

## A New Description of Equilateral Five-membered Rings during Pseudorotation

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An equilateral, five-membered ring is described based on the assumption that during conformational changes the distances between atoms and the position of the centre of mass of the ring remain constant. In the course of pseudorotation the atoms move along three-dimensional curves whose projections on a mean plane of the ring are circles. Assuming a certain bond length of the initial regular pentagon, the model allows for high accuracy calculations of the geometry of any future conformations of the ring. Our model is compared with those of Kilpatrick and Adams.

Five-membered rings are very common in nature, especially furanose rings in nucleic acids or the D ring in steroids. The conformational flexibility of such rings plays an important role in the structure of entire molecules and consequently in their biological activity.

Models for pseudorotation of five-membered rings may be divided into three groups. The first is that of Kilpatrick *et al.*,<sup>1</sup> also used by other authors (Cremer and Pople<sup>2</sup>), the second is that formulated by Altona *et al.*,<sup>3-6</sup> and the third is that postulated by Adams *et al.*,<sup>7</sup> hereafter referred to as AGB. Each model has its disadvantages. In Kilpatrick's model, the endocyclic bond lengths do not remain constant during pseudorotation (Figure 1). Altona has contributed a description rather than a model of pseudorotation; it contains merely mathematical relationships between the valency or torsional angles and the pseudorotation parameters  $P$  and  $\tau_m$ , which are adequate only for small values of the puckering amplitude  $q$ ; no attention is given to the nature of the phenomenon. The AGB model, though lacking the disadvantages of Kilpatrick's model, is based on strict requirements limiting its application to small out-of-plane displacements of the atoms (see Discussion section) and does not fulfill Eckart's conditions.

**Definition of the Model.**—The approach to pseudorotation presented in this work is based on the following assumptions: (a) the  $z_j$  co-ordinate of atoms in the five-membered ring satisfies Kilpatrick's<sup>1</sup> relation (1) where  $q$  is the puckering amplitude

$$z_j = \sqrt{2/5}q \cos(\varphi + 4\pi j/5) \quad (1)$$

and  $\varphi$  the phase angle and (b) the  $x_j$  and  $y_j$  co-ordinates depend on the pseudorotation parameters of equations (2) and (3) and

$$x_j = F_j(q, \varphi) \quad (2)$$

$$y_j = G_j(q, \varphi) \quad (3)$$

functions  $F$  and  $G$  are periodic with respect to  $\varphi$ .

Thus we propose a parametric dependence on  $\varphi$  of the planar curve describing the motion of the projection of an atom on the plane of the unpuckered ring. In Kilpatrick's model this curve is a point; in the AGB model an interval, while in this paper it is postulated to be an ellipse. This ellipse is easily described by an equation in a local reference system of the atom concerned. The local reference system is determined by the two characteristic directions of the unpuckered five-membered ring, *i.e.* the direction defined by the geometrical centre of the planar ring and the atom in question, and the direction perpendicular to it. The first direction defines the  $p_j$

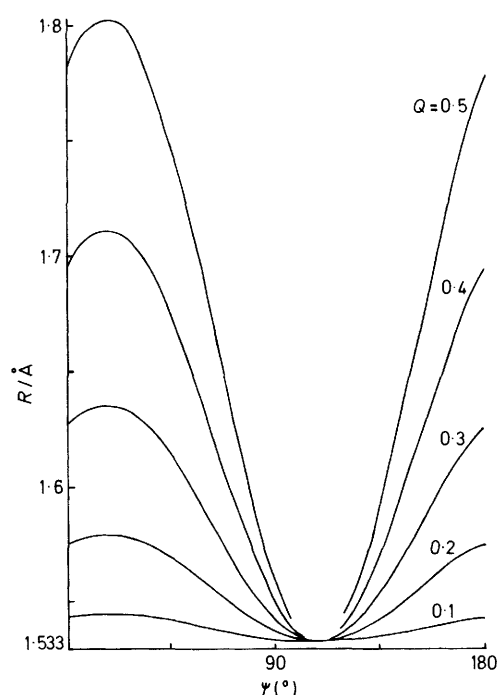


Figure 1. Distance between first and fifth atom of an equilateral ring of bond length 1.533 Å versus phase angle  $\varphi$  for several values of  $Q = \sqrt{2/5}q$ , calculated according to Kilpatrick *et al.*<sup>1</sup>

