# Photoelectron Spectra of Group 5 Compounds. Part IV.<sup>1</sup> A Study of the E-X Bond in $R_3EX$ Compounds (R = Me or F; E = N or P; X = O or S)

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The He(I) photoelectron (p.e.) spectra of the isoelectronic  $C_{3\nu}$  compounds Me<sub>3</sub>NO, Me<sub>3</sub>PO, Me<sub>3</sub>PS, and F<sub>3</sub>PS are discussed and assigned by comparison with the p.e. spectra of the parent bases R<sub>3</sub>P and R<sub>3</sub>N and with reference to CNDO and EHMO calculations. Bond descriptions for the E–X multiple bonds ( $\sigma$  donation,  $\rho_{\pi}$ - $d_{\pi}$  back donation) have been critically revised taking into account the variation of the central atom E, of the acceptor X, and of the substituent R, and symmetry arguments. Inductive, hyperconjugative, and d-orbital effects have been determined by the CNDO method.

VARIOUS models have been proposed to describe multiple bonding in R<sub>3</sub>EX systems in order to reflect differences in E-X bond orders, bond distances, or nucleophilicities. The models range from (a) to (d). While the equivalent

$R_2 E - \bar{X}$	R₃E→X	R3E=X	R₃E≡X
(a)	(b)	(c)	(d)

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formulations (a) and (b) are generally accepted for amine oxides, (c) and (d) are used to indicate  $p_{\pi}-d_{\pi}$  interactions<sup>2</sup> as was reported for phosphine complexes in numerous ab initio (e.g. Me<sub>3</sub>PO <sup>3,4</sup> and F<sub>3</sub>PX <sup>4-9</sup>) and semiempirical calculations (R<sub>3</sub>EO <sup>10-12a</sup> and R<sub>3</sub>PX<sup>'13,14</sup>).

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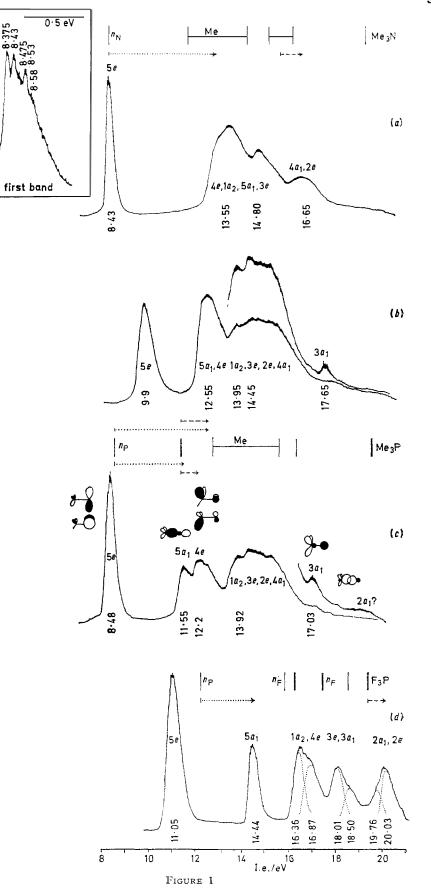
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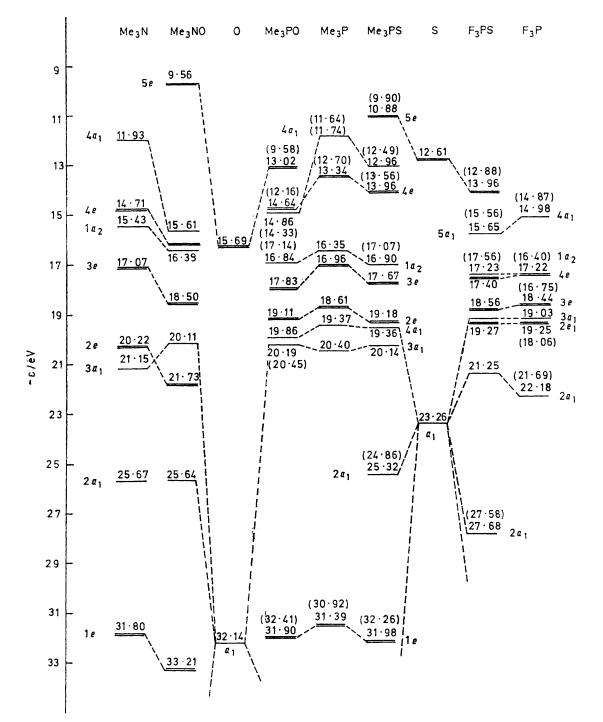


