# Molecular Structures of Titanium(IV) and Vanadium(IV) Amides and Alkoxides $\dagger$ 

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#### Abstract

Gas-phase electron diffraction data for $\mathrm{Ti}\left(\mathrm{NMe}_{2}\right)_{4}, \mathrm{~V}\left(\mathrm{NMe}_{2}\right)_{4}$ and $\mathrm{V}(\mathrm{OBu})_{4}$ obtained with nozzle temperatures of $130-156{ }^{\circ} \mathrm{C}$ are consistent with molecular models of $S_{4}$ symmetry. The metal to ligand bond distances $\left(r_{0}\right)$ are $\mathrm{Ti}-\mathrm{N} 191.7(3), \mathrm{V}-\mathrm{N} 187.9(4)$ and $\mathrm{V}-\mathrm{O} 177.9(6) \mathrm{pm}$, respectively. A set of bonding radii for $\mathrm{Ti}^{\mathrm{Iv}}, \mathrm{V}^{\mathrm{v}}, \mathrm{Cr}^{\prime v}$ and $\mathrm{W}^{\mathbf{v}}$ for use with the modified Schomaker-Stevensen rule is proposed. The valence angle bisected by the $S_{4}$ axis in the titanium amide, $\mathrm{N}-\mathrm{Ti}-\mathrm{N} 114.2(17)^{\circ}$, indicates that the co-ordination tetrahedron is slightly flattened. The corresponding angles in the vanadium amide and alkoxide are $\mathrm{N}-\mathrm{V}-\mathrm{N} 100.6(5)^{\circ}$ and $\mathrm{O}-\mathrm{V}-\mathrm{O} 115.1(19)^{\circ}$, indicating that the co-ordination polyhedron in the former is significantly elongated, while in the latter it is significantly flattened. The shapes of the coordination tetrahedra in these and related molecules are discussed in terms of donation of $\pi$ electrons from N or O into low-lying d orbitals on the metal atom.


Relatively few studies have been made of a series of homoleptic molecules $\mathbf{M X}_{n}$ where only one variable, the metal, the ligand or the oxidation number, is changed in a systematic manner. The present study is part of a project aimed at the structure determination of homoleptic transition-metal compounds $\mathbf{M X} \mathbf{X}_{n}$, where $M$ is a Group 4, 5 or 6 metal, the ligand $X$ is an alkyl, dimethylamino or tert-butoxy group, and the co-ordination and oxidation numbers $n$ are 4. An X-ray study of $\mathrm{Mo}\left(\mathrm{NMe}_{2}\right)_{4}{ }^{1}$ and a gas-phase electron diffraction (GED) study of $\mathrm{Zr}\left(\mathrm{NMe}_{2}\right)_{4}{ }^{2}$ have been published by others. We have previously investigated the GED structure of $\mathrm{Cr}\left(\mathrm{OBu}^{1}\right)_{4}{ }^{3}$ and the molecular structure of $\mathrm{Cr}\left(\mathrm{CH}_{2} \mathrm{Me}_{3}\right)_{4}{ }^{4}$ has been published recently. We have, however, been unsuccessful in our attempts at recording GED data for $\mathrm{Ti}\left(\mathrm{CH}_{2} \mathrm{SiMe}_{3}\right)_{4}, \mathrm{Ti}\left(\mathrm{OBu}^{1}\right)_{4}$ and $\mathrm{Zr}\left(\mathrm{CH}_{2} \mathrm{CMe}_{3}\right)_{4}$.

In this paper we report the gas-phase thermal average molecular structures of $\mathrm{Ti}\left(\mathrm{NMe}_{2}\right)_{4}, \mathrm{~V}\left(\mathrm{NMe}_{2}\right)_{4}$ and $\mathrm{V}\left(\mathrm{OBu}^{1}\right)_{4}$. In our discussion of the structures determined to date we emphasize metal-ligand bond distances as well as deviation from perfect tetrahedral co-ordination geometry about the metal atom. Both the amides and the alkoxide have been investigated by various spectroscopic methods (IR, NMR, PS and ESR). ${ }^{5}{ }^{7}$ These studies provided some evidence for distortion of the $\mathbf{M N}_{4}$ or $\mathbf{M O}_{4}$ skeletons, particularly when $\mathbf{M}=\mathrm{V}$.

