

## A Mössbauer Study of Rare-earth Metal Pyrochlores $A_2FeSbO_7$ (A = Lu, Er, Tb, or Sm) using the 37.2-keV $\gamma$ -Ray of Antimony-121

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Isomer shifts and quadrupole splittings have been obtained from the  $^{121}\text{Sb}$  Mössbauer spectra of four members of the pyrochlore series  $A_2FeSbO_7$  (A = Lu, Er, Tb, or Sm). The isomer shift (i.s.) decreases linearly with  $r(A^{3+})$ , the radius of the rare-earth metal ion, and is in the opposite sense to that previously found for the  $^{57}\text{Fe}$  i.s. in the same series. This result is interpreted in terms of a decrease in the covalent character of the Sb–O bonds as the distortion of the  $\text{SbO}_6$  octahedron increases. The quadrupole-splitting results are of insufficient precision to show any clear trends.

THE crystallographic and  $^{57}\text{Fe}$  Mössbauer parameters of the pyrochlore series  $A_2FeSbO_7$  (A = rare-earth metal) have been recorded and discussed in an earlier paper.<sup>1</sup> In the pyrochlore structure the Fe and Sb atoms randomly occupy trigonally distorted octahedral B sites with the distortion decreasing with increasing radius of the rare-earth metal ion,  $r(A^{3+})$ . This work was extended by that of Ayasse<sup>2</sup> who studied also the stannate series  $A_2Sn_2O_7$  by  $^{119m}\text{Sn}$  Mössbauer spectroscopy. Mössbauer data on the stannates has also been

published by Belyaev *et al.*<sup>3</sup> and by Loebenstein *et al.*<sup>4</sup> Several correlations have been established between the quadrupole splitting (q.s.) and isomer shifts (i.s.) in these compounds and  $r(A^{3+})$  or the positional parameter  $x(\text{O}_2)$  of the 48(f) oxygen anions.  $x(\text{O}_2)$  and the cubic lattice parameter  $a_0$  are sufficient to define all the ionic positions in the pyrochlore structure (space group  $Fd\bar{3}m$ ).

The present work was undertaken to see whether

<sup>1</sup> O. Knop, F. Brisse, R. E. Meads, and J. Bainbridge, *Canad. J. Chem.*, 1968, **46**, 3829.

<sup>2</sup> C. Ayasse, Ph.D. Thesis, Dalhousie University, Halifax, Nova Scotia, 1972.

<sup>3</sup> L. M. Belyaev, I. S. Lyubutin, L. M. Dem'yanets, T. V. Dmitrieva, and L. P. Matina, *Soviet Phys. (Solid State)*, 1969, **11**, 424.

<sup>4</sup> H. M. Loebenstein, R. L. Zilber, and H. Zmora, *Phys. Letters*, 1970, **A33**, 453.

similar correlations exist between the q.s. and i.s. of  $^{121}\text{Sb}$  in the iron antimonates and the radius of the rare-earth metal  $r(\text{A}^{3+})$ . The experiments were carried out on the compounds for which  $\text{A} = \text{Lu}, \text{Er}, \text{Tb}, \text{and Sm}$ .

#### EXPERIMENTAL

The method of preparation of the iron antimonates has already been given in detail.<sup>1,2</sup> Mössbauer absorbers were made by enclosing the powder in Perspex holders, each

