

Rare Earth Transition Metal Sulfides, $LnMS_3$ *

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Ternary rare earth transition metal sulfides $LnMS_3$ with $Ln = La, Nd, \text{ and } Gd$, and $M = V \text{ and } Cr$; as well as $Ln = La$ and $M = Mn, Fe, Co, \text{ and } Ni$ have been prepared and characterized. The vanadium and chromium sulfides crystallize in a monoclinic layer structure isotypic with $LaCrS_3$, while the other $LnMS_3$ sulfides crystallize in a hexagonal structure. Chemical shifts of the metal K -absorption edge and XPS binding energies of core levels indicate that the transition metal is trivalent in the V and Cr sulfides, while it is divalent in the Mn, Fe, Co, and Ni sulfides. Electrical and magnetic properties of the sulfides are discussed in terms of their structures and the electronic configurations of the transition metal ions.

1. Introduction

A large number of perovskite oxides of the general formula $LnMO_3$ ($Ln = \text{lanthanide metal and } M = \text{transition metal}$) exhibiting a variety of structures and electronic properties have been investigated in detail (1). In view of the limited information available on the analogous ternary sulfur compounds, we have taken up a systematic study of several members of the $LnMS_3$ family, where M is a first-row transition metal ion. In the literature, the preparation of only a few of the $LnMS_3$ members has been reported (2), while their characterization has been limited to a structural study of $LaCrS_3$ (3). In the present study, we have investigated the structures as well as electronic and magnetic and related properties

of $LnMS_3$ compounds where M is V, Cr, Mn, Fe, Co, and Ni.

2. Experimental

The ternary sulfides $LnMS_3$ were prepared from the corresponding ternary oxides by reaction with H_2S or CS_2 vapor at elevated temperatures. $LnMS_3$ compounds with $M = V$ and Cr were obtained by the H_2S reaction at 1200 and 1000°C, respectively (for a total period of ~150 hr). $LaMS_3$ with $M = Mn, Fe, Co, \text{ and } Ni$ were synthesized by passing dry oxygen-free nitrogen saturated with CS_2 vapor over the oxides kept at 1100°C. Each sample was repeatedly heated in the stream of CS_2/N_2 with intermittent grindings until the reaction was complete as indicated by its constant weight. It generally took around 150 hr of total reaction time for the reaction to be complete. Chemical analysis for $Ln, M, \text{ and } S$, carried out by standard methods,

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showed that the compositions of the sulfides are very close to $LnMS_3$ in all the cases studied. We must mention here that great care had to be exercised in ensuring the completion of sulfide formation and to establish exact conditions for the preparations. We have taken over a year to prepare these compounds reliably. The compounds were stored in evacuated sealed tubes.

X-Ray powder patterns recorded with a Philips diffractometer (vertical goniometer type PW 1050/70 employing Ni-filtered $CuK\alpha$ radiation) showed the compounds to be monophasic with unique crystal structures (Table I).

Experimental procedures for the study of metal K -absorption edges and X-ray photoelectron spectra (XPS) have been described elsewhere (4, 5). Electrical resistivities reported here are from four-probe measurements on pressed pellets carried out in an atmosphere of dry oxygen-free nitrogen in the temperature range 300–600 K; for the iron compound, resistivity was measured up to 800 K. The pellets were presintered under the same conditions employed for the preparation. Seebeck coefficients were also measured on the same pellets (6). Magnetic susceptibility measurements were carried out down to 100 K under vacuum by the Faraday method employing a Cahn electrobalance.

TABLE I
UNIT CELL PARAMETERS OF $LnMS_3$

Compound	a (Å)	b (Å)	c (Å)	β (°)	Unit cell volume, V (Å ³)
LaVS ₃	5.985	16.854	11.158	89.7	1125
NdVS ₃	5.886	16.800	10.958	89.3	1083
GdVS ₃	5.847	16.372	10.866	89.4	1040
LaCrS ₃	5.970	17.110	11.033	87.5	1126
NdCrS ₃	5.889	16.446	10.859	87.8	1051
GdCrS ₃	5.838	16.205	10.703	86.9	1011
LaMnS ₃	10.370	—	5.772	—	537.5
LaFeS ₃	10.314	—	5.806	—	534.9
LaCoS ₃	10.318	—	5.776	—	532.5
LaNiS ₃	10.295	—	5.762	—	528.9

