

## Infrared and Electronic Spectra of Matrix-isolated Tetrafluoro- and Tetrachloro-oxo-molybdenum(vi) and -oxotungsten(vi) and Tetrabromo-oxotungsten(vi)

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The prominent i.r. bands of the title compounds isolated as monomers in low-temperature matrices have been obtained. Some anomalies in the literature regarding the fundamentals of  $\text{WOCl}_4$  have been resolved, and in particular, it is shown *via* high-resolution studies, that whilst argon matrices generally yield multiple trapping sites, nitrogen matrices typically produce only one site. Electronic spectra in  $\text{N}_2$  matrices have been recorded for all five compounds and compared with the predictions of  $X\alpha$  calculations. The matrix electronic spectra are shown to be superior to gas-phase spectra in terms of resolution, the observation of vibrational fine structure, and the absence of any complications due to the possible presence of decomposition products.

FIVE oxide tetrahalides of molybdenum(vi) and tungsten(vi) have been prepared and structurally characterised in the solid state, *viz.*  $\text{MoOF}_4$ ,  $\text{MoOCl}_4$ ,  $\text{WOF}_4$ ,  $\text{WOCl}_4$ , and  $\text{WOBBr}_4$ . All contain a distorted six-coordinate arrangement about the metal. In  $\text{MoOF}_4$  the structure consists of infinite chains with *cis*-fluorine bridges,<sup>1</sup> while  $\text{WOF}_4$  is a fluorine-bridged tetramer<sup>2</sup> not an oxygen-bridged molecule as originally reported.<sup>3</sup> Molybdenum(vi) oxide tetrachloride contains square-pyramidal molecules weakly associated as dimeric units *via* chlorine bridges.<sup>4</sup> In contrast,  $\text{WOCl}_4$  and  $\text{WOBBr}_4$  are *trans*-oxygen-bridged polymers.<sup>5</sup> In the gas phase, electron diffraction,<sup>6-8</sup> i.r. and Raman spectroscopy,<sup>9-11</sup> and mass spectrometry<sup>10,12</sup> have shown that the major species present are five-co-ordinate square-pyramidal monomers with  $C_{4v}$  symmetry. Evidence for the existence of  $\text{MoOBr}_4$  and the oxide tetraiodides is tenuous, although Ward and Stafford<sup>9</sup> suggested that new i.r. bands observed in the vapours over  $\text{MO}_3\text{-LiX}$  ( $M = \text{Mo}$  or  $\text{W}$ ;  $X = \text{halide}$ ) melts could be due to  $\text{MoOBr}_4$  and  $\text{WOI}_4$ .

There is considerable recent interest in the existence of  $\text{MOX}_4$  species in the gas phase, not least because of the presence of such species in lamps in which traces of halogens and/or halocarbons have been added to the noble-gas lamp atmospheres. As part of an investigation of simple monomeric transition-metal oxohalides using matrix-isolation spectroscopy, *e.g.*  $\text{CrOCl}_3$ ,<sup>13</sup> we report here a study of the monomeric  $\text{MOX}_4$  species, isolated in noble-gas matrices, employing high-resolution i.r. spectroscopy and electronic spectroscopy. Such studies<sup>14</sup> have a number of important advantages over gas- and solution-phase studies: (a) the matrix sampling technique results in well isolated monomeric molecules with no significant interference from dimers or impurities, (b) narrow linewidths in the i.r. spectra allow the structure of the species to be deduced unambiguously because isotopic splitting patterns can be resolved, and (c) the low temperatures and the isolation of molecules are the most favourable for the observation of vibrational fine structure in electronic spectroscopy.

During the course of this work two reports appeared

describing the analysis of the electronic spectra of some  $\text{MOX}_4$  species in high-temperature vapours on the basis of SCF- $X\alpha$ -SW calculations using the overlapping atomic sphere model.<sup>15,16</sup> In particular, assignments for  $\text{MoOCl}_4$ ,<sup>15</sup>  $\text{WOCl}_4$ ,<sup>15,16</sup> and  $\text{WOBBr}_4$ <sup>15,16</sup> were made and may be compared with the results described in our experiments. An electronic spectrum of  $\text{MoOCl}_4$  in  $\text{CCl}_4$  solution has also been reported.<sup>17</sup>

### EXPERIMENTAL

*Preparation of Compounds.*—The compounds  $\text{MoOF}_4$  and  $\text{WOF}_4$  were prepared by the reaction of  $\text{MO}_3$  and  $\text{F}_2$  in a flow system.<sup>18</sup>

The compounds  $\text{MoOCl}_4$  [%Cl: 55.6 (found), 55.9 (calc.)] and  $\text{WOCl}_4$  [%Cl: 41.0 (found), 41.5 (calc.)] were prepared by the reaction of  $\text{MO}_3$  with thionyl chloride under reflux followed by vacuum sublimation.<sup>19</sup> The compound  $\text{WOCl}_4$  was also conveniently prepared<sup>20</sup> from  $\text{WO}_3$  and  $\text{BCl}_3$  in a manner analogous to that used to prepare  $\text{WOBBr}_4$  (see below).

$\text{WOBBr}_4$ . Purified  $\text{WO}_3$  (*ca.* 3 g) was placed in an ampoule, attached to a vacuum line, and evacuated. A small quantity of  $\text{BBr}_3$  was condensed onto the trioxide using liquid nitrogen and the ampoule allowed to warm slowly in an ice-bath.† As the ampoule warmed a vigorous exothermic reaction occurred and a black solid was produced. This procedure was repeated until some  $\text{BBr}_3$  remained unreacted, then the ampoule was evacuated, sealed, and the products vacuum sublimed. The product,  $\text{WOBBr}_4$  [%Br: 61.5 (found), 61.5 (calc.)], was a black lustrous solid.

