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Crystal structure of Fe2Os(CO)12, a structural analog of Fe3(CO)12 with no required site symmetry and only a ~12:1 "star of david" disorder pattern. Example of an acentric, but close to centric, structure. Some thoughts on the solid-state structure of Fe3(CO)12

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ygen atom of the carboxyl group. By analogy, it is known that tertiary phosphines will add to the CH, group of σ - π -coordinated alkenyl groups to form the zwitterionic ligands $R_3P^+CH_2C^-R$, in which the negatively charged carbon atom bridges two metal atoms.¹⁵

(15) Churchill, M. R.; DeBoer, B. G.; Shapley, J. R.; Keister, J. B. *J. Am. Chem.* **SOC. 1976, 98, 2357.**

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Supplementary Material Available: Tables of positional parameters for hydrogen atoms and anisotropic thermal parameters for all three crystal structure analyses (10 pages); tables of structure factor amplitudes **(55** Pages). Ordering information is given on any current masthead page.

Crystal Structure of Fe₂Os(CO)₁₂, a Structural Analogue of $Fe₃(CO)₁₂$ with No Required Site Symmetry and Only a \sim 12:1 **"Star of David" Disorder Pattern: Example of an Acentric, but Close to Centric, Structure. Some Thoughts on the Solid-state Structure of Fe₃(CO)₁₂**

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The heteroatomic trinuclear cluster $Fe₂Os(CO)₁₂$ crystallizes in the noncentrosymmetric monoclinic space group *Pn* (variation of *Pc*; C_5^2 , No. 7) with $a = 8.377$ (2) Å, $b = 22.715$ (8) Å, $c = 8.953$ (2) Å, $\beta = 96.510$ $(18)^\circ$, $V = 1692.6$ (8) \AA^3 , and $Z = 4$. The molecular packing is close to, but not exactly, that required for space group P_{21}/n (cf. $Fe_3(CO)_{12}$, P_{21}/n , $a = 8.359$ (2) Å, $b = 11.309$ (2) Å, $c = 8.862$ (2) Å, $\beta = 97.00$ (2)°, $V = 831.5$ (3) Å³, and $Z = 2$, i.e., similar except that, in the Fe₂Os(CO)₁₂ structure, 2, a absent and *b* is doubled). Diffraction data for $\pm h, +k, \pm l$ were collected for $2\theta = 4.5-50.0^{\circ}$ with Mo K α radiation. The structure was solved by Patterson and difference-Fourier methods, and refinement converged radiation. The structure was solved by Patterson and difference-Fourier methods, and refinement converged
with $R_F = 4.9\%$, $R_{\rm wF} = 4.8\%$, and GOF = 1.772 for all 5993 point-group-unique reflections ($R_F = 4.0\%$, $R_{\text{wF}} = 4.6\%$ for those 5220 data with $|F_0| > 6\sigma(|F_0|)$. There are two crystallographically distinct $\text{Fe}_2\text{Os}(\text{CO})_{12}$ molecules in the unit cell. Each suffers an approximately 12:l "star of David" disorder of the metal atoms. The major components are unambiguously defined and contain an $Os(CO)₄$ system linked to two $Fe(CO)₃$ units, which, in turn, are bridged by two asymmetrically bridging carbonyl ligands-i.e., a connectivity similar to that in $Fe_3(CO)_{12}$. Bond lengths of note (in Å) are as follows: molecule A, Os(A)-Fe(1A) = 2.735 (2), Os(A)-Fe(2A) = 2.740 (3), Fe(1A)-Fe(2A) = 2.589 (4), Fe(1A)-C(BR1A) = 2.245 (16), Fe(2A)-C(BR1A) = 1.939 (16), Fe(1A)–C(BR2A) = 1.918 (17), Fe(2A)–C(BR2A) = 2.224 (16); molecule B, Os(B)–Fe(1B) = 2.745 (3), Os(B)–Fe(2B) = 2.746 (3), Fe(1B)–Fe(2B) = 2.594 (4), Fe(1B)–C(BR1B) = 2.084 (20), Fe- $(2B)-C(BR1B) = 2.047 (19), F_e(1B)-C(BR2B) = 1.937 (17), F_e(2B)-C(BR2B) = 2.123 (17).$

Introduction

The saga of the structural characterization of triiron dodecacarbonyl, $Fe₃(CO)₁₂$, has been vividly recounted by Wei and Dahl.¹ Further refinement of the crystal structure was later carried out by Cotton and Troup.² It is established that *in the macroscopic solid state* the molecule has structure **1** and that the molecules are statistically disordered about sites of $\overline{1}$ (C_i) symmetry in space group $P2_1/n$.

The heteronuclear species $Fe₂Os(CO)₁₂$ has been synthesized from $Fe_2(CO)_9$ and $H_2O_8(CO)_4$.³ The substitu-

tional reactivity of this material **has** recently been studied by Shojaie and Atwood.4 **As** a prelude to investigating some products of substitution, we have carried out a single-crystal X-ray diffraction study on the parent compound, $Fe₂Os(CO)₁₂$. Our results, reported below, may also have relevance toward understanding the crystal structure of $Fe₃(CO)₁₂$.

