Comparative Solid-State NMR Study of the Molecular Dynamics of Si(SiMe3)4 and C(SiMe3)4

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Variable-temperature 13C and 29Si one- and two-dimensional MAS NMR experiments are reported for the low-temperature phases of Si(SiMe3)4 (**1**) and C(SiMe3)4 (**2**). Quantitative determination of exchange rate constants from 13C and 29Si NMR spectra of **1** and **2** proves considerable dissimilarities of the molecular solid-state dynamic properties of **1** and **2**, despite the chemically and structurally homologous nature of these two compounds. In the lowtemperature phase of 1, internal SiMe₃ and whole-molecule reorientation occurs, while for the low-temperature phase of **2** no internal SiMe₃ reorientation but only whole-molecule reorientation is observed.

Introduction

Compounds $E(SiMe₃)₄$ (with $E = Si$, C; Me = methyl), which may be viewed as higher homologues of methane, CH4, belong to a class of solids with nearly spherically shaped and nonpolar molecular units. Such molecular solids are known to undergo structural solid-state phase transitions, often including a plastically crystalline state ("rotator phase") at elevated temperatures. These structural phase transitions are closely related to thermally activated solid-state molecular dynamics such as internal modes of reorientation or whole-molecule reorientational modes.¹

Solid Si(SiMe₃)₄, **1**, and C(SiMe₃)₄, **2**, exist as strongly orientationally disordered phases at ambient temperatures and undergo structural phase transitions upon cooling. The low-temperature phase of **1** exists at temperatures $T < 238-241$ K;² for **2** the low-temperature phase is stable up to $T = 213 \text{ K}^3$. Neither for 1 nor for **2** have single-crystal X-ray diffraction investigations been reported; the gas-phase structures of **1** and **2** have been examined by electron diffraction.4,5 Empirical force field calculations6,7 have indicated that **1** and **2** exhibit molecular *T* symmetry in the gas phase. On the basis of previous variable-temperature 13 C and 29Si CP/MAS NMR results, molecular point group symmetry *C*³ has been proposed for the low-temperature phases of **1** and **2**; similar dynamic solid-state properties of **1** and **2**, that is, the occurrence of both internal SiMe3 reorientation and reorientation of the entire molecule around preferred axes, have been inferred from these NMR results.^{8,9}

Here we will focus on the quantitative analysis of variable-temperature one- and two-dimensional 13C and 29Si MAS NMR experiments on the low-temperature phases of **1** and **2**. It will be shown that quantitative determination of exchange rate constants as a function of temperature from 13C and 29Si MAS NMR experiments on **1** and **2** is necessary for unambiguous identification of different reorientational modes. Quantification of exchange rate constants in addition reveals distinct differences in the thermally activated molecular solid-state dynamics of the low-temperature phases of **1** and **2**.

Results and Discussion

In the following we will concentrate on the molecular dynamic properties of the low-temperature phases of compounds $Si(SiMe₃)₄$, **1**, and $C(SiMe₃)₄$, **2**, as viewed by solid-state ¹³C and ²⁹Si NMR. The low-temperature phases of **1** and **2** have recently been shown, by means of synchrotron powder diffraction techniques, to crystallize in space group *P*213,10 confirming molecular *C*³ point group symmetry. The possible modes of reorientation of a molecule $E(SiMe₃)₄$ with molecular $C₃$ symmetry (apart from internal reorientation of the methyl groups themselves) are illustrated schematically in Chart 1. Two different reorientational modes have to be taken into account: internal reorientation of the SiMe_3 groups around the respective $E-Si$ bond $(E = Si, C)$ directions, and reorientational jumps of the entire molecule around

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Chart 1. Schematic Representation of Molecular Reorientational Modes in the Low-Temperature Phases of E(SiMe₃)₄, E = C, Si, with Molecular C_3

 $a \, k$ and K denote exchange rate constants for internal SiMe₃ (*k*) and for whole-molecule (*k*′) reorientation. The numbering scheme $Me^{(1)}$ to $Me^{(4)}$ and $Si^{(1)}$, $Si^{(2)}$ denotes symmetryequivalent/inequivalent sites. See text for detailed description.

preferred molecular axes, also coincident with the E-Si bond directions.