All reagents ( $\text{SOCl}_2$ ,  $\text{BCl}_3$ ,  $\text{BBr}_3$ ,  $\text{Na}_2[\text{WO}_4]$ , and  $[\text{NH}_4]_2[\text{MoO}_4]$ ) were obtained from B.D.H. Chemicals Ltd. The boron halides were distilled under vacuum before use. Commercial samples of  $\text{MoO}_3$  and  $\text{WO}_3$  were found to be relatively unreactive and so fresh samples were prepared as follows. The compound  $\text{MoO}_3$  was prepared by the thermal decomposition of  $[\text{NH}_4]_2[\text{MoO}_4]$  at *ca.* 200 °C. For  $\text{WO}_3$ , a solution of  $\text{Na}_2[\text{WO}_4]$  was acidified with HCl and the resultant precipitate was filtered off, washed with distilled water, and dried in an oven at 120 °C.

*Matrix-isolation Spectroscopy.*—All samples were vacuum resublimed immediately before use and were vaporised

† **CAUTION:** The use of large amounts of  $\text{BBr}_3$  or the omission of the ice-bath can result in an explosion.

from Pyrex break-seal ampoules, where necessary equipped with a poly(tetrafluoroethylene) stopcock [Youngs (Acton) Ltd.]. The vapours were co-condensed with the appropriate matrix gas (B.O.C. Special Gases Ltd.) on to a CsI window for i.r. studies, or a LiF window for u.v.-visible studies, attached to a closed-cycle liquid helium refrigerator (Air Products and Chemicals Ltd., CSW-202 Displex) in a high-vacuum ( $10^{-6}$  Torr)\* cryostat. Deposition times varied in the range 5–30 min. For i.r. studies the spray-on temperatures of the samples ( $\text{MoOF}_4$  and  $\text{WOF}_4$ , 25 °C;  $\text{MoOCl}_4$ , 25–40 °C;  $\text{WOCl}_4$ , 60 °C;  $\text{WOBr}_4$ , 300 °C) were higher than those required for u.v.-visible studies (20, 20, 40–50, 120–160 °C) because of the lower absorption coefficients of the i.r. bands. Perkin-Elmer models 225 and 580B spectrometers, for i.r. studies, and a Pye-Unicam SP 1800B spectrometer, for u.v.-visible studies (200–850 nm), were regularly calibrated with the appropriate standard reference materials.

[Figure 1(a)]. However, the degenerate bending mode  $\nu_8$  appeared as a distinct doublet at 309 and 304  $\text{cm}^{-1}$ , and although there was no evidence for dimerisation or sample decomposition, it would appear that the nitrogen matrix environment produced a slight distortion from  $C_{4v}$  symmetry.

The results for  $\text{MoOF}_4$  were also in good agreement with the literature, and entirely consistent with the isolation of monomers. Table 1 compares our matrix frequencies with the gas-phase fundamentals reported by Beattie and co-workers,<sup>11</sup> and the only significant perturbation (under low resolution) is the splitting again observed for  $\nu_8$  in a nitrogen matrix. For this system, however, isotope effects due to the central metal atom are easier to resolve, and the nature of possible 'matrix effects' thus probed in more detail.

TABLE 1

Fundamental modes	Infrared bands ( $\text{cm}^{-1}$ ) and assignments for monomeric $\text{MOX}_4$ species											
	$\text{WOF}_4$			$\text{MoOF}_4$			$\text{WOCl}_4$		$\text{MoOCl}_4$		$\text{WOBr}_4$	
	$\text{N}_2^a$	$\text{Ar}^a$	gas <sup>a</sup>	$\text{N}_2^a$	$\text{Ar}^a$	gas <sup>b</sup>	$\text{N}_2^a$	gas <sup>d</sup>	$\text{N}_2^a$	gas <sup>c</sup>	$\text{N}_2^a$	gas <sup>e</sup>
$\nu_1(A_1)$	1 058 <sup>f</sup>	1 053 <sup>f</sup>	1 055	1 050 <sup>f</sup>	1 045 <sup>f</sup>	1 048	1 032 <sup>f</sup>	1 027	1 017 <sup>f</sup>	1 015	1 025	1 004
$\nu_2(A_1)$	726	729	733	714	719	714	400	405				
$\nu_3(A_1)$	254	245	248	267	258	264						
$\nu_7(E)$	686	690	698	708 <sup>f</sup>	712 <sup>f</sup>	720	380 <sup>f</sup>	383	395 <sup>f</sup>	404	264	265
$\nu_8(E)$	309/304	301	298	309/304	301	294	260 <sup>g</sup>	260 <sup>g</sup>				
$\nu_9(E)$	236	232	236	238	232	236		145 <sup>g</sup>				

Isotope fine structure for $\text{MoOF}_4$ <sup>h</sup>							
Band	<sup>92</sup> Mo	<sup>94</sup> Mo	<sup>95</sup> Mo	<sup>96</sup> Mo	<sup>97</sup> Mo	<sup>98</sup> Mo	<sup>100</sup> Mo
$\nu_1(\text{Mo=O})$ , $\text{N}_2$ matrix	1 052.9	1 051.1	1 050.3	1 049.4	1 048.6	1 047.8	1 046.2
$\nu_1$ , Ar matrix, site A	1 047.9	1 046.0	1 045.2	1 044.4	1 043.6	1 042.8	1 041.3
$\nu_7(\text{Mo-F})$ , site A	716.5	714.4	713.3	712.3	711.3	710.3	708.4

<sup>a</sup> This work: estimated matrix ratios *ca.* 1:1 000. <sup>b</sup> Ref. 11. <sup>c</sup> Refs. 11, 21, and 23. <sup>d</sup> Refs. 24 and 25. <sup>e</sup> Refs. 9 and 23. <sup>f</sup> Centre of complex absorption. <sup>g</sup> Tentative assignment. <sup>h</sup> These isotope frequencies, together with the other observed stretching modes, may be fitted to better than 0.1  $\text{cm}^{-1}$  using several force fields. In particular, the calculated line diagram accompanying Figure 1(b) employs three independent parameters  $F_R = 8.813$ ,  $F_r = 4.318$ , and  $F_{r_{\text{rcis}}} = 0.485$   $\text{mdyn } \text{Å}^{-1}$ , together with the constraints  $F_{r_{\text{trns}}} = 0.5 F_{r_{\text{rcis}}}$  and  $F_{Rr} = -0.1 F_{r_{\text{rcis}}}$  ( $R = \text{Mo-O}$ ,  $r = \text{Mo-F}$ ) ( $1 \text{ dyn} = 10^{-5} \text{ N}$ ).

