

Influence of Cis and Trans Ligands in Platinum(II) Complexes on the Ability of the Platinum Center to Activate C–H Bonds. A Density Functional Theory Study

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We have studied the influence of different ligands X (X = F, Cl, Br, I, NO₂, and CN) on the C–H bond activation of CH₄ in *trans*-PtCl₂X(CH₄)[−], **1**, and *trans*-PtClX₂(CH₄)[−], **2**, where X is either *trans* (**1**) or *cis* (**2**) to methane. For **1** with X in the *trans* position, the *trans*-PtCl₂X[−] fragment interacts with CH₄ through the overlap between the empty d_σ-based orbital 2a₁ pointing along the Pt–X direction and σ_{CH} on CH₄. An interaction also takes place between an occupied d_σ-based orbital 1b₁ and the empty σ*_{CH} orbital on CH₄, where the d_π metal orbital is positioned perpendicular to the PtCl₂X[−] plane. The d_σ metal orbital contribution in 2a₁ is antibonding with respect to σ_x on X, whereas d_π in 1b₁ is antibonding with respect to π_x. Through the series F, Cl, Br, I, NO₂, and CN, the energies of σ_x and π_x increase. This is mostly an electronegativity effect. The increase in energy causes an increase in the contribution from σ_x and π_x to 2a₁ and 1b₁, respectively. As a consequence, the bonding overlaps ⟨σ_{CH}|2a₁⟩ and ⟨σ*_{CH}|1b₁⟩ will diminish, as only the d-component in 2a₁ and 1b₁ contributes to the overlap. As a result of the decreasing bonding overlaps, the Pt–CH₄ bond strength will decline. It is thus shown that the experimentally established order of *trans*-labilizing power for the series of ligands X studied here, F < Cl < Br < I < NO₂ < CN, can be related to the orbital energies of σ_x and π_x and the electronegativity of the elements that are involved in these orbitals. The labilization of the Pt–CH₄ bond in the C–H activation transition state is even larger than in the adduct **1**, leading to an increase in the C–H activation barrier along the series F < Cl < Br < I < NO₂ < CN. For **2** with X in the *cis* position, solvation has the largest influence on trends in the Pt–CH₄ bond for both **2** and the transition state. However C–H activation barriers are quite similar for different X.

1. Introduction

In the Shilov reaction,^{1,2} methane replaces H₂O in *cis*-PtCl₂(H₂O)₂ to form *cis*-PtCl₂(H₂O)(CH₄) followed by C–H activation. We have found in a recent theoretical study³ that the rate-determining step of methane activation by PtCl₂(H₂O)₂ is the uptake of CH₄ via an associative S_N2 substitution reaction rather than the C–H activation step. Shteinman et al.⁴ have extended the scope of the original Shilov reaction to involve the activation of methane by *cis*-PtX₂(H₂O)₂ with X = F, Cl, Br, I, NO₂, and CN. We find in a theoretical analysis⁵ of the Shteinman experiment that the uptake of CH₄ to form *cis*-PtX₂(H₂O)(CH₄) rather than C–H activation again is rate-determining for methane activation, independent of X. It was however found that the barrier of C–H activation in *cis*-PtX₂(H₂O)(CH₄) increases through the series F < Cl < Br < I < NO₂ < CN. We shall in the current investigation examine in more detail how the different ligand, X, influences the C–H activation barrier when placed either *cis* or *trans* to methane. Our analysis will be based on a bond energy decomposition

scheme that takes into account both steric and electron interactions between CH₄ and the platinum center with its co-ligands including X. A similar bond decomposition analysis will be performed to understand how X influences the barrier in the S_N2 substitution reaction that replaces H₂O with CH₄. As an outcome of this bond decomposition analysis, we shall provide new insight into the factors that make X a strong or weak *trans*-directing ligand in the ground state as well as the transition state.