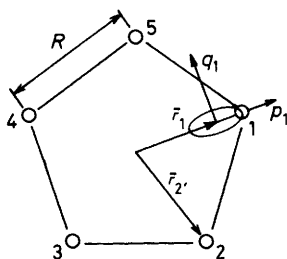
axis of the new system, the second the  $q_j$  axis (Figure 2). The planar ring is described by the points of intersection of the trajectories of the atoms with the mean plane. In our model the geometry of this planar ring is required for establishing these trajectories. The model is limited to an equilateral five-membered ring. Hence, we assume that the planar ring is a regular pentagon and is described unequivocally by its bond length  $R$ . The origin of the system of the  $j$ th atom is given by

$$\vec{r}_j = \vec{r}_j' - \vec{a}_j \quad (4)$$

the vector in equation (4) where  $\vec{r}_j'$  is the radius vector of the  $j$ th atom in the planar ring, and  $\vec{a}_j = (\vec{r}_j'/r_j)a$ , where  $a$  is the major axis of the ellipse. The ellipse described by the projection of the  $j$ th atom during its motion is given by equations (5)

$$p_j = a \cos[t(\varphi)] \quad (5)$$

$$q_j = b \cos[t(\varphi)] \quad (6)$$



**Figure 2.** Elliptical trajectories of the atomic projections on the  $xy$  mean plane and local systems of co-ordinates  $(p_j, q_j)$ . Small circles show the positions of the atoms in the planar ring.  $R$  is the bond length

and (6) where  $b$  is the minor axis of the ellipse. Constants  $a$  and  $b$  must be independent of  $\varphi$  and dependent on  $q$ . They also must be different for different starting geometries of the planar ring, defined by the bond length  $R$ .

The function  $t(\varphi)$  was found on the basis of assumptions (7) and (8). It is given by equation (9). Hence substituting

$$\text{if } |z_j| = \max = \sqrt{2/5}q \text{ then } p_j = -a \quad (7)$$

$$\text{if } |z_j| = \min = 0 \quad \text{then } p_j = +a \quad (8)$$

$$t(\varphi) = 2(\varphi + 4\pi j/5) + \pi \quad (9)$$

equation (9) in equations (5) and (6), one can obtain a three-dimensional parametric curve described by the atom during pseudorotation in the local reference system  $(p_j, q_j, z_j)$  [equations (10)].

$$p_j = -a \cos 2(\varphi + 4\pi j/5) \quad (10a)$$

$$q_j = -b \sin 2(\varphi + 4\pi j/5) \quad (10b)$$

$$z_j = \sqrt{2/5}q \cos(\varphi + 4\pi j/5) \quad (10c)$$

Using the geometrical centre of the planar ring as the reference for the co-ordinate system, equations (10) became (11) where

$$x(j) = -a \cos(2\varphi_j) \cos \alpha_j + b \sin(2\varphi_j) \sin \alpha_j + r_j^x \quad (11a)$$

$$y(j) = -a \cos(2\varphi_j) \sin \alpha_j - b \sin(2\varphi_j) \cos \alpha_j + r_j^y \quad (11b)$$

$$z(j) = Q \cos \varphi_j \quad (11c)$$

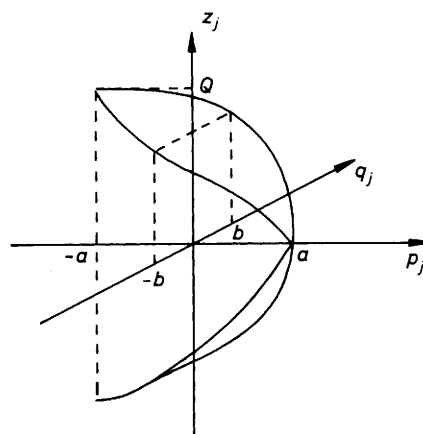
$\varphi_j = \varphi + 4\pi j/5$ ,  $Q = \sqrt{2/5}q$ , and is the maximum possible deviation of atoms from the mean plane of the ring,  $\alpha_j$  is the angle between vector  $\vec{r}_j$  and the  $x$ -axis, and  $r_j^x$  and  $r_j^y$  are the  $x$  and  $y$  co-ordinates of the origin of the  $(p_j, q_j, z_j)$  system. The values of  $r_j^x$  and  $r_j^y$  are given by equations (12).

