

Location of Cation Sites in a Hydrated Copper-exchanged Chabazite by Extended X-Ray Absorption Fine Structure†

Gianni Antonioli

Dipartimento di Fisica, Università di Parma, 43100 Parma, Italy

Gilberto Vlaic

Dipartimento di Scienze Chimiche, Università di Cagliari, 09124 Cagliari, Italy

Giorgio Nardin and Lucio Randaccio*

Dipartimento di Scienze Chimiche, Università di Trieste, 34127 Trieste, Italy

The local structure around copper ions in a hydrated copper-exchanged chabazite has been characterized by means of extended X-ray absorption fine structure spectroscopy. The model proposed on the basis of the experimental results indicates that copper ions occupy two sites, one with a distorted octahedral co-ordination, the other with a square-pyramidal geometry. The previously suggested hypothesis that copper ions are located in the large cavity of the chabazite is confirmed, although a different co-ordination for one site was detected.

Location of cations in hydrated zeolites is generally a hard task often because of fractional occupancies of cations and water molecules. As a consequence, X-ray analysis sometimes gives controversial results. A typical example is furnished by the crystal structure analysis of a hydrated Cu²⁺-exchanged chabazite, reported some years ago by Smith and co-workers.¹ The interpretation of the electron density in the large cage of the chabazite was uncertain, showing however that all Cu²⁺ and H₂O molecules are present in those cages. These authors concluded that further speculation is futile, but the placement of Cu²⁺ in two sites of approximately equal occupancies, one in a near-octahedral environment, the other in a pentagonal bipyramid, is not crystal-chemically unreasonable.¹ More reliable information about the local co-ordination geometry of cation sites could be obtained by extended X-ray absorption fine structure (EXAFS) measurements. This technique provides a description of the short-range order of selected atomic species in terms of the number of neighbours, distances, and thermal and static disorder within a range of those distances without the need for long-range order.

Experimental

Sample Preparation.—A sample of a natural calcium-rich chabazite, of starting composition Ca_{1.4}Sr_{0.3}Al_{3.8}Si_{8.3}O₂₄·13H₂O from N.E. Azerbaijan, Iran,² was ion exchanged in 1 mol dm⁻³ NaCl solution and then with 0.01 mol dm⁻³ CuCl₂ solution, following the procedure of ref. 1. Chemical analysis of the pale blue-green powder showed complete exchange by copper.

EXAFS Spectroscopy and Data Reduction.—The EXAFS experiment consisted in the measurement of the X-ray absorption coefficient as a function of the X-ray energy. Structural information was deduced by relating the experimental modulation of μx (where μ is the absorption co-efficient and x the thickness of the sample) above the absorption edge E_0 , $\chi(k)$, to a theoretical description in the single-electron, single-scattering, plane-wave approximation (1),³ where k is the wave

$$\chi(k) = \frac{1}{k} \sum_i \frac{N_i}{R_i^2} F_i(k) e^{-2\sigma_i^2 k^2} \sin[2R_i + \Phi_i(k)] \quad (1)$$

vector of the ejected photoelectron, N_i the number of

neighbours at a distance R_i from the absorbing atomic species, $F_i(k)$ the amplitude of back scattering, σ_i a Debye-Waller-like term, and $\Phi_i(k)$ the total phase shift experienced by the photoelectron for each co-ordination shell i ; $\Phi_i(k)$ is characteristic of a particular combination of absorbing atom and back-scattering shell, whereas $F_i(k)$ is almost entirely dependent on the back-scattering shell. The last two terms mentioned can be calculated theoretically⁴ or derived experimentally from reference compounds of known crystal structure. The latter procedure was used here. Analysis of the EXAFS data was carried out according to a standard procedure of Fourier filtering and fitting described elsewhere.⁵

The EXAFS spectra of CuO, CuSO₄·5H₂O, and the chabazite were measured. Known amounts of finely powdered samples were deposited on nucleopore membranes in order to obtain a very homogeneous film of thickness x and to have $\mu x \approx 2$ after the photoelectron excitation. The measurements were performed at LURE (French synchrotron national facility, Orsay) at the station EXAFS-I, installed on the DCI storage ring, working at 1.85 GeV, a typical current being 200 mA, in the classic transmission mode. A Si 311 channel-cut monochromator was used in order to reject harmonics. All EXAFS spectra were collected in the energy range 8 700–9 800 eV, $\Delta E = 2$ eV. Each point was measured for 1.5 s. All spectra were measured three times, in air at room temperature.