2. Computational Details

Results were obtained from DFT calculations based on the Becke–Perdew exchange–correlation functional,^{6–8} using the Amsterdam Density Functional (ADF) program.⁹ A standard double-ζ STO basis with one set of polarization functions was applied for the H, N, F, Cl, Br, C, and O atoms, while a standard triple-ζ basis set was employed for the I and Pt atoms. The 1s electrons of N, C, F, and O, as well as the 1s–2p electrons of Cl, 1s–3d electrons of Br, 1s–4d electrons of I, and 1s–4f electrons of Pt, were treated as a frozen core. A standard set of auxiliary s, p, d, f, and g STO functions, centered on each

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nucleus, was used to fit the electron density and calculate the Coulomb and exchange potentials in each SCF cycle. The reported energies include first-order scalar relativistic corrections.¹⁰ Gas-phase electronic enthalpies were calculated from the Kohn–Sham energies. For all free energy values reported, the electronic entropy was neglected and standard expressions¹¹ were used to calculate the remaining gas-phase enthalpic and entropic contributions at nonzero temperature, as well as the zero-point vibrational contribution.

Thermodynamic parameters for the solvation of the halogen ions were obtained from a recent compilation of experimental values.¹² The remaining solvation enthalpies were obtained using the COSMO method¹³ as implemented in ADF.¹⁴ The solvent excluding surface was used along with an epsilon value of 78.5 for the dielectric constant of water as the solvent. Atomic radii employed were 1.39, 1.8, 1.16, 1.4, 1.3, 2.3, 1.96, 1.33, and 2.2 Å for Pt, Cl, H, N, O, C, Br, F, and I, respectively. Although the Born energy reported by the COSMO model is, strictly speaking, a free energy, the entropic contribution amounts to perhaps 2% of the total energy.¹⁵ The solvation enthalpy was therefore taken as the difference between the gas-phase energy and that calculated using the COSMO solvation model.

For the purpose of calculating the remaining solvation entropies, the solvation process was broken up into three steps, following Wertz.¹⁶ Here, the solute in the gas phase is first compressed to the molar volume of the solvent. The compressed solute gas then loses the same fraction of its entropy as would be lost by the solvent in going from gas (at its liquid-phase density) to liquid. Finally, the solute gas is extended to the concentration of the desired solution (i.e., 1.0 mol/L).

The solute entropy change for the first and third steps, which are strictly changes in molar volume, is given by $\Delta S = R \ln(V_{m,f}/V_{m,i})$, where $V_{m,f}$ is the final solute molar volume and $V_{m,i}$ is the initial solute value. The entropy fraction α lost in the second step can be determined from the absolute entropies of the solvent in its gas (S_{gas}°) and liquid (S_{liq}°) phases, as shown in eq 1.

$$\alpha = \frac{S^\circ(S^\circ + R \ln V_{m,\text{liq}}/V_{m,\text{gas}})}{(S^\circ + R \ln V_{m,\text{liq}}/V_{m,\text{gas}})} \quad (1)$$

Substituting the appropriate parameters for water¹⁷ affords the value of $\alpha = -0.46$. The sum of the entropy changes accompanying each of the three steps then gives the total solvation entropy; at a temperature of 298.15 K, we have (again, for water) eq 2.

$$\Delta S_{\text{sol}} = (-14.3 \text{ cal} \cdot \text{mol}^{-1} \cdot \text{K}^{-1}) - 0.46(S_{\text{gas}}^\circ - 14.3 \text{ cal} \cdot \text{mol}^{-1} \cdot \text{K}^{-1}) + (7.98 \text{ cal} \cdot \text{mol}^{-1} \cdot \text{K}^{-1}) \quad (2)$$

It is convenient that (by chance) the constant terms independent of S° in eq 2 nearly cancel on expansion, and the solvation entropy in water can therefore be approximated in more