## RESULTS AND DISCUSSION

**Infrared Studies.**—(a)  $\text{MoOF}_4$  and  $\text{WOF}_4$ . The vibrational spectra previously obtained for these species leave little doubt that in the vapour phase the predominant species is the monomer, with symmetry  $C_{4v}$ . For this structure,  $\Gamma_{\text{vib.}} = 3A_1$  (i.r., R) +  $2B_1$  (R) +  $B_2$  (R) +  $3E$  (i.r., R) and Beattie and co-workers<sup>11</sup> in particular have assigned the *i.r.-active* bands in vapour-phase  $\text{WOF}_4$  as  $A_1$  1 055, 733, and 248  $\text{cm}^{-1}$  and  $E$  698, 298, and 236  $\text{cm}^{-1}$ . These authors also studied  $\text{WOF}_4$  in low-temperature matrices under medium resolution, and noted corresponding bands ( $\text{cm}^{-1}$ ) at 1 057 ( $\text{N}_2$ ), 727 ( $\text{N}_2$ ), 246 (Ar), 691 (Ar), 302 (Ar), and 233 (Ar).

Our spectra in nitrogen and argon matrices were very similar (Table 1), but under high resolution several of our argon-matrix bands were accompanied by weak shoulders separated by up to 5  $\text{cm}^{-1}$  from the principal absorption. In nitrogen, the bands were generally much cleaner, and a tungsten isotope pattern was just discernible on the  $\text{W=O}$  stretching mode at 1 058  $\text{cm}^{-1}$

Figure 1(b) shows a nitrogen-matrix spectrum of the  $\text{Mo=O}$  stretching mode under high resolution. The seven prominent isotopes of molybdenum are well resolved, and the accompanying line diagram indicates the extent of spectral agreement with calculations based on a simple force field (Table 1). In contrast, Figure 1(c) shows the same mode in argon, and it is evident that more than one trapping site must be invoked to interpret this pattern. A similar effect is also observed for the intense  $\text{Mo-F}$  stretching mode at *ca.* 710  $\text{cm}^{-1}$  ( $\nu_7$ ,  $E$ ). Figure 1(d) shows the relevant spectral region under high resolution, and again the observed spectrum is readily interpreted as two overlapping septets. The relative intensities of the two patterns are similar to those associated with the  $\text{Mo=O}$  mode, and this provides very strong evidence for multiple site trapping.<sup>13,14</sup> As the lower-frequency mode is doubly degenerate, these spectra also indicate that both these trapping sites (A and B) have relatively high symmetry, as no splitting of the degeneracy could be detected. In contrast, the nitrogen-matrix spectrum of this region showed a broad band for which we were unable adequately to resolve isotope fine structure. We believe that this is because of a slight

\* Throughout this paper: 1 Torr = 101 325/760 Pa.

lifting of degeneracy, the effects of which are not as pronounced as for  $\nu_6$ , but are sufficient to destroy resolution.

For both  $\text{WOF}_4$  and  $\text{MoOF}_4$  we therefore conclude that in nitrogen matrices there is essentially *one* trapping site of rather low symmetry, and that in argon there are at least two sites, neither of which causes a detectable lifting of degeneracy.

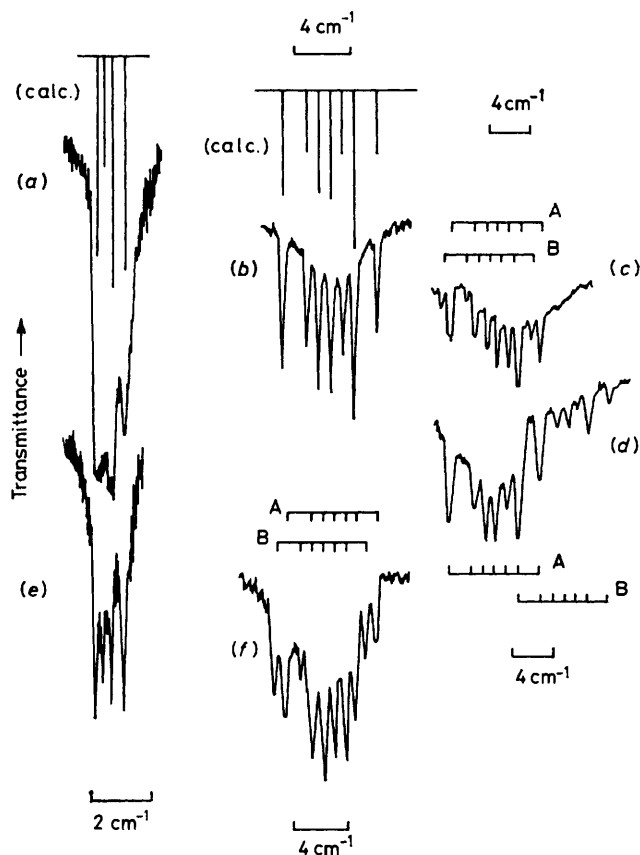


FIGURE 1 Observed and calculated isotope patterns in the i.r. spectra of matrix isolated  $\text{MOX}_4$  species: (a)  $\nu(\text{W}=\text{O})$  in  $\text{WOF}_4$  ( $\text{N}_2$  matrix); (b)  $\nu(\text{Mo}=\text{O})$  in  $\text{MoOF}_4$  ( $\text{N}_2$  matrix); (c)  $\nu(\text{Mo}=\text{O})$  in  $\text{MoOF}_4$  (Ar matrix) showing sites A and B; (d)  $\nu(\text{Mo}-\text{F})$  in  $\text{MoOF}_4$  (Ar matrix) showing sites A and B; (e)  $\nu(\text{W}=\text{O})$  in  $\text{WOF}_4$  ( $\text{N}_2$  matrix); (f)  $\nu(\text{Mo}=\text{O})$  in  $\text{MoOCl}_4$  (Ar matrix) showing sites A and B