Experimental Section

Collection of X-ray Diffraction Data for $Fe₂Os(CO)₁₂$. Crystals were provided by Shojaie and Atwood.' **A** black crystal with approximate dimensions $0.30 \times 0.30 \times 0.33$ mm was inserted into a glass capillary with its extended direction parallel to the wall of the capillary. It was then mounted in a eucentric goniometer and accurately aligned and centered on a Syntex P2₁ automated four-circle diffractometer. All further setup operations (Le. determination of the crystal's unit cell parameters and orientation matrix) and collection of the intensity data were performed **as** described previously.6 Details appear in Table I. The

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⁽¹⁾ Wei, C. H.; Dahl, L. F. *J. Am. Chem. SOC.* **1969,91,1351.** *See* **also: (2)** Cotton, F. A.; Troup, J. M. *J. Am. Chem. SOC.* **1974, 96, 4155.** Desiderato, R., Jr.; Dobson, *G.* R. *J. Chem. Educ.* **1982, 59, 752.**

Table 1. Experimental Data for the X-ray Diffraction Study of $Fe₂Os(CO)₁₀(\mu-CO)₂$

 (4) U (4) U (1)

(B) Collection of X-ray Diffraction Data

diffractometer: Syntex P21

radiation: Mo K $\bar{\alpha}$ ($\bar{\lambda}$ = 0.710730 Å)

- monochromator: highly oriented (pyrolytic) graphite; equatorial mode with $2\theta(m) = 12.160^{\circ}$; assumed to be 50% perfect/50% ideally mosaic for polarizn cor
- rflns measd: $\pm h, \pm k, \pm l$ for $2\theta = 4.5 50.0^{\circ}$, yielding 5993

point-group-unique data

scan type: coupled θ (cryst)-2 θ (counter)

scan width: $[2\dot{\theta}(K_{\alpha_1}) - 0.8] - [2\theta(K_{\alpha_2}) + 0.8]$ °

scan speed: $4.0^{\circ}/\text{min}$ (20)

- bkgds: stationary cryst, stationary counter at the two extremes of the 2 θ scan; each for $\frac{1}{4}$ of the total scan time
- std rflns: 3 mutually orthogonal rflns collected before each set of 97 data points; no decay obsd
- abs cor: $\mu(\text{Mo K}\bar{\alpha}) = 97.3 \text{ cm}^{-1}$, corrected empirically by interpolation (in 2θ and ϕ) between 6 close-to-axial (ψ -scan) rflns; rflns used were 162 ($2\theta = 15.30^{\circ}$, $T_{\min}/T_{\max} = 0.693$), 282 (20.20', 0.699), 392 (24.43', 0.728), 2,11,3 (26.72', 0.7181, 3,11,3 (29.22', 0.744), and 4,12,3 (33.50°, 0.771)

Table II. Intensity Statistics for $Fe₂Os(CO)₁₀(\mu-CO)₂$

		ideal		
function	obsd	acentric	centric	
$\langle E \rangle$	0.873	0.886	0.798	
	1.000	1.000	1.000	
$\ket{1}$ E^2	0.747	0.736	0.968	
> 1.0, % ĿΊ	36.72	36.79	31.73	
> 1.2, % E	23.36	23.69	23.01	
$> 1.4, \%$ EI	13.79	14.09	16.15	
> 1.6, % Εl	7.94	7.73	10.96	
$> 1.8, \%$ E	4.05	3.92	7.19	
> 2.0, % E	2.33	1.89	4.55	
$2.5, \, \%$ El ⋗	0.37	0.19	1.24	
$3.0, \%$	0.03	0.01	0.27	

final cell parameters are based upon a least-squares analysis of the setting angles $(2\theta, \omega, \chi)$ of 25 automatically centered reflections, well dispersed in reciprocal space, with $2\theta = 20-30^{\circ}$ (Mo K $\bar{\alpha}$) radiation).

The crystal was found to exhibit the systematic absences *h01* for $h + l = 2n + 1$, consistent with either the centrosymmetric monoclinic space group $P2/n$ (a nonstandard setting of $P2/c$; C_{2h}^4 ; No. 13, having the equipoints $\pm(x, y, z; ^{1}/_{2} + x, \bar{y}, ^{1}/_{2} + z))$ or the noncentrosymmetric space group *Pn* (a nonstandard setting of Pc; C_8^2 ; No. 7, having the equipoints *x*, *y*, *z*; $\frac{1}{2}$ + *x*, \bar{y} , $\frac{1}{2}$ + *z*). Data were therefore collected for the octants $\pm h, +k, \pm l$. This provides two equivalent forms for space group $P2/n$ $(+h, +k, +l) = -h, +k, -l$ and $-h, +k, +l = +h, +k, -l$ or a single point-groupunique form for space group *Pn* (thereby allowing the determination of crystal polarity via anomalous dispersion). The completed structural analysis revealed that the true space group is the noncentrosymmetric *Pn.* (See Table **I1** for intensity statistics, which are uniformly indicative of the acentric case.)