## Experimental

Synheses.--V(OBu') ${ }^{\text {. }}$. Vanadium tetra-tert-butoxide was prepared by a synthetic method analogous to that used to prepare $\mathrm{Cr}\left(\mathrm{OBu}^{1}\right)_{4} \cdot{ }^{8}$ Solid $\mathrm{VCl}_{3}(\text { thf })_{3}$ (thf $=$ tetrahydrofuran) ( $16.7 \mathrm{~g}, 0.035 \mathrm{~mol}$ ) was added slowly over 24 h to $\mathrm{LiOBu}{ }^{1}(11.3 \mathrm{~g}$, 0.14 mol ) in hexane ( $125 \mathrm{~cm}^{3}$ ) and thf ( $50 \mathrm{~cm}^{3}$ ). Solid CuCl ( 3.5 $\mathrm{g}, 0.035 \mathrm{~mol}$ ) was added to the blue-purple suspension and the mixture was refluxed for 3 h . The blue suspension was allowed to cool to room temperature, settle and filtered. The filtrate was evaporated to dryness and the blue oil was distilled at $85-90$ C at 0.1 mmHg (lit., ${ }^{9} 60-70 \mathrm{C}$ at 0.1 mmHg ) giving $8.3 \mathrm{~g}\left(63^{\circ}\right.$. yield) of $\mathrm{V}\left(\mathrm{OBu}^{1}\right)_{4}$. UV (pentane): 720 and 850 (sh) nm [lit.: ${ }^{7} 719$ and 914(sh) nm (benzene)].
$\mathrm{V}\left(\mathrm{NMe}_{2}\right)_{4}$. Vanadyl chloride, $\mathrm{VOCl}_{3}(4.15 \mathrm{~g}, 0.024 \mathrm{~mol})$, was added dropwise to $\mathrm{LiNMe}_{2}(6.3 \mathrm{~g}, 0.12 \mathrm{~mol})$ suspended in pentane ( $120 \mathrm{~cm}^{3}$ ) at 0 C . The green suspension was warmed to

[^0]Table 1 Nozzle temperatures, number of exposed plates and $s$ limits for GED intensity curves of $\mathrm{Ti}\left(\mathrm{NMe}_{2}\right)_{4}, \mathrm{~V}\left(\mathrm{NMe}_{2}\right)_{4}$ and $\mathrm{V}\left(\mathrm{OBu}^{\prime}\right)_{4}$

|  | $\mathrm{Ti}\left(\mathrm{NMe}_{2}\right)_{4}$ | $\mathrm{V}\left(\mathrm{NMe}_{2}\right)_{4}$ | $\mathrm{V}(\mathrm{OBu})_{4}$ |
| :---: | :---: | :---: | :---: |
| Nozzle temperature ${ }^{\circ} \mathrm{C}$ | 152-155 | 130155 | 152 156 |
| Number of plates |  |  |  |
| 50 cm | 5 | 5 | 6 |
| 25 cm | 5 | 10 | 4 |
| $s$ limits $/ \mathrm{nm}^{-1}$ |  |  |  |
| $\begin{aligned} & 50 \mathrm{~cm} \\ & \left(\Delta s=1.25 \mathrm{~nm}^{1}\right) \end{aligned}$ | 25.00-148.75 | 20.00-135.00 | 30.00-148.75 |
| $\begin{aligned} & 25 \mathrm{~cm} \\ & \left(\Delta s=2.50 \mathrm{~nm}^{-1}\right) \end{aligned}$ | $25.00-280.00$ | 40.00-265.00 | 70.00-245.00 |

room temperature and stirred for 12 h . It was then filtered and the filtrate evaporated to dryness. The green semi-solid residue was distilled at $50-60 \mathrm{C}$ at $10^{-2} \mathrm{mmHg}$ to give $2.0 \mathrm{~g}\left(37^{\circ}\right.$, yield $)$ of green, solid $\mathrm{V}\left(\mathrm{NMe}_{2}\right)_{4}$. UV (pentane) 559 and 750 nm [lit.:? 559 and 752 nm (benzene)]

The compound $\mathrm{Ti}\left(\mathrm{NMe}_{2}\right)_{4}$ was prepared as described by Andersen et al. ${ }^{10}$