From the schematic representation of a molecule $E(SiMe₃)₄$ with molecular $C₃$ symmetry in Chart 1, we expect the following ${}^{13}C$ and ${}^{29}Si$ MAS NMR features for the low-temperature phases of **1** and **2**. At sufficiently low temperatures, in the so-called slowexchange regime (on the one-dimensional ¹³C and ²⁹Si MAS NMR time scale), we expect two ²⁹Si resonances in a 1:3 relative intensity ratio for the $\rm Si^{(1)}Me_3$ and $\rm Si^{(2)}$ -Me3 groups. The 12 methyl groups will give rise to four 13C resonances in a relative 1:1:1:1 intensity ratio for the four crystallographically distinct types of methylcarbon sites $Me^{(1)}$ to $Me^{(4)}$ (see Chart 1). For reasons of symmetry, exchange rate constants $k_1 = k_2 = k_3$ for internal $Si^{(2)}Me_3$ reorientation and $K_1 = K_2 = K_3$ for whole-molecule reorientation, while exchange rate constants k_4 and k'_4 for internal and whole-molecule reorientation around the molecular 3-fold axis of symmetry may differ from k_1 , k_2 , k_3 and k'_1 , k'_2 , k'_3 since the $Si^{(1)}Me^{(4)}$ ₃ group is crystallographically distinct from the remaining three $Si^{(2)}Me^{(1)}Me^{(2)}Me^{(3)}$ groups. In fact, unequal exchange rate constants for internal reorientation of chemically equivalent but crystallographically inequivalent sites ER_3 in solid organosilicon compounds have previously been reported for moieties $ER_3 = \frac{B1^{11}}{B1^{11}}$
and $ER_2 = \text{SiMe}_2^{-12,13}$. Under conditions of fast magicand $ER_3 = SiMe_3$.^{12,13} Under conditions of fast magic-
angle spinning, exchange rate constants k_t and k' , for angle spinning, exchange rate constants *k*⁴ and *k*′⁴ for **1** and **2** cannot be measured: reorientations around the proper 3-fold axis of symmetry lead to mutual exchange of equivalent sites (mutual interchange of $Me⁽⁴⁾$ sites for k_4 , mutual interchange of Me⁽⁴⁾ sites and of Si⁽²⁾Me⁽¹⁾- $Me^{(2)}Me^{(3)}$ groups for K_4) and hence do not contribute to one- and two-dimensional ¹³C and ²⁹Si NMR spectral line shapes. In principle, it would be possible to also

Figure 1. 2D 29Si EXSY experiments on **1** and **2**: (a) **1**, *T* = 160 K, mixing time τ_{mix} = 300 ms, only the SiMe₃ region is shown; (b) **2**, $T = 170$ K, mixing time $\tau_{mix} = 75$ ms. For long mixing times relative intensities corresponding to complete exchange in a four-site exchange (diagonal 1:9 and off-diagonal 3:3) are observed for **1** and **2**.

determine k_4 and k'_4 from MAS NMR experiments recorded under slow magic-angle spinning conditions. For instance, the ODESSA experiment, 14 or a timereversed version thereof,¹⁵ would be a suitable approach. In practice, however, it turns out impossible to determine k_4 and k'_4 because for **1** and **2** very slow magicangle spinning would be necessary ($\omega_{\rm rot}/2\pi < 200$ Hz), resulting in severe spectral overlap problems.

Under fast magic-angle spinning conditions, for **1** and **2** we can only determine the exchange rate constants *k* k_1, k_2, k_3 and $k' = k'_1, k'_2, k'_3$ (see Chart 1) as these exchange processes lead to mutual exchange of inequivalent sites. We have to be able to distinguish internal SiMe₃ reorientation (*k*) from whole-molecule reorientation (*k*′). Under fast-spinning MAS conditions, the two ²⁹Si resonances of the SiMe₃ groups selectively

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b.

