



A SIMPLE DERIVATION OF THE ACOUSTIC BOUNDARY  
CONDITION IN THE PRESENCE OF FLOW

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(Received 21 December 1998)

In modelling the noise propagation in a duct of a ducted fan engine, the authors needed the boundary condition satisfied by the acoustic pressure on a lined duct wall in the presence of flow. The result in the form used by the authors was derived by M. K. Myers [1] who presented a formal derivation of a related result by K. Taylor [2]. A good review of the subject can be found in reference [3]. Here, a brief and simple derivation of the acoustic boundary condition of Myers is given. The main difference in the derivation from that in reference [1] is that the Gaussian co-ordinate system  $(q^1, q^2, q^3(q^1, q^2, t))$  is used to specify the instantaneous position, i.e., Lagrangian variable, of the point on the mean position of the wall with curvilinear (Gaussian) co-ordinates  $(q^1, q^2)$ . Myers uses a locally orthogonal coordinate system which is somewhat less specific than what is used here.

One assumes that one has a base or background flow with velocity  $\vec{u}_0(\vec{x})$  which is time independent perturbed by a small velocity distribution  $\varepsilon\vec{u}_1(\vec{x}, t)$  where  $0 < \varepsilon \ll 1$ . One also assumes that the wall boundary's mean position  $S_0$  is independent of time and specified by the position vector  $\vec{x}_0(q^1, q^2)$  where  $(q^1, q^2)$  are the curvilinear (Gaussian) co-ordinates on the boundary surface. The position of the time dependent boundary  $S$  is given by

$$\vec{x}(q^1, q^2, q^3, t) = \vec{x}_0(q^1, q^2) + \varepsilon q^3(q^1, q^2, t)\vec{n}_0(q^1, q^2), \quad (1)$$

where  $\vec{n}_0$  is the local unit normal to  $S_0$  and  $\varepsilon q^3$  is the distance along the normal  $\vec{n}_0$  from  $S_0$  to  $S$  at  $(q^1, q^2, t)$ . The fundamental physical requirement at the boundary is

$$(\vec{u}_0 + \varepsilon\vec{u}_1) \cdot \vec{n}_0 = \varepsilon \partial q^3 / \partial t. \quad (2)$$

This is the instantaneous equality of the normal fluid velocity and the surface velocity. Note that the symbol  $\varepsilon$  will be retained for order of magnitude comparison for now.

The left side of equation (2) will now be expanded as follows and equated to the right side:

$$\begin{aligned} & [\vec{u}_0(\vec{x}_0 + \varepsilon q^3 \vec{n}_0) + \varepsilon \vec{u}_1(\vec{x}_0 + \varepsilon q^3 \vec{n}_0)] \cdot \vec{n}_0 \\ &= \vec{u}_0(\vec{x}_0) \cdot \vec{n}_0 + \varepsilon [q^3 \vec{n}_0 \cdot \nabla \vec{u}_0(\vec{x}_0) + \vec{u}_1(\vec{x}_0)] \cdot \vec{n}_0 = \varepsilon \partial q^3 / \partial t. \end{aligned} \quad (3)$$

The equality of the zeroth and first order terms from both sides gives

$$\vec{u}_0(\vec{x}_0) \cdot \vec{n}_0 = 0 \quad (4)$$

and

$$q^3 \vec{n}_0 \cdot [\vec{n}_0 \cdot \nabla \vec{u}_0(\vec{x}_0)] + \vec{u}_1(\vec{x}_0) \cdot \vec{n}_0 = \partial q^3 / \partial t. \quad (5)$$

Equation (4) tells us that the normal velocity based on the mean flow is zero on the mean surface. Equation (5) can be written as

$$\vec{u}_1(\vec{x}_0) \cdot \vec{n}_0 = \partial q^3 / \partial t - q^3 \vec{n}_0 \cdot [\vec{n}_0 \cdot \nabla \vec{u}_0(\vec{x}_0)]. \quad (6)$$

One now assumes  $\varepsilon = 1$  and thus  $|\vec{u}_1| \ll |\vec{u}_0|$  and  $q^3$  is the local normal distance between  $S_0$  and  $S$  as a function of time. Note that  $\partial q^3 / \partial t$  is the local normal velocity of  $S$  in terms of the Lagrangian variables  $(q^1, q^2)$ . Using Eulerian variables  $\vec{x}_0$ , one can define a new function  $g(\vec{x}_0, t)$  such that  $q^3 = g(\vec{x}_0, t)$ . Note that since  $q^3$  is the normal distance between  $S_0$  and  $S$ , one has  $|\nabla g| = 1$ . One notes that

$$(\partial q^3 / \partial t)(q^1, q^2) = \partial g / \partial t + \vec{u}_0(\vec{x}_0) \cdot \nabla g. \quad (7)$$

Using this result in equation (6) gives equation (11) of Myers [1]:

$$\vec{u}_1(\vec{x}_0) \cdot \vec{n}_0 = \partial g / \partial t + \vec{u}_0(\vec{x}_0) \cdot \nabla g - g \vec{n}_0 \cdot [\vec{n}_0 \cdot \nabla \vec{u}_0(\vec{x}_0)], \quad (8)$$

which is the condition that the perturbation velocity  $\vec{u}_1$  must satisfy on the mean surface  $S_0$ .

One now derives the liner boundary condition based on equation (8). For a time harmonic disturbance proportional to  $e^{i\omega t}$ , the complex acoustic pressure  $p$  and  $g$  are related to each other by the following relation on  $S_0$ :

$$g = -p/i\omega Z, \quad (9)$$

where  $Z$  is the complex normal impedance. Using this result in equation (8) gives

$$\vec{u}_1 \cdot \vec{n}_0 = -p/Z - (1/i\omega) \vec{u}_0 \cdot \nabla(p/Z) + (p/i\omega Z) \vec{n}_0 \cdot (\vec{n}_0 \cdot \nabla \vec{u}_0), \quad (10)$$

which is the liner boundary condition, equation (15), in Myers [1]. Equation (10) is implemented in a ducted fan noise prediction code developed for NASA Langley Research Center by the authors [4].

#### REFERENCES

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