Ground-State and Excited-State C≡C Stretching **Frequencies of Dimetallapolyyne Redox Congeners.** Direct Evidence for Metal \rightarrow Alkynyl π -Back-Bonding

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Summary: Paramagnetic $[K(crypt-222)][Mo_2(C \equiv CSiMe_3)_4 (PMe_3)_4$] ([K(crypt-222)]1) is prepared by one-electron reduction of $Mo_2(C \equiv CSiMe_3)_4(PMe_3)_4$ (1) by $K(C_{10}H_8)$ in the presence of crypt-222. The $v(C \equiv C)$ frequencies for **1** and $[K(crypt-222)]\mathbf{1}$ in their ground states and $\delta \rightarrow \delta^*$ electronic excited states are controlled by the occupancy of the $\delta(M_2)$ and $\delta^*(M_2)$ orbitals, lying in the order **1** $(\delta^2) > \mathbf{1}^* (\delta^1 \delta^{*1}) > [K(crypt-222)]\mathbf{1} (\delta^2 \delta^{*1}) > [K(crypt$ $(222)|\mathbf{1}^*(\delta^1\delta^{*2});$ this order indicates that $M \rightarrow C \equiv CSiMe_3$ π -back-bonding is of increasing importance across this series.

A general description of the metal-carbon bond in transition-metal-alkynyl compounds has not yet been developed.^{1,2} This is despite the fact that understanding the nature of this bond is a logical prerequisite for the rational design of electronic materials composed of metal-alkynyl building blocks³⁻⁷ and for interpreting their physical properties. Particular uncertainty surrounds the contribution of M \rightarrow CCR π -back-bonding to the metal-carbon interaction. Although the degree to

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U. H. F. Angew. Chem., Int. Ed. Engl. 1996, 35, 969. (e) Bartik, T.;
Bartik, B.; Brady, M.; Dembinski, R.; Gladysz, J. A. Angew. Chem. Bartik, B.; Brady, M.; Dembinski, R.; Gladysz, J. A. Angew. Chem., Int. Ed. Engl. 1996, 35, 414. (f) Coat, F.; Lapinte, C. Organometallics 1996, 15, 477. (g) Yam, V. W.-W.; Lau, V. C.-Y.; Cheung, K.-K. Organometallics 1996, 15, 1740. (h) Weng, W.; Bartik, T.; Brady, M.; Bartik, B.; Ramsden, J. A.; Arif, A. M.; Gladysz, J. A. J. Am. Chem. Soc. 1995, 117, 11922. (i) Le Narvor, N.; Toupet, L.; Lapinte, C. J. Am. Chem. Soc. 1995, 117, 7129. (j) Brady, M.; Weng, W.; Gladysz, J. A. J. Chem. Soc. Chem. Commun. 1994, 2655. (k) Lang, H. Angew. Chem., Int. Ed. Engl. 1994, 33, 547. For a review of M-C₂-M compounds, see: Beck, W.; Niemer, B.; Wieser, M. Angew. Chem., Int. Ed. Engl. 1993, 32, 923. which alkynyl ligands act as π -acceptors has often been inferred from comparisons among C=C stretching frequencies of $ML_m(C \equiv CR)_n$ compounds, such interpretations are problematic because $\nu(C \equiv C)$ typically varies so little (a few tens of cm⁻¹) as a function of oxidation state and ancillary ligand that electronic effects other than M–CCR π interactions cannot be excluded from consideration.¹ Moreover, Lichtenberger and co-workers have recently reported definitive evidence from photoelectron-spectroscopic studies that alkynyl ligands are π -donors in FeCp(C=CR)(CO)₂ complexes and that their π -acceptor character is negligible,^{2,8} although, as they noted, more electron-rich compounds might exhibit M \rightarrow CCR π -back-bonding. Herein, we report the synthesis of an electron-rich compound, [Mo₂(C≡CSiMe₃)₄-(PMe₃)₄]⁻, that together with its one-electron-oxidized congener $Mo_2(C \equiv CSiMe_3)_4(PMe_3)_4$ (1)^{9,10} possess unique electronic attributes for testing this hypothesis: their $\delta(M_2)$ and $\delta^*(M_2)$ orbitals are of π -symmetry with respect to the $\pi^*(C \equiv CR)$ orbitals, so changing the δ, δ^* occupancy chemically and photochemically allows broad and systematic variation of the M–CCR π -electron configuration. The ground-state and $\delta(M_2) \rightarrow \delta^*(M_2)$ excited-state C=C stretching frequencies of these compounds provide the first unambiguous evidence for M \rightarrow CCR π -back-bonding.