Results and Discussion

Figure 1 shows the EXAFS spectra for the reference compound CuO, CuSO₄·5H₂O, and for the chabazite. The structure of CuO is characterized by a first shell of four oxygen atoms around each Cu²⁺ ion, nearly in a plane at the corners of a square at a distance of 1.95 Å.⁶ The k^3 multiplied spectra were then Fourier-transformed in the range 48–650 eV after application of a Gaussian window. The moduli which represent the radial distribution function in R space are shown in Figure 2, where it should be noted that the peaks are shifted towards lower values of R than the real distances because of the k dependency of the function $\Phi(k)$. Only one peak appears in the Fourier transform of both CuSO₄·5H₂O and the chabazite, which is easily attributed to Cu–O atom pairs. The amplitude and phase functions were calculated from the Fourier-filtered first peak of CuO.

† Non-S.I. unit employed: eV $\approx 1.60 \times 10^{-19}$ J.

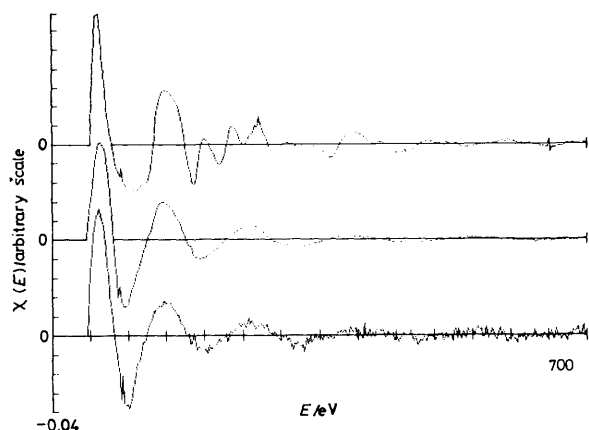


Figure 1. Experimental EXAFS signals for (top to bottom) CuO, $\text{CuSO}_4 \cdot 5\text{H}_2\text{O}$, and the chabazite as $\chi(E)$ versus E

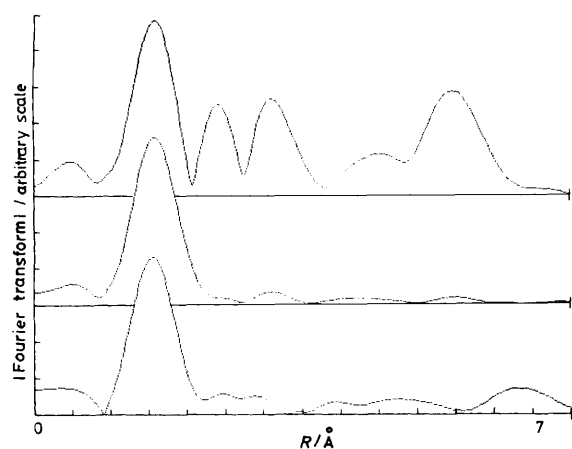


Figure 2. Moduli of the Fourier transforms for (top to bottom) CuO, $\text{CuSO}_4 \cdot 5\text{H}_2\text{O}$, and the chabazite

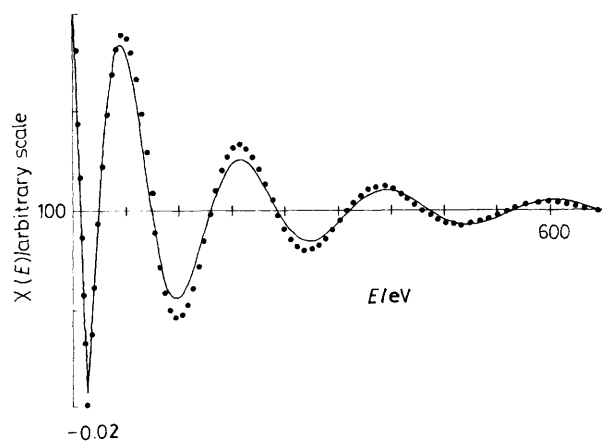


Figure 3. One-shell fit (continuous line) of the Fourier-filtered peak (dotted line) for $\text{CuSO}_4 \cdot 5\text{H}_2\text{O}$

