiH<sub>4</sub> beam direction. This result is consistent with the result of previous crossed beam studies of the same reaction using a universal molecular machine.<sup>22</sup> Similarly, the raw images for the SiH<sub>3</sub> product at two other collision energies, 1.25 and 8.17 kcal/mol, are also recorded.

**3.3. Product Angular Distributions.** In order to obtain the product velocity and angular distributions, the density-to-flux transformation is performed through an analysis program developed earlier.<sup>27</sup> Product angular distributions were obtained by integrating signals at the different angles in the density-to-flux transformed images. Figure 4 shows the total product angular distributions at three collision energies. At all three collision energies, the SiH<sub>3</sub> product is mainly forward scattered relative to the SiH<sub>4</sub> direction. This is also true for the SiH<sub>3</sub> product at different vibrational states. The angular distribution of the HF product at  $v' = 3$  and 4 states are also clearly forward scattered, relative to the F-atom beam. The degree of the forwardness for the SiH<sub>3</sub> product is also different at different collision energies. At the collision energy of 1.25 kcal/mol, the reaction product is mainly forward scattered. However, there is a significant amount of product at the sideway and backward scattering directions. As the collision energy increases, the backward/sideway scattered product flux becomes smaller relative to that for the forward scattering direction. At the collision energy of 8.17 kcal/mol, backward scattering product

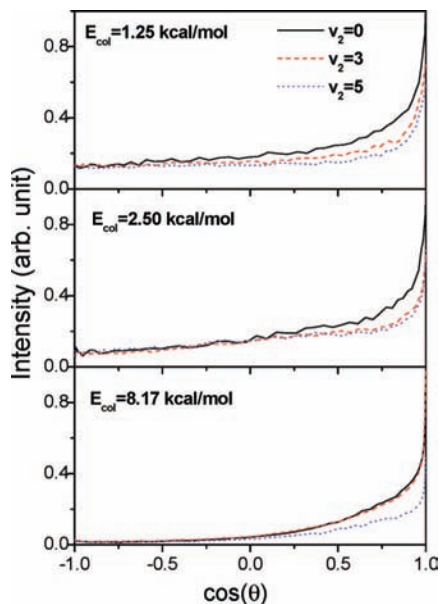


**Figure 3.** Raw images of the SiH<sub>3</sub> product at the collision energy of 2.5 kcal/mol: (a) SiH<sub>3</sub>( $v_2 = 0$ ); (b) SiH<sub>3</sub>( $v_2 = 3$ ); (c) SiH<sub>3</sub>( $v_2 = 5$ ).

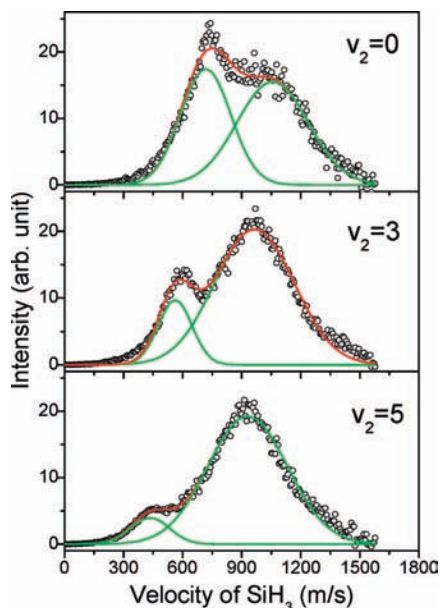
becomes very small relative to the forward scattering product and the total product angular distribution is exclusively forward scattered. These results suggest that as the collision energy increases the product angular distribution becomes more forward scattered. For the SiH<sub>3</sub> product at different vibrational states, the angular distributions are also different. It seems that if the vibrational excitation of SiH<sub>3</sub> is smaller, the SiH<sub>3</sub> product is overall more forward scattered.