(b)  $\text{WOCl}_4$  and  $\text{MoOCl}_4$ . By comparison with the oxyfluorides of tungsten and molybdenum, vibrational data on monomeric  $\text{WOCl}_4$  and  $\text{MoOCl}_4$  are considerably less extensive, although a  $C_{4v}$  structure has been established for both species. Beattie and co-workers<sup>11</sup> studied the Raman and i.r. spectra of the vapours at 180–200 °C, and although  $\text{MoOCl}_4$  was too intensely coloured to yield results, they report Raman bands for  $\text{WOCl}_4$  at 1 027 and 402  $\text{cm}^{-1}$  and i.r. bands at 1 027, 404vw, 383s, and 358w  $\text{cm}^{-1}$ . Bues *et al.*<sup>21</sup> have obtained Raman data for  $\text{WOCl}_4$  in all three phases, and their assignments were subsequently used by Brockner and Hovdan<sup>22</sup> to calculate the vibration frequencies of the various isotopic  $\text{WOCl}_4$  monomers. Unfortunately,

these calculations were based on fundamentals obtained from different phases, and this led to a tabulation of i.r.-active ( $A_1 + E$ ) modes for  $\text{WOCl}_4$  of  $A_1$  1 028.0, 407.6, and 161.7  $\text{cm}^{-1}$  and  $E$  337.4, 265.4, and 149.9  $\text{cm}^{-1}$ . These frequencies agree quite well with the four vapour-phase bands reported by Bues, at 1 025w, 407s, ca. 260w, and ca. 145w  $\text{cm}^{-1}$ , but the prediction of a highest  $E$  mode at 337.4  $\text{cm}^{-1}$  is considerably lower than the value (ca. 383  $\text{cm}^{-1}$ ) found experimentally by Beattie.<sup>11</sup> The position of this mode ( $\nu_7$ ) has subsequently been confirmed by Mal'tsev and co-workers,<sup>23</sup> in an independent i.r. study of the vapour, where they found, in addition, absorptions at 305 and 254  $\text{cm}^{-1}$  which they attribute also to  $\text{WOCl}_4$ . A combination of these results thus leads to the following assignment for the i.r.-active fundamentals of  $\text{WOCl}_4$ :  $A_1$  ca. 1 027 and ca. 405  $\text{cm}^{-1}$ ;  $E$  383  $\text{cm}^{-1}$ , with the weaker features at ca. 260 and 145  $\text{cm}^{-1}$  being assigned either as the two remaining  $E$  modes or as  $A_1$  and  $E$ .

Our nitrogen-matrix spectra for this system (Table 1) gave strong bands centred at ca. 1 032 and ca. 380  $\text{cm}^{-1}$ , both of which showed fine structure, and two much weaker features at ca. 400 and ca. 260  $\text{cm}^{-1}$ , whilst in argon a similar pattern of fundamentals was observed with additional site splittings. Figure 1(e) shows the multiplet structure on the band at 1 032  $\text{cm}^{-1}$  ( $\text{N}_2$  matrix) which corresponds closely to the pattern expected for a tungsten isotope effect. (The effect of different *chlorine* isotopes on this mode has been shown to be negligible.<sup>22</sup>) In the region of 380  $\text{cm}^{-1}$ , the nitrogen-matrix spectrum shows four prominent bands under high resolution, at 384.2, 379.3, 375.2, and 371.4  $\text{cm}^{-1}$ . However, the separation between these bands is too large to be attributed solely to a chlorine isotope effect,<sup>21</sup> and we believe that this fine structure probably arises from the removal of the degeneracy. The absence of other spectral features leaves little doubt that monomeric  $\text{WOCl}_4$  has been isolated.

There are even less published data on monomeric  $\text{MoOCl}_4$ . Barraclough and Kew<sup>24</sup> report  $\nu(\text{Mo}=\text{O})$  ( $A_1$ ) at 1 015  $\text{cm}^{-1}$ , whilst unpublished observations by Bukovsky<sup>25</sup> confirm this value and locate  $\nu_7(E)$  at ca. 404  $\text{cm}^{-1}$  for the vapour-phase molecule. Our matrix experiments yielded a deceptively simple spectrum under low resolution, and in nitrogen we observed two strong features at ca. 1 017 and 395  $\text{cm}^{-1}$  with corresponding bands in argon at ca. 1 012 and ca. 399  $\text{cm}^{-1}$ . However, under high resolution, it was evident that neither matrix provided a unique trapping site. Figure 1(f) shows the  $\text{Mo}=\text{O}$  stretching mode under high resolution in an argon matrix and the complicated pattern observed is rationalised by invoking two overlapping molybdenum isotope septets. The intense lower-frequency feature was similarly complex under high resolution, and typically comprised ca. 10 poorly resolved components separated by ca. 2  $\text{cm}^{-1}$ . No attempt was made to analyse this pattern in view of the abundance of molybdenum and chlorine isotopomers, and the existence of multiple sites, but despite these complications,

we are confident that the spectra show that monomeric  $\text{MoOCl}_4$  has been isolated in Ar and  $\text{N}_2$  matrices.

(c)  $\text{WOBBr}_4$ . Spectroscopic data on monomeric  $\text{WOBBr}_4$  are also sparse. Ward and Stafford,<sup>9</sup> working in the KBr region ( $>400\text{ cm}^{-1}$ ), assigned a prominent band at  $1\,004\text{ cm}^{-1}$  to this molecule in their study of high-temperature oxohalides of molybdenum and tungsten, whilst Mal'tsev and co-workers<sup>23</sup> reported an intense low-frequency band at  $265\text{ cm}^{-1}$  in their far-i.r. studies ( $200\text{--}500\text{ cm}^{-1}$ ). Our nitrogen-matrix spectra obtained from purified samples of  $\text{WOBBr}_4$  gave only two bands (Table 1). The most intense feature occurred at  $264\text{ cm}^{-1}$  and the second band was observed at  $1\,025\text{ cm}^{-1}$ . Both bands were relatively narrow (*ca.*  $3\text{ cm}^{-1}$ ) and there were no obvious site effects. Attempts to resolve the tungsten isotope pattern on the higher-frequency band resulted in spectra of slightly poorer quality than those previously obtained for  $\text{WOF}_4$  (and  $\text{WOCl}_4$ ). However, there seems little doubt that monomeric  $\text{WOBBr}_4$  has also been satisfactorily isolated.