Gas-phase Electron Diffraction.-Data were recorded on a Balzers Eldigraph K DG-2 instrument. ${ }^{11.12}$ Two data sets were recorded for each molecule, with nozzle-to-plate distances of 50 and 25 cm , respectively. After photometry of the photographic plates the data were processed by standard procedures. ${ }^{13}$ Backgrounds were computer drawn and an average curve was calculated for each data set. Atomic scattering factors were taken from ref. 14. Information about the experiment (nozzle temperature, number of exposed plates) and $s$ limits in the final refinements is given in Table 1. Intensity curves for $\mathrm{V}\left(\mathrm{NMe}_{2}\right)_{4}$ and $\mathrm{V}\left(\mathrm{OBu}^{1}\right)_{4}$ are shown in Figs. 1 and 2, respectively: those of $\mathrm{Ti}\left(\mathrm{NMe}_{2}\right)_{4}$ cannot be visually distinguished from the ones of the vanadium analogue.

Model Choice and Refinements.- $\mathrm{M}\left(\mathrm{NMe}_{2}\right)_{4}(\mathrm{M}=\mathrm{Ti}$ or V$)$. Structure refinements were based on models with overall $S_{4}$ symmetry, $\mathrm{NCH}_{3}$ fragments of $C_{3 r}$ symmetry and $\mathrm{MN}\left(\mathrm{CH}_{3}\right)_{2}$ fragments of $C_{2}$ symmetry. The latter assumption implies that the co-ordination sphere around the N atoms is planar as found in previous studies of $\mathrm{Mo}\left(\mathrm{NMe}_{2}\right)_{4}$ and $\mathrm{Zr}\left(\mathrm{NMe}_{2}\right)_{4}{ }^{1.2}$


Fig. 1 Calculated (full line) and experimental ( $)$ intensity curves for $\mathrm{Ti}\left(\mathrm{NMe}_{2}\right)_{4}$ with the difference curves shown below


Fig. 2 Calculated (full line) and experimental () intensity curves for $\mathrm{V}\left(\mathrm{OBu}^{\mathbf{t}}\right)_{4}$ with the difference curves shown below

Exploratory refinements with non-planar co-ordination of the nitrogen atoms did not improve the fit to the experimental data. A molecular model of $\mathrm{Ti}\left(\mathrm{NMe}_{2}\right)_{4}$ is shown in Fig. 3. This model is described by three bond distances, $\mathrm{M}-\mathrm{N}, \mathrm{N}-\mathrm{C}$ and $\mathrm{C}-\mathrm{H}$, three bond angles, $\mathrm{N}^{*}-\mathrm{M}-\mathrm{N}$ (see Fig. 3), $\mathrm{M}-\mathrm{N}-\mathrm{C}$ and $\mathrm{N}-\mathrm{C}-\mathrm{H}$, and two dihedral angles $\varphi\left(\mathrm{N}^{*} \mathrm{MNC}\right)$ and $\varphi(\mathrm{MNCH})$. When $\varphi\left(\mathrm{N}^{*} \mathrm{MNC}\right)$ is 0 (i.e. when the two $\mathrm{NC}_{2}$ groups are coplanar) or $90^{\circ}$ the molecular symmetry is $D_{2 d}$.

The eight independent parameters and eight root-meansquare vibrational amplitudes were refined in a least-squares fit of a calculated to the experimental intensity curve. For both amides the absolute value of all correlation coefficients was less than $68 \%$.
In their study of $\mathrm{Zr}\left(\mathrm{NMe}_{2}\right)_{4}$, which is pseudo-isoelectronic with $\mathrm{Ti}\left(\mathrm{NMe}_{2}\right)_{4}$, Rice and co-workers ${ }^{2}$ found the gas-phase electron diffraction data to be equally compatible with a $D_{2 d}$ model with $\varphi\left(\mathbf{N}^{*} \mathrm{MNC}\right)=0$ and an $S_{4}$ model with $\varphi\left(\mathrm{N}^{*} \mathrm{MNC}\right)=66^{\circ}$. Initial refinements of $\mathrm{Ti}\left(\mathrm{NMe}_{2}\right)_{4}$ with $\varphi\left(\mathbf{N}^{*} \mathbf{M N C}\right)$ fixed at 0 yielded an $R$ factor of 0.111 . When $\varphi\left(\mathrm{N}^{*} \mathrm{MNC}\right)$ was refined it increased to $51(1)^{\prime}$ and the $R$ factor decreased to 0.041 . Refinements with starting values for $\varphi\left(\mathrm{N}^{*} \mathrm{MNC}\right)$ spaced at 10 intervals between 0 and 90 all converged to this least-squares minimum. We conclude that a $D_{2 d}$ model of $\mathrm{Ti}\left(\mathrm{NMe}_{2}\right)_{4}$ is inconsistent with our data.
$\mathrm{V}\left(\mathrm{OBu}^{1}\right)_{4}$. In our study of $\mathrm{Cr}\left(\mathrm{OBu}^{1}\right)_{4}$ we investigated four