**3.4. Product Velocity Distributions and State-to-State Correlation between HF and SiH<sub>3</sub>.** Through the density-to-flux transformation of the raw images, velocity distributions of the SiH<sub>3</sub> product can also be obtained. Figure 5 shows the velocity distributions of the SiH<sub>3</sub> product at  $v_2 = 0, 3, 5$  obtained from the images at the collision energy of 2.5 kcal/mol. The curve contains two features, which can be assigned to the HF  $v' = 3$  and  $v' = 4$  levels. The relative branching ratios of HF( $v' = 3$ ) and HF( $v' = 4$ ) were determined by fitting the distributions using two components. The relative branching ratios obtained are listed Table 1. The relative branching ratios of the HF( $v' =$





**Figure 4.** Angular distributions of the  $\text{SiH}_3$  product with respect to the  $\text{SiH}_4$  beam direction at 1.25, 2.5, and 8.17 kcal/mol. The black solid line is the result for  $\text{SiH}_3(v_2 = 0)$ , the red dashed line is for  $\text{SiH}_3(v_2 = 3)$ , and the blue dotted line is for  $\text{SiH}_3(v_2 = 5)$ .



**Figure 5.** Velocity distributions of the  $\text{SiH}_3$  product at 2.5 kcal/mol for  $\text{SiH}_3$  at  $v_2 = 0, 3,$  and  $5$ . The open circles in the figure are the experiment results, the green lines are simulated results for the correlated  $v' = 3$  and  $v' = 4$  HF products, and the red line is the total simulated result.

3) and  $\text{HF}(v' = 4)$  products vary considerably for different vibrationally excited  $\text{SiH}_3$  products. For the  $\text{SiH}_3$  product at  $v_2 = 5$ , nearly all HF product is populated in the  $v' = 3$  state, with very little in  $v' = 4$  state (9%), whereas for the  $\text{SiH}_3$  product at  $v_2 = 0$ , there is a considerable increase in the  $\text{HF}(v' = 4)$  population (44%). From these results, there is clearly vibrational state-to-state correlation between the  $\text{SiH}_3$  and HF products. The  $\text{SiH}_3$  and HF products seem to be anticorrelated vibrationally. The higher the  $\text{SiH}_3$  vibrational excitation, the lower the HF vibrational excitation. From the correlated images, we can also conclude that the rotational excitation of the  $\text{HF}(v' = 3)$  product is significantly higher than that of the  $\text{HF}(v' = 4)$  product. This is largely due to the different available energies of the two

**TABLE 1: Averaged  $\cos \theta$ , Product Energy Disposals, and HF Vibrational Branching Ratios (All Energies in kcal/mol)**

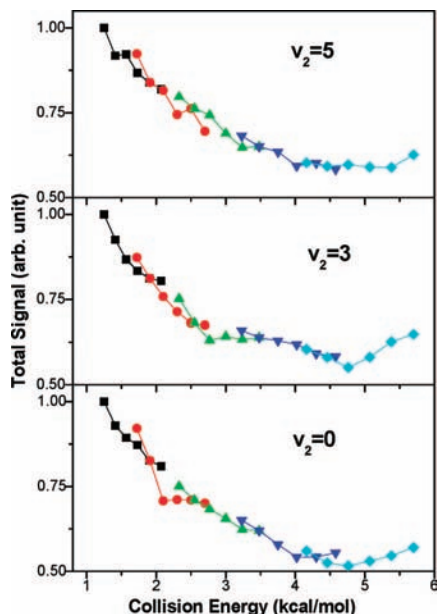
	$v_{\text{SiH}_3} = 0$	$v_{\text{SiH}_3} = 3$	$v_{\text{SiH}_3} = 5$
$E_C = 1.25$			
$\langle \cos \theta \rangle$	0.23	0.18	0.15
$\langle E_T \rangle_{\text{total}}$	5.6	6.7	7.1
$\langle E_{v'} \rangle_{\text{HF}}$	36.3	33.8	32.7
$\langle E_{v'} \rangle_{\text{SiH}_3}$	0	2.1	4.1
$\langle E_{v'} \rangle_{\text{total}}$	36.3	35.9	36.8
$\nu_{\text{HF}} = 3$ (%)	62	87	100
$\nu_{\text{HF}} = 4$ (%)	38	13	0
$E_C = 2.5$			
$\langle \cos \theta \rangle$	0.31	0.25	0.21
$\langle E_T \rangle_{\text{total}}$	8.2	8.2	8.2
$\langle E_{v'} \rangle_{\text{HF}}$	36.9	34.4	33.4
$\langle E_{v'} \rangle_{\text{SiH}_3}$	0	2.1	4.1
$\langle E_{v'} \rangle_{\text{total}}$	36.9	36.5	37.5
$\nu_{\text{HF}} = 3$ (%)	56	81	91
$\nu_{\text{HF}} = 4$ (%)	44	19	09
$E_C = 8.17$			
$\langle \cos \theta \rangle$	0.56	0.54	0.48
$\langle E_T \rangle_{\text{total}}$	11.2	11.0	11.0
$\langle E_{v'} \rangle_{\text{HF}}$	39.2	36.6	34.6
$\langle E_{v'} \rangle_{\text{SiH}_3}$	0	2.1	4.1
$\langle E_{v'} \rangle_{\text{total}}$	39.2	38.7	38.7
$\nu_{\text{HF}} = 3$ (%)	31	59	78
$\nu_{\text{HF}} = 4$ (%)	69	41	22