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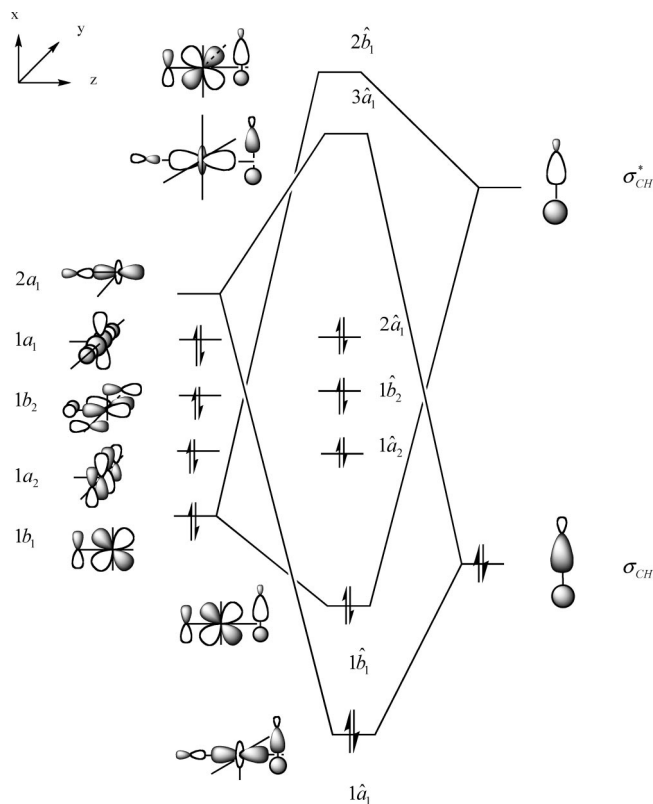


Figure 1. Orbital diagram of the formation of $\text{trans-PtCl}_2\text{X}(\text{CH}_4)^-$ from $\text{trans-PtCl}_2\text{X}^-$ and CH_4 fragments.

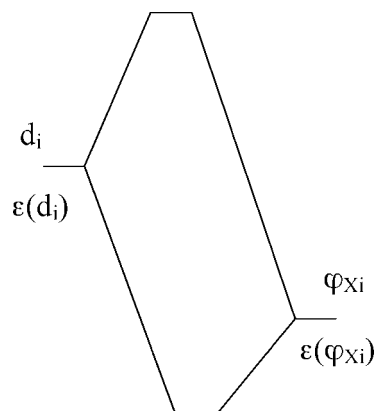


Figure 2. Orbital interaction between d_i metal orbital and φ_{Xi} ligand orbital.

qualitative discussions as 50% of the gas-phase entropy S_{gas}° , with the opposite sign.

3. Results and Discussion

We shall now study the influence of the ligand X on the strength of the Pt–CH₄ bond in $\text{trans-PtCl}_2\text{X}(\text{CH}_4)^-$ when methane is trans to X as well as $\text{trans-PtX}_2\text{Cl}(\text{CH}_4)^-$ when CH₄ is cis to X. Our analysis will consider the orbital interaction between $\text{trans-PtCl}_2\text{X}^-$ and CH₄ in the trans case and the interaction between $\text{trans-PtX}_2\text{Cl}^-$ and CH₄ in the cis case. We shall begin with the interaction between $\text{trans-PtCl}_2\text{X}^-$ and CH₄. To this end, we need first to analyze the frontier orbitals of $\text{trans-PtCl}_2\text{X}^-$.

3.1. Orbital Interaction in $\text{trans-PtCl}_2\text{X}^-$. Figure 1 illustrates the orbital interaction diagram between the $\text{trans-PtCl}_2\text{X}^-$ and CH₄ fragments. On the left-hand side, we display

Table 1. Percent Contribution^a from Ligand Combinations to the Frontier Orbitals of *trans*-PtCl₂X⁻

ligand X	percent contribution from φ_{Xi}	
	2a ₁ (d _{z²)^b}	1b ₁ (d _{xz}) ^c
F	31.4	20.8
Cl	40.1	29.5
Br	42.6	32.4
I	47.7	34.5
NO ₂	45.5	31.4
CN	48.3	35.5

^a $C_i^2 \times 10^2$ where C_i is defined in eq 4. ^b Percent contribution from σ_x to the 2a₁ orbital of *trans*-PtCl₂X⁻. ^c Percent contribution from π_x to 1b₁ of *trans*-PtCl₂X⁻.