Fig. 3 Molecular model of $\mathrm{Ti}\left(\mathrm{NMe}_{2}\right)_{4}$ of $S_{4}$ symmetry viewed down the $S_{4}$ axis


Fig. 4 Molecular model of $\mathrm{V}\left(\mathrm{OBu}^{\prime}\right)_{4}$ of $S_{4}$ symmetry viewed down the $S_{4}$ axis
models of $D_{2 d}, D_{2}$ or $S_{4}$ symmetry. ${ }^{3}$ The only model which could be brought into satisfactory agreement with the data was an $S_{4}$ model similar to the one shown in Fig. 4. Since the radial distribution curves of the two molecules are indistinguishable we based our structure refinement of $\mathrm{V}\left(\mathrm{OBu}^{\prime}\right)_{4}$ on the same model. The $\mathrm{C}\left(\mathrm{CH}_{3}\right)_{3}$ and $\mathrm{CCH}_{3}$ fragments were assumed to have local $C_{3}$ symmetry. The model is determined by four bond distances [ $\mathrm{V}-\mathrm{O}, \mathrm{O}-\mathrm{C}, \mathrm{C}-\mathrm{C}$ (mean) and $\mathrm{C}-\mathrm{H}$ ], four valence angles ( $\mathrm{O}-\mathrm{V}-\mathrm{O}^{\prime}, \mathrm{V}-\mathrm{O}-\mathrm{C}, \mathrm{O}-\mathrm{C}-\mathrm{C}$ and $\mathrm{C}-\mathrm{C}-\mathrm{H}$ ), two dihedral angles, $\varphi\left(\mathrm{O}^{\prime} \mathrm{VOC}\right)$ and $\varphi(\mathrm{VOCC})$, and a tilt angle. This tilt angle, $i . e$. the angle between the three-fold symmetry axis of the $\mathrm{C}\left(\mathrm{CH}_{3}\right)_{3}$ group and the $\mathrm{O}-\mathrm{C}$ bond, was defined as positive when the non-bonded $V \cdots C(1)$ and $V \cdots C(2)$ distances were increased. Refinements of the tilt angle failed to converge, and after some trial and error it was fixed at 0.6 which led to marginally better agreement than at 0 .
The remaining ten structure parameters and eleven vibrational amplitudes were refined in a least-squares fit of a calculated intensity curve to the experimental one. The correlation coefficients between the parameters $r(\mathrm{O}-\mathrm{C}), r(\mathrm{C}-\mathrm{C})$,

Table 2 Selected structure parameters, root-mean-square vibrational amplitudes (l) and $R$ factors determined for the best model of $\mathrm{Ti}\left(\mathrm{NMe}_{2}\right)_{4}$ and $\mathrm{V}\left(\mathrm{NMe}_{2}\right)_{4}$. Estimated standard deviations in parentheses in units of the last digit. $R=\left[\Sigma \mathrm{w}\left(I_{\mathrm{obs}}-I_{\text {catc }}\right)^{2} / \Sigma w I_{\mathrm{obs}}\right]^{2}$