microchannels. Obviously, the  $\text{HF}(v' = 3)$  product has more product rotational states for energy disposal and the  $\text{HF}(v' = 4)$  product has fewer.

As collision energy increases, the  $\text{HF}(v' = 4)$  product channel becomes more and more important than the  $\text{HF}(v' = 3)$  channel, while the  $\text{HF}(v' = 3)$  channel is less important. This seems to be true for all measured  $\text{SiH}_3$  vibrational channels. At the collision energy of 1.25 kcal/mol, the correlated HF product are all in the  $v' = 3$  state with no population at all in the  $v' = 4$  state for the  $\text{SiH}_3$  ( $v_2 = 5$ ) product channel. At the collision energy of 8.17 kcal/mol, the  $\text{HF}(v' = 4)$  channel already accounts for 22% of all HF product for the same  $\text{SiH}_3$  ( $v_2 = 5$ ) channel. For the  $\text{SiH}_3$  ( $v_2 = 0$ ) channel, the branching ratio of the  $\text{HF}(v' = 4)$  increases from 38% at 1.25 kcal/mol to 69% at 8.17 kcal/mol.

A rather remarkable observation is the large proportion of the available energy deposited into the vibration of the HF product in such a polyatomic reaction. From Table 1, more than 70% of the available energy is partitioned into the HF vibration. This means that the majority of the available energy in this reaction is deposited to the HF vibration. This extreme vibrational excitation is very likely produced via a peculiar extraction mechanism, in which the F atom abstracts an H atom from the silane molecule at a relatively large distance and then flies away. Such reaction mechanism can also explain the strong forward scattering of the reaction products, which is produced likely via abstraction at large impact parameters.

**3.5. Collision Energy Dependence of the Total REMPI Signal.** Integral cross sections (ICS) for the  $\text{SiH}_3$  product at specific vibrational states were also measured for the title reaction. Figure 6 shows the measured total REMPI signal for  $\text{SiH}_3$  at  $v_2 = 0, 3, 5$  versus the collision energy. These data were measured by integrating the total  $\text{SiH}_3$  REMPI ion signals when the REMPI laser was parked at a specific transition. The collision energy variation is achieved by both changing the crossing angle between the two beams from  $65^\circ$  to  $90^\circ$  and varying the mixing ratio of the 5%  $\text{F}_2/\text{He}$  gas and the  $\text{NF}_3$  gas in a premixed gas sample. By varying this mixing ratio, the



**Figure 6.** Relative integral cross section (ICS) as a function of the collision energy in the  $F + SiH_4 \rightarrow HF + SiH_3$  reaction for different vibrationally excited  $SiH_3$  products.