FIGURE 2 Correlation of CNDO eigenvalues of  $R_3EX$  and  $R_3E$  molecules and singlet X atoms (R = Me or F; E = N or P; X = O or S). Values in parentheses refer to calculations without d orbitals in the basis set

Caption to Figure 1

FIGURE 1 P.e. spectra of Me<sub>3</sub>NO (a), Me<sub>3</sub>PO (b), Me<sub>3</sub>PS (c), and F<sub>3</sub>PS (d) and the p.e. energies of the parent molecules Me<sub>3</sub>N,<sup>1</sup> Me<sub>3</sub>P,<sup>1</sup> and F<sub>3</sub>P.<sup>16,18,19</sup> Composite p.e. bands of R<sub>3</sub>E are reproduced as their half-widths drawn as solid lines. The stabilization of the original E-R  $\sigma$ -bonding orbitals of R<sub>3</sub>E (4e or 2e) is shown by broken arrows; lone-pair orbital stabilization ( $n_{\rm E} \rightarrow \sigma_{\rm E-x}$ ) is indicated by dotted arrows. M.o. pictures refer to 5e, 5a<sub>1</sub>, 4e, 3a<sub>1</sub>, and 2a<sub>1</sub> CNDO eigenfunctions of Me<sub>3</sub>PS

Restrictions were made with respect to the weaker and less polar P-S bond compared to P-O.9 In contrast to alkylamine oxides, (c) and (d) appeared to give a better fit to the photoelectron (p.e.) spectroscopic <sup>12</sup> and structural <sup>15</sup> peculiarities of the F<sub>3</sub>NO molecule. Symmetry properties of R<sub>3</sub>EX compounds require a cylindrical E-X bonding type as suggested by (a), (b), or (d) and no fundamental differences in P-X and N-X bonding are expected. The satisfactory agreement between calculated and experimental ionization potentials (i.p.s) has encouraged the discussion of p.e. band shifts on coordination in terms of  $\sigma$  donation,  $\pi$  back donation,<sup>14a,16</sup> and even the Lewis acidities of the acceptors X.17

In order to examine the various effects of changing substituents R, central atoms E, and acceptors X on the co-ordinate bond, we have studied a series of representative, isoelectronic, and isostructural compounds, R<sub>3</sub>EX, by p.e. spectroscopy and semiempirical calculations.

#### RESULTS AND DISCUSSION

The He(I) p.e. spectra of Me<sub>3</sub>NO, Me<sub>3</sub>PO, Me<sub>3</sub>PS, and  $F_3PS$  shown in Figure 1 are assigned by comparison with the p.e. data of the free bases (Me<sub>3</sub>E<sup>1</sup> and  $F_{3}P^{16,18,19}$ ) and molecular-orbital (m.o.) calculations (Figure 2). Numerical values of the vertical i.p.s of the resolved p.e. band maxima of  $R_3 EX$  are included in Figure 1. Here, too, the vertical i.p.s of the parent phosphines and trimethylamine are inserted. (Composite bands are approximated by showing their half-widths as horizontal solid lines.) Bands are labelled and numbered according to the  $C_{3v}$  point group and to valence orbitals, respectively.