*Ultraviolet-Visible Studies.*—(a) *Spectra.* Previous attempts<sup>15,16</sup> to obtain electronic spectra of  $\text{MOX}_4$  compounds in the gas phase have been fraught with uncertainties concerning disproportionation because of the high temperatures ( $>80\text{ }^\circ\text{C}$ ) needed to achieve a sufficient vapour pressure of the compounds to measure the spectra. For example, although mass spectrometric studies indicate that  $\text{MoOCl}_4$ ,<sup>12,24</sup>  $\text{WOCl}_4$ ,<sup>12,24</sup> and  $\text{WOBBr}_4$ <sup>12,25,26</sup> vaporise as monomeric molecules, other studies indicate that  $\text{MoOCl}_4$ <sup>27,28</sup> dissociates as  $2\text{MoOCl}_4 \rightarrow 2\text{MoOCl}_3 + \text{Cl}_2$ , and that  $\text{WOCl}_4$ <sup>29-32</sup> disproportionates in the vapour phase above  $200\text{ }^\circ\text{C}$  with the formation of  $\text{WO}_2\text{Cl}_2$  and  $\text{WCl}_6$ . Some of the bands appearing in the gas-phase spectra above  $\text{MoOCl}_4$  and  $\text{WOCl}_4$  were therefore assigned<sup>15,16</sup> to  $\text{Cl}_2$  ( $\text{MoOCl}_4$ ) and  $\text{WCl}_6$  ( $\text{WOCl}_4$ ). In view of the weak association of  $\text{MoOCl}_4$  molecules as dimers *via* chlorine bridges,<sup>4</sup> it seemed possible that the reported<sup>17</sup> solution spectrum of  $\text{MoOCl}_4$  was not of monomeric  $\text{MoOCl}_4$  molecules. Experiments to obtain the electronic spectra of  $\text{MOX}_4$  molecules using the matrix-isolation technique do not require high temperature for vaporisation (see Experimental section); in fact lower vaporisation temperatures were required than for matrix i.r. spectra where isolation of monomeric species has been unambiguously established (see above).

The electronic spectra of matrix-isolated monomeric  $\text{MOX}_4$  molecules [ $\text{N}_2$  matrices,\* *ca.*  $1:2\,000$ ; Figure 2(a)—(e)] show a large number of bands, some with resolved vibrational fine structure. Band maxima positions † and vibrational spacings are given in Table 2. Unlike  $\text{CrO}_2\text{Cl}_2$ ,<sup>33-35</sup>  $\text{VOCl}_3$ ,<sup>35</sup> and  $\text{MnO}_3\text{Cl}$ <sup>35</sup> for which luminescence spectra have been obtained for the matrix-isolated molecules, no luminescence was detected on

\* Nitrogen was used as a matrix because of the i.r. observations of single trapping sites.

† Absorption coefficients could not be measured because it was not possible to determine the amount of substrate in the beam of the spectrometer.

exciting any of the  $\text{MOX}_4$  molecules in any of their electronic absorption bands.

(b) *Assignment of spectra.* Assignment of the electronic spectra can be attempted using molecular-orbital calculations<sup>15,16,36</sup> and from calculations of charge-transfer energies *via* optical electronegativities.<sup>37</sup>

TABLE 2

Observed <sup>a</sup> band positions [ $\lambda_{\text{max.}}/\text{cm}^{-1}$  (nm)] for  $\text{MOX}_4$  compounds isolated in  $\text{N}_2$  matrices at 12 K

Compound	Band position	Tentative assignment <sup>b</sup>	
$\text{MoOF}_4$	44 250 (226.0)		
	39 150 (255.4) <sup>c</sup>	c.t. (F→Mo)	
$\text{WOF}_4$	39 560 (252.8) <sup>d</sup>	c.t. (F→W)	
	48 780 (205.0)	c.t. (Cl→Mo)	
$\text{MoOCl}_4$ <sup>e</sup>	36 630 (273.0)	c.t. (Cl→Mo)	
	31 250 (320.0)	c.t. (O→Mo)	
	20 790 (481.0)	c.t. (Cl→Mo)	
	15 770 (634.0) <sup>f</sup>	c.t. (Cl→Mo)	
	15 380 (650.0) <sup>f</sup>	c.t. (Cl→Mo)	
	$\text{WOCl}_4$ <sup>g</sup>	44 840 (223.0)	c.t. (Cl→W)
		38 870 (257.3) <sup>h</sup>	c.t. (Cl→W)
		37 040 (270.0)	c.t. (O→W)
		33 330 (300.0)	c.t. (Cl→W)
	$\text{WOBBr}_4$ <sup>i</sup>	28 270 (353.7)	c.t. (Cl→W)
26 320 (380.0)		c.t. (Cl→W)	
22 400 (446.5)		c.t. (Cl→W)	
43 480 (230.0)		c.t. (Br→W)	
37 880 (264.0)		c.t. (O→W)	
35 460 (282.0)		c.t. (Br→W)	
33 110 (302.0)		c.t. (Br→W)	
31 250 (230.0)		c.t. (Br→W)	
29 410 (340.0)		c.t. (Br→W)	
25 970 (385.0)		c.t. (Br→W)	
$\text{WOBBr}_4$ <sup>i</sup>	20 830 (480.0)	c.t. (Br→W)	
	17 830 (561.0)	c.t. (Br→W)	
	15 110 (662.0)	c.t. (Br→W)	