3. Results and Discussion

3.1 Crystal Chemistry

X-Ray powder diffraction data show that the $LnMS_3$ compounds studied by us crystallize in two different structures (Table I). The vanadium and chromium sulfides (of La, Nd, and Gd) crystallize in a monoclinic structure isotypic with LaCrS₃ (2), wherein alternating layers of LaS and CrS₂ are present. Other $LnMS_3$ compounds with $Ln = La$ and $M = Mn, Fe, Co,$ or Ni are isostructural, forming in a hexagonal structure with six formula units per cell. Powder diffraction data of two typical compounds, LaMnS₃ and LaNiS₃, are given in Table II. The hexagonal cell appears to be in a way related to the $2H$ hexagonal structure of CsNiCl₃ (1). In order to determine the oxidation state of the transition metal and the nature of sulfur in these sulfides, we have measured the chemical shifts in the K -absorption edge of the transition metal as well as the valence- and core-level binding energies using XPS. In Table III, we have listed the chemical shifts of the transition metal K -absorption edges in $LaMS_3$ compounds. The corresponding values of binary transition metal sulfides are also shown in this table. We readily see that, in the case of LaCrS₃, the shift is close to that of Cr₂S₃ indicating that chromium in LaCrS₃ is in the trivalent state. In the other $LaMS_3$ sulfides, the chemical shifts of the transition metal are close to those of the monosulfides, indicating that the oxidation state of transition metal is likely to be 2+.

Further evidence for the oxidation state of M and the nature of sulfur is provided by XPS of $LaMS_3$ compounds. XPS binding energies of valence and core levels are listed in Table IV. For purposes of comparison, we have listed the binding energies of the corresponding ternary oxides (7) in the same table. We see that the binding energies of the transition metal levels are slightly lower in the sulfides than in the

TABLE II
X-RAY POWDER DIFFRACTION DATA OF LaMnS_3 AND LaNiS_3 ^a

<i>hkl</i>	LaMnS_3			LaNiS_3		
	<i>I</i> / <i>I</i> ₀	<i>d</i> _{obs} (Å)	<i>d</i> _{calc} (Å)	<i>I</i> / <i>I</i> ₀	<i>d</i> _{obs} (Å)	<i>d</i> _{calc} (Å)
101	25	4.87	4.86	30	4.84	4.84
200	25	4.51	4.49	25	4.46	4.46
111	50	3.870	3.857	60	3.839	3.839
201	100	3.555	3.544	100	3.527	3.526
210	40	3.401	3.394	45	3.373	3.370
211	10	2.929	2.926	15	2.912	2.909
002	20	2.885	2.886	20	2.876	2.881
301	35	2.659	2.657	20	2.638	2.641
220	30	2.595	2.592	35	2.574	2.574
112	25	2.522	2.522	30	2.514	2.514
310	50	2.493	2.491	65	2.475	2.473
202	40	2.428	2.428	40	2.417	2.420
311	65	2.288	2.287	60	2.272	2.272
212	40	2.201	2.199	30	2.189	2.190
{401	—	—	—	—	—	2.079
{302	—	—	—	8	2.078	2.069
312	30	1.883	1.886	30	1.873	1.876
411	50	1.853	1.856	30	1.842	1.843
311	20	1.656	1.656	15	1.645	1.644

^aUnit cell parameters are given in Table I.

oxides due to the higher covalency of the former compounds. The lanthanum levels in LaMS_3 are similar to those in LaMO_3 showing 4*d* and 3*d* peaks with characteristic satellites (7). A comparison of the core level binding energies of transition metals in the LnMS_3 compounds with the values

TABLE III
CHEMICAL SHIFTS OF TRANSITION METAL (*M*)
K-ABSORPTION EDGES IN LaMS_3

Compound	<i>K</i> -absorption edge energy (±0.5 eV)	Chemical shift ^a ΔE (eV)
LaCrS_3	5996.25	7.4 ^b (5.0)
LaMnS_3	6543.80	5.80 (5.0)
LaFeS_3	7118.49	7.0 (6.5)
LaCoS_3	7717.68	8.0 (7.4)
LaNiS_3	8337.10	5.9 (5.3)

^a ΔE values for the corresponding transition metal monosulfides are shown in parentheses.