A common feature of the spectra is the overall stabilization of R<sub>3</sub>E energies on co-ordination. This is in accordance with the simple picture of donor-acceptor bond formation according to bonding types (a) or (b), and is due to charge migration from  $(R_3)E$  to X. The known p.e. spectra of halides Y<sub>3</sub>PX <sup>14,19</sup> were interpreted correspondingly, and their assignments were supported by ab initio calculations.<sup>3-9</sup> In contrast to F<sub>3</sub>NO <sup>19a</sup> R<sub>3</sub>E band stabilization is also valid for Me<sub>3</sub>NO.

The dotted arrows in Figure 1 indicate the shifts of the original lone-pair orbitals  $n_{\rm E}$  (4a<sub>1</sub>) on formation of the dative bond  $\sigma_{E-X}$  ( $\sigma$  donation). The E-X  $\sigma$ -bonding orbital is attributed to the  $5a_1$  bands of  $R_3EX$  (Figure 1), in agreement with corresponding experimental and ab *initio* results for  $Y_3P-Y_3PX$  pairs (Y = halogeno),

\* It has been pointed out already (e.g. refs. 6 and 8) that E-R  $\sigma$  bonding is delocalized over a number of orbitals. A qualitative m.o. model based on the  $D_{3h}$  point group (see e.g. refs. 19a and 12b) shows that 4e and 2e orbitals (e' in  $D_{3h}$ ) should mix, whereas 3e (e'' in  $D_{3h}$ ; substituent group orbitals orientated perpendicular to the R<sub>3</sub> plane) should be less affected. Although this approach appears to be rather artificial on adaption to  $C_{3v}$  systems, it is supported, for example, by the various spin-orbit splittings of degenerate  $n_{\rm X}$  orbitals of  ${\rm Y}_3{\rm E}$  and  ${\rm Y}_3{\rm EX}$ compounds (e.g. ref. 14) due to the differing steric arrangements of substituent group orbitals (cf. K. Wittel and R. Manne, Theor. Chim. Acta, 1974, 33, 347 and refs. therein).  $\dagger 1 {\rm eV} \approx 1.60 \times 10^{-19} {\rm J}.$ \* It has been pointed out already (e.g. refs. 6 and 8) that

Complex vibrational fine structure was observed in the first band of Me<sub>3</sub>NO (5e, Figure 1) which could be caused either by two overlapping progressions or by Jahn-Teller distortion.

except for  $Cl_3PO.^8$  The *p*-type E-X  $\sigma$ -bonding m.o.  $(5a_1)$  has contributions from  $n_X$  lone-pair orbitals as well as R orbitals due to symmetry  $(a_1)$  and energy effects [i.e.  $(p_{\rm X}) \prec \rightarrow$  i.e.  $(p_{\rm R})$  where i.e. = ionization energy]. The character of the three highest  $a_1$  orbitals is illustrated in Figure 1 with CNDO eigenfunctions of Me<sub>3</sub>PS.

Assuming a singlet state for the acceptors X during complex formation according to simple Lewis structures, their orbitals transform as  $e(p_X)$  and  $a_1(s_X)$ . On coordination the  $a_1$  orbital of X is involved mainly in stype  $E-X \sigma$  bonding and ionizes at higher energies. The degenerate lone-pair orbitals,  $n_{\rm X}$ , undergo hyperconjugative interaction with the  $E-R \sigma$ -bonding m.o.s of  $R_{a}E$  (4e or 2e, respectively), thus producing RE-X  $\pi$ -antibonding (5e) and RE-X  $\pi$ -bonding m.o.s (4e or 2e, respectively), if E-X  $p_{\pi}$ - $d_{\pi}$  feedback is ignored. The corresponding 5e and 4e CNDO eigenfunctions are portrayed in Figure 1 for the case of Me<sub>3</sub>PS. The broken arrows (Figure 1) indicate the shifts of the degenerate E-R  $\sigma$ -bonding orbitals of R<sub>3</sub>E on co-ordination, including the main  $\pi$ -back-donation effects. From calculations, the degenerate E-Y and N-Me σ-bonding orbitals are assigned to the 2e orbitals of  $R_3E$ , as is commonly accepted (Me<sub>3</sub>N  $^{1,20a}$  and  $Y_{3}P^{19}$ ), whereas the P-Me  $\sigma$ m.o.s are mainly attributed to the 4e orbital of Me<sub>3</sub>P.<sup>1,20b,\*</sup> For the alkyl derivatives this approach is supported by experimental evidence, since bands due to orbitals containing a large amount of central-atom character are expected to be separated from those due to Me. This is fulfilled for the second band of  $Me_3P$  (4e, 11.34 eV) and the third band of Me<sub>3</sub>N (2e, 15.7 eV).<sup>†</sup>