Table 2. Energies^a of $\epsilon(\varphi_{\sigma_{Xi}})$ and $\epsilon(\varphi_{\pi_{Xi}})$ for Different Ligands X (X = F, Cl, Br, I, NO₂, and CN)

X ^b	$\epsilon(\varphi_{\sigma_{Xi}})$	$\epsilon(\varphi_{\pi_{Xi}})$
F	-30.309	-11.523
Cl	-21.047	-8.809
Br	-20.381	-8.000
I	-17.689	-7.171
NO ₂	-16.920	-6.397
CN	-12.217	-9.421

^a ϵ in eV. ^b Orbital energies of neutral ligands.

a simplified orbital diagram for the *trans*-PtCl₂X⁻ fragment, containing the five antibonding d-based frontier orbitals. The LUMO is an antibonding combination between d_{z²} and the trans σ_x orbital on X, whereas the HOMO is an out-of-phase combination between d_{x²-y²} and two cis σ_{Cl} orbitals on the two chlorines. We have in addition three d π orbitals destabilized by antibonding interactions with π -ligand orbitals. The interaction of the Pt orbital and the ligand combination can be explained by perturbation theory (see Figure 2).

The upper valence levels in Figure 2 are represented by d orbitals destabilized by out-of-phase antibonding interactions with the ligands. We can write the destabilized d orbitals as

$$\Psi_i = d_i - C_i \varphi_{Xi} (i = 1, 5) \quad (3)$$

where φ_{Xi} is the ligand orbital combination that interacts with d_i. We have from perturbation theory that C_i of eq 3 is given by

$$C_i = \frac{\langle d_i | H | \varphi_{Xi} \rangle}{\epsilon(d_i) - \epsilon(\varphi_{Xi})} \quad (4)$$

where $\epsilon(d_i)$ and $\epsilon(\varphi_{Xi})$ are the orbital energies of d_i and φ_{Xi} , respectively, as shown in Figure 2. It is clear that the larger the C_i in absolute terms, the larger the percent contribution of φ_{Xi} to Ψ_i . We present in Table 1 $C_i^2 \times 10^2$ for 2a₁ and 1b₁ of Figure 1 for different X. It follows for both the d_{z²}-based LUMO and the d_{xz}-based HOMO-3 that the percent contribution from φ_{Xi} to these orbitals with different trans ligands X follows the order F < Cl < Br < I < NO₂ < CN.

An analysis of the trend in Table 1 shows that it is set primarily by $\epsilon(\varphi_{Xi})$ in the denominator for the expression for C_i in eq 4. Thus, the more electronegative the ligand X is, the larger $-\epsilon(\varphi_{Xi})$ will be, Table 2, and the smaller the C_i² percent contribution from φ_{Xi} .

3.2. Orbital Interaction in PtCl₂X(CH₄)⁻. Having analyzed the frontier orbitals of PtCl₂X⁻, we now turn toward the interaction between PtCl₂X⁻ and methane to form the methane adduct PtCl₂X(CH₄)⁻. It can be seen from Figure 1 that the interaction diagram for the formation of PtCl₂X(CH₄)⁻ involves the occupied σ_{CH} orbital of CH₄ and the LUMO 2a₁ of PtCl₂X⁻

Table 3. Composition^{a,b} of the Bonding Orbitals 1 \hat{a}_1 and 1 \hat{b}_1 of PtCl₂X(CH₄)⁻ in Terms of CH₄ and PtCl₂X⁻ Fragment Contributions

ligand	composition	
	$C_2^2 \times 10^2$ in 1 \hat{a}_1	$C_3^2 \times 10^2$ in 1 \hat{b}_1
F	4.3	6.2
Cl	4.2	4.3
Br	4.0	3.7
I	4.0	1.4
NO ₂	2.6	1.0
CN	2.2	0

^a C₂ is defined in eq 6 as the contribution to 1 \hat{a}_1 from $\psi(2a_1)$. ^b C₃ is defined in eq 8 as the contribution from σ_{CH}^* to 1 \hat{b}_1 .