The reliability of these functions was checked by transferring them to the Fourier-filtered peak of the $\text{CuSO}_4 \cdot 5\text{H}_2\text{O}$ sample, considered as a structural unknown. In this compound, copper is surrounded by a tetragonally distorted octahedron of oxygen atoms⁷ from water molecules (four oxygens) and from sulphate ions (distances: 4×1.97 and 2×2.40 Å). Since Joyner,⁸ studying the same sample, concluded that EXAFS is not able to

Table. Results of the one- (a) and two-shells (b) fits for the three samples; N is the number of O atoms at a distance R from Cu

Sample	N	R (Å)	$\Delta\sigma$ (Å)	E_0 (eV)	ΔE_0 (eV)
CuO	4	1.95	—	8 988	—
$\text{CuSO}_4 \cdot 5\text{H}_2\text{O}$	(a) 3.5	1.95	0.000	8 991	-3.0
	(b) { 4.2 1.8	{ 1.96 2.30	{ 0.020 0.075	8 991	-2.5
Chabazite	(a) 3.2	1.95	0.002	8 988	-3.6
	(b) { 3.9 1.5	{ 1.96 2.25	{ 0.002 0.080	8 988	-1.4

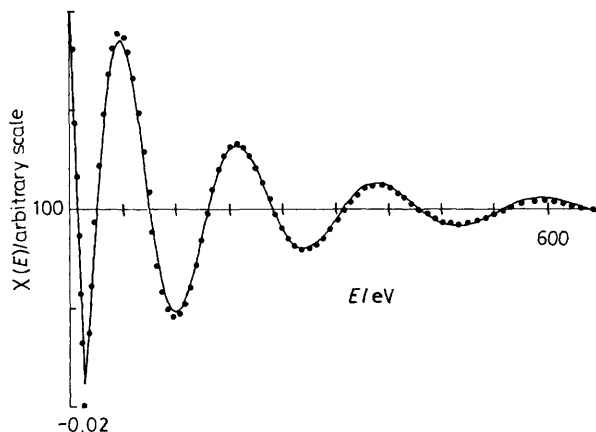


Figure 4. Two-shells fit (continuous line) of the Fourier-filtered peak (dotted line) for $\text{CuSO}_4 \cdot 5\text{H}_2\text{O}$

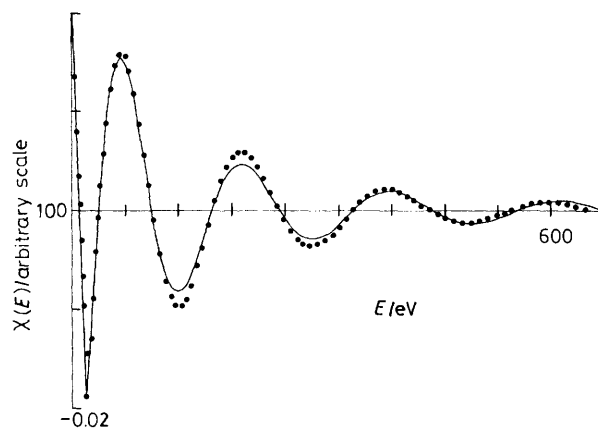


Figure 5. One-shell fit (continuous line) of the Fourier-filtered peak (dotted line) for the chabazite

reveal the longest distance, a least-squares fitting of the signal assuming only one detectable distance was carried out.