$$r_j^x = \left[ \frac{R}{2\sin(\pi/5)} - a \right] \cos \alpha_j \quad (12a)$$

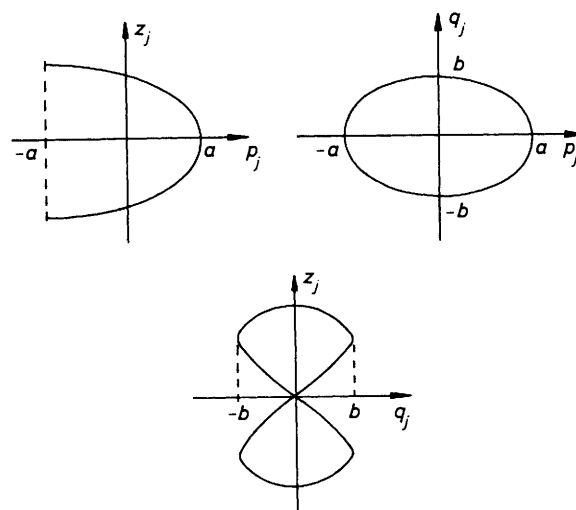
$$r_j^y = \left[ \frac{R}{2\sin(\pi/5)} - a \right] \sin \alpha_j \quad (12b)$$

Curve (10) is shown in Figure 3 and its projections on the planes of the  $(p_j, q_j, z_j)$  reference system are shown in Figure 4.

The description of curve (11) is incomplete without knowledge of the values of constants  $a$  and  $b$ . These values were calculated assuming that the bond lengths in the ring are constant on the pseudorotation path, that is by solving



**Figure 3.** Trajectory of the  $j$ th atom given by equation (10) in the local reference system  $(p_j, q_j, z_j)$



**Figure 4.** Projections of the trajectory of the  $j$ th atom on: (a) the  $p_j, z_j$  plane; (b) the  $p_j, q_j$  plane; (c) the  $q_j, z_j$  plane, in the  $(p_j, q_j, z_j)$  reference system

$$[x(j+1) - x(j)]^2 + [y(j+1) - y(j)]^2 + [z(j+1) - z(j)]^2 = R^2 \quad (13)$$

equation (13) with respect to  $a$  and  $b$ , provided that bond length  $R$  is independent of  $\varphi$  and  $q$ .  $x(j), y(j)$ , and  $z(j)$  were substituted from equations (11).

Calculated values of  $a$  and  $b$  for various  $Q$  and  $R$  values are given in Table 1;  $a$  and  $b$  are unequivocally defined for all values of  $Q$  and  $R$ . Thus there exists an unequivocal relationship between  $a$  and  $b$ . It was found to be very simple [equation (14)]. It was also found that equation (15) holds. The depen-

$$b = 0.2361a \quad (14)$$

$$a/R = f(Q/R) \quad (15)$$

dence of  $a$  on  $Q$  and  $R$  was found using the least-squares method and is in the form of a polynomial equation (16). The

$$a = AR - BQ + CQ^2/R \quad (16)$$

**Table 1.** Constants  $a$  and  $b$  calculated from equation (13) for various values of  $R$  and  $Q$ 

$R/\text{nm}$	$Q/\text{nm}$				
	0.01	0.02	0.03	0.04	0.05
0.140	$10^{-5} a =$	224	517	956	1 585
	$55$				
0.145	$10^{-5} b =$	52	122	225	374
	$13$				
0.150		53	216	498	917
		13	51	117	216
0.155		52	209	480	882
		12	49	113	208
0.160		50	202	463	849
		12	48	109	200
		48	195	448	819
		11	46	106	193

$a$  and  $b$  in nm.

**Table 2.** Dependence of major axis of the ellipse on  $Q$  and  $R$  values found using the least-squares method, according to the equation  $a = AR - QB + CQ^2/R$ 

$Q/R$ range		$A$	$B$	$C$	Accuracy of fit (nm)
from	to				
0.0	0.1	0.0	0.0015	0.7937	$10^{-5}$
0.1	0.2	0.0010	0.0195	0.8708	$10^{-5}$
0.2	0.3	0.0115	0.1224	1.1239	$10^{-5}$
0.3	0.4	0.0820	0.5800	1.8693	$2 \times 10^{-5}$
0.4	0.5	0.9086	4.5205	6.5832	$2 \times 10^{-4}$

constants  $A$ – $C$  are collected in Table 2. Substituting co-ordinates (11) and parameters  $a$  and  $b$  [equations (14), (16)] in equation (13), we found that on the pseudorotation path the bond lengths are constant with an accuracy of  $10^{-5}$  nm.