The cyclic voltammogram of $Mo_2(C \equiv CSiMe_3)_4(PMe_3)_4$ exhibits a reversible, one-electron reduction wave at

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-2.13 V (vs FeCp₂/FeCp₂⁺) in tetrahydrofuran (THF),¹¹ which suggested that [Mo₂(C≡CSiMe₃)₄(PMe₃)₄]⁻ could be synthesized and isolated. Accordingly, addition of a THF solution of $K(C_{10}H_8)$ (1 equiv) to an equimolar mixture of royal blue, diamagnetic 1 and 4,7,13,16,21,-24-hexaoxa-1,10-diazabicyclo[8.8.8]hexacosane (crypt-222) in THF at -78 °C results in a red-brown solution, from which paramagnetic, orange-brown [K(crypt-222)]- $[Mo_2(C \equiv CSiMe_3)_4(PMe_3)_4]$ ([K(crypt-222)]1) can be isolated in ca. 65% yield upon removal of solvent and extraction of naphthalene with pentane.¹² [K(crypt-222)]1 reacts instantly with air in both the solid state and solution to form a dark blue product that consists primarily of 1, as indicated by UV-vis and infrared spectroscopy. The infrared spectrum of [K(crypt-222)]-**1** exhibits a single, strong band in the $\nu(C \equiv C)$ region at 1954 cm^{-1} ; for 1, the corresponding band appears at 1991 cm^{-1,10,13} We expect that the structure of $1^$ possesses idealized D_{2d} symmetry, like that of **1**, based on the fact that the D_{2d} symmetry of related compounds of the type $[Re_2Cl_4(PR_3)_4]^{2+}$ is maintained upon both one- and two-electron reduction.14

Frequencies for ν (C=C) in the $\delta \rightarrow \delta^*$ electronic excited states of 1 and [K(crypt-222)]1 were determined from electronic absorption spectra of the compounds at 10 K. These spectra exhibit vibronically structured bands attributable to the ${}^{1}(\delta \rightarrow \delta^{*})$ and ${}^{2}(\delta \rightarrow \delta^{*})$ transitions of 1 (Figure 1) and [K(crypt-222)]1 (Figure 2), respectively, as the lowest-energy features. Aside from the fact that the ${}^{2}(\delta \rightarrow \delta^{*})$ transition of [K(crypt-222)]1 is red shifted by ca. 7000 cm⁻¹ from the $(\delta \rightarrow \delta)$ δ^*) transition of **1**, as a result of the larger spin-pairing energy contributions to the latter,¹⁵ the two bands are qualitatively similar: both are dominated by a 350-370 cm⁻¹ vibronic progression that is attributable to nominal ν (Mo–Mo) $/\nu$ (Mo–C) modes¹⁶ and, additionally, exhibit weaker, lower frequency vibronic origins (ca.



Figure 1. Electronic absorption spectrum of a polycrystalline film of 1 at ca. 10 K. Vibronic spacings are given in cm⁻¹. The trace in the upper right corner is an expansion of the ν (C=C) origin and associated features.



Figure 2. Electronic absorption spectrum of a polycrystalline film of [K(crypt-222)]1 at ca. 10 K. Vibronic spacings are given in cm⁻¹. The trace in the upper right corner is an expansion of the ν (C=C) origin and associated features. The dip in intensity at 7850 cm^{-1} , marked by an asterisk, is an instrumental artifact.

160/170, 240 cm⁻¹). Near the high-energy edge of each band-beyond the point at which features attributable to combinations of the lower-frequency fundamentals are still clearly resolved-an additional, weak vibronic

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⁽¹²⁾ Anal. Calcd (found) for $C_{50}H_{108}KMo_2N_2O_6P_4Si_4$: C, 46.18 (46.46); H, 8.37 (8.06); N, 2.15 (2.09). IR, ν (C=C) region (neat film): 1954 cm⁻¹ (s). EPR (solid, 295 K): g = 2.05.

⁽¹³⁾ The C=C stretching coordinates of D_{2d} -symmetric M₂(CCR)₄- $(PMe_3)_4$ transform according to $a_1 + b_2 + e_2$ of which b_2 and e are allowed in the infrared spectrum. The infrared bands of 1 and [K(crypt-222)]1 at 1991 and 1954 cm⁻¹, respectively, are assigned to e-symmetric C=C stretches on the basis that these should be more intense than the b₂-symmetric modes; a vector-sum model predicts $[I(e)/I(b_2)] = \frac{1}{2}$ - $(\tan^2 \theta) \cong 2$ (where θ is the angle between the Mo–Mo and the C=C bonds, taken from the structure of Mo₂(CCH)₄(PMe₃)₄ in ref 9c). The FT-Raman spectrum of 1 exhibits a band at 1999 cm⁻¹ attributable to the a₁-symmetric ν (C=C) mode. Repeated attempts to measure the ν -(C=C) frequency for [K(crypt-222)]1 by Raman spectroscopy have been (14) Cotton, F. A.; Dunbar, K. R.; Falvello, L. R.; Tomas, M.; Walton, R. A. J. Am. Chem. Soc. 1983, 105, 4950–4954.