Solution **and** Refinement of the Structure. All crystallographic calculations were performed on a Syntex **XTL** system, which incorporated the use of the SUNY-Buffalo modifiedversion of the **XTL** interactive crystallographic program package.

The analytical form of the scattering factor for the appropriate neutral atoms^{6a} was used in calculating F_c values; corrections were made for both the real $(\Delta f')$ and imaginary $(\Delta f'')$ components of anomalous dispersion.^{6b}

Table 111. Final Atomic Parameters **for** Fe,Oe(CO)lz

			z	, $B(\text{iso})$, \AA^2
atom	x	у		
		Molecule A		
Os(A)	0.00000(0)	$-0.13493(3)$	0.00000(0)	
Os(A')	0.0965 (20)	$-0.11341(47)$	0.3374 (19)	4.62(25)
Fe(1A)	0.22549(31)	–0.13714 (11)	0.24599 (26)	
Fe(2A)	$-0.05702(29)$	$-0.09640(13)$	0.27907 (28)	
Fe(1A')	0.1580(21)	–0.15860 (83)	0.0659 (19)	0.79(28)
Fe(2A')	$-0.1090(30)$	–0.1046 (11)	0.0869(28)	2.88(47)
C(1A)	$-0.2126(23)$	$-0.13030(65)$	–0.1154 (21)	4.33 (33)
C(2A)	0.1250(21)	$-0.15977(72)$	$-0.1575(21)$	4.31 (33)
C(3A)	0.0449(16)	$-0.05296(62)$	$-0.0386(16)$	2.95(26)
C(4A)	$-0.0364(21)$	–0.21778 (83)	0.0595(20)	4.29 (35)
C(11A)	0.3596(20)	$-0.12474(66)$	0.4171 (19)	4.37 (30)
C(12A)	0.3404(21)	$-0.09781(73)$	0.1155 (19)	4.24 (33)
C(13A)	0.2957(19)	$-0.20828(72)$	0.1971(18)	4.19 (30)
C(21A)	$-0.0792(21)$	$-0.08091(71)$	0.4747(20)	4.34 (33)
C(22A)	–0.1657 (20)	$-0.03552(75)$	0.2015(19)	4.55 (32)
C(23A)	$-0.2207(22)$	$-0.14544(82)$	0.2445(20)	4.74 (36)
C(BR1A)	0.1324(18)	$-0.04684(67)$	0.2899(17)	3.37(29)
C(BR2A)	0.0752(20)	$-0.17784(66)$ $-0.13025(56)$	0.3554(20)	3.87(29)
0(1A) O(2A)	$-0.3295(18)$ 0.1921(22)	$-0.17091(76)$	$-0.1823(20)$ $-0.2496(18)$	
O(3A) O(4A)	0.0722(16) $-0.0618(18)$	$-0.00574(52)$ $-0.26381(52)$	$-0.0688(15)$ 0.0792(16)	
0(11A)	0.4540(20)	$-0.11667(52)$	0.5197(17)	
O(12A)	0.4267(18)	$-0.07401(65)$	0.0503(18)	
O(13A)	0.3433(19)	$-0.25435(62)$	0.1748(17)	
O(21A)	$-0.0922(20)$	$-0.07080(69)$	0.5991(15)	
O(22A)	-0.2342 (20)	0.00645(72)	0.1599(17)	
O(23A)	$-0.3318(18)$	$-0.17801(65)$	0.2451(19)	
O(BR1A)	0.1888(15)	$-0.00306(54)$	0.3012(15)	
O(BR2A)	0.0485(17)	$-0.21933(59)$	0.4213(15)	
		Molecule B		
Os(B)	$-0.41083(9)$	$-0.38736(3)$	$-0.15249(7)$	
Os(B')	$-0.5281(12)$	-0.36324 (37)	$-0.4867(13)$	2.89 (20)
Fe(1B)	$-0.65092(31)$	$-0.37452(12)$	$-0.38559(30)$	
Fe(2B)	$-0.35689(31)$	–0.35676 (12)	–0.43921 (29)	
Fe(1B')	$-0.5715(29)$	-0.3855 (11)	–0.1926 (27)	2.74(44)
Fe(2B')	$-0.2781(29)$	$-0.3824(12)$	$-0.2540(27)$	2.77(42)
C(1B)	$-0.3922(19)$	–0.46780 (72)	$-0.2098(17)$	3.65(29)
C(2B)	$-0.4286(19)$	$-0.30233(72)$	–0.1087 (17)	3.84 (32)
C(3B)	$-0.1981(23)$	-0.39117 (68)	$-0.0467(22)$	4.62 (35)
C(4B)	$-0.5382(21)$	$-0.40487(79)$	0.0093(22)	4.70 (36)
C(11B)	$-0.7984(19)$	-0.36307 (64)	$-0.5477(18)$	4.13(29)
C(12B)	$-0.7595(25)$	-0.33443 (87)	$-0.2523(23)$	5.35 (40)
C(13B)	$-0.7190(21)$	$-0.44675(80)$	$-0.3371(20)$	4.95 (35)
C(21B) C(22B)	$-0.3439(19)$	$-0.33584(71)$	$-0.6338(19)$	3.92(32)
	$-0.2391(22)$	$-0.29676(84)$	$-0.3572(21)$	5.31 (39)
C(23B) C(BR1B)	$-0.1929(23)$ $-0.5050(24)$	$-0.40941(83)$ –0.42563 (76)	$-0.4234(20)$	4.68 (35) 4.88 (34)
C(BR2B)	$-0.5539(20)$		$-0.5078(23)$ $-0.4228(19)$	
O(1B)	$-0.3806(17)$	$-0.29928(75)$ $-0.51696(50)$	$-0.2406(16)$	4.27 (32)
O(2B)	-0.4424 (16)	$-0.25561(48)$	$-0.0873(13)$	
O(3B)	$-0.0726(18)$	–0.39435 (59)	0.0222 (19)	
O(4B)	$-0.6128(21)$	–0.41759 (83)	0.1060(17)	
0(11B)	$-0.8919(18)$	–0.35451 (55)	$-0.6428(18)$	
O(12B)	$-0.8300(18)$	–0.31039 (63)	$-0.1763(19)$	
O(13B)	$-0.7655(18)$	-0.49216 (65)	$-0.3092(20)$	
O(21B)	$-0.3222(20)$	$-0.32549(75)$	$-0.7447(15)$	
O(22B)	$-0.1575(22)$	$-0.25814(73)$	$-0.3160(18)$	
O(23B)	$-0.0951(19)$	$-0.44471(74)$	$-0.4187(23)$	
0(BR1B)	$-0.5122(18)$	–0.46721 (58)	$-0.5817(14)$	
O(BR2B)	$-0.5703(18)$	-0.24882 (56)	$-0.4455(16)$	