<sup>a</sup> Wavenumber positions obtained from spectra measured in wavelength ( $\pm 0.2\text{ nm}$ ). <sup>b</sup> See text, Table 3, and Figure 4; c.t. = charge transfer. <sup>c</sup> Vibrational progression: 41 390, 40 540, 39 870, 39 150, 38 580, 37 950, 37 370, 36 710, 36 100, 35 390, 34 690, 33 990, and 33 330  $\text{cm}^{-1}$ ; *i.e.*  $\nu_{\text{av.}} = 670 \pm 70\text{ cm}^{-1}$ . <sup>d</sup> Vibrational progression: 46 320, 45 600, 44 980, 44 390, 43 800, 43 100, 42 430, 41 700, 41 030, 40 180, 39 530, 38 790, and 38 120  $\text{cm}^{-1}$ ; *i.e.*  $\nu_{\text{av.}} = 685 \pm 70\text{ cm}^{-1}$ . <sup>e</sup> Band positions in  $\text{CCl}_4$  solution: 37 200, 20 800, and 14 400  $\text{cm}^{-1}$ ; data from ref. 17. Band positions in gas phase: 37 040, 21 280, and *ca.* 15 400  $\text{cm}^{-1}$ ; data from ref. 15. <sup>f</sup> Treating this splitting (see text) as two resolved peaks of a vibrational progression would give a value of  $\nu' = 440 \pm 50\text{ cm}^{-1}$ . <sup>g</sup> Band positions in gas phase: 44 400, 27 800, and 20 830  $\text{cm}^{-1}$ ; data from ref. 15. <sup>h</sup> Vibrational progression: 39 450, 39 100, 38 870, 38 540, 38 240, 37 910, and 37 660  $\text{cm}^{-1}$ ; *i.e.*  $\nu' = 300 \pm 70\text{ cm}^{-1}$ . <sup>i</sup> Band positions in gas phase: 35 000, 22 000, 17 000, and 15 000  $\text{cm}^{-1}$ ; data from ref. 16.

The  $X\alpha$  calculations<sup>36</sup> of the electronic energy levels in the  $d^1$  ions  $[\text{CrOCl}_4]^-$ ,  $[\text{MoOCl}_4]^-$ , and  $[\text{MoOBr}_4]^-$  suggest that the highest occupied orbital is the predominantly metal  $d_{xy}$  orbital ( $2b_2$ ) and that below this are several levels ( $1a_2$ ,  $6e$ ,  $3b_1$ ,  $5e$ ,  $7a_1$ ) which are essentially non-bonding since they are localised on the halogens. For the  $d^0$   $\text{MOX}_4$  compounds the  $2b_2$  level will be empty and all the absorption bands must involve ligand→metal charge-transfer transitions. The SCF- $X\alpha$ -SW calculation<sup>15,16</sup> of the orbital populations and orbital energies for  $\text{MoOCl}_4$ ,  $\text{WOCl}_4$ , and  $\text{WOBBr}_4$  is in agreement with the relative ordering of energy levels for the  $d^1$  ions, although the  $C_{4v}$  symmetry designation of levels is different.

Figure 3 gives a tentative assignment of energy levels.

A consideration<sup>37</sup> of the charge-transfer energies and optical electronegativities for a range of related oxo-compounds of vanadium, molybdenum, tungsten, and

progressions on the lowest-energy charge-transfer absorption bands and emission bands of  $\text{CrO}_2\text{Cl}_2$ .<sup>33-35</sup>

The lowest-energy charge-transfer bands for  $\text{MoOF}_4$  and  $\text{WOF}_4$  (Figure 2, Table 2) are predicted to be  $\text{O} \rightarrow \text{M}$  bands on the basis of optical electronegativities<sup>37</sup> (Table

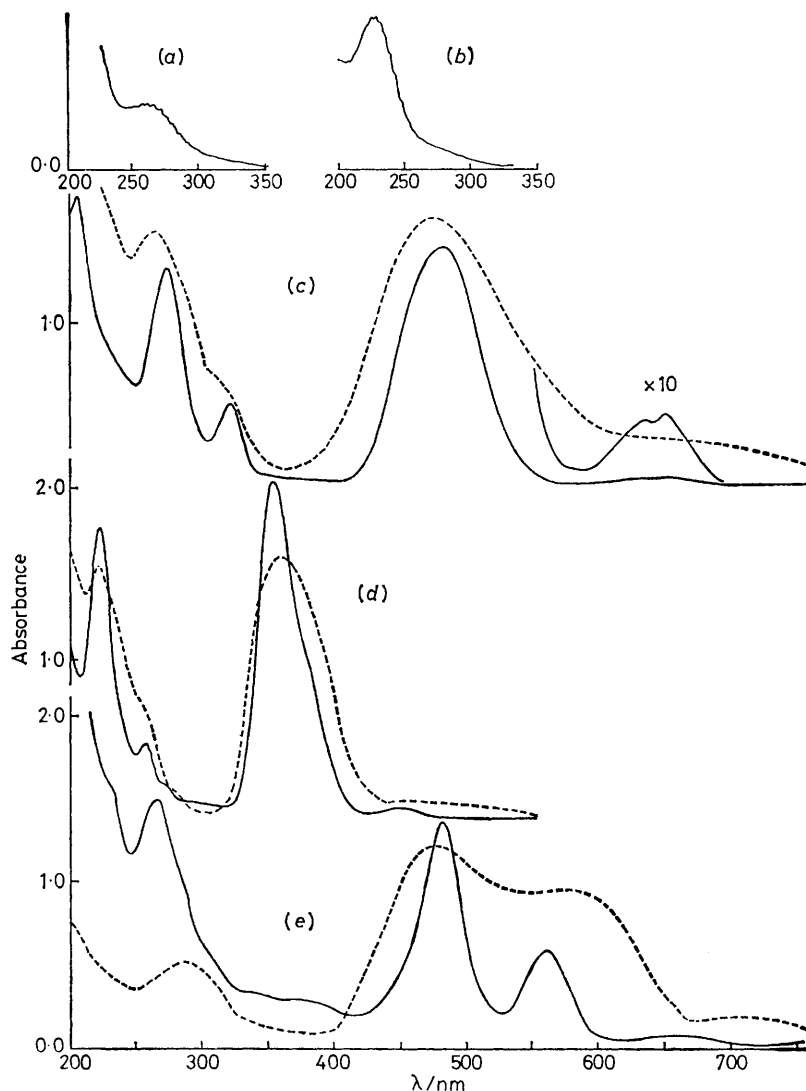


FIGURE 2. Electronic absorption spectra (—) for monomeric  $\text{MOX}_4$  compounds isolated in  $\text{N}_2$  matrices at 12 K: (a)  $\text{MoOF}_4$ , (b)  $\text{WOF}_4$ , (c)  $\text{MoOCl}_4$ , (d)  $\text{WOCl}_4$ , and (e)  $\text{WOBr}_4$ . Previously obtained gas-phase spectra (---) are included where these have been obtained (refs. 15 and 16)

niobium allows the calculation of the lowest-energy  $\text{O} \rightarrow \text{M}$  and  $\text{X} \rightarrow \text{M}$  ( $\text{X} = \text{halogen}$ ) charge-transfer bands (Table 3). The charge-transfer bands would be expected to show vibrational progressions because the transfer of charge will alter the equilibrium geometry of the excited state with respect to the ground state. These progressions should be particularly informative in assigning bands because they will involve the excited state  $A_1$  metal-halogen [ $\nu(\text{M-X})'$ ] and metal-oxygen [ $\nu(\text{M-O})'$ ] stretching modes and also the  $A_1$  metal-halogen [ $\delta(\text{M-X})'$ ] deformation modes, cf. the  $\nu(\text{Cr-Cl})$  and  $\delta(\text{Cr-Cl})$