|  | $\mathrm{Ti}\left(\mathrm{NMe}_{2}\right)_{4}$ |  | $\mathrm{V}\left(\mathrm{NMe}_{2}\right)_{4}$ |  |
| :---: | :---: | :---: | :---: | :---: |
| Bond distances/pm | $r_{a}$ | $l$ | $r_{\text {a }}$ | $l$ |
| M-N | 191.7(3) | 5.9(3) | 187.9(4) | 5.8(4) |
| $\mathrm{N}-\mathrm{C}$ | 146.1(2) | 4.0(3) | 145.7(3) | 5.1(4) |
| C-H | 112.0(4) | 6.9(4) | 111.1(5) | 7.4(5) |
| Bond angles/" |  |  |  |  |
| $\mathrm{M}-\mathrm{N}-\mathrm{C}$ | 124.3(3) |  | 123.2(3) |  |
| $\mathrm{N}-\mathrm{C}-\mathrm{H}$ | 109.3(7) |  | 112.5(9) |  |
| $\mathrm{N}-\mathrm{M}-\mathrm{N}^{*}$ | 114.2(17) |  | 100.6(5) |  |
| $\mathbf{N}-\mathbf{M}-\mathbf{N}^{* *}$ | 107.2(9) |  | 114.1(1) |  |
| Torsional angles/ ${ }^{\circ}$ |  |  |  |  |
| $\varphi\left(\mathrm{N}^{*} \mathrm{MNC}\right)$ | 51(1) |  | 71(2) |  |
| $\varphi$ (MNCH) | 20(2) |  | -28(2) |  |
| Non-bonded distances/pm | $r_{\text {a }}$ | $l$ | $r_{\text {a }}$ | 1 |
| M . . C | 299.4(4) | 12.4(3) | 294.1(3) | 10.9(4) |
| N... ${ }^{*}$ | 322(3) | 22(3) ${ }^{\text {a }}$ | 289.2(10) | 14(3) ${ }^{\text {a }}$ |
| N... $\mathrm{N}^{* *}$ | 309(2) | 22(3) ${ }^{\text {a }}$ | 315.3(6) | 14(3) ${ }^{\text {a }}$ |
| N...H | 211(1) | 10.1(10) | 214(1) | 9.7(10) |
| C... $\mathrm{C}^{\prime}$ | 241.3(8) | 8.1(10) | 243.8(9) | 7.6(13) |
| $R$ factor of best model | 0.041 |  | 0.049 |  |

${ }^{a}$ Refined with equal shifts

Table 3 Selected structure parameters, root-mean-square vibrational amplitudes ( $l$ ) and $R$ factors determined from the best model for $\mathrm{V}\left(\mathrm{OBu}^{\prime}\right)_{4}$. See Table 2 for definition of uncertainties and $R$ factor

| Bond distances/pm | $r_{\text {a }}$ | 1 |
| :---: | :---: | :---: |
| $\mathrm{V}-\mathrm{O}$ | 177.9(6) | 6.9(4) |
| O-C | 143.3(15) | 4.6(21) |
| C-C | 152.7(7) | 3.4(12) |
| $\mathrm{C}-\mathrm{H}$ | 110.4(4) | 6.1(5) |
| Bond angles |  |  |
| V -O-C | 138.6(18) |  |
| $\mathrm{O}-\mathrm{C}-\mathrm{C}$ | 113.1(15) |  |
| $\mathrm{C}-\mathrm{C}-\mathrm{H}$ | 109.7(12) |  |
| $\mathrm{O}-\mathrm{V}-\mathrm{O}^{\prime}$ | 115.1(19) |  |
| $\mathrm{O}-\mathrm{V}-\mathrm{O}^{\prime \prime}$ | 106.7(11) |  |
| Torsional angles/ |  |  |
| $\varphi\left(\mathrm{O}^{\prime} \mathrm{VOC}\right)$ | 50(2) |  |
| $\varphi[\mathrm{VOCC}(2)]$ | 37(3) |  |
| Non-bonded distances/pm | $r_{\text {a }}$ | 1 |
| V...C | 301(2) | 7.5(11) |
| $\mathrm{V} \cdot . \cdot \mathrm{C}(1)$ | 340(2) | 17.5(22) |
| V...C(2) | 388(4) | 51(20) |
| $\mathrm{V} \cdot \mathrm{C}$ C(3) | 417(1) | 15.6(17) |
| $\mathrm{C}(1) \cdots \mathrm{C}(2)$ | 243(3) | $9.7(5)^{\text {a }}$ |
| O...C(1) | 248(2) | $9.7(5)^{a}$ |
| $\mathrm{O} \cdot . . \mathrm{C}(3)$ | 246(2) | $9.7(5)^{\text {a }}$ |
| $\mathrm{C} \cdot \mathrm{CH}$ | 216(2) | 10.0(13) |
| O... $\mathrm{O}^{\prime}$ | 300(3) | $6.1(25)^{\text {b }}$ |
| $\mathrm{O} \cdot \mathrm{CO}$ | 285(2) | $6.1(25)^{b}$ |
| $R$ factor of best model | 0.055 |  |

" Refined with equal shifts. ${ }^{6}$ Refined with equal shifts.
$l(\mathrm{O}-\mathrm{C})$ and $/(\mathrm{C}-\mathrm{C})$ have absolute values in the range $74-91 \%$. All other correlation coefficients are less than $66 \%$.