The model described by equation (11) does not fulfill Eckart's conditions. During pseudorotation the centre of mass of the ring rotates in the  $xy$  plane about the  $z$  axis in the direction opposite to the direction of the motion of atomic projections on the  $xy$  plane. The motion discussed is an artifact resulting from the fact that the trajectory of an atom intersects the mean plane at the apex of the planar pentagon [equation (8)]. Hence, when the atom crosses the mean plane it 'pulls' the whole ring towards the point of intersection. The motion of the centre of mass is eliminated by adding to equations (11a and b) the corrections  $x_c$  and  $y_c$ , respectively [equations (17) where

$$x_c = R_{\text{cm}} \sin 2\varphi \quad (17a)$$

$$y_c = R_{\text{cm}} \cos 2\varphi \quad (17b)$$

$R_{\text{cm}}$  is the radius of rotation of the centre of mass]. The radius  $R_{\text{cm}}$  was found to depend on the constants  $a$  and  $b$  [equation (18)]. Hence, substituting equation (14) into (18) one obtains

$$R_{\text{cm}} = (a-b)/2 \quad (18)$$

$$R_{\text{cm}} = 0.3820a \quad (19)$$

(19). When correction (17) is taken into consideration, the trajectory of the  $j$ th atom is given by equations (20).

$$x(j) = -a \cos 2(\varphi + 4\pi j/5) \cos \alpha_j + b \sin 2(\varphi + 4\pi j/5) \sin \alpha_j + r_j^x + \left(\frac{a-b}{2}\right) \sin 2\varphi \quad (20a)$$

**Table 3.** Dependence of the radius of the circle  $r_c$  on  $Q$  and  $R$  values found using the least-squares method, according to the equation  $r_c = AR - BQ + CQ^2/R$ 

$Q/R$ range		$A$	$B$	$C$	Accuracy of fit (nm)
from	to				
0.0	0.1	0.0	0.0	0.485	$10^{-5}$
0.1	0.2	0.0008	0.0145	0.5479	$10^{-5}$
0.2	0.3	0.0082	0.0843	0.7116	$10^{-5}$
0.3	0.4	0.0534	0.3741	1.1773	$2 \times 10^{-5}$

$$y(j) = -a \cos 2(\varphi + 4\pi j/5) \sin \alpha_j - b \sin 2(\varphi + 4\pi j/5) \cos \alpha_j + r_j^y + \left(\frac{a-b}{2}\right) \cos 2\varphi \quad (20b)$$

$$z(j) = \sqrt{2/5} q \cos(\varphi + 4\pi j/5) \quad (20c)$$

One can thus unequivocally obtain the co-ordinates of the  $j$ th atom in the ring for all given values of parameters  $\varphi$ ,  $q$ , and  $R$  from equations (20) using equations (12), (14), and (16). Correction (17) introduces an additional rotation of the whole ring, and therefore does not affect the geometry of the ring, *i.e.* the interatomic distances and torsional and bond angles.

It was also found that co-ordinates  $x(j)$  and  $y(j)$  satisfy equation (21). It can be seen that the planar curve, describing

$$[x(j) - r_j^x]^2 + [y(j) - r_j^y]^2 = \left(\frac{a+b}{2}\right)^2 \quad (21)$$

the motion of the projection of the  $j$ th atom on the mean plane, proves to be a circle rather than an ellipse. The circle and the ellipse have a common centre and the length of the radius of the circle is the average of the lengths of the major and minor axes of the ellipse. Equations (20) can thus be simplified to (22)

$$x(j) = r_c [\sin 2(\varphi + 4\pi j/5) \sin \alpha_j - \cos 2(\varphi + 4\pi j/5) \cos \alpha_j] + r_j^x \quad (22a)$$

$$y(j) = -r_c [\sin 2(\varphi + 4\pi j/5) \cos \alpha_j + \cos 2(\varphi + 4\pi j/5) \sin \alpha_j] + r_j^y \quad (22b)$$

$$z(j) = \sqrt{2/5} q \cos(\varphi + 4\pi j/5) \quad (22c)$$

where  $r_c$  is the radius of the circle described by the projection of the  $j$ th atom on the mean plane. The dependence of the radius  $r_c$  on  $Q$  and  $R$  is similar to the dependence of  $a$  on  $Q$  and  $R$  [equation (23)]. The constants  $A$ – $C$  are collected in Table 3.

$$r_c = AR - BQ + CQ^2/R \quad (23)$$

An algorithm for cyclopentane geometry calculation with respect to bond length  $R$  and pseudorotation parameters  $q$  and  $\varphi$  is presented in the Appendix.