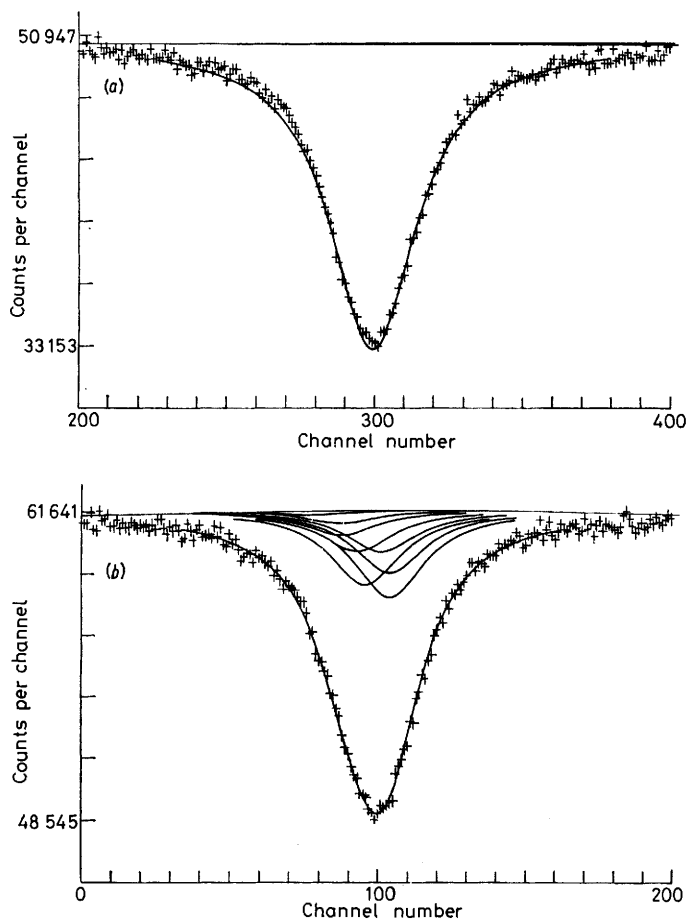


FIGURE 1. Antimony-121 Mössbauer spectra at 77 K of (a)  $\text{Sm}_2\text{FeSbO}_7$  fitted to a single Lorentzian line, (b)  $\text{Er}_2\text{FeSbO}_7$  fitted by the eight-line model

absorber having a thickness of *ca.*  $10 \text{ mg cm}^{-2}$  of antimony. Spectra were recorded using a conventional constant-acceleration Mössbauer spectrometer under conditions of constant source-absorber-detector geometry. Both source and absorber were kept at 77 K. The source used was  $^{121m}\text{Sn}$  in a matrix of  $\text{CaSnO}_3$  with a nominal strength of 1 mCi (New England Nuclear Corporation). All the spectra show a single resonance-absorption dip in the centre of the spectrum with an isomer shift (i.s.) close to that of  $\text{Sb}^{5+}$  in an oxide environment, as expected. A representative spectrum of the series, that of  $\text{Sm}_2\text{FeSbO}_7$ , is shown in Figure 1. All the spectra were least-squares fitted to two models: (a) a single line and (b) eight lines corresponding to the eight lines expected in a pure quadrupole spectrum of  $^{121}\text{Sb}$ . In the fits to model (b) the line intensities were

constrained to be proportional to the squares of the appropriate Clebsch-Gordan coefficients and the linewidths were constrained to be equal. The ratio of the quadrupole moments of the excited and ground states was taken as 1.38 : 1. Both positive and negative starting values of the quadrupole coupling constant  $eQ_gV_{zz}$  were tried and, although both fits converged, the negative value always gave the lower  $\chi^2$  values. The asymmetry parameter  $\eta$  was assumed to be zero. [See the results of point-charge calculations of electric-field gradient (e.f.g.) due to the distortion of the nearest-neighbour  $\text{O}^{2-}$  octahedron given in the following section.]

The  $\chi^2$  values obtained for both fits to models (a) and (b), with negative  $eQ_gV_{zz}$  values, are given in the Table.

Antimony-121 Mössbauer parameters <sup>a</sup> of  $\text{A}_2\text{FeSbO}_7$  at 77 K

A	$\delta$	$eQ_gV_{zz}$	$\Gamma$	$\chi^2$ <sup>b</sup>	
				(a)	(b)
Lu	0.132		3.29	203; 216	179; 205
Er	0.102	-6.7	3.49	225; 255	184; 214
Tb	0.058	-6.5	3.54	250; 273	208; 210
Sm	0.014	-6.5	3.39	259; 252	197; 197

<sup>a</sup>  $\delta$  = Isomer shift relative to  $\text{CaSnO}_3$  at 77 K ( $\pm 0.05 \text{ mm s}^{-1}$ );  $eQ_gV_{zz}$  = ground-state quadrupole-coupling constant ( $\pm 1.0 \text{ mm s}^{-1}$ );  $\Gamma$  = linewidth ( $\pm 0.05 \text{ mm s}^{-1}$ ). <sup>b</sup> Refer to fits to models (a) and (b) respectively (see text). A negative value of  $e^2qQ_g$  was used in each type (b) fit quoted. A fit using positive  $e^2qQ_g$  of the  $\text{Sm}_2\text{FeSbO}_7$  spectrum gave  $\chi^2 = 254$  and 230.