All structure refinements were carried out with diagonal weight matrices. The formal estimated standard deviations calculated by the program, $\sigma_{\mathrm{LS}}$, therefore include neither the uncertainty due to data correlation nor to non-refined parameters. The estimated standard deviations quoted in


Fig. 5 Radial distribution curves for $\mathrm{Ti}\left(\mathrm{NMe}_{2}\right)_{4}$, experimental ( $)$ and theoretical (full line). The artificial damping factor $k=30 \mathrm{pm}^{2}$. The most important peaks are indicated by bars of height approximately proportional to the weight of the distance in the theoretical intensity curve. The difference curve is shown below

Tables 2 and 3 have been obtained from $\sigma_{\mathrm{LS}}$ by multiplication with a factor of three. They have also been expanded to include an estimated scale uncertainty of $0.1 \%$.

## Discussion

The parameters determined from the best model for the amides and the alkoxide are listed in Tables 2 and 3, respectively. The radial distribution curves for $\mathrm{Ti}\left(\mathrm{NMe}_{2}\right)_{4}, \mathrm{~V}\left(\mathrm{NMe}_{2}\right)_{4}$ and $\mathrm{V}\left(\mathrm{OBu}^{1}\right)_{4}$ are shown in Figs. 5, 6 and 7, respectively.
No significant differences are found between the structures of the amide ligands in $\mathrm{Ti}\left(\mathrm{NMe}_{2}\right)_{4}$ and $\mathrm{V}\left(\mathrm{NMe}_{2}\right)_{4}$ : in both compounds the co-ordination sphere of the N atom is trigonal planar.

Under the $S_{4}$ point group the four occupied $p_{\pi}$ orbitals on the N atoms may be combined to form symmetry orbitals of A, B and E symmetry. The 3d orbitals on the metal atom transform like A, B (twice) and E. The two low-lying d orbitals in a
tetrahedral field, viz. $3 \mathrm{~d}_{z^{2}}$ and $3 \mathrm{~d}_{x^{2}-y^{2}}$, transform like A and B respectively. All $p_{\pi}$ electrons on N may therefore be stabilized by


Fig. 6 Radial distribution curves for $\mathrm{V}\left(\mathrm{NMe}_{2}\right)_{4}$. Details as in Fig. 5


Fig. 7 Radial distribution curves for $\mathrm{V}\left(\mathrm{OBu}^{1}\right)_{4}$. Details as in Fig. 5
interactions with metal $d$ orbitals, leading to partial $\pi$ bonding between the N and metal atoms. The strongest interaction is however expected with the $3 \mathrm{~d}_{z^{2}}$ and $3 \mathrm{~d}_{x^{2}-y^{2}}$ orbitals.

Since the observed planar co-ordination of the N atoms may be due to repulsion between the methyl groups and the large metal atoms, we do not believe that this observation in itself is sufficient to establish that such $\pi$ interactions are significant.

If the dihedral angle $\varphi\left(\mathrm{N}^{*} \mathrm{MNC}\right)$ is $90^{\circ}$ the non-bonding $\pi$ electron pairs on the N atoms interact only with the $3 \mathrm{~d}_{z^{2}}$ orbital on the metal atom. If the dihedral angle is $0^{\circ}$ they may interact with $3 \mathrm{~d}_{x^{2}-y^{2}}$. The observed dihedral angle in the $\mathrm{d}^{0}$ titanium amide, $51(1)^{\circ}$, is close to $45^{\circ}$, thus indicating that nitrogen $\pi$ electrons are donated equally into both low-lying $d$ orbitals. The $\mathrm{N}-\mathrm{Ti}-\mathrm{N}$ valence angle, $114(2)^{\circ}$, indicates that the coordination tetrahedron of Ti may be slightly flattened. The large value obtained for the $\mathrm{N} \cdots \mathrm{N}$ vibrational amplitudes, 22(3) pm, suggests, however, that the $\mathrm{TiN}_{4}$ core undergoes relatively large-amplitude deformation vibrations.