The structure was solved by the successful location of two independent osmium atoms from a Patterson map in space group *Pn* with $Z = 4$; two independent molecules therefore define the asymmetric unit. Least-squares refinement (three cycles full matrix) led to $R_F = 29.3\%$ and $R_{wF} = 35.2\%$.⁷ A difference-Fourier synthesis was now performed, yielding the positions of the four iron atoms. Three cycles of full-matrix least-squares refinement led to $R_F = 15.2\%$, $R_{\text{wF}} = 19.8\%$. Subsequent difference-Fourier syntheses led to the location of all atoms of the carbonyl groups $(R_F = 7.4\%)$. At this point a further difference-Fourier synthesis was performed. Six independent peaks

⁽⁵⁾ Churchill, M. R.; Lashewycz, R. A.; Rotella, F. J. *Inorg. Chem.* **1977**, 16, 265.

^{(6) (}a) *International Tables /or X-ray Crystallography;* Kynoch Press Birmingham, England, 1974; Vol IV, pp 99-101. (b) *Ibid.,* pp 149-150.

⁽⁷⁾ $R_F(\%) = 100[(\sum ||F_0| - |F_0||)/\sum |F_0|]; R_{wF}(\%) = 100[(\sum w(|F_0| - |F_0|)^2)/\sum w|F_0|^2]^{1/2}$, where $1/w = [\sigma(|F_0|)]^2 + [0.015|F_0]|^2$; GOF = $[\sum w(|F_0| - |F_0|)^2/(NO - NV)]^{1/2}$, where NO = number of observations and NV = number of variables.

Figure 1. ORTEP-II diagram for molecule A of $Fe₂Os(CO)₁₀(\mu$ -CO)₂. The *b* direction is vertical, and a is horizontal. The minor-component atoms are artificially reduced.

Figure 2. ORTEP-II diagram for molecule **B** of $Fe₂Os(CO)₁₀(\mu$ -CO)₂, with *b* vertical and a horizontal (i.e., the same as for Figure 1). This is the molecule related to molecule A by a pseudo inversion center. Note that the major $Fe₂Os$ core is tipped from the horizontal by a lesser degree than that in Figure 1.

with heights of up to \sim 6 e/Å³ indicated that the structure possessed some form of disorder. These peaks were found to lie in the planes of the three metal atoms in each of the two independent molecules. Within each molecule, the three metal atoms and three residual peaks formed a "star of David" pattern indicative of disorder involving either (i) a 180' rotational disorder of a major and minor molecular orientation or (ii) an inversion disorder of the major and minor molecular orientation. Since the independent molecule has C_2 (and approximate $C_{2\nu}$) symmetry, these alternatives are indistinguishable in the present case. The secondary osmium atoms were distinguished from iron atoms on the basis of (a) peak heights on the difference-Fourier map and (b) interatomic distances $(Os-Fe > Fe-Fe)$.