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⁽¹⁶⁾ Normal-coordinate calculations on 1 and the isotopomers Mo₂- $(C \equiv CSiMe_3 - d_3)_4$ (PMe₃)₄ and Mo₂(¹³C \equiv ¹³CSiMe₃)₄(PMe₃)₄ (John, K. D.; Miskowski, V. M.; Vance, M. A.; Dallinger, R. F.; Wang, L. C.; Geib, S. J.; Hopkins, M. D. Manuscript in preparation) indicate that there is significant coupling between the symmetric ν (Mo–Mo) and ν (Mo–C) stretching modes, but that ν (C=C) is uncoupled from other modes (>90% C=C stretching in character).

origin is observed, the spacing of which relative to the 0–0 band is 1970 cm⁻¹ for **1** and 1890 cm⁻¹ for [K(crypt-222)]**1**. Progressions built on these origins arising from the 350/370 cm⁻¹ mode are resolved, as are features attributable to the lower-frequency modes. Isotopic labeling of the alkynyl ligands of **1** with ¹³C allows the 1970 cm⁻¹ origin of this compound to be unambiguously assigned to one quantum of the ν (C=C) stretching mode, as this origin is shifted to 1900 cm⁻¹ in the spectrum of Mo₂(¹³C=¹³CSiMe₃)₄(PMe₃)₄; this frequency is in close agreement with what is expected from the diatomicoscillator approximation (1890 cm⁻¹). By analogy, the 1890 cm⁻¹ origin observed in the spectrum of [K(crypt-222)]**1** is also assigned to ν (C=C).

The relationship between the frequency of ν (C=C) and the $[\delta, \delta^*(M_2)]$ electron configuration of **1** and [K(crypt-222)]1 provides direct evidence for $M \rightarrow CCR \pi$ -backbonding. The ground states and $\delta \rightarrow \delta^*$ excited states (1*, [K(crypt-222)]1*) of the two compounds are described by four different configurations of the δ and δ^* orbitals: **1**, $[\sigma^2 \pi^4 \delta^2]$; **1***, $[\sigma^2 \pi^4 \delta^1 \delta^{*1}]$; [K(crypt-222)]**1**, $[\sigma^2 \pi^4 \delta^2 \delta^{*1}]$; [K(crypt-222)]**1***, $[\sigma^2 \pi^4 \delta^1 \delta^{*2}]$. The b₁-symmetry δ and a₂-symmetry δ^* orbitals are of π symmetry with respect to the ligands and, hence, are capable of interacting with the $\pi(CCR)$ and $\pi^*(CCR)$ orbitals (a₂ + b₁ + e), with the δ^* orbital being closer in energy to the $\pi^*(CCR)$ levels than is the δ orbital. Importantly, increasing the occupancy of the δ^* orbital and raising the energy of δ and δ^* by reduction of Mo₂⁴⁺ to Mo₂³⁺ results in a significant, systematic reduction in the frequency of ν (C=C), indicating that M→CCSiMe₃ π -back-bonding increases concomitantly. Context for



these frequencies and the differences among them is provided by the ca. 80 previously reported $\nu(C=C)$ frequencies for $[ML_m(C=CSiMe_3)_n]^{z+}$ complexes (M = group 4 to group 10, d^0-d^8 ; z = +1 to -2; Figure 3), all but three of which lie in a narrow 80 cm⁻¹ range (1990– 2070 cm⁻¹). In contrast, the $\nu(C=C)$ frequencies of **1** and [K(crypt-222)]**1*** differ by >100 cm⁻¹, with that for [K(crypt-222)]**1*** being, by >50 cm⁻¹, the lowest yet reported; the latter frequency is suggestive of significant double-bond character in the alkynyl C–C linkages.



Figure 3. Distribution of ν (C=C) stretching frequencies for ML_m(C=CSiMe₃)_n complexes. A listing of compounds, frequencies, and references is available as Supporting Information. Ground-state and excited-state frequencies for **1** and [K(crypt-222)]**1** are shown by unshaded bars.

In summary, we have demonstrated that $M\rightarrow CCR \pi$ -back-bonding can contribute substantially to the metal-carbon interaction in metal-alkynyl compounds. It is important to note that the present data do not rule out contributions from $M\leftarrow CCR \pi$ -bonding in these compounds; unfortunately, attempts to prepare $[Mo_2-(C\equiv CSiMe_3)_4(PMe_3)_4]^+$, for which such interactions could be particularly important, have not been successful. The observation that the extent of $M\rightarrow CCR \pi$ -back-bonding is markedly different in the ground and excited states of **1** and **1**⁻ has important implications for metal-alkynyl electronic materials, as this opens the possibility of photochemically switching those properties that are electronically governed by the $M-CCR \pi$ interactions.

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Supporting Information Available: Tables of ν (C=C) frequencies for ML_m(C=CSiMe₃)_n complexes, with references (6 pages). Ordering information is given on any current masthead page.

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