If the central atoms are allowed to have d character, 5e correlates with the bonding combination of the acceptor e orbital and the first virtual 5e orbital of  $R_3E$ which is believed to have appreciable d character.<sup>3</sup> No decision can be made from p.e. spectra as to whether the 5e m.o.s of  $R_3EX$  are destabilized by suitable empty orbitals on the central atom (e.g. d and  $\sigma_{E-R}^*$ ). RE-X  $\pi$ -bonding character might be reflected in the greater half-width of the first band of Me<sub>3</sub>PO (5e) compared to Me<sub>2</sub>PS, or its rather high ionization energy in relation to the amine oxide ‡ (Figure 1). A consistent interpretation can still be given by neglecting d-orbital effects, since the 5e energies parallel the relative  $p_{\rm X}$ valence ionization potentials (v.i.p.s) as well as E-X distances and X electronegativities within the series  $R_3PX [X = 0 \text{ or } S (Se^{21})].$  Furthermore, an analogous trend in half-widths of  $n_{\mathbf{X}}$  bands was found in the furan

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further details can be given by the authors.

series  $(CH_2)_4 X$  (X = O to Te)<sup>22</sup> and was discussed in terms of increased hyperconjugation on going from the Te to the O derivative.

Neither the spectrum of Me<sub>3</sub>NO nor the calculations (Figures 1 and 2) show any fundamental differences from Me<sub>3</sub>PO. The same is true for F<sub>3</sub>PX. The lower first i.p. (Me<sub>3</sub>NO, 8.43; Me<sub>3</sub>PO, 9.95 eV) may be a consequence of enhanced hyperconjugative *e-e* orbital interaction due to a smaller E-X bond distance and/or a strong inductive effect. This is compatible with the relatively larger average stabilization of *R* levels (Figure 1), with  $\Delta a_2$  data from CNDO calculations (Figure 2) and with p.e. band shifts of analogous phosphine and amine complexes suggesting greater inductive effects on N-bonding.<sup>23</sup> The average stabilization of *R* levels ( $n_{\rm X}$ , Me) is *ca.* 0.75 eV in Me<sub>3</sub>NO, *ca.* 0.5 eV in Me<sub>3</sub>PX, and

concerning 5e can be deduced from the p.e. spectra of  $R_3E$  and  $R_3EX$ .

Although no conclusions about inductive and hyperconjugative interactions can be drawn from the experimental evidence, a ' composite-molecule ' approach based on the method of Baird <sup>24</sup> may be helpful in gaining some insight into the dimensions of formal inductive (I) and pseudo- $\pi$  effects during the hypothetical reaction  $R_3E +$  $X \longrightarrow R_3EX$ . Figure 3 shows the stepwise coupling procedure for *e* orbitals of the moieties  $R_3E$  and X which has been simulated by CNDO eigenvalues of  $H_3P$ , O and S (singlet states),  $H_3PX$ , and their fictitious fragments.<sup>24</sup> Hydrogen derivatives were taken for clarity, although the inductive R standard  $1a_2$  is missing here. (This omission is not serious, since the relative  $\Delta a_2$  values from Figure 2 should still hold.) Figure 3 can give an