Careful refinement of the occupancies of the major and minor components yielded average normalized values of 0.925 and 0.075 (an approximately 121 ratio). Inversion of coordinates, with appropriate changes in B_{12} and B_{23} values, now gave lower R_F and R_{wF} values; the crystal polarity was thus redefined at this stage of the analysis. Two cycles of refinement gave the residuals R_F = 5.1% and R_{WF} = 5.0%. An extinction correction was now

performed; the correction applied was very small, with $g = 1.245$ \times 10⁻⁷ in the expression $|F_{\text{o,corr}}| = |F_{\text{o,uncorr}}|(1.0 + gI_o)$. Convergence of the model was achieved with the major components of the metal atoms and the oxygen atoms assigned anisotropic thermal parameters and **all** other atoms **having** isotropic thermal parameters. The residuals for all 5993 reflections were $R_F = 4.9\%$, $R_{wF} = 4.8\%$, and GOF = 1.772. These were reduced to $R_F = 4.3\%$, $R_{WF} = 4.7\%$, and GOF = 1.813 for those 5585 reflections with $|F_o| > 3\sigma(|F_o|)$ and to $R_F = 4.0\%$, $R_{wF} = 4.6\%$, and GOF = 1.841 for those 5220 reflections with $|F_o| > 6\sigma(|F_o|)$. A final difference-Fourier synthesis was now performed, revealing only peaks of height $\langle 1 \rangle$ e/Å³.

If we assume that the disorder pattern follows that of $Fe₃(C 0)_{12}$ ¹ our model omits secondary carbon atoms (and, possibly, secondary oxygen sites²) of occupancy 0.075. This would correspond to "pseudoatoms" containing 0.45 e (0.075 **X** C) and 0.60 e (0.075×0) . We feel justified in this omission.

Final atomic positions are listed in Table III. (Anisotropic thermal parameters appear as supplementary material.)

Description of the Crystal Structure of $Fe₂Os(CO)_{12}$

Crystals are composed of discrete molecular units of composition $Fe₂Os(CO)₁₀(\mu$ -CO)₂. There are two crystallographically independent units in the asymmetric unit. The labeling of molecule A is shown in Figure 1; that of molecule B is shown in Figure 2. Interatomic distances and angles are compiled in Tables **IV** and **V,** respectively.

Each of the molecular sites is subject to a 0.925:0.075 $(\sim 12:1)$ disorder, with the same "star of David" pattern of disorder for the Fe₂Os cores as is observed in Fe₃(C- O_{12} ^{1,2} However, for $Fe₂Os(CO)₁₀(\mu$ -CO₂, the molecules are in general positions and the connectivity of the major components can be observed directly. Crystalline $Fe₃(C O_{12}$ contains units of Fe₃(CO)₁₀(μ -CO)₂ that are statistically disordered about sites of C_i (1) symmetry, and elucidation of the *molecular* geometry must be inferred.^{1,2} Here, the two $Fe₃$ units form a "star of David" pattern with a required 0.5:0.5 occupancy about crystal-averaged inversion centers. Further comparisons of the *crystal structures* of $Fe₂Os(CO)₁₀(\mu$ -CO $)₂$ and $Fe₃(CO)₁₀(\mu$ -CO $)₂$ will be deferred to later.

We now come to the details of the molecular geometry of $Fe₂Os(CO)₁₀(\mu$ -CO)₂. Parameters for both the major and minor components are available in Tables **IV** and **V.** Within the text we will describe the geometry of the major components only. In molecule A, metal-metal distances are Os(A)-Fe(1A) = 2.735 (2) **A,** Os(A)-Fe(2A) = 2.740 (3) \AA , and $\text{Fe}(1\text{A})-\text{Fe}(2\text{A}) = 2.589$ (4) \AA ; in molecule B, Os-(B)-Fe(1B) = 2.745 (3) **A,** Os(B)-Fe(2B) = 2.746 (3) **A,** and $Fe(1B)-Fe(2B) = 2.594$ (4) Å. Average metal-metal bond **f** 0.004 **A.** (The comparable dibridged Fe-Fe distance in $Fe₃(CO)₁₀(\mu$ -CO)₂ is 2.558 (1) Å, while the nonbridged distances are 2.677 (2) and 2.683 (1) Å; the dibridged Fe-Fe linkage in $Fe_3(CO)_{10}(\mu$ -CO)₂ being shorter than that in $Fe₂Os(CO)₁₀(\mu$ -CO)₂ probably is a geometric result of shorter nonbridged Fe-Fe, vis à vis Os-Fe, linkages.) lengths are $Os-Fe = 2.742 \pm 0.005$ Å and Fe-Fe = 2.592

Osmium-carbonyl distances range from 1.909 (16) through 1.989 (19) **A,** averaging 1.943 **A;** the precision of the structural determination is not sufficient to enable us to distinguish between axial and equatorial Os-CO dis $tances$ (cf. Os-CO(axial) = 1.946 \pm 0.006 Å vs Os-CO- $(equatorial) = 1.912 \pm 0.007$ Å in $Os_3(CO)_{12}$.⁸ Terminal iron-carbonyl distances range from 1.755 (17) through A in $\text{Fe}_3(\text{CO})_{10}(\mu\text{-\text{CO}})_2$.² Terminal C-O distances are 1.056 (22)-1.190 (24) **A** (average 1.134 **A).** 1.828 (17) **A,** averaging 1.804 *1* (cf. Fe-CO(average) = 1.82