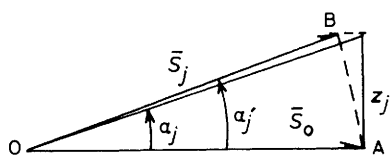
The point of intersection of the trajectory of the  $j$ th atom with the mean plane is described by the radius vector, equation

$$\vec{r}_j^{\text{int}} = \vec{r}_j' - \left(\frac{a-b}{2}\right) (\vec{r}_j'/r_j') \quad (24)$$

(24), and lies closer to the geometrical centre of the ring than the  $j$ th apex of the planar pentagon by  $R_{\text{cm}}$ . The foregoing eliminates the arbitrary assumption discussed beforehand.

**Table 4.** Comparison of torsion angle values and sums of valency angles in cyclopentane obtained with AGB and HR models for envelope and twist conformations ( $R$  0.1546 nm,  $q$   $4.35 \times 10^{-2}$  nm)

	$\varphi$ 0° (envelope)		$\varphi$ 18° (twist)	
	AGB	HR	AGB	HR
$\angle C(2)-[C(3)-C(4)]-C(5)$ (°)	25.01	26.04	13.16	13.71
$\angle C(3)-[C(4)-C(5)]-C(1)$	40.27	41.99	34.34	35.77
$\angle C(4)-[C(5)-C(1)]-C(2)$	40.27	41.99	42.29	44.13
$\angle C(5)-[C(1)-C(2)]-C(3)$	25.01	26.04	34.34	35.77
$\angle C(1)-[C(2)-C(3)]-C(4)$	0.0	0.0	13.16	13.71
$\Sigma\theta_i$	522.30	520.33	522.30	520.33



**Figure 5.** Description of the position of the  $j$ th atom according to Adams *et al.*<sup>7</sup> (AGB). Points O and A are in the  $xy$  plane. The  $xy$  plane is perpendicular to the plane of the diagram. O is the centre of the planar ring; A is the position of the  $j$ th atom in the planar ring; B is the momentary position of the  $j$ th atom in the puckered ring.

## Discussion

Any geometrical features of the pentagon, such as atomic distances, valence angles, and torsional angles may be analytically established for all values of the pseudorotation parameters  $\varphi$ ,  $Q$ , and  $R$ . Table 4 presents a comparison of some angle values obtained with the AGB model<sup>7</sup> and the model described in the present paper, referred to as HR. They were calculated using  $R$  0.1546 and  $Q$   $2.75 \times 10^{-2}$  nm ( $q$   $4.35 \times 10^{-2}$  nm) which correspond to the values used by Adams *et al.*<sup>7</sup> Table 4 shows that the absolute values of torsional angles are larger in the HR model, while the sum of valence angles is smaller. This result may be interpreted as a tendency of the HR model to increase the amplitude of puckering  $q$  or a tendency of the AGB model to diminish it. A more detailed analysis confirms the latter hypothesis.

The basic quantity describing the displacement of the  $j$ th atom from the  $xy$  plane in the AGB model is the angular displacement of the vector  $\vec{S}_j$  from this plane (Figure 5). This displacement is given by equation (25).<sup>7</sup> This approximation is

$$\alpha_j = \text{arctg}(z_j/S_0) \quad (25)$$

valid only for small values of the angle  $\alpha_j$ . The angle  $\alpha_j$  described by equation (25) is smaller than the actual angle  $\alpha_j'$  (see Figure 5); the amplitude of puckering  $q$  is thus diminished.

All distances in the five-membered ring independent of  $\varphi$ , such as  $a$  and  $b$  or the mean non-bonded atom-atom distance, calculated in the frame of the HR model, satisfy relationship

$$X/R = F(Q/R) \quad (26)$$

(26) where  $X$  is any distance independent of  $\varphi$ . All the angular quantities independent of  $\varphi$ , such as the sum of the valence angles or the sum of the squares of the torsional angles, satisfy the following relationship:  $X = f(Q/R)$ . The relationships permit an easy approximation of the quantities  $X$  by means of the least-squares fit. They also show that puckering of the ring

**Table 5.** Dependence of sum of valence angles on  $Q$  and  $R$ , according to the equation  $\Sigma\theta_i$  (°) =  $A + BQ/R - CQ^2/R^2$