3) and the ionisation and appearance potentials of ions observed by mass spectrometry.<sup>12</sup> However, the lowest-energy bands for  $\text{MoOF}_4$  and  $\text{WOF}_4$  show vibrational progressions ( $\nu_{\text{Mo-O}}' = 670 \pm 70$ ,  $\nu_{\text{W-O}}' = 685 \pm 70$   $\text{cm}^{-1}$ ) which are more comparable with  $A_1 \nu(\text{Mo-F})$  (714  $\text{cm}^{-1}$ ) and  $A_1 \nu(\text{W-F})$  (726  $\text{cm}^{-1}$ ) than with  $A_1 \nu(\text{Mo-O})$  (1 050  $\text{cm}^{-1}$ ) and  $A_1 \nu(\text{W-O})$  (1 058  $\text{cm}^{-1}$ ) vibrational modes. Furthermore, the absorption spectrum of  $\text{MoF}_6$  shows a vibrational progression of 639  $\text{cm}^{-1}$  which has been assigned to the  $A_{1g} \nu(\text{Mo-F})'$  vibration in the excited state corresponding to a wavenumber of 736  $\text{cm}^{-1}$  in the

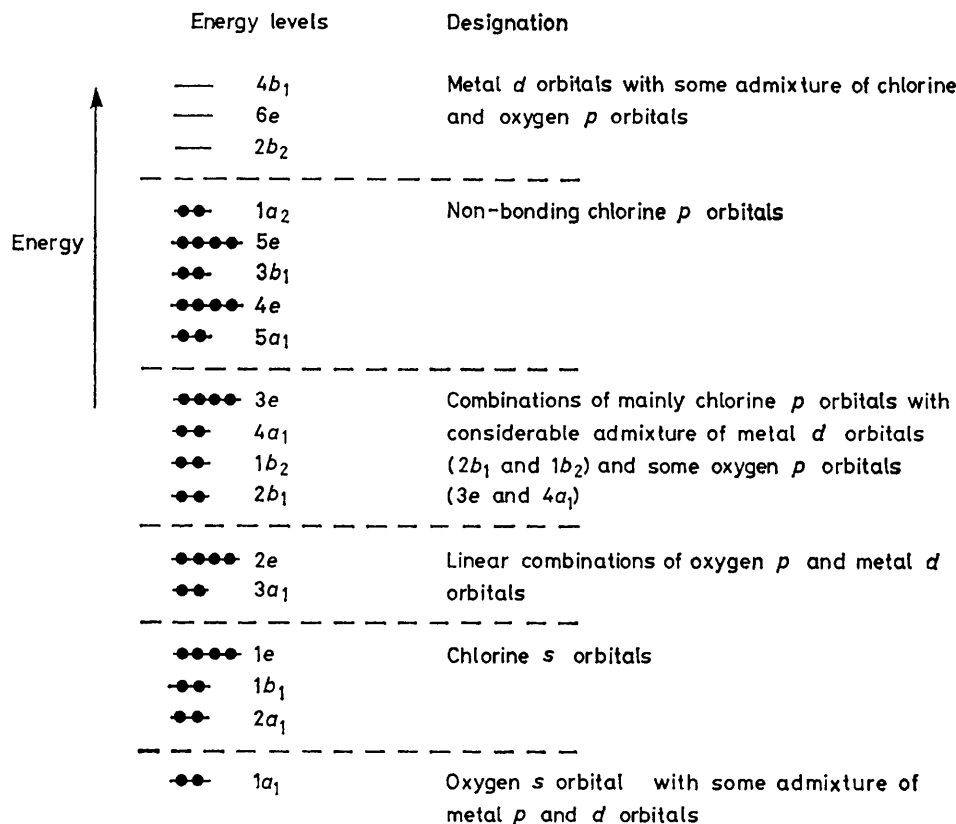


FIGURE 3 Schematic molecular-orbital energy diagram (based on refs. 15 and 16) for  $\text{MOX}_4$  molecules showing occupancy of energy levels (● = electron) and their nature and  $C_{4v}$  symmetry designations

TABLE 3

Calculated <sup>a</sup> band positions ( $\text{cm}^{-1}$ ) for the lowest-energy  $\text{O} \rightarrow \text{M}$  and  $\text{X} \rightarrow \text{M}$  charge-transfer (c.t.) transitions of  $\text{MOX}_4$  compounds

Compound	c.t. ( $\text{O} \rightarrow \text{M}$ )	c.t. ( $\text{X} \rightarrow \text{M}$ )
$\text{MoOF}_4$ <sup>b</sup>	$31\,000 \pm 2\,000$	$42\,000 \pm 2\,000$
$\text{WOF}_4$ <sup>b</sup>	$38\,000 \pm 2\,000$	$49\,000 \pm 2\,000$
$\text{MoOCl}_4$	$31\,000 \pm 2\,000$	$15\,000 \pm 2\,000$
$\text{WOCl}_4$	$38\,000 \pm 2\,000$	$22\,000 \pm 2\,000$
$\text{WOBr}_4$	$38\,000 \pm 2\,000$	$14\,000 \pm 2\,000$

<sup>a</sup> Ref. 37. <sup>b</sup> Band positions obtained by using an optical electronegativity of 3.9 for  $\text{F}^-$  (C. K. Jorgensen, *Struct. Bonding (Berlin)*, 1966, 1, 3) in the expressions given in ref. 37.