We now come to the $Fe(\mu$ -CO)₂Fe bridges within the two $Fe₂Os(CO)₁₀(\mu$ -CO)₂ molecules. In molecule A, the short Fe(lA)-C(BR2A) bond is counterbalanced by the long Fe(2A)-C(BR2A) interaction (1.918 (17) vs 2.224 (16) **A)** and the short Fe(2A)-C(BRlA) bond is counterbalanced by the long Fe(1A)-C(BR1A) interaction (1.939 (16) vs 2.245 (16) **A);** the differences of 0.306 and 0.306 A are identical, and the idealized total molecular symmetry of C_2 (as determined for $Fe_3(CO)_{10}(\mu$ -CO)₂^{1,2}) is confirmed. Appropriate M-C-O angles are $Fe(IA)-C(BR2A)-O(BR2A) = 146.7 (14)°$ vs $Fe(2A)-C(BR2A)-O(BR2A) =$ $(136.2 \ (14)^{\circ}$ and $Fe(2A)$ -C(BR1A)-O(BR1A) = 150.4 (14)^o vs $Fe(1A)-C(BR1A)-O(BR1A) = 133.4 (13)°$.

Unfortunately, the situation for molecule B is not quite so clear-cut. Thus, we have $Fe(1B)-C(BR2B) = 1.937(17)$ \AA vs Fe(2B)-C(BR2B) = 2.123 (17) \AA and Fe(2B)-C-(BRlB) = 2.047 (19) **A** vs Fe(1B)-C(BR1B) = 2.084 (20) A-i.e., lesser asymmetries of 0.186 and 0.037 **A** only. The

Figure 3. Positioning of the $Fe₂Os$ units (major components only) in the packing of $Fe₂Os(CO)₁₀(\mu-CO)₂$. Type A molecules are shaded. The basic unit cell is outlined *(b* vertical, *a* horizontal). The *n*-glide planes are at $y = 0$, $\frac{1}{2}$, 1.

Figure 4. Positioning of Fe₃ units in the packing of disordered Figure 4. Positioning of Fe₃ units in the packing of disordered $Fe_3(CO)_{10}(\mu$ -CO)₂. The double unit cell 2b (vertical) and *a* (horizontal) is outlined. The *n*-glide planes are at $y = \frac{1}{4}$, $\frac{3}{4}$, $\frac{1}{4}$, $\frac{1}{2}$, $\frac{1}{4}$, $\frac{1}{4}$, the 2₁ axes lie along y at $x = \frac{1}{4}$, $\frac{3}{4}$ and $z = 0$, $\frac{1}{2}$, $\frac{1}{1}$; $\frac{1}{2}$, $\frac{1}{2}$, $\frac{1}{4}$, $\frac{1}{2}$, $\frac{1}{2}$, $\frac{1}{2}$, $\frac{1}{2}$, $\frac{1}{2}$, $\frac{1}{2}$, \frac $= 0, \frac{1}{2}, 1.$

most probable explanation of this is that the energy profile from the asymmetrically bridging case to the symmetrically bridging case is very flat and is in the same energetic range as weak intermolecular forces.

⁽⁸⁾ Churchill, M. R.; DeBoer, B. G. *Inorg. Chem.* **1977,16, 878.**

Unit Cells of $Fe₂Os(CO)₁₀(\mu-CO)₂$ and
 $Fe₃(CO)₁₀(\mu-CO)₂$

The unit cell dimensions of the $Fe₂Os$ complex are $a =$ 8.377 (2) Å, $b = 22.715$ (8) Å, $c = 8.953$ (2) Å, and $\beta =$ The exact ratios (Fe₂Os:Fe₃) are 1.002:1 for a, 2.009:1 for b, 1.010:1 for c, and 0.995:1 for β . This correspondence led 96.510 (18)^o; those for the Fe₃ complex are² $a = 8.359$ (2) \AA , $b = 11.309$ (2) \AA , $c = 8.862$ (2) \AA , and $\beta = 97.00$ (2)^o. us to hope that we might be able to determine whether there could be any microscopic order as opposed to macroscopic disorder in the crystal structure of $Fe_3(CO)_{10}$ - $(\mu$ -CO)₂.

Figures 3 and 4 show the orientations of $Fe₂Os$ and $Fe₃$ triangles in $Fe₂Os(CO)₁₀(\mu$ -CO)₂ and $Fe₃(CO)₁₀(\mu$ -CO)₂.

Let us first consider the crystal packing of $Fe₂Os (CO)_{10}(\mu$ -CO)₂ in detail. As is clearly shown in Figure 3, there are four orientations of molecules, arranged in layers, up the b axis. These are labeled as follows (starting from the bottom of the figure): B (molecule B, whose coordinates are listed in Table 111), A (molecule A from Table III), $A(n)$ (related to molecule A by an *n* glide at $y = 0$), and $B(n)$ (related to molecule B by an n glide at $y = 0$); this pattern then repeats by virture of a b translation.