^b ΔE of Cr_2S_3 is 7.6 eV.

in the corresponding transition metal monosulfides (8) indicates that the metal is trivalent in LaVS_3 and LaCrS_3 and divalent in the other LnMS_3 compounds. From the XPS data we also see that LaVS_3 and LaCrS_3 show single sharp peaks for the sulfur (3*s*) and (2*s*) levels, whereas the other LaMS_3 (*M* = Mn, Fe, Co, Ni) compounds show two S(3*s*) peaks and broad S(2*s*) peaks suggesting presence of two types of sulfur anions. We have similarly found presence of two sulfur anions, S^{2-} and S_2^{2-} , in the case of CuS (8). The results from XPS and *K*-absorption edge measurements are consistent with the following formulation of the LnMS_3 compounds: (i) $\text{Ln}^{3+}\text{M}^{3+}3\text{S}^{2-}$ in the case of *M* = V and Cr and (ii) $\text{La}^{3+}\text{M}^{2+}2\text{S}^{2-} \cdot 0.5(\text{S}_2^{2-})$ in the case of *M* = Mn, Fe, Co, and Ni.

3.2 Electron Transport Properties

Results of our measurements of electrical

TABLE IV
XPS BINDING ENERGIES (eV) OF LaMS_3^a

Compound	<i>M</i> levels		<i>S</i> levels		La levels	
LaVS ₃	3 <i>d</i>	0.1 (1.0)	3 <i>p</i>	4.5	5 <i>p</i>	17.8
	3 <i>p</i>	38.9 (40.7)	3 <i>s</i>	13.2	4 <i>d</i> _{5/2}	102.1 (102.7)
	2 <i>p</i> _{3/2}	514.6 (516.6)	2 <i>p</i>	161.3	3 <i>d</i> _{5/2}	836.5 (834.5)
	2 <i>p</i> _{1/2}	522.1 (524.1)	2 <i>s</i>	225.3		
LaCrS ₃	3 <i>d</i>	2.0 (2.4)	3 <i>p</i>	3.8, 5.0	5 <i>p</i>	17.6
	3 <i>p</i>	43.0 (43.2)	3 <i>s</i>	14.0	4 <i>d</i> _{5/2}	102.3 (101.1)
	2 <i>p</i> _{3/2}	575.8 (576.0)	2 <i>p</i>	161.2	3 <i>d</i> _{5/2}	836.4 (834.4)
	2 <i>p</i> _{1/2}	585.3 (586.1)	2 <i>s</i>	225.1		
LaMnS ₃	3 <i>d</i>	1.2 (1.0, 3.5)	3 <i>p</i>	3.5, 7.8	5 <i>p</i>	17.6
	3 <i>p</i>	— (48.3)	3 <i>s</i>	11.0–13.5	4 <i>d</i> _{5/2}	102.3 (101.5)
	2 <i>p</i> _{3/2}	639.2 (640.9)	2 <i>p</i>	161.4	3 <i>d</i> _{5/2}	836.7 (834.1)
	2 <i>p</i> _{1/2}	650.0 (652.9)	2 <i>s</i>	225.4		
LaFeS ₃	3 <i>d</i>	0.3, 3.1 (2.0, 4.0)	3 <i>p</i>	4.6	5 <i>p</i>	17.8
	3 <i>p</i>	— (55.2)	3 <i>s</i>	10.5, 13.0	4 <i>d</i> _{5/2}	102.6 (101.2)
	2 <i>p</i> _{3/2}	706.6 (710.5)	2 <i>p</i>	161.4	3 <i>d</i> _{5/2}	836.7 (833.7)
	2 <i>p</i> _{1/2}	719.9 (724.0)	2 <i>s</i>	225.4		
LaCoS ₃	3 <i>d</i>	0.2, 1.5 (—)	3 <i>p</i>	3.8, 7.7	5 <i>p</i>	17.8
	3 <i>p</i>	— (60.6)	3 <i>s</i>	10.5, 13.7	4 <i>d</i> _{5/2}	102.3 (101.5)
	2 <i>p</i> _{3/2}	778.3 (779.6)	2 <i>p</i>	161.3	3 <i>d</i> _{5/2}	836.7 (833.6)
	2 <i>p</i> _{1/2}	793.4 (795.2)	2 <i>s</i>	225.6		
LaNiS ₃	3 <i>d</i>	1.1, 2.8 (1.0, 2.8)	3 <i>p</i>	4.5, 7.4	5 <i>p</i>	17.7
	3 <i>p</i>	67.1 (67.1)	3 <i>s</i>	11.5, 14.0	4 <i>d</i> _{5/2}	102.2 (101.1)
	2 <i>p</i> _{3/2}	852.7 (—)	2 <i>p</i>	161.1	3 <i>d</i> _{5/2}	836.8 (833.8)
	2 <i>p</i> _{1/2}	869.8 (—)	2 <i>s</i>	225.5		