In the $d^{1}$ vanadium amide the dihedral angle $\varphi\left(N^{*} \mathbf{M N C}\right)$ has increased to $71(2)^{\circ}$. This suggests that the additional electron on the metal atom occupies the $3 \mathrm{~d}_{x^{2}-y^{2}}$ orbital and that the increased dihedral angle reflects a preference for donation of the nitrogen $\pi$ electrons into the empty $3 \mathrm{~d}_{z}=$ orbital. Preferential donation of such electrons into the metal $3 \mathrm{~d}_{z^{2}}$ orbitals would be expected to decrease the $\mathrm{N}^{*}-\mathrm{M}-\mathrm{N}$ angle, and $\mathrm{N}^{*}-\mathrm{V}-\mathrm{N}$ is indeed found to be significantly smaller than tetrahedral, $100.6(5)^{\circ}$. Measurements of spectroscopic and magnetic properties have led to the same conclusion. ${ }^{7}$

The structure of $\mathrm{Cr}\left(\mathrm{NMe}_{2}\right)_{4}$ is unknown, but that of $\mathrm{Mo}\left(\mathrm{NMe}_{2}\right)_{4}$ which is diamagnetic has been determined by X-ray crystallography. ${ }^{1}$ The crystal contains two molecules in the asymmetric unit, both close to $D_{2 d}$ symmetry with $\varphi\left(\mathrm{N}^{*} \mathrm{MoNC}\right)=90^{\circ}$, indicating that the nitrogen $\pi$ electrons are donated into the $4 \mathrm{~d}_{z^{2}}$ orbital exclusively. The two d electrons presumably occupy the $4 \mathrm{~d}_{x^{2}-y^{2}}$ orbital. ${ }^{1}$ The $\mathrm{N}^{*}-\mathrm{M}-\mathrm{N}$ angles in crystalline $\mathrm{Mo}\left(\mathrm{NMe}_{2}\right)_{4}$ would consequently be expected to be smaller than tetrahedral. The $\mathrm{N}-\mathrm{Mo}-\mathrm{N}$ angles in the two crystallographically independent molecules are, however, found to vary in an apparently random manner between 107.3 and $112.5^{\circ}$; it appears that in the crystalline phase the angles are determined by intermolecular forces.
The difference between the $\mathrm{M}-\mathrm{O}$ bond distances in the tertbutoxides of V and $\mathrm{Cr}, 177.9(6)$ and $175.1(7) \mathrm{pm},{ }^{3}$ is of the expected magnitude, but of marginal statistical significance. The $\mathrm{O}-\mathrm{M}-\mathrm{O}$ angle in the $\mathrm{d}^{1}$ vanadium alkoxide is significantly greater than tetrahedral, corresponding to a flattened co-

Table 4 Observed metal-ligand bond distances ( pm ) in compounds $\mathrm{MX}_{4}$ where $\mathrm{M}=\mathrm{Ti}, \mathrm{V}$ or Cr and $\mathrm{WX} \mathrm{X}_{6}$ where $\mathrm{X}=$ alkyl, $\mathrm{NMe}_{2}$, alkoxy or F . Distances calculated from the metal bonding radii with those from the modified Schomaker Stevenson rule ${ }^{a}$ in square brackets