Table 4. Bonding Overlaps in 1 \hat{a}_1 and 1 \hat{b}_1 of PtCl₂X(CH₄)⁻ between CH₄ and PtCl₂X⁻ Fragment Orbitals

ligands	overlap	
	$\langle \sigma_{CH} \psi(2a_1) \rangle^a$	$\langle \sigma_{CH}^* \psi(1b_1) \rangle^b$
F	3.17×10^{-1}	1.69×10^{-1}
Cl	2.29×10^{-1}	1.54×10^{-1}
Br	2.25×10^{-1}	1.49×10^{-1}
I	2.17×10^{-1}	1.46×10^{-1}
NO ₂	1.80×10^{-1}	1.30×10^{-1}
CN	1.56×10^{-1}	0.78×10^{-1}

^a Bonding overlap in 1 \hat{a}_1 . ^b Bonding overlap in 1 \hat{b}_1 .

leading to the bonding orbital 1 \hat{a}_1 as well as the empty σ_{CH}^* of CH₄ and the occupied 1b₁ orbital of PtCl₂X⁻ leading to the bonding orbital 1 \hat{b}_1 . We can write

$$\psi(1\hat{a}_1) = \sigma_{CH} + C_2 \psi(2a_1) \quad (5)$$

where

$$C_2 = \frac{\langle \sigma_{CH} | H | \psi(2a_1) \rangle}{\epsilon(2a_1) - \epsilon(\sigma_{CH})} \quad (6)$$

and

$$\psi(1\hat{b}_1) = \psi(1b_1) + C_3 \sigma_{CH}^* \quad (7)$$

with

$$C_3 = \frac{\langle \sigma_{CH}^* | H | \psi(1b_1) \rangle}{\epsilon(\sigma_{CH}^*) - \epsilon(1b_1)} \quad (8)$$

We display in Table 3 the contributions from $\psi(2a_1)$ to $\psi(1\hat{a}_1)$ and from σ_{CH}^* to $\psi(1\hat{b}_1)$ for different trans ligands X. We find in both cases the order F > Cl > Br > I > NO₂ > CN. This trend correlates with the corresponding order for the overlaps $\langle \sigma_{CH} | \psi(2a_1) \rangle$ and $\langle \sigma_{CH}^* | \psi(1b_1) \rangle$ given by F > Cl > Br > I > NO₂ > CN as shown in Table 4.

The trend in the overlaps is in turn a consequence of the d orbital participation in $\psi(1b_1)$ and $\psi(2a_1)$. The more electronegative the X ligand, the larger the d participation in $\psi(1b_1)$ and $\psi(2a_1)$, and the larger the interacting overlaps of $\langle \sigma_{CH} | \psi(2a_1) \rangle$ and $\langle \sigma_{CH}^* | \psi(1b_1) \rangle$.

For square-planar platinum complexes *trans*-PtXCl₂L⁻, it is well established experimentally that the ability of X to destabilize the Pt-L bond trans to it follows the order¹⁸ F < Cl < Br < I < NO₂ < CN. It seems that the order of this "trans influence"¹⁹ at the most fundamental level is related to the electronegativity of X, according to our analysis in Sections

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Table 5. Decomposition of the Bonding Energy^{a,d} between the *trans*-PtCl₂X⁻ Fragment and the CH₄ Fragment in *trans*-PtCl₂X(CH₄)⁻ for Different X (X = F, Cl, Br, I, NO₂, and CN)