For the fit a variation of ΔE_0 was introduced, to take into account any approximation in the chemical transferability of amplitude and phase functions (E_0 was determined as the inflection point of the edge jump). Furthermore, σ has only a relative value $\Delta\sigma$, compared to that of the reference compound, where the term σ remains convoluted with the $F(k)$ function during the extraction of the amplitude function. The fit is reported in Figure 3, and the values obtained are listed in the Table (a). As can be seen, the fit is of poor quality: the frequency is reproduced, but not the amplitude. At this point, a second shell was introduced in the calculation. The results are summarized in the Table (b) and the fit is shown in Figure 4. A significant improvement is apparent and the numerical values

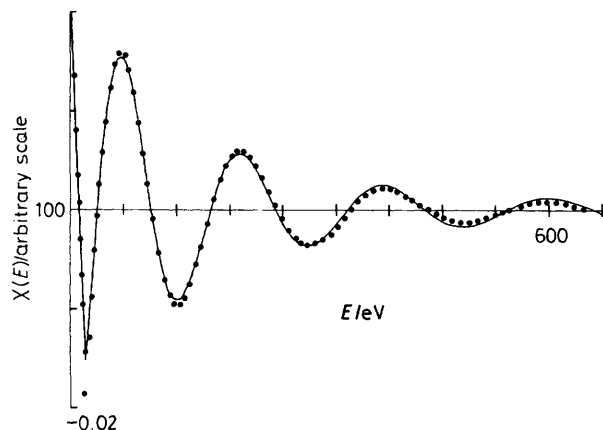


Figure 6. Two-shells fit (continuous line) of the Fourier-filtered peak (dotted line) for the chabazite

are very consistent with known crystallographic data. A shortening of the second distance (about 0.1 Å) deduced by EXAFS, as compared with the crystallographic one, is observed. This is in line with the observations of Crozier and Seary⁹ about the shortening of distances seen by EXAFS with respect to crystallographic ones when the parameter σ is high, as in our case ($\Delta\sigma = 0.075$ Å).

Similar calculations were performed on the zeolite signals. The one-shell fit is reported in Figure 5, while a two-shell fit is shown in Figure 6. The values are summarized in the Table. The analysis clearly shows that the experimental results are in agreement with a model where the Cu^{2+} ions are placed in two

different environments, one represented by a tetragonally distorted octahedron, the other by an axially elongated square pyramid. The ratio between these two environments is about 1:1. This model may be interpreted by assuming that Cu^{2+} ions occupy one site, half of them with a distorted octahedral geometry, the remainder with a pyramidal geometry. Alternatively, Cu^{2+} ions occupy two different sites each with one of the two observed co-ordination geometries. Comparison of these results with those of the X-ray crystal structure analysis¹ strongly support the latter hypothesis. The previously suggested¹ pentagonal bipyramidal co-ordination for one site should be dismissed in favour of a square-pyramidal one. However, the location of cation sites in the large cavity is definitely confirmed.

References

- 1 J. J. Pluth, V. J. Smith, and W. J. Mortier, *Mater. Res. Bull.*, 1977, **12**, 1001.
- 2 P. Comin-Chiaramonti, O. Pongilupi, and G. Vezzalini, *Bull. Soc. Fr. Minér. Cryst.*, 1979, **102**, 386.
- 3 E. A. Stern, D. E. Sayers, and F. W. Lytle, *Phys. Rev. B*, 1975, **11**, 4836.
- 4 B. K. Teo, in 'EXAFS Spectroscopy,' eds. B. K. Teo and D. C. Jay, Plenum, New York, 1981.
- 5 G. Vlaic and J. C. J. Bart, *Recl. Trav. Chim.*, 1982, **101**, 171.
- 6 S. Asbrink and L. J. Norrby, *Acta Crystallogr., Sect. B*, 1970, **26**, 8.
- 7 G. E. Bacon and N. A. Curry, *Proc. R. Soc. London, Sect. A*, 1962, **266**, 95.
- 8 R. W. Joyner, *Chem. Phys. Lett.*, 1980, **72**, 162.
- 9 E. D. Crozier and A. J. Seary, *Can. J. Phys.*, 1980, **58**, 1388.

Received 7th June 1989; Paper 9/02411E