^a Values in parentheses are those found in the corresponding oxides, LaMO_3 (from Ref. (7)).

properties of LnMS_3 compounds are summarized in Table V. Among the isostructural LnVS_3 and LnCrS_3 compounds possessing layer structures, the vanadium compounds show low resistivities and small positive Seebeck coefficients. While LaVS_3 clearly shows metal-like temperature dependence of resistivity, NdVS_3 and GdVS_3 do not show this behavior (Fig. 1). Electron transport in LnVS_3 may arise from holes in the $\text{S}(3p)$ band created by the partial overlap of $\text{V}(3d)$ states with the top of $\text{S}(3p)$ bonding states; this model is similar to the type (e) band model proposed by Jellinek (9) for transition metal chalcogenides. In contrast to the behaviour of LnVS_3 , the isostructural LnCrS_3 compounds are insula-

tors with resistivities of the order of 10^6 ohm · cm at 300 K. The activation energy, E_a , for electrical conduction varies in the order $\text{Gd} > \text{Nd} > \text{La}$ just as in the oxides (6). The large negative Seebeck coefficient of LnCrS_3 at 300 K suggests the presence of a small sulfur deficiency which can be shown as $\text{LnCr}_{1-x}^{3+}\text{Cr}_x^{2+}\text{S}_{3-x/2}$. The sharp change in the electron transport properties of LnMS_3 compounds on passing from $M = \text{V}$ to $M = \text{Cr}$ is reminiscent of the behavior in LaMO_3 perovskites (1), where LaTiO_3 is a metal, LaVO_3 is a semiconductor with a small activation energy, and LaCrO_3 is an insulator. In LnMO_3 the changes have been explained by Goodenough (10) on the basis of the varying M -O covalent interaction

TABLE V
ELECTRICAL AND MAGNETIC PROPERTIES OF $LnMS_3$

Compound	μ_{eff}^a (μ_B)		Θ^b (K)	Resistivity at 300 K ρ (ohm · cm)	Seebeck coefficient at 300 K, α ($\mu\text{V}/\text{deg}$)	Activation energy E_a (eV)
	Expt.	Calcd.				
	Pauli para-					
LaVS ₃	magnetic	2.83	—	3×10^{-3}	+9	Metallic
NdVS ₃	4.24	4.59	-180	5×10^{-3}	+10	—
GdVS ₃	8.37	8.43	-50	10×10^{-3}	+10	—
LaCrS ₃	3.70	3.87	-350	5×10^6	-800	0.48
NdCrS ₃	5.00 ^c	5.30	-140 ^c	6×10^6	-300	0.50
GdCrS ₃	9.40	8.83	-50	7×10^6	-200	0.56
LaMnS ₃	5.50	5.92	-300	1.0	+22	0.05
LaFeS ₃	4.00	4.90	+450	0.3	+19	0.08
LaCoS ₃	2.32	1.73	-110	4.0×10^{-2}	-10	Metallic
LaNiS ₃	1.35	2.83	25	1.0×10^{-2}	-11	Metallic

^a μ_{eff} for the transition metal, M .

^b Paramagnetic Curie temperature.

^c μ_{eff} from Cr³⁺ contribution alone is 3.55 in agreement with the spin-only value of Cr³⁺; the corresponding Θ is -300 K.

across the transition metal series. It is also to be noted that in the isostructural $NaMS_2$

series, $NaTiS_2$ is a metal, $NaVS_2$ a good conductor, and $NaCrS_2$ an insulator (11).