Since in each experimental run the spectrum was recorded twice, corresponding to positive and negative values of the relative acceleration between source and absorber, the  $\chi^2$  values are quoted in pairs. As a lower value of  $\chi^2$  was consistently obtained for fits to model (b), it may be reasonably concluded that there is a significant, albeit small, quadrupole splitting. The  $^{121}\text{Sb}$  Mössbauer parameters obtained from fitting of model (b) to the spectra are also given in the Table. The isomer shifts are quoted relative to calcium stannate at 77 K; to convert into InSb add  $+8.4 \text{ mm s}^{-1}$ .

#### RESULTS AND DISCUSSION

Although the precision of values for the  $^{121}\text{Sb}$  quadrupole splitting (q.s.) obtained in the present experiments is insufficient to show any definite trend with the radius of the rare-earth metal  $r(\text{A}^{3+})$ , it seems firmly established that the sign of  $e^2qQ_g$  is negative. Because of the negative value ( $-0.28$ ) for  $Q_g$  in  $^{121}\text{Sb}$ , the result requires a positive value for the  $V_{zz}$  component of the nuclear e.f.g. tensor. Using a point-charge model, the e.f.g. at a B site in the pyrochlore structure due to the nearest-neighbour octahedron of  $\text{O}^{2-}$  ions can be calculated in terms of the  $x(\text{O}_2)$  parameter  $x$  and the cubic lattice parameter  $a_0$ . The results are (1) and (2). The charge

$$V_{zz} = \frac{-3(16x - 7)}{32\pi\epsilon_0 a_0^3 [(\frac{1}{2} - x)^2 + \frac{1}{32}]^{\frac{5}{2}}} \quad (1)$$

$$V_{xx} = V_{yy} = \frac{3(16x - 7)}{64\pi\epsilon_0 a_0^3 [(\frac{1}{2} - x)^2 + \frac{1}{32}]^{\frac{5}{2}}} \quad (2)$$

on each  $\text{O}^{2-}$  has been assumed to be  $-2e$  and the polarizability of the ions has been ignored. A positive

sign for the  $V_{zz}$  component of the nuclear e.f.g. tensor would be expected from this calculation provided that: (a) the nuclear e.f.g., including the valence contribution, has the same sign as the e.f.g. due to the external charges; and (b) that  $x$  is less than  $\frac{7}{18}$ . Although extended point-charge calculations by Ayasse<sup>2</sup> have shown that the effects of more remote ions on  $V_{zz}$  may not be neglected, the nearest-neighbour contribution is expected to dominate. Values of  $x(\text{O}_2)$  have been obtained experimentally in a limited number of cases for stannates and iron antimonates by refinement of crystallographic data,<sup>2,5,6</sup> and are indeed found to be significantly lower than  $0.4375(\frac{7}{18})$ ; for example, the value in  $\text{Lu}_2\text{FeSbO}_7$  is 0.4216. The values of  $x(\text{O}_2)$  required to give agreement between measured values and point-charge calculations of the q.s. of  $^{57}\text{Fe}$  in the iron antimonates are in good agreement with the crystallographically determined values. In the stannates it is possible to establish proportionality between the measured quadrupole splittings and the calculated values.<sup>2</sup> No direct experimental evidence is available for the sign of  $V_{zz}$  at  $^{119}\text{mSn}$  in the stannates.