F-atom beam velocity can be changed between 1000–2000 m/s. In this experiment, we have used a total of five different mixtures, in which the percentage of the  $NF_3$  sample in the mixture are 50%, 25%, 16%, 10%, 0%, respectively. The collision energy for the  $F + SiH_4$  reaction can be varied from 1.25 to about 5.7 kcal/mol with the beam crossing angle changing from  $65^\circ$  to  $90^\circ$  in addition to the gas mixture changes. The five measurements using different gas mixtures are put together by scaling the different sets of data to a single overlapping point. This allows us to put together the whole total REMPI signal plots from 1.25 to 5.7 kcal/mol for specific vibrationally excited  $SiH_3$  products at  $v_2 = 0, 3, 5$ . The total REMPI signal curves of the  $SiH_3$  product in different vibrational state are very much similar. All curves decrease monotonically as the collision energy increases. Since the effects of the density-flux transformation and the relative velocity are somewhat similar at different collision energies, the trend of the total REMPI plots should be similar to that of the integral cross sections. Thus, we can conclude that the integral cross section for all reaction products should also decrease monotonically as the collision energy increases. This means that the  $F + SiH_4 \rightarrow HF + SiH_3$  reaction is likely a barrierless reaction governed by an attractive potential, which is quite consistent with the picture obtained from the product energy partitions.

#### 4. Conclusions

In this work, we have investigated the dynamics of the  $F + SiH_4 \rightarrow HF + SiH_3$  reaction using the time-sliced velocity map ion imaging technique with  $(2 + 1)$  REMPI detection of the  $SiH_3$  product. From the REMPI spectra of the  $SiH_3$  product, we found that the  $SiH_3$  product is vibrationally excited only in the umbrella vibrational mode. Sliced raw images of the  $SiH_3$  product at  $v_2 = 0, 3, 5$  levels were acquired. After density-to-flux transformation of these raw images, both the velocity and angular distributions of the product  $SiH_3$  at  $v_2 = 0, 3, 5$  were obtained at three collision energies: 1.25, 2.50, and 8.17 kcal/mol. From the experimental results, we found that the majority of the available reaction energy is partitioned into the vibration degrees of freedom in the HF and  $SiH_3$  products. Vibrational

state distributions of the correlated HF product can be determined for the  $SiH_3$  product at various vibrationally excited states. Therefore, information on the vibrational state-to-state correlation between the  $SiH_3$  and HF products can be derived. From these results, vibrational state distributions for the  $SiH_3$  and HF products are clearly anticorrelated. In addition, we found that the  $SiH_3$  product at  $v_2 = 0, 3, 5$  vibrational states are clearly forward scattered. At the low collision energy, the  $SiH_3$  product has a large forward scattering peak, and backward and sideways scatterings are also significant. As the collision increases, however, the product angular distribution becomes more forward scattered. The integral cross sections for the  $SiH_3$  product at  $v_2 = 0, 3, 5$  states have been measured and obviously decrease monotonically as the collision energy increase. These results suggest that the  $F + SiH_4 \rightarrow HF + SiH_3$  reaction proceeds via a direct abstraction mechanism with little or no barrier.