Among the $LaMS_3$ sulfides with $M = Mn, Fe, Co,$ or Ni , the cobalt and nickel compounds are definitely metallic with small negative Seebeck coefficients (Fig. 1). The manganese and iron compounds, on the other hand, are p-type semiconductors with small activation energies (Fig. 1). It is interesting to compare the electronic properties of pyrite disulfides of $Mn, Fe, Co,$ and Ni (12) with those of the $LaMS_3$ compounds. MnS_2 shows localized electron behavior similar to $LaMnS_3$, which is characteristic of octahedral-site high-spin $Mn^{2+}:3d^5$. In the other pyrites, covalent interaction between the metal and sulfur is stronger than in the manganese compound creating a narrow σ^* band from the metal e_g orbitals. Electronic properties of $FeS_2, CoS_2,$ and NiS_2 have been explained in terms of electron correlation and occupancy of the σ^* band. Although $LaFeS_3$ is a semiconductor similar to FeS_2 , the magnetic properties are different (as can be

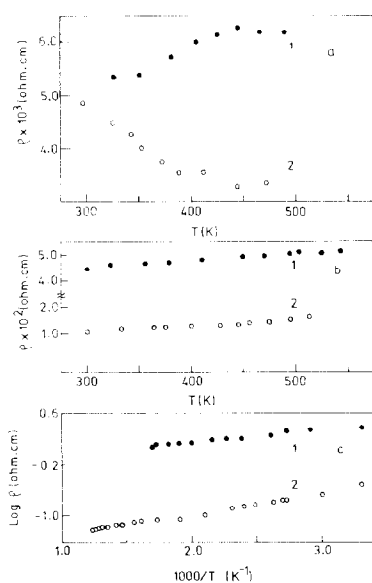


FIG. 1. Temperature dependence of electrical resistivity of $LnMS_3$ sulfides: (a) 1, LaVS₃ and 2, NdVS₃; (b) 1, LaCoS₃ and 2, LaNiS₃; (c) 1, LaMnS₃ and 2, LaFeS₃.

seen in the next section); Fe^{2+} is in the high-spin state in LaFeS_3 unlike in FeS_2 . The result seems to indicate that covalent interaction between iron and sulfur in LaFeS_3 is not strong enough to give rise to the low-spin state of Fe^{2+} . A definite interpretation of electronic properties of this isostructural LaMS_3 series must await determination of their detailed crystal structures.

3.3 Magnetic Properties

Magnetic properties of LnMS_3 compounds are summarized in Table V. LaVS_3 (Fig. 2) has a small molar magnetic susceptibility, χ_m ($\sim 350 \times 10^{-6}$ cgs units) which is nearly independent of temperature down to 200 K; at low temperatures, the susceptibility develops a "Curie tail," similar to that found in VO_x (13) and VS_x (14). In order to see whether the magnetic properties of NdVS_3 and GdVS_3 are similar to LaVS_3 , we subtracted the rare earth ion contribution

from the measured susceptibility and plotted the reciprocal of the difference susceptibility, $\Delta\chi_m^{-1}$, against temperature (Fig. 2). We find the behavior to be similar to LaVS_3 except that $\Delta\chi_m$ is higher ($\sim 820 \times 10^{-6}$ cgs units for NdVS_3) indicating a greater degree of localization of V(3d) electrons in NdVS_3 and GdVS_3 than in LaVS_3 . This behavior is consistent with electrical properties of these vanadium sulfides.

LnCrS_3 compounds show Curie-Weiss paramagnetism at high temperatures in all the three cases studied (Fig. 3). The effective magnetic moments, μ_{eff} , obtained from the linear region of the $\chi_m^{-1}-T$ plots are in fair agreement with the values calculated assuming spin-only contribution for Cr^{3+} and spin and orbital contribution for Ln^{3+} (Table V). The paramagnetic Curie temperatures, Θ , are negative, the absolute values varying in the order $\text{La} > \text{Nd} > \text{Gd}$. The $\chi_m^{-1}-T$ plots do not, however, show a Néel temperature; instead, the plots deviate from the Curie-Weiss law at low tempera-

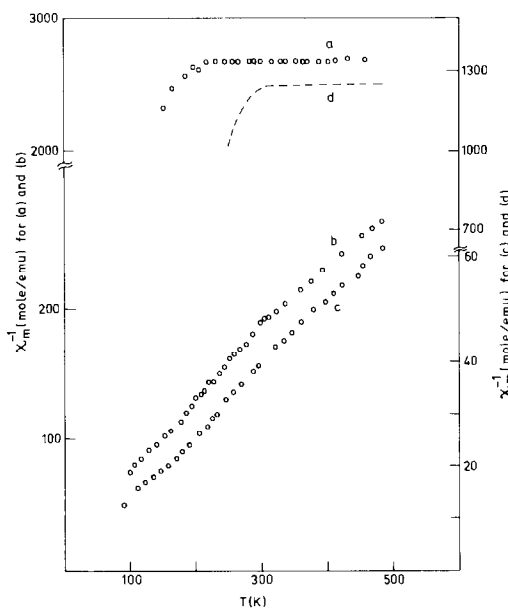


FIG. 2. $\chi_m^{-1}-T$ plots of (a) LaVS_3 , (b) NdVS_3 , and (c) GdVS_3 ; plot of vanadium contribution in the case of NdVS_3 is shown by the broken curve (d).