The results obtained by Ayasse<sup>2</sup> show a linear decrease of  $^{57}\text{Fe}$  q.s. with increasing  $r(\text{A}^{3+})$  in the iron antimonate series at both room temperature and 77 K. At room temperature the values range from  $1.01 \text{ mm s}^{-1}$  for the lutetium compound to  $0.79 \text{ mm s}^{-1}$  for that of Eu. The corresponding range at 77 K is from  $1.04$  to  $0.84 \text{ mm s}^{-1}$ . The behaviour of the iron antimonates having  $\text{A}_2 = \text{Nd}_2, \text{Pr}_2, \text{ or LaPr}$  deviated from the above linear variation. This anomaly is explained by the presence of an impurity orthoferrite phase. Ayasse<sup>2</sup> found also a linear variation with  $r(\text{A}^{3+})$  of the q.s. at  $^{119}\text{mSn}$  sites in the stannates. At room temperature the measured range is from  $0.79 \text{ mm s}^{-1}$  at Lu to  $0.44 \text{ mm s}^{-1}$  at La. At room temperature the results range from  $0.82 \text{ mm s}^{-1}$  at Yb to  $0.46 \text{ mm s}^{-1}$  at La. The q.s. data at 77 K for the stannates obtained by Loebenstein *et al.*<sup>4</sup> show consistently higher values. The reason for this is not clear but may be due to the use of rather thick absorbers.<sup>7</sup> At 77 K, quadrupole splittings for both pyrochlore series extrapolate to zero at a value of  $r(\text{A}^{3+})$  of *ca.*  $1.33 \text{ \AA}$ . This decrease of the quadrupole splittings with increasing  $r(\text{A}^{3+})$  is consistent with the decreasing trigonal distortion from octahedral symmetry of the B site.

Figure 2(a) shows a plot of the  $^{121}\text{Sb}$  i.s. against  $r(\text{A}^{3+})$  for the four pyrochlores studied, and shows convincingly that the i.s. decreases with  $r(\text{A}^{3+})$ . This is quite unlike the variation obtained at both room temperature and 77 K for  $^{57}\text{Fe}$  in the iron antimonates<sup>2</sup> where the i.s. increases with  $r(\text{A}^{3+})$ , as shown in Figure 2(b). Now an increase in the  $^{57}\text{Fe}$  i.s. corresponds to a decrease in  $s$ -electron density at the iron nucleus. Ayasse<sup>2</sup> argued that this decrease of  $s$ -electron density with  $r(\text{A}^{3+})$  is

consistent with what would be expected from the simple  $\sigma$ -bonding covalent model: increase in the Fe-O separation with increasing  $r(\text{A}^{3+})$  will lead to an increase in the ionicity of the bond, and a decrease in the covalent donation to the  $4s$  orbital of the Fe atom, resulting in a decrease in  $s$ -electron density at the Fe nucleus, as

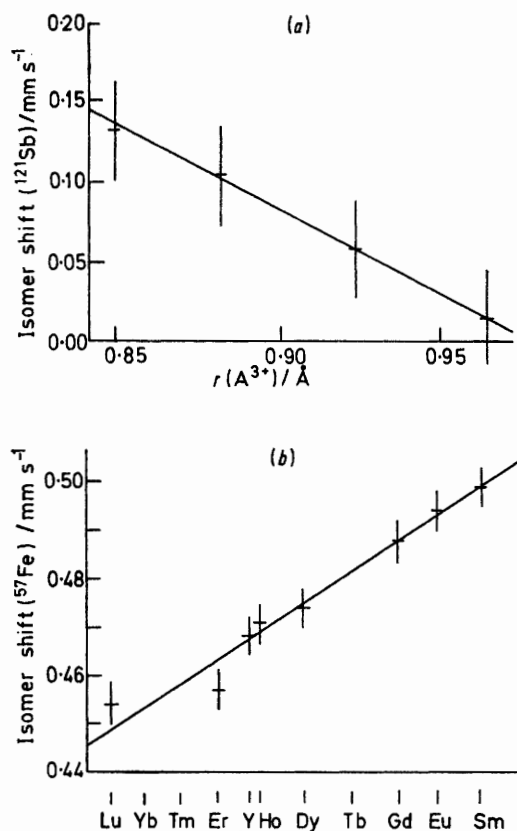


FIGURE 2 Variations of isomer shift with  $r(\text{A}^{3+})$  in  $\text{A}_2\text{FeSbO}_7$ : (a)  $^{121}\text{Sb}$ ; (b)  $^{57}\text{Fe}$