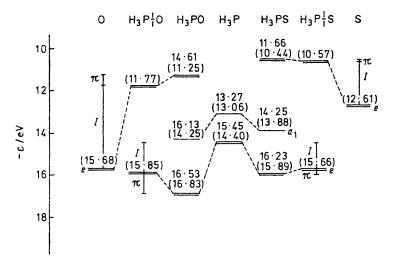


FIGURE 3 Inductive (I) and hyperconjugative ( $\pi$ ) effects on co-ordination  $R_3E + X \rightarrow R_3EX$ , simulated by CNDO eigenvalues in a 'composite-molecule' model. Values in parentheses refer to calculations without d orbitals in the basis set

ca. 0.4 eV in  $R_3PS$  (most clearly in the latter case due to band separation), while E-R  $\sigma$ -bonding orbitals are shifted to greater extents.

All R<sub>3</sub>E and R<sub>3</sub>EX molecules contain a unique  $a_2$  orbital. Although entirely located in the R<sub>3</sub> skeleton due to symmetry, its energy changes with charge transfer to (or from) X and thus is a quantitative measure of the inductive effect (neglecting geometry changes on coordination). The inductive lowering  $\Delta a_2$  is most obvious in the p.e. spectra of Y<sub>3</sub>P and Y<sub>3</sub>PX, where  $1a_2$  bands are better resolved (Cl<sub>3</sub>P and Cl<sub>3</sub>PO, 0.70; <sup>14, 19a</sup> Cl<sub>3</sub>P and Cl<sub>3</sub>PS, 0.32; <sup>14, 19a</sup> F<sub>3</sub>P and F<sub>3</sub>PS, 0.47 eV), and is assumed to parallel the average stabilization of Me levels of Me<sub>3</sub>E. A simple electrostatic model considering electron-density distribution <sup>21</sup> shows that the inductive stabilization for e and  $a_1$  orbitals is of course not the same as  $\Delta a_2$ . Thus the inductive contributions to shifts  $\Delta 4e$  or  $\Delta 2e$  remain undetermined (Figure 1), and neither hyperconjugative destabilization nor  $p_{\pi}$ - $d_{\pi}$  back donation

impression of the complexity of the inductive and hyperconjugative perturbations on co-ordination. An analogous approach is shown for the present compounds in Figure 2.

In order to study the effect of  $\sigma$  transfer, indicated by the dotted arrows in Figure 1, the  $a_1$ - $a_1$  interactions have to be analyzed. The transformation of the orginal  $n_{\rm E}$ lone-pair orbitals on complex formation has already been mentioned. The p-type E-X  $\sigma$ -bonding orbitals correlating with  $n_{\rm E}$  have been assigned to the  $5a_1$  bands in each case. These bands are separated only for R<sub>3</sub>PS. The  $\sigma_{\rm P-O}$  level is expected to coincide with  $\sigma_{\rm P-Me}$  energies (12.5 eV) according to *ab initio*<sup>3</sup> and CNDO eigenvalues (Figure 2). The  $\sigma_{\rm N-O}$  energy is obscured by the Me bands centred at 13.6 eV (Figure 1), but is assigned to the low-energy shoulder of this band in agreement with our CNDO eigenvalues (Figure 2).

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<sup>&</sup>lt;sup>24</sup> N. C. Baird, Theor. Chim. Acta, 1970, **16**, 239; C. Müller, A. Schweig, and H. Vermeer, Angew. Chem., 1974, **86**, 275; W. Enßlin, H. Bock, and G. Becker, J. Amer. Chem. Soc., 1974, **96**, 2757. Extinction of  $\pi$  (and  $\sigma$ ) interactions is achieved by setting the off-diagonal elements [A(i,j) in subroutine SCFCLO] between  $\pi$ -type orbitals P-H and  $n_{\mathbf{X}}$  in the Hartree-Fock matrix to zero before diagonalization.

