

Twenty-seven-layer barium chromium oxide,  $\text{BaCrO}_3$ , was prepared at  $1200^\circ\text{C}$  and 60–65 kbar by reaction of  $\text{CrO}_2$  and  $\text{Ba}_2\text{CrO}_4$ . The compound crystallizes in space group  $R\bar{3}m$  with hexagonal cell dimensions  $a = 5.652 \pm 0.002 \text{ \AA}$ ,  $c = 62.75 \pm 0.02 \text{ \AA}$ . The structure was determined from 1070 independent reflections, of which 811 were considered observed, collected by automated counter methods and refined by least-squares methods to a conventional  $R$  value of 4.1%. The structure consists of a 27-layer stacking sequence of close-packed  $\text{BaO}_3$  layers, Zhdanov notation for the sequence (3)2(2)2, with all of the  $\text{O}_6$  octahedral sites occupied by Cr. Strings of four pairs of face-sharing octahedra in which the pairs are linked to each other by corner sharing are in turn joined to each other by octahedra sharing only corners. The structure is thus closely related to the four-layer and six-layer polytypes of  $\text{BaCrO}_3$ . Bariums in two of the five crystallographic positions showed large thermal anisotropy and are better described in terms of half-atom occupancy in split sites.

*Neutron Diffraction Experiments on  $\text{CsCrI}_3$  at 300, 77, and  $1.2^\circ\text{K}$ .* H. W. ZANDBERGEN AND D. J. W. IJDO. Gorlaeus Laboratories, University of Leiden, P.O. Box 9502, 2300 RA Leiden, The Netherlands.  $\text{CsCrI}_3$  has been investigated by neutron powder diffraction at room temperature,  $77^\circ\text{K}$ , and  $1.2^\circ\text{K}$ . It undergoes a phase transition at  $150^\circ\text{K}$  due to the cooperative Jahn–Teller effect. The high-temperature form,  $\alpha\text{-CsCrI}_3$  (hexagonal, space group  $P6_3/mmc$ ,  $a = 8.127(1) \text{ \AA}$ ,  $c = 6.944(1) \text{ \AA}$ ,  $Z=2$ ), adopts the  $\text{BaNiO}_3$  structure with a local Jahn–Teller distortion. The low-temperature form  $\beta\text{-CsCrI}_3$  (orthorhombic, space group  $Pbcn$ ,  $a = 8.102(1) \text{ \AA}$ ,  $b = 13.792(1) \text{ \AA}$ ,  $c = 6.900(1) \text{ \AA}$ ,  $Z=4$ ), has a structure not yet reported for a Jahn–Teller distorted  $\text{BaNiO}_3$  structure. It is shown that the low-temperature form can be derived from the  $\text{BaNiO}_3$  structure by means of canting of triangles, formed by the three common  $\text{I}^-$  ions of two adjacent  $\text{CrI}_6^{4-}$  octahedra. The magnetic structure of  $\beta\text{-CsCrI}_3$  at  $1.2^\circ\text{K}$  is found to consist of an antiparallel sequence of ferromagnetic (001) planes with a magnetic moment in the  $|100|$  direction of  $3.26 \mu_B$ .

*X-Ray Study of the Ordering of the Alkali Ions in the Intercalation Compounds  $\text{Na}_x\text{TiS}_2$  and  $\text{Li}_x\text{TiS}_2$ .* TJPKE HIBMA. Brown Boveri Research Center, CH-5405 Baden-Dättwil, Switzerland.

The arrangement of the alkali ions in electrochemically intercalated  $\text{TiS}_2$  crystals was studied by diffuse X-ray techniques. In  $\text{Na}_x\text{TiS}_2$  a stage 3 phase was discovered in addition to the known stage 1 and 2 phases. Three types of three-dimensionally ordered superstructures were observed: a  $(3^{1/2} \times 3^{1/2})$  superstructure for stage 2 and 3 phases, a  $(2 \times 2)$  and a  $(2 \times 3^{1/2})$  superstructure for stage 1 and 2 phases. The appearance of these superstructures is consistent with a screened Coulomb interaction between the sodium ions. In the single-phase region above  $x = \frac{1}{2}$ , diffuse rings show up in addition to the  $(2 \times 3^{1/2})$  superlattice. These rings are caused by local rearrangements of the ions to accommodate the excess sodium ions. The  $(3^{1/2} \times 3^{1/2})$  and  $(2 \times 2)$  superstructures were also observed in  $\text{Li}_x\text{TiS}_2$  crystals.

*Studies of Layered Uranium(VI) Compounds. III. Structural Investigations of Hydrogen Uranyl Phosphate and Arsenate Tetrahydrates below the Respective Transition Temperatures of  $274$  and  $301^\circ\text{K}$ .* MARK G. SHILTON AND ARTHUR T. HOWE. Department of Inorganic and Structural Chemistry, University of Leeds, Leeds LS2 9JT, England.

The layered hydrates  $\text{HUO}_2\text{PO}_4 \cdot 4\text{H}_2\text{O}$  (HUP), and  $\text{HUO}_2\text{AsO}_4 \cdot 4\text{H}_2\text{O}$  (HUAs), which are proton-conducting solid electrolytes above the conductivity transitions at  $274$  and  $301^\circ\text{K}$  respectively, have been shown, using powder X-ray diffraction, to change from tetragonal to orthorhombic symmetry below these temperatures. For HUP the unit-cell dimensions were  $a = 6.985(5)$  and  $c = 17.45(1) \text{ \AA}$  at  $290^\circ\text{K}$ , and  $a = 6.966(5)$ ,  $b = 7.004(5)$ , and  $c = 17.43(1) \text{ \AA}$  at  $260^\circ\text{K}$ . The values for HUAs were  $a = 7.150(2)$  and  $c = 17.608(5) \text{ \AA}$  at  $305^\circ\text{K}$ , and  $a = 7.128(2)$ ,  $b = 7.168(2)$ , and  $c = 17.613(5) \text{ \AA}$  at  $293^\circ\text{K}$ . The enthalpies of these displacive-type transitions were found from differential scanning calorimetry to be less than  $0.5 \text{ kJ}$  per mole of water for both compounds. Such a small value indicates that the rigid-like water lattices existing below the transitions do not become liquid-like above the transitions. The infrared spectra of HUP and HUAs both above the transitions and down to  $80^\circ\text{K}$  showed clear evidence of the presence of  $\text{H}_3\text{O}^+$  ions, showing that the conductivity transitions are not caused by a loss of carriers. Rather, the antiferroelectric ordering, known to exist for HUAs, would appear to cause the conductivity drop. Upon this indication of ordering within the water layers, two possible related H-bond ordered structures have