Figure 2. Variable-temperature ²⁹Si MAS NMR spectra of 1 (a) and 2 (b); only the SiMe₃ region is shown. Temperatures are indicated and $*$ denotes the presence of resonances from the coexisting high-temperature phase. Upper traces represent experimental spectra; lower traces are the corresponding best-fit calculated spectra.

monitor the exchange rate constants *k*′ for wholemolecule reorientation, that is, interchanging $Si^{(1)}$ with Si⁽²⁾. Formally, this exchange process may be described as a four-site exchange with three degenerate sites. 13C one- and two-dimensional MAS NMR spectra will be affected by both molecular reorientational modes. Internal $Si^{(2)}Me_3$ reorientation only interchanges the Me⁽¹⁾, $Me^{(2)}$, Me⁽³⁾ sites with rate constant *k*, while wholemolecule reorientation with rate constant *k*′ leads to mutual interchange of Me⁽¹⁾, Me⁽²⁾, Me⁽³⁾ sites, as well as to interchange of $(Me^{(1)}, Me^{(2)}, Me^{(3)})$ sites with the $Me⁽⁴⁾$ sites. Only reorientation of the entire molecule can lead to exchange between all four 13 C-methyl resonances. Exchange rate constants *k*, *k*′ for two simultaneously occurring processes are additive;¹⁶ hence, the simultaneous presence of both reorientational modes would lead to, for example, exchange peaks of unequal integrated intensities in a two-dimensional 13C exchange MAS NMR spectrum (2D EXSY). Since the 13C resonance of the $Me^{(4)}$ sites is only affected by wholemolecule reorientation, it would always be possible to determine *k*′ directly from the relative integrated intensities of those exchange peaks involving the ^{13}C resonance of the $Me^{(4)}$ sites. Alternatively, the observation of exchange between all four 13C-methyl resonances with *all* exchange peaks of identical integrated intensities is consistent only with the presence of wholemolecule reorientation in the absence of additional internal $Si^{(2)}$ Me₃ reorientation.

We may now inspect more closely the various one- and two-dimensional, variable-temperature ¹³C and ²⁹Si NMR results on the low-temperature phases of **1** and **2**. In accord with the predictions for molecular *C*³ symmetry from the above discussion of Chart 1, ²⁹Si MAS NMR spectra of **1** and **2** at low temperatures display two 29Si resonances in a relative intensity ratio of 1:3 for the $Si^{(1)}Me_3$ and $Si^{(2)}Me_3$ groups. This slowexchange limit on the one-dimensional 29Si NMR time scale is reached at temperatures $T \le 190$ K for 1 and at $T \le 170$ K for **2**. ²⁹Si 2D EXSY experiments on **1** and **2** in the respective 29Si slow-exchange limit provide proof that for both compounds, even at low temperatures, whole-molecule reorientation does take place: in both cases off-diagonal exchange peaks connecting the two 29Si resonances are observed (see Figure 1). Exchange rate constants *k*′ for whole-molecule reorientation as a function of temperature in this slow-exchange regime for **1** and **2** are obtained from a series of 29Si 2D EXSY experiments, recorded at different temperatures and with different mixing times. At temperatures $T \geq$ 190 K (for **1**) or $T \ge 170$ K (for **2**) the line shapes of one-dimensional 29Si MAS NMR spectra become affected by whole-molecule reorientation, as can be seen in Figure 2. Spectral line shape simulations of the exchange-broadened 29Si resonances yield exchange rate constants *k*′ for whole-molecule reorientation of **1** and **2** up to the temperature ranges where the onset of the phase transition to the high(er)-temperature phase starts to interfer by the coexistence of high- and lowtemperature phase. For **1**, the one-dimensional 29Si spectral line shapes can be exploited to determine exchange rate constants in the temperature range $T \approx$ ¹⁹⁰-215 K; for **²**, the corresponding informative temperature range is $T \approx 170-195$ K (see Figure 2).

In contrast to the one- and two-dimensional variabletemperature 29Si MAS NMR experiments, selectively monitoring exchange rate constants *k*′ of whole-molecule

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Figure 3. 2D ¹³C EXSY experiments on **1** and **2**; only the methyl region is shown: (a) **1**, $T = 154$ K, mixing time $\tau_{\text{mix}} =$ 500 ms; (b) **2**, $T = 164$ K, mixing time $\tau_{mix} = 250$ ms; (c) 1, $T = 190$ K, mixing time $\tau_{mix} = 1$ s.