Although the lengths of the arrows could reflect the expected  $\sigma$ -donation abilities as predicted by theory, Figure 1 and particularly Figure 2 show the remarkable influence of the relative s v.i.p.s of the acceptors X on all original  $a_1$  orbitals of the free bases. Even though the  $s_{\mathbf{X}}$  orbital interferes mainly with  $3a_1$  (the  $3a_1$  orbital of F<sub>3</sub>P, e.g. at 18.57 eV, remaining nearly constant in  $F_3PS$ ,  $4a_1$  at ca. 18.50 eV) and less with  $4a_1$  ( $n_E$ ) due to energy effects, the rather small  $n \rightarrow \sigma$  shifts for sulphur and borane complexes do not necessarily mean weak  $\sigma$ donation ( $\Delta n_{\rm E} \rightarrow 5a_1$ : Me<sub>3</sub>P and Me<sub>3</sub>PO, ca. 3.95; Me<sub>3</sub>P and  $Me_3PS$ , 2.96;  $Me_3P$  and  $Me_3P\cdot BH_3$ , ca. 2.0; <sup>21</sup>  $Me_3N$  and  $Me_3NO$ , ca. 4.65;  $Me_3N$  and  $Me_3N\cdot BH_3$ , 2.86; <sup>16</sup>  $F_3P$  and  $F_3PO$ , 3.4; <sup>19</sup>  $F_3P$  and  $F_3PS$ , 2.15 eV). The shifts reflect the effect of nearby  $s_8$  or  $B_8H_3$  orbitals  $(a_1)$ that counteract the inductive stabilization on complex formation. The remarkable amount of  $s_{\mathbf{X}}$  character, even in the  $5a_1$  m.o., as a function of the  $s_X \leftrightarrow n_E$ energy gap is confirmed by ab initio calculations (e.g.  $F_3PO$ , 14;  $F_3PS$ , 20%)<sup>4</sup> and our CNDO results.<sup>21</sup>

Obviously, concepts such as ' $\sigma$  donation ' or ' $\pi$  back donation ' cannot be substantiated by p.e. spectroscopy and their application therefore seems vague and only relative. In contrast, theoretical studies can be used to determine differing bond orders or charge effects as suggested by the models (a)—(d) of the E-X co-ordinate bonding which were mentioned in the Introduction. For

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comparison, CNDO and EHMO bond orders, net atomic charges, experimental stretching force constants, and observed and calculated dipole moments of the present compounds (and partially of their parent bases  $R_3E$ ) are summarized in Supplementary Publication No. SUP 21743 (2 pp.).\* The sequence of CNDO bond orders calculated without d orbitals <sup>25</sup> for analogous  $C_{3v}$  phosphine and amine complexes <sup>21</sup> is:  $_{\rm HB}$ -P<sub>F</sub> < S-P<sub>Me</sub>  $\approx$   $_{\rm HB}$ -P<sub>Me</sub> < O-N<sub>Me</sub>  $\approx$   $_{\rm HB}$ -N<sub>Me</sub>  $\leqslant$  S-P<sub>F</sub>  $\approx$  S-P<sub>Cl</sub> < O-P<sub>Me</sub>  $\approx$  O-N<sub>F</sub> < P-P<sub>Cl</sub> = O-P<sub>F</sub>. Apart from reversals in the order due to the inclusion of d orbitals with respect to N-X and P-X, the sequence is well reproduced by results of Choplin and Kaufmann <sup>10</sup> for comparative data.

### EXPERIMENTAL

The He(I) (584 Å) spectra were recorded on a Perkin-Elmer PES 16 spectrometer. Calibration was with argon and xenon. The estimated error in the ionization potentials is  $\pm 30-40$  meV. The resolution ranged between 25 and 40 meV.

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\* For details see Notices to Authors No. 7, *J.C.S. Dalton*, 1975 Index issue (items less than 10 pp. are supplied as full-size copies) <sup>25</sup> R. J. Boyd, *Canad. J. Chem.*, 1973, **51**, 1151.