We solved this structure in the noncentrosymmetric space group *Pn*, rather than its potential centrosymmetric supergroup $P2/n$. This choice is clearly correct—there is, indeed, only one orientation of molecules at any one level along the *b* axis. (In space group $P2/n$, there must be a 2-fold axis about *b,* two molecular orientations in each layer, and only two such layers in the cell.) Nevertheless, the structure has definite indications of having approximate (but not exact) inversion centers. (Note that the *^x* and *z* coordinates are not specifically determined in space group *Pn*; we defined an origin by fixing $x = 0$ and $z = 0$ for $Os(1A)$.) Molecule A is related to molecule B by an approximate inversion center at $x = -0.2084$, $y = -0.2478$, $z = -0.0754$. The metal atoms, taken in pairs, yield averaged coordinates as follows: $\text{Os}(A) + \text{Os}(\overline{B})$, $x = -0.2054$, $y = -0.2611$, $z = -0.0763$; Fe(1A) + Fe(1B), $x = -0.2127$, $y = -0.2558$, $z = -0.0698$; $Fe(2A) + Fe(2B)$, $x = -0.2070$, $y = -0.2266$, $z = -0.0801$. The correspondence between the lighter atoms (whose positions are, in essence, dictated by the orientations of the metal clusters) is not so precise. Thus, the coordinates $O(BR1A) + O(BR1B)$ average to *x* $= -0.16165$, $y = -0.2351$, $z = -0.1403$ and $O(BR2A)$ + O(BR2B) average to $x = -0.2610$, $y = -0.2341$, $z = -0.0121$.

Note that, if the inversion center is precise and is at y $= -1/4$, then all four molecules are interrelated by sym-Note that, if the inversion center is precise and is at y
= $-1/4$, then all four molecules are interrelated by sym-
metry elements: A \rightarrow B (inversion center at y = $-1/4$), A $=$ $-1/4$, then all four molecules are interrelated by symmetry elements: A \rightarrow B (inversion center at $y = -1/4$), A \rightarrow A(n) (n glide at $y = 0$), and A \rightarrow B(n) (2₁ axis); in this case, the space group will then be the centrosymmetric $P2₁/n$. (Cell dimensions are unchanged, and a standard set of coordinates can be obtained by moving the inversion

Table VI. $0k0$ Data for $\text{Fe}_2\text{Os}(\text{CO})_{10}(\mu\text{-CO})_2$ ($|F_o|$, $\sigma(|F_o|)$, $|F_o|$),

וט –ו							
rfln <i>hkl</i>	$ F_{\rm o} $ $(\sigma(F_{\rm o}))$	$\left F_{\rm e}\right $	rfln <i>hkl</i>	$ F_{\rm o} $ $(\sigma(F_{\rm o}))$	$ \mathit{F}_\mathit{c} $		
030	102(10)	66	0,16,0	834 (15)	857		
040	4260 (64)	4074	0.17.0	645 (13)	658		
050	207(11)	300	0,18,0	138 (26)	130		
060	0(27)	27	0.19.0	606 (13)	609		
070	655 (11)	660	0.20.0	43 (56)	62		
080	3337 (50)	3172	0,21,0	510 (13)	516		
090	603 (12)	635	0.22.0	0(40)	87		
0.10.0	0(30)	44	0.23,0	800 (16)	798		
0.11.0	489 (10)	517	0,24,0	(31) 0	20		
0,12,0	1159 (18)	1203	0.25.0	339 (17)	402		
0,13,0	982 (16)	969	0,26,0	138 (32)	82		
0.14,0	0(25)	37	0.27,0	379 (18)	392		
0,15,0	940 (16)	928					

center to 0, 0, 0 from \sim -0.2084, $-1/4$ (required), \sim -0.0754.)

The alert reader will now wonder why we have not made this transformation in order (a) to correctly define the structure and (b) to avoid being "Marshed"? The answer is simple. The raw data negate this possibility. We have taken the liberty of reproducing our *OkO* data in Table VI. Here it can clearly be seen that we do not have the required systematic absences $0k0$ for $k = 2n + 1$ necessary for a crystallographically precise $2₁$ axis along b. The profiles for the $0(2n + 1)0$ reflections are normal-i.e., similar to all other data. Furthermore, the final agreement between $|F_{o}|$ and $|F_{c}|$ for this group of reflections is excellent! **(An** interesting, but unrelated, observation is that *Ok0* reflections with $k = 4n$ are intense and those with $k = 4n + 2$ are very weak; this, of course, results from the four almost equally spaced and close-to-parallel layers of metal atoms along b giving rise to constructive and destructive interference, respectively, for these two groups of reflections.)

Thus, the structure of $\rm OsFe_2(CO)_{10}(\mu\text{-}CO)_2$ is correctly described in space group Pn but is close to the related supergroup $P2_1/n$. (Overall intensity statistics also predict the acentric case (vide supra).)