reorientation in the low-temperature phase of **1** and **2**, 13C MAS NMR experiments can report on the kinetics of internal $Si^{(2)}Me_3$ reorientation as well as on wholemolecule reorientation. In agreement with the requirements of molecular *C*³ symmetry for the low-temperature phases of **1** and **2**, four 13C-methyl resonances of relative intensities 1:1:1:1 are observed for both compounds in 13C CP/MAS NMR spectra, recorded in the slow-exchange temperature regimes on the one-dimensional 13 C NMR time scale. This temperature regime is reached at $T \le 150$ K for **1** and at $T \le 170$ K for **2**. 2D 13C EXSY experiments on **1** and **2** at low temperatures reveal major differences in the molecular dynamics of **1** and **2**. As illustrated in Figure 3, for **1** only exchange between three methyl sites, $Me^{(1)}$, $Me^{(2)}$, $Me^{(3)}$, is observed at low temperatures, while exchange between all methyl sites $Me^{(1)}$ to $Me^{(4)}$ occurs for 2. Obviously, for 1 at low temperatures only internal $Si^{(2)}$ -Me3 reorientation takes place with nonnegligible exchange rate constants *k*; at higher temperatures additional whole-molecule reorientation occurs (see Figure 3c). In solid **2** the only low-temperature dynamic process occurring with substantial exchange rate constants *k*′ is whole-molecule reorientation. The integrated intensities of all 12 exchange peaks in the 2D 13C EXSY spectra of **2** are equal within experimental error and thus exclude the simultaneous presence of additional internal $Si^{(2)}Me₃$ reorientation. The dissimilarities in the solid-state molecular dynamics of **1** and **2** are also evident in the temperature range where one-dimensional 13C MAS NMR spectra display temperature-dependent exchange broadening. Variabletemperature 13C MAS NMR spectra of **1** and **2** in this temperature range are shown in Figure 4. For **1** there are two different coalescence regimes: one occurs in the $T = 160 - 180$ K temperature range and affects only the three ¹³C-methyl resonances for sites $Me^{(1)}$, $Me^{(2)}$, $Me^{(3)}$ (internal $Si^{(2)}$ Me₃ reorientation); a second coalescence regime is reached in the temperature range $T = 180-$ 210 K, brought about by the additional onset of wholemolecule reorientation (see Figure 4a). For **2** the evolution of the one-dimensional ¹³C MAS NMR spectral line shape as a function of temperature is quite different (see Figure 4b). Over the temperature range $T = 160-$ 190 K all four 13C-methyl resonances are part of the emerging coalescence pattern. From qualitative inspection of this coalescence pattern we cannot safely exclude that at somewhat elevated temperatures also for **2** additional internal Si⁽²⁾Me₃ reorientation may take place as well.

However, if we take into consideration *all* exchange rate constants *k*, *k*′ determined from 2D 13C and 29Si

b.

Figure 4. Variable-temperature 13C MAS NMR spectra of **1** (a) and **2** (b); in b only the methyl region is shown. Temperatures are indicated, and $*$ denotes the presence of resonances from the coexisting high-temperature phase. Upper traces represent experimental spectra; lower traces are the corresponding best-fit calculated spectra.

Figure 5. Semilogarithmic plots of exchange rate constants *k*, *k*′ for **1** (a) and **2** (b) versus 1/*T*. The symbols refer to exchange rate constants determined from ²⁹Si MAS (*), ²⁹Si 2D EXSY (\times), ¹³C MAS (\diamond), and ¹³C 2D EXSY (\triangle) NMR experiments, respectively.

EXSY experiments and from iterative fitting of the spectral line shapes of variable-temperature onedimensional 13C and 29Si MAS NMR spectra, the following picture is obtained for the low-temperature phases of **1** and **2**. Figure 5 displays semilogarithmic plots of all exchange rate constants *k*, *k*′ versus 1/*T* for the low-temperature phases of **1** and **2**. The Arrhenius plot for **1** (Figure 5a) clearly shows the two exchange rate constants *k* for internal Si⁽²⁾Me₃ reorientation and *k*′ for whole-molecule reorientation to differ considerably. Note that the exchange rate constant *k* for internal $Si^{(2)}$ Me₃ reorientation in the low-temperature phase of **1** is always higher than *k*′ and that *k* exhibits