|  | X |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  | $\begin{aligned} & \mathrm{CH}_{2} \mathrm{CMe}_{3} \\ & R(\mathrm{C})=79^{a} \\ & \chi(\mathrm{C})=2.50^{h} \end{aligned}$ | $\mathrm{NMe}_{2}$ $\begin{aligned} & R(\mathrm{~N})=73^{a} \\ & \chi(\mathrm{~N})=3.07^{b} \end{aligned}$ | $\begin{aligned} & \mathrm{OCMe}_{3} \\ & R(\mathrm{O})=72^{a} \\ & \chi(\mathrm{O})=3.50^{b} \end{aligned}$ | $\begin{aligned} & \mathrm{F} \\ & R(\mathrm{~F})=74^{\prime \prime} \\ & \chi(\mathrm{F})=4.10^{\prime \prime} \end{aligned}$ |
| $\begin{aligned} & \mathrm{Ti}^{\mathrm{IV}} \\ & R=137^{c} \end{aligned}$ | - | $191.7(3)^{d}$ [191] | --- | $175.4(3)^{\circ}$ |
| $\begin{aligned} & \chi=1.32^{b} \\ & \mathbf{V}^{\text {IV }} \\ & R=130^{c} \end{aligned}$ | - | $\begin{aligned} & 187.9(4)^{d} \\ & {[186]} \end{aligned}$ | $\begin{aligned} & 177.9(6)^{d} \\ & {[179]} \end{aligned}$ | [175] |
| $\underset{\mathrm{Cr}^{\text {IV }}}{\chi}=1.45^{b}$ | $203.8(4)^{\text {f }}$ | [ | $175.1(7)^{4}$ | 170.6(2) ${ }^{n}$ |
| $\begin{aligned} & R=128^{c} \\ & \chi=1.56^{b} \end{aligned}$ | [199] $\mathrm{CH}_{3}$ | $\mathrm{NMe}_{2}$ | [179] OMc | $[171]$ F |
| $\begin{aligned} & \mathbf{W}^{\mathrm{v}_{1}} \\ & R=145^{\mathrm{c}} \\ & \chi=1.40^{h} \end{aligned}$ | $\begin{aligned} & 214.6(3)^{i} \\ & {[214]} \end{aligned}$ | $\begin{aligned} & 203.5(5)^{j} \\ & {[201]} \end{aligned}$ | $\begin{aligned} & 190.2(3)^{k} \\ & {[193]} \end{aligned}$ | $\begin{aligned} & 183.2(3)^{\prime} \\ & {[185]} \end{aligned}$ |


ordination tetrahedron, indicating preferential donation of oxygen $\pi$ electrons into the $3 \mathrm{~d}_{x^{2}-y^{2}}$ orbital and that the unpaired d electron occupies the $3 \mathrm{~d}_{z^{2}}$ orbital. The dihedral angle $\varphi\left(\mathrm{O}^{\prime}\right.$ VOC $)=50(2)^{\circ}$, however, does not favour donation into either of the two orbitals. It is suggested that the dihedral angle is determined by steric repulsion between the very bulky ligands. The $\mathrm{d}^{2}$ chromium alkoxide ${ }^{3}$ is paramagnetic with a magnetic moment corresponding to two unpaired electrons, $\left(3 \mathrm{~d}_{z^{2}}\right)^{1}\left(3 \mathrm{~d}_{x^{2}-y^{2}}\right)^{1}$, and as expected the $\mathrm{O}^{\prime}-\mathrm{Cr}-\mathrm{O}$ angle, $112(3)^{\circ}$, is not significantly greater than tetrahedral.

To what extent do the bond distances indicate that they are shortened by partial $\pi$ bonding? We have found that in simple molecules single bond distances between the eighteen maingroup elements in Groups 14-17 and H can be predicted with an average error of 3.2 pm from the modified SchomakerStevensen (MSS) rule [equation (1)], where $R$ is the bonding

$$
\begin{equation*}
R(\mathrm{~A}-\mathrm{B})=R(\mathrm{~A})+R(\mathrm{~B})-c|\chi(\mathrm{~A})-\chi(\mathrm{B})|^{1.4} \tag{1}
\end{equation*}
$$

radius, $\chi$ is the Allred-Rochow electronegativity and the constant $c=8.5 \mathrm{pm} .{ }^{15}$ Since the bonding radii of $\mathrm{N}, \mathrm{O}$ and F have been adjusted to reproduce single bond distances to atoms without low-lying acceptor orbitals, we believe them to be suitable for estimation of single $\sigma$ bonds to transition metals. In Table 4 we compare experimental and calculated bond distances obtained after adjustment of the bonding radii of the transition-metal atoms: $R=137,130,128$ and 145 pm for $\mathrm{Ti}^{\mathrm{iv}}$, $\mathrm{V}^{\mathrm{IV}}, \mathrm{Cr}^{1 \mathrm{~V}}$ and $\mathrm{W}^{\mathrm{V1}}$, respectively. The average difference between calculated and observed bond distances is about 2 pm . Thus, although dative $\pi$ bonding may influence the valence angles at the metal and the orientations of the ligands, its effect on $M-X$ bond distances is uncertain.

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[^0]:    $\dagger$ Non-SI unit employed: $\mathrm{mmHg} \approx 133 \mathrm{~Pa}$.