We now come to the crystal structure of $Fe₃(CO)₁₀(\mu$ -CO), (Figure **4).** The macroscopic structure has been treated as a disordered ensemble in space group $P2_1/n^{1,2}$ This structure has two distinct layers of molecules up the b axis. Arranged about the origin are the three atoms **of** the defined $Fe₃$ cluster; superposed on these is the inversion-related Fe3 cluster, forming **an** overall "star of David" pattern. The second layer is centered about $1/2$, $1/2$, and contains mutually opposed triangular Fe₃ units related (by an *n*-glide plane and a $2₁$ axis) to the defined Fe₃ cluster. The contents of the double cell *(a,* 2b, **c)** is the

⁽⁹⁾ This is a colloquialism heard recently and based upon Dr. R. E. rected by the inclusion of further elements of symmetry. See, for exam-
ple: Marsh, R. E. Organometallics **1989**, 8, 1583.

same as that for the idealized $(P2_1/n)$ cell of Fe₂Os- $(CO)_{10}(\mu$ -CO)₂ (after a change of origin (vide supra)), except that each Fe3 unit is **0.50:0.50** disordered, whereas the Fe₂Os units are only 0.925:0.075 disordered. The observed $Fe₃(CO)₁₀(\mu$ -CO)₂ structure can arise by two routes.

(1) As described by Wei and Dahl,' the exterior of the molecule (an icosahedron of 0 atoms) may be invariant to inversion of the molecule. In this case, each site may be occupied, on a statistical basis, by "normal" and "inverted" $Fe₃(CO)₁₀(\mu$ -CO)₂ molecules. However, Cotton and Troup2 found not only two sites for each carbon atom but also *two sites for each oxygen atom.* This indicates that the **12** exterior oxygen atoms of the "normal" and "inverted" molecules do not overlap precisely and that the polyhedron defined by the **12** oxygen atoms does not possess a perfect inversion center. This brings us to a second possible model, based upon the $Fe₂Os(CO)₁₀(\mu$ -CO)₂ structure.

(2) The microscopic packing of $Fe_3(CO)_{10}(\mu$ -CO)₂ molecules could be based on the double cell *a, 2b, c* in which the individual sites are ordered. **A** perfectly ordered domain is then possible if molecules stack as in the $Fe₂Os (CO)_{10}(\mu$ -CO)₂ structure. The overall composite *disordered macrostructure* can now be formed by twinning of these microdomains along the *b* axis.

Either of these (or a combination of the two) is possible. The first requires fewer assumptions (and thus is favored by Occham's razor); the second allows for the additional observation that the O_{12} periphery of $Fe_3(CO)_{10}(\mu$ -CO)₂ does not have a perfect inversion center and that a precise 0.50:0.50 occupancy is therefore not favored. It would be interesting to survey $Fe_3(CO)_{10}(\mu$ -CO)₂ crystals and search for a possible superlattice in their diffraction pattern, particularly at extremely low temperature, where the possibility of intramolecular rearrangement **is** reduced. We have previously observed one example of a structure with a well-defined subcell but with a diffuse superlattice doubling the axis length in one direction (i.e. $[(Me_6C_6)_3Nb_3Cl_3]Cl$).¹⁰

Thermal Parameters of the Iron Atoms in $Fe₂Os(CO)₁₀(\mu$ -CO)₂

It is clear from Figures 1 and **2** that iron atoms (but not the osmium atoms) have thermal parameters extended in the *b* direction. Thus, in molecule A, $B_{22}(Fe) = 5.87-6.28$ A^2 and $B_{22}(Os) = 2.91$ A^2 ; in molecule B, $B_{22}(Fe) =$ **4.98-5.41** \bar{A}^2 and $B_{22}(\text{Os}) = 2.65 \text{ Å}^2$. A similar situation also occurs in $Fe_3(C\ddot{\tilde{O}})_{10}(\mu$ -CO)₂.² Here (in different units) β_{22} (Fe, bridged) = 0.0140-0.0162 whereas β_{22} (Fe, non $bridgeed) = 0.0067$.

This latter fact was pointed out to M.R.C. by Dr. B. F. *G.* Johnson in a private conversation. He regarded it as important evidence for his proposed model for stereochemical nonrigidity of $Fe_3(CO)_{10}(\mu$ -CO)₂ involving the interchange of carbonyl ligands by a C_2 rotation¹¹ of the Fe₃ triangle as opposed to a C_3 rotation^{2,12,13} of the Fe₃ triangle relative to the external carbonyl groups. **A** recent variable-temperature Mössbauer-effect study on Fe₃- $(CO)_{10}(\mu$ -CO)₂¹⁴ has been reported and also indicates that a "low-temperature single-crystal X-ray study of $Fe_3(CO)_{12}$ " is being undertaken. The results of this study appear to be inconsistent with either a C_2 or C_3 rotation of the Fe₃ triangle.

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Supplementary Material Available: A listing of **anisotropic thermal parameters (1 page); a** listing **of observed and calculated structure factors (29 pages). Ordering information is given on any current masthead page.**

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