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#### References and Notes

- (1) Ding, L.; Marshall, P. J. *J. Chem. Phys.* **1993**, *98* (11), 8545.
- (2) Lin, J. J.; Lee, Y. T.; Yang, X. J. *J. Chem. Phys.* **2000**, *113* (5), 1831.
- (3) Skodje, R. T.; Skoouteris, V. D.; Manolopoulos, D. E.; Lee, S.-H.; Dong, F.; Liu, K. J. *J. Chem. Phys.* **2000**, *112*, 4536; *Phys. Rev. Lett.* **2000**, *85*, 1206.
- (4) Neumark, D. M.; Wodtke, A. M.; Robinson, G. N.; Hayden, C. C.; Lee, Y. T. *Phys. Rev. Lett.* **1984**, *53*, 226; *J. Chem. Phys.* **1985**, *82*, 3045.
- (5) Qiu, M.; Ren, Z.; Che, L.; Dai, D.; Harich, S. A.; Wang, X.; Yang, X. *Chin. J. Chem. Phys.* **2006**, *20*, 1.
- (6) Qiu, M.; Ren, Z.; Che, L.; Dai, D.; Harich, S.; Wang, X.; Yang, X.; Xu, C.; Qian, D.; Gustafsson, M.; Skodje, R. T.; Sun, Z.; Zhang, D. H. *Science* **2006**, *311*, 1440.
- (7) Qiu, M.; Ren, Z.; Che, L.; Dai, D.; Harich, S. A.; Wang, X.; Yang, X.; Xu, C.; Xie, D.; Gustafsson, M.; Skodje, R. T.; Sun, Z.; Zhang, D. H. *Science* **2006**, *311* (5766), 1440.
- (8) Che, L.; Ren, Z. F.; Wang, X. G.; Dong, W. R.; Dai, D. X.; Wang, X. Y.; Zhang, D. H.; Yang, X. M.; Sheng, L. S.; Li, G. L.; Werner, H. J.; Lique, F.; Alexander, M. H. *Science* **2007**, *317* (5841), 1061.
- (9) Yang, X.; Zhang, D. H. *Acc. Chem. Res.* **2008**, *41* (8), 981.
- (10) Harper, W. W.; Nizkorodov, S. A.; Nesbitt, D. J. *J. Chem. Phys.* **2000**, *113*, 3670.
- (11) Harper, W. W.; Nizkorodov, S. A.; Nesbitt, D. J. *J. Chem. Phys. Lett.* **2001**, *335*, 381.
- (12) Lin, J. J.; Zhou, J.; Shiu, W.; Liu, K. *Science* **2003**, *300*, 966.
- (13) Zhou, J.; Lin, J. J.; Shiu, W.; Pu, S.-C.; Liu, K. *J. Chem. Phys.* **2003**, *119*, 4997.
- (14) Parson, J. M.; Lee, Y. T. *J. Chem. Phys.* **1972**, *56*, 4658.
- (15) Robinson, G. N.; Continetti, R. E.; Lee, Y. T. *J. Chem. Phys.* **1990**, *92*, 275.
- (16) Ran, Q.; Yang, C.-H.; Shen, G.; Lee, Y. T.; Yang, X. *J. Chem. Phys.* **2005**, *122*, 044307.
- (17) Ran, Q.; Yang, C.-H.; Shen, G.; Lee, Y. T.; Yang, X. *J. Chem. Phys.* **2004**, *121*, 6302.
- (18) Shobatake, K.; Lee, Y. T.; Rice, S. A. *J. Chem. Phys.* **1973**, *59*, 1453.
- (19) Shobatake, K.; Parson, J. M.; Lee, Y. T.; Rice, S. A. *J. Chem. Phys.* **1973**, *59*, 1416.
- (20) Parson, J. M.; Shobatake, K.; Lee, Y. T.; Rice, S. A. *J. Chem. Phys.* **1973**, *59*, 1402.
- (21) Smith, D. J.; Setser, D. W.; Kim, K. C.; Bogan, D. J. *J. Phys. Chem.* **1977**, *81* (9), 898.
- (22) Shen, G.; Yang, X.; Shu, J.; Yang, C.-H.; Lee, Y. T. *J. Chem. Phys.* **2006**, *125* (13), 133103.
- (23) Lin, J. J.; Zhou, J.; Shiu, W.; Liu, K. *Rev. Sci. Instrum.* **2003**, *74* (4), 2495.
- (24) Lin, J. J.; Zhou, J. G.; Shiu, W. C.; Liu, K. P. *Science* **2003**, *300* (5621), 966.
- (25) Huang, C.; Li, W.; Suits, A. G. *J. Chem. Phys.* **2006**, *125* (13), 133107.
- (26) Johnson, R. D., III; Tsai, B. P.; Hudgens, J. W. *J. Chem. Phys.* **1989**, *91* (6), 3340.

(27) Wu, G.; Zhang, W.; Pan, H.; Shuai, Q.; Jiang, B.; Dai, D.; Yang, X. *Rev. Sci. Instrum.* **2008**, 79 (9), 094104.

(28) Ren, Z.; Qiu, M.; Che, L.; Dai, D.; Wang, X.; Yang, X. *Rev. Sci. Instrum.* **2006**, 77 (1), 016102.

(29) Shiu, W.; Lin, J. J.; Liu, K.; Wu, M.; Parker, D. H. *J. Chem. Phys.* **2004**, 120 (1), 117.

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