



COMMENTS ON “NATURAL FREQUENCIES OF RECTANGULAR PLATES
USING A SET OF STATIC BEAM FUNCTIONS IN RAYLEIGH–RITZ METHOD”

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The author is to be congratulated for developing co-ordinate functions which are combinations of sinusoidal terms and polynomials. The coefficients of the polynomial expressions are determined by the boundary conditions of the plate under analysis. Excellent accuracy is achieved [1].

It is quite surprising to the writer that the author does not refer at all to an alternative approach developed extensively by the writer and coworkers, which consists of generating approximations to the displacement function of the plate using polynomials which satisfy, at least, the essential boundary conditions.

Admittedly, the analysis has been restricted in many instances to the lower modes of vibrations, and, in a great majority of the cases, to the fundamental mode of vibration [2–18]. Among the situations treated, one may mention edges elastically restrained against rotation and in-plane stresses [3, 4], concentrated and elastically mounted masses [5, 6], non-uniform thickness [7], orthotropic materials [8], etc.

The approach has been used in the case of forced vibration problems [9, 10] and extended to circular plates [11, 12] and plates of complicated boundary shape using the conformal mapping method [13].

Furthermore, polynomial approximations have been employed in the case of plates with inner supports [14] and the technique has also been used in the case of portal frames [15], arches [16] and rings [17].

As a matter of fact, it was shown in reference [2] that using a simple one-term polynomial approximation one obtains $\lambda_1 = 9.00$ for the fundamental eigenvalue of a clamped square plate while the author, using $p, q = 2, 2$ (Table 1 of reference [1]) obtains $\lambda_1 = 9.012$, the “exact” result being $\lambda_1 = 8.996$ [1]. The methodology has also been applied in the case of vibrating plates with openings having free edges [18].

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AUTHOR'S REPLY

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I would like to thank Professor Laura for his calling attention to additional literature [1] pertinent to the title problem. I am sorry that I did not pay much attention to Professor Laura's papers, partly owing to the lack of opportunities to obtain the total literature in the field when preparing my paper [2]. Besides, readers may find that in most cases, Laura's papers were mainly aimed at the calculation of the fundamental mode of vibration of

plates, beams, etc., while my paper focuses on higher modes in addition to the fundamental modes of vibration of rectangular plates. My method has been extended to vibration of elastically restrained rectangular plates with eight different rotational and translational flexural coefficients [3] and line supported rectangular plates in one and two ways [4]. Good accuracy was also achieved.

In reply to his last comment, the point to note is that $\lambda_1 = 9.012$ is actually obtained by using $p, q = 1, 1$ (Table 1 of reference [2]) in my paper. Apparently the structures are symmetric about the central axes for the clamped rectangular plates and the basis functions $Y_i(\zeta)$ in my paper are symmetric for $i = 1, 3, 5, \dots$ and antisymmetric for $i = 2, 4, 6, \dots$. In such cases, the accuracy of taking $p, q = 1, 1$ is the same as that of $p, q = 2, 2$ for λ_1 ; the antisymmetric function $Y_2(\zeta)$ does not affect the symmetric modes of vibration. If one wants, one may take advantage of the symmetry about the central axes of the plates by taking i and j to be 1, 3, 5, ... for full-symmetric, i, j to be 2, 4, 6, ... for full-antisymmetric, i to be 1, 3, 5, ... and j to be 2, 4, 6, ... or i to be 2, 4, 6, ... and j to be 1, 3, 5, ... for symmetric-antisymmetric or antisymmetric-symmetric modes of vibration in the eigenvalue equation of clamped rectangular plates; furthermore, the computational cost will be reduced.

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