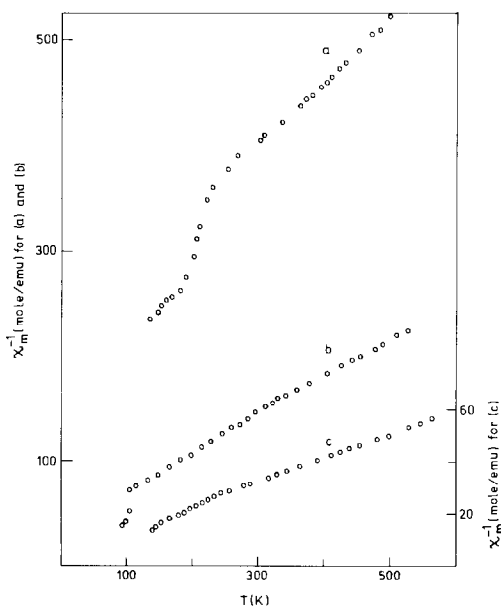


FIG. 3. $\chi_m^{-1}-T$ plots of (a) LaCrS_3 , (b) NdCrS_3 , and (c) GdCrS_3 .

tures towards larger susceptibility values. In the LaCrS_3 structure consisting of alternating rocksalt-like LaS and CdI_2 -like CrS_2 layers (3), two types of magnetic interaction between chromium ions are possible: antiferromagnetic Cr–Cr interaction and a ferromagnetic 90° Cr–S–Cr interaction between the edge-shared CrS_6 octahedra in the CrS_2 layers. While the negative Weiss constant suggests antiferromagnetic interaction, the deviation from the Curie–Weiss law at low temperatures may arise from a short-range ferromagnetic interaction within the CrS_2 layers. Similar weak ferromagnetism is shown by LnCrSe_3 crystallizing in the $(\text{NH}_4)\text{CdCl}_3$ structure (15).

The χ_m^{-1} - T plots (Fig. 4) of LaMnS_3 , LaCoS_3 , and LaNiS_3 show Curie–Weiss law behavior with negative Weiss constants. A μ_{eff} value of $5.5 \mu_B$ found for the manganese compound is consistent with the $2+$ oxidation state of manganese ($3d^5$). The μ_{eff} value ($2.32 \mu_B$) found in LaCoS_3 is slightly higher than the value expected for

octahedral-site low-spin Co^{2+} ion. The CoS_2 – CoSe_2 system where divalent cobalt is in the low-spin state, also has a similar μ_{eff} (16). The susceptibility of LaNiS_3 at room temperature is quite small ($\sim 690 \times 10^{-6}$ cgs units) and shows a weak temperature dependence, the behavior being consistent with its metallic nature. The magnetic property of LaNiS_3 is roughly comparable to that of NiSe_2 (17).

LaFeS_3 shows a sharp magnetic transition around 620 K (Fig. 4). Above the transition temperature, the susceptibility seems to follow Curie–Weiss law behavior with μ_{eff} of about $4 \mu_B$ which is close to that expected from high-spin Fe^{2+} ions. Because of thermal hysteresis it was not possible to obtain an exact value of μ_{eff} . At the transition temperature, the susceptibility increases by more than an order of magnitude and thereafter remains almost constant down to 150 K. Similar transitions to weak ferromagnetic state occur in $\alpha\text{-Fe}_2\text{O}_3$ and LaFeO_3 (18). We are unable to ascribe the observed behavior of LaFeS_3 to oxide impurities since LaFeS_3 prepared by us was monophasic with a distinct crystal structure and exhibited very low resistivity and a relatively small Seebeck coefficient unlike Fe_2O_3 or LaFeO_3 .