a significant deviation from Arrhenius behavior at elevated temperatures. From the data displayed in Figure 5a, the activation barrier for internal $Si^{(2)}Me₃$ reorientation in 1 at temperatures up to $T \approx 185$ K is determined as $E_a = 53.3 \pm 6.0 \text{ kJ} \text{ mol}^{-1}$. The exchange rate constants *k*′ accessible for whole-molecule reorientation of **1** permit determination of the activation barrier for this process as $E_a = 49.3 \pm 3.0 \text{ kJ} \text{ mol}^{-1}$. A different situation is found in the Arrhenius plot for **2** (see Figure 5b). Over the entire accessible temperature range, the exchange rate constants at a given temperature, as obtained from 13C *and* 29Si NMR data, are identical within experimental error, leading to the conclusion that

the only dynamic process occurring in the low-temperature phase of **2** is whole-molecule reorientation with an activation barrier $E_a = 44.7 \pm 3.0 \text{ kJ mol}^{-1}$.

We need to answer the question why there should be such differences in the intramolecular (internal $Si^{(2)}Me₃$ reorientation) solid-state dynamics of **1** and **2**. Obviously, the difference in intramolecular steric crowding in $E(SiMe₃)₄$ on going from $E = Si (1)$ to $E = C (2)$, caused by different bond lengths E-Si (236 pm (**1**) and 192 pm (**2**)10), has a significant impact on the molecular dynamic properties of the two compounds. It renders internal $Si^{(2)}$ Me₃ reorientation in $C(SiMe₃)₄$, **2**, a highly unfavorable process, while "conventional" solid-state dynamic properties (where internal Si⁽²⁾Me₃ reorientation is thermally activated at lower temperatures, followed by the onset of global reorientation at more elevated temperatures) are found for the low-temperature phase of $Si(SiMe₃)₄$, **1**. It appears reasonable to consider intramolecular steric strain as the dominating cause of hindrance to internal $Si^{(2)}Me₃$ reorientation in solid **1** and **2**, although crystal-packing effects should not be completely ignored. With the crystal structure of the low-temperature phases of **1** and **2** determined,10 it becomes feasible to compare ab initio and force-field calculations based on the gas-phase and crystal structures in order to understand better the molecular dynamic properties of **1** and **2**, as well as the relationship between solid-state molecular dynamics and structural phase transitions. This work is in progress and will be reported elsewhere.¹⁰

Experimental Section

Compound 1, Si(SiMe₃)₄, is commercially available (ABCR Chemikalien, Karlsruhe) and has been used without further purification. Compound 2, C(SiMe₃)₄, was kindly donated by G. Fritz, Karlsruhe.

All 13C and 29Si NMR experiments were carried out on a Bruker MSL 300 NMR spectrometer, corresponding to Larmor frequencies of 75.5 (13 C) and 59.6 (29 Si) MHz, respectively. Standard 4 and 7 mm $ZrO₂$ rotors and double-bearing probes were used. For variable-temperature MAS NMR experiments, purified N_2 gas served as the drive and bearing gas. MAS frequencies were in the range $2-3$ kHz, ¹H $\pi/2$ pulse durations were 3-⁵ *^µ*s, recycle delays had to be 3-5 s, and the contact times for Hartmann-Hahn cross-polarization (¹³C and ²⁹Si) were 5 ms. 2D exchange spectroscopy (EXSY)¹⁶⁻¹⁹ employed phase cycling according to the TPPI method.²⁰ Spectral line shape simulations of the exchange-broadened 13C and 29Si MAS NMR spectra of **1** and **2** were carried out as described elsewhere;^{11,21} iterative fitting employed the MATLAB²² simplex routine for nonlinear least-squares minimization.

Isotropic chemical shifts δ (¹³C) and δ (²⁹Si) are quoted with respect to external SiMe4. Selected data are as follows. **1**: ¹³C, 3.5 ppm ($T = 297$ K), and 4.6, 3.0, 2.4, 1.5 ppm ($T = 154$) K); ²⁹Si, -9.9 (SiMe₃), -135.6 ppm ($T = 297$ K), and -9.3, -10.0 (SiMe₃), -140.9 ppm (*T* = 168 K). **2**: ¹³C, 6.4 (Me), -1.2 ppm (*T* = 297 K), and 7.1, 5.9, 5.0, 4.5 (Me), -2.6 ppm (*T* = 164 K); ²⁹Si, -1.5 ppm (*T* = 297 K), and -0.9, -1.9 ppm (*T* = 163 K).

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