ground state.<sup>38</sup> It seems reasonable, therefore, to assign the lowest-energy charge-transfer band to the  $5e \rightarrow 2b_2$  fluorine  $\rightarrow$  metal transition (Figure 3), analogous to the other  $\text{MOX}_4$  compounds (see below), rather than to an oxygen  $\rightarrow$  metal transition. Some inconsistencies have been noted previously for oxyfluorides,<sup>37</sup> e.g. while the lowest charge-transfer band of  $\text{VOF}_3$  ( $25\,700\text{ cm}^{-1}$ )<sup>39</sup> agrees with the calculated value ( $26\,000\text{ cm}^{-1}$ ),  $[\text{VOF}_5]^{2-}$  shows the lowest-energy band at a considerably higher energy ( $31\,000\text{ cm}^{-1}$ ).<sup>39</sup> Further work on other oxyfluorides seems to be called for.

The lowest-energy bands for  $\text{MoOCl}_4$ ,  $\text{WOCl}_4$ , and  $\text{WOBr}_4$  (Figure 2, Table 2) show a good correlation with the gas-phase bands [Figure 2 (c)–(e)] and the positions

of the bands predicted for charge transfer ( $\text{X} \rightarrow \text{M}$ ) using optical electronegativities (Table 3). The set of bands for  $\text{MoOCl}_4$  in a  $\text{N}_2$  matrix at 12 K also shows reasonable agreement with those obtained previously for the compound in a  $\text{CCl}_4$  solution,<sup>17</sup> and hence confirms that isolated  $\text{MoOCl}_4$  molecules occur in  $\text{CCl}_4$  solution. This solvent shifts the lowest-energy band to higher wavenumbers and may cause band broadening so that matrix bands observed at  $36\,630$  and  $31\,250\text{ cm}^{-1}$  appear as a single band at  $37\,200\text{ cm}^{-1}$  for  $\text{CCl}_4$  solution. With the exception of the band of  $\text{WOCl}_4$  at  $38\,870\text{ cm}^{-1}$  and possibly the band \* of  $\text{MoOCl}_4$ , at  $15\,500\text{ cm}^{-1}$ , the lack of vibrational fine structure on the bands prevents definitive assignments. However, an assignment of the bands should be possible using the ordering of energy levels in Figure 3 and the resulting transitions.<sup>15,16</sup> Surprisingly, the SCF- $X\alpha$ -SW calculations predict the first pure  $\text{O} \rightarrow \text{M}$  charge-transfer band to occur at  $>50\,000\text{ cm}^{-1}$  (from  $2e$  and  $3a_1$ ) while the lowest-energy band which *might* involve orbitals with some oxygen  $p$  character ( $3e$  and  $4a_1$ ) is assigned at *ca.*  $22\,000\text{ cm}^{-1}$  for  $\text{MoOCl}_4$  ( $3e \rightarrow 2b_2$ )<sup>15</sup> and *ca.*  $28\,000\text{ cm}^{-1}$  for  $\text{WOCl}_4$  ( $3e \rightarrow 2b_2$ ).<sup>15</sup> The large number of bands predicted in the gas-phase studies and their assignment in clusters to broad bands is dis-

\* The doublet splitting could be a vibrational progression (Table 2) or could arise from two new degenerate transitions or spin-orbit coupling.

appointingly uninformative. On the basis of optical electronegativities, the lowest-energy O→M charge-transfer bands are tentatively assigned at 31 250 cm<sup>-1</sup> for MoOCl<sub>4</sub> (31 000 ± 2 000 cm<sup>-1</sup>, calculated), 37 040 cm<sup>-1</sup> for WOCl<sub>4</sub> (38 000 ± 2 000, calculated), and 37 880 cm<sup>-1</sup> for WOBr<sub>4</sub> (38 000 ± 2 000, calculated). Some confidence in the optical electronegativity approach can be drawn from the close agreement<sup>37</sup> of the lowest-energy absorption bands of [MOX<sub>4</sub>]<sup>-</sup> ions, e.g. [MoOCl<sub>4</sub>]<sup>-</sup> [27 000 (calculated) and 26 000–26 700 cm<sup>-1</sup> (observed)], [MoOBr<sub>4</sub>]<sup>-</sup> [19 000 (calculated) and 20 600 cm<sup>-1</sup> (observed)], and [WOBr<sub>4</sub>]<sup>-</sup> [24 000 (calculated) and 24 700 cm<sup>-1</sup> (observed)]. A firm assignment of the band of WOCl<sub>4</sub> at 38 870 cm<sup>-1</sup> as a charge-transfer (Cl→W) band can be made on the basis of the vibrational progression (300 ± 70 cm<sup>-1</sup>) which may be correlated with the ground-state (*E*) mode (260 cm<sup>-1</sup>), or alternatively, with the A<sub>1</sub> ν(W-Cl) vibrational mode (400 cm<sup>-1</sup>). All other assignments (Table 2) must be tentative. Considering the large number of X→M charge-transfer bands it is perhaps not surprising that excitation does not lead to luminescence.

**Conclusions.**—The i.r. spectra of the matrix-isolated monomer MOX<sub>4</sub> species show no detectable amounts of decomposition products or polymers. The literature assignments of the vibrational fundamentals of these monomers have been generally confirmed and certain details clarified. High-resolution i.r. spectra, utilizing the isotopic fine structure, have been used to probe matrix site effects, and it has been shown that whilst argon often produces two trapping sites, nitrogen produces one site of rather low symmetry.

The electronic spectra of the matrix-isolated MOX<sub>4</sub> in nitrogen have been recorded, the results being superior to gas-phase spectra in terms of resolution, greater ease of observation of vibrational fine structure, and most importantly when taken in *parallel* with the i.r. studies, the absence of any doubts about the presence of decomposition products. A comparison of the experimental data with the results of recent *Xα* calculations shows that, although these in principle give information about the number and relative energies of the orbitals, the numerical agreement is poor, and the much simpler concept of optical electronegativity provides a better estimate of the energies of the lowest O→M and X→M charge-transfer bands.

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