$Q/R$ range		$A$	$B$	$C$	Accuracy of fit (°)
from	to				
0.0	0.1	540.00	0.08	610.11	$10^{-2}$
0.1	0.2	539.75	2.94	-620.91	$10^{-1}$
0.2	0.3	579.48	-314.67	0.00	1
0.3	0.4	503.27	242.52	1 019.55	1

**Table 6.** Dependence of mean non-bonded atom-atom distance on  $Q$  and  $R$  values according to the equation  $R_{NB} = AR + BQ - CQ^2/R$

$Q/R$ range		$A$	$B$	$C$	Accuracy of fit (nm)
from	to				
0.0	0.1	1.6180	0.0020	1.2782	$10^{-5}$
0.1	0.2	1.6161	0.0337	1.4178	$10^{-5}$
0.2	0.3	1.5964	0.2152	1.8601	$10^{-4}$
0.3	0.4	1.4911	0.9268	3.0389	$2 \times 10^{-4}$

is better described by the ratio  $Q/R$  than by the puckering amplitude  $q$ . Hence, the ratio  $Q/R$  may be called the degree of puckering of the ring.

Our approximations for the sum of valence angles in the ring and the mean non-bonded atom-atom distance are presented in Tables 5 and 6. Our approximation for the sum of valence angles is more accurate than the formula proposed by Dunitz.<sup>8</sup>

## Conclusions

The model proposed in this work describes most accurately pseudorotation in an equilateral five-membered ring. It ensures the constancy of the bond length for all values of the pseudorotation parameters  $\varphi$  and  $q$ . Unlike the Adams model,<sup>7</sup> ours fulfills Eckart's condition of immobility of the centre of mass. Furthermore, its application is not limited to small values of the puckering amplitude  $q$ . Our model permits a rapid calculation of the co-ordinates of the atoms in the ring for all given values of parameters  $\varphi$ ,  $q$ , and  $R$ .

The model is being adapted to irregular five-membered rings.

## Appendix

An algorithm for cyclopentane geometry calculation with respect to bond length  $R$  and pseudorotational parameters  $q$  and  $\varphi$  is as follows:

$$\begin{aligned} j &= \text{atom number} \\ \varphi &= \text{phase angle (}^\circ\text{)} \\ q &= \text{amplitude of puckering (nm)} \\ R &= \text{bond length (nm)} \\ \varphi_j &= \pi\varphi/180 + 4\pi j/5 \\ \alpha_j &= \pi/10 + (6-j)4\pi/5 \\ Q &= \sqrt{2/5}q \\ r_c &= R[A - BQ/R + C(Q/R)^2] \quad (A-C \text{ are listed in Table 3)} \\ r_j^x &= \{R/[2\sin(\pi/5)] - 1.6179r_c\}\cos\alpha_j \\ r_j^y &= \{R/[2\sin(\pi/5)] - 1.6179r_c\}\sin\alpha_j \\ x(j) &= +r_c[\sin(2\varphi_j)\sin\alpha_j - \cos(2\varphi_j)\cos\alpha_j] + r_j^x \\ y(j) &= -r_c[\cos(2\varphi_j)\sin\alpha_j + \sin(2\varphi_j)\cos\alpha_j] + r_j^y \\ z(j) &= Q\cos(2\varphi_j) \end{aligned}$$

$x(j)$ ,  $y(j)$ , and  $z(j)$  are cartesian co-ordinates of the  $j$ th atom (nm).

**References**

- 1 J. E. Kilpatrick, K. S. Pitzer, and R. Spitzer, *J. Am. Chem. Soc.*, 1947, **69**, 2483.
- 2 D. Cremer and J. A. Pople, *J. Am. Chem. Soc.*, 1975, **97**, 1354.
- 3 H. J. Geise, C. Altona, and C. Romers, *Tetrahedron Lett.*, 1967, 1383.
- 4 C. Altona, H. J. Geise, and C. Romers, *Tetrahedron*, 1968, **24**, 13.
- 5 C. Altona and M. Sundaralingam, *J. Am. Chem. Soc.*, 1972, **94**, 8205.
- 6 H. P. M. de Leeuw, C. A. G. Haasnoot, and C. Altona, *Isr. J. Chem.*, 1980, **20**, 108.
- 7 W. J. Adams, H. J. Geise, and L. S. Bartell, *J. Am. Chem. Soc.*, 1970, **92**, 5013.
- 8 J. D. Dunitz, *Tetrahedron*, 1972, **28**, 5459.

*Received 29th September 1981; Paper 1/1508*