observed. Since the sign of the change in the nuclear radius  $\Delta R$  between the ground and the first excited state is the same for both  $^{57}\text{Fe}$  and  $^{121}\text{Sb}$ , this simple  $\sigma$ -bonding model predicts the variation of the  $^{57}\text{Fe}$  and  $^{121}\text{Sb}$  isomer shifts with  $r(\text{A}^{3+})$  to be similar since Fe and Sb occupy identical sites in the  $\text{A}_2\text{FeSbO}_7$  lattice. This is inevitably an oversimplified view since the charge-density calculations of Ruby *et al.*<sup>8</sup> show that, for antimony, whilst the addition of  $5s$  electrons increases the  $s$ -electron density at the nucleus, addition of  $5p$  electrons will decrease it. However, the change in  $5p$  density must be much more marked if it is to produce the same effect as a change in  $5s$  density, so that an increase in  $\sigma$  covalency will still generally cause a decrease in the i.s. at  $^{121}\text{Sb}$ . This has also been demonstrated by Brukhanov *et al.*<sup>9</sup> who showed that in a series of octahedrally co-ordinated antimony(v) compounds the

<sup>5</sup> O. Knop, F. Brisse, L. Castelliz, and Sutarno, *Canad. J. Chem.*, 1965, **43**, 2812.

<sup>6</sup> F. Brisse and O. Knop, *Canad. J. Chem.*, 1968, **46**, 859.

<sup>7</sup> R. E. Meads, B. M. Place, F. W. D. Woodhams, and R. C. Clark, *Nuclear Instr. Methods*, 1972, **98**, 29.

<sup>8</sup> S. L. Ruby, G. M. Kalvius, G. B. Beard, and R. E. Snyder, *Phys. Rev.*, 1967, **159**, 239.

<sup>9</sup> V. A. Brukhanov, B. Z. Iofa, V. Kotkhekar, S. I. Semenov, and V. S. Shpinel, *Zhur. eksp. teor. Fiz.*, 1967, **53**, 1582.

i.s. at  $^{121}\text{Sb}$  decreases with decreasing electronegativity difference between Sb and the ligand atoms. Moreover, in the present series of compounds, the relative constancy of the Sb q.s. implies that the  $5p$  density remains reasonably constant through the series.

Thus the fact that the  $^{121}\text{Sb}$  i.s. trend is in the opposite sense to that of  $^{57}\text{Fe}$  indicates that there must be at least one other contribution to the  $s$ -electron density which varies in the opposite sense to the  $\sigma$ -bonding contribution through the series. In Sb this second contribution must dominate the  $\sigma$ -bonding contribution so that the  $s$ -electron density increases through the series with  $r(\text{A}^{3+})$ . This second contribution could arise in the following manner. So far we have not found it necessary to consider the trigonal distortion of the B-site oxygen octahedron. The q.s. of the high-spin  $\text{Fe}^{3+}$  ion (and to a limited extent the  $\text{Sn}^{4+}$  ion) is of course an accurate measure of this distortion which decreases with increasing  $r(\text{A}^{3+})$ . Now as the distortion of this octahedron is reduced as  $r(\text{A}^{3+})$  changes there will be increased overlap of the Sb and O orbitals which will lead to a greater degree of covalency. This in turn will give rise to an increase in  $5s$  density on the Sb ion.

If this increase outweighs the decrease expected from the increase in bond length, then the i.s. will decrease with increasing  $r(\text{A}^{3+})$ , as observed.

Covalency effects of Sn-O bonds are expected to be intermediate between those of Fe-O and Sb-O bonds, since the polarizing power of the metal ions increases in the order  $\text{Fe} < \text{Sn} < \text{Sb}$ , so that it is tempting to conclude that the relative constancy of the  $^{119\text{m}}\text{Sn}$  i.s. arises from a cancellation of the effects due to bond-length and distortion changes. Very precise measurements of the  $^{119\text{m}}\text{Sn}$  isomer shifts would be required to decide with certainty whether the gradient of the variation with  $r(\text{A}^{3+})$  is positive or negative.

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