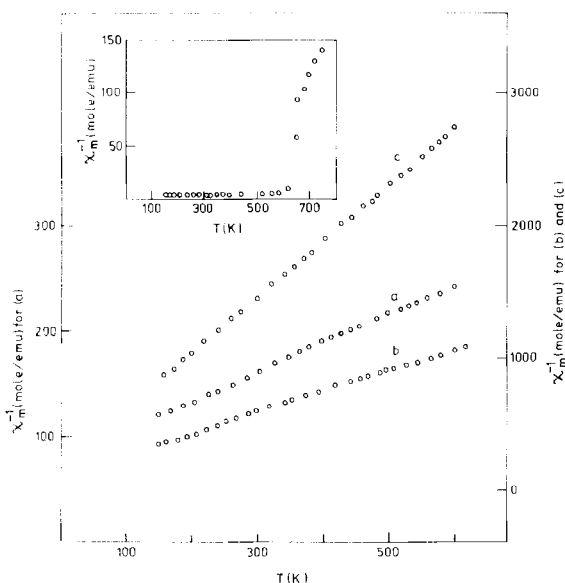


FIG. 4. χ_m^{-1} - T plots of (a) LaMnS_3 , (b) LaCoS_3 , and (c) LaNiS_3 . The χ_m^{-1} - T plot of LaFeS_3 is shown in the inset.

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References

- (a) J. B. GOODENOUGH AND J. M. LONGO, "Crystallographic and Magnetic Properties of Perovskite and Perovskite-Related Compounds," Landolt-Börnstein Tabellen, New Series, III/4a, p. 126, Springer-Verlag, Berlin (1970). (b) S. NOMURA, Landolt-Börnstein Tabellen, III/12a, p. 368, Springer-Verlag, Berlin (1978).
- T. TAKAHASHI, T. OKA, O. YAMADA, AND K. AMETANI, *Mater. Res. Bull.* **6**, 173 (1971).

3. K. KATO, I. KAWADA, AND T. TAKAHASHI, *Acta Crystallogr. Sect. B* **33**, 3437 (1977).
4. T. MURUGESAN, P. R. SARODE, J. GOPALAKRISHNAN, AND C. N. R. RAO, *J. Chem. Soc. Dalton*, 837 (1980).
5. C. N. R. RAO, D. D. SARMA, S. VASUDEVAN, AND M. S. HEGDE, *Proc. Roy. Soc. London A* **367**, 239 (1979).
6. P. GANGULY, OM PARKASH, AND C. N. R. RAO, *Phys. Status Solidi A* **36**, 669 (1976).
7. W. H. MADHUSUDAN, S. KOLLALI, P. R. SARODE, M. S. HEGDE, P. GANGULY, AND C. N. R. RAO, *Pramana* **12**, 317 (1979).
8. J. GOPALAKRISHNAN, T. MURUGESAN, M. S. HEGDE, AND C. N. R. RAO, *J. Phys. C*, **12**, 5255 (1979).
9. F. JELLINEK, "MTP International Review of Science," Series One, Inorganic Chemistry, Vol. 5, p. 343, University Park Press, Baltimore (1972).
10. J. B. GOODENOUGH, "Progress in Solid State Chemistry" (H. Reiss, Ed.), Vol. 5, Pergamon, New York (1971).
11. W. BRONGER, "Crystallography and Crystal Chemistry of Materials with Layered Structures" (F. Levy, Ed.), p. 106, Reidel, Dordrecht (1976).
12. J. B. GOODENOUGH, "Solid State Chemistry" (C. N. R. Rao, Ed.), Dekker, New York (1974).
13. M. D. BANUS AND T. B. REED, "The Chemistry of Extended Defects in Nonmetallic Solids" (L. Eyring and M. O'Keeffe, Eds.), p. 488, North-Holland, Amsterdam (1970).
14. A. B. DE VRIES AND C. HAAS, *J. Phys. Chem. Solids* **34**, 651 (1973).
15. O. GOROCHOV AND H. MCKINZE, *J. Solid State Chem.* **7**, 400 (1973).
16. K. ADACHI, M. MATSUI, AND M. KAWAI, *J. Phys. Soc. Japan* **46**, 1474 (1979).
17. F. GRONVOLD AND E. JACOBSON, *Acta Chem. Scand.* **10**, 1440 (1956).
18. H. WATANABE, *J. Phys. Soc. Japan* **14**, 511 (1959).