X	ΔE_{dist}	$\Delta E_{\text{steric}}^c$	$\Delta E_{\text{orbital}}$	χ^b	$\Delta E_{\text{bond}}^{\text{int}}$	ΔE_{solv}	$\Delta E_{\text{bond}}^{\text{d,e}}$
F	18.6	45.9	-77.7	1.20	-13.2	15.4	2.2
Cl	9.1	35.1	-51.6	1.17	-7.4	14.2	6.8
Br	8.6	34.2	-49.0	1.14	-6.2	12.9	6.7
I	6.8	32.1	-43.8	1.13	-5.0	11.6	6.6
NO ₂	5.4	28.4	-33.1	0.98	0.7	12.2	12.9
CN	3.7	20.4	-25.9	1.07	-1.8	24.1	22.3

^a Energy in kcal/mol. ^b $\chi = |\Delta E_{\text{orbital}}| / (\Delta E_{\text{dist}} + \Delta E_{\text{steric}})$. ^c $\Delta E_{\text{steric}} = \Delta E_{\text{elstat}} + \Delta E_{\text{Pauli}}$. ^d $\Delta E_{\text{bond}} = \Delta E_{\text{bond}}^{\text{int}} + \Delta E_{\text{solv}}$. ^e Bond energy in solution. ^f Bond energy in the gas phase, see eq 9.

3.1 and 3.2. Thus the most electronegative X has the weakest trans influence, for reasons explained above. We are not aware that such a simple explanation of the “trans influence” has been given before for the series of X ligands discussed here.

3.3. Energy Decomposition Analysis of the Bond between CH₄ and *trans*-PtCl₂X⁻ in *trans*-PtCl₂X(CH₄)⁻. Up to now, we have analyzed how different trans ligands X affect the overlap between orbitals on *trans*-PtCl₂X⁻ and the CH₄ fragment when a methane adduct is formed. We shall now study the interaction between the two fragments in more detail using the extended transition state (ETS) energy decomposition scheme (eq 9).^{20–24} The total interaction energy between PtCl₂X⁻ and CH₄ in the gas phase can be written as

$$\begin{aligned} \Delta E_{\text{bond}}^{\text{tot}} &= \Delta E_{\text{dist}} + [\Delta E_{\text{elstat}} + \Delta E_{\text{Pauli}}] + \Delta E_{\text{orbital}} \\ &= \Delta E_{\text{dist}} + \Delta E_{\text{steric}} + \Delta E_{\text{orbital}} \end{aligned} \quad (9)$$

The first component is referred to as the distortion term ΔE_{dist} , which is the amount of energy taken to promote the two separate fragments, PtCl₂X⁻ and CH₄, from their equilibrium structures to the structures they will take up in the combined molecule. The second component is the steric repulsion term ΔE_{steric} , which is the sum of ΔE_{elstat} and ΔE_{Pauli} where ΔE_{elstat} is the classical electrostatic interaction between the promoted fragments, and ΔE_{Pauli} represents the repulsive Pauli interaction between occupied orbitals on the two fragments. Finally, the last term, $\Delta E_{\text{orbital}}$, stands for the interactions between the occupied molecular orbitals on one fragment and the unoccupied molecular orbitals of the other fragment as well as the mixing of occupied and virtual orbitals within the same fragment (intrafragment polarization). The terms in eq 9 have been evaluated as described elsewhere.^{20–22}

Table 5 displays the total interaction energy $\Delta E_{\text{bond}}^{\text{int}}$ in PtCl₂X(CH₄)⁻ decomposed into the various contributions according to eq 9 for different X. We note that the interaction energy $\Delta E_{\text{bond}}^{\text{int}}$ in absolute terms decreases in the order F > Cl > Br > I > NO₂ > CN as the overlaps between σ_{CH} and 2a₁ as well as σ_{CH}^* and 1b₁ decrease for reasons explained previously.

It is not surprising that the system (X = F) with the strongest interaction energy $\Delta E_{\text{bond}}^{\text{int}}$ also has the largest distortion energy, Table 5. Thus a strong interaction $\Delta E_{\text{bond}}^{\text{int}}$ will lead to a considerable donation of charge from σ_{CH} to 2a₁ as well as a sizable back-donation from 1b₁ to σ_{CH}^* ; see Table 6. Both

Table 6. Occupations of Fragment Orbitals in PtCl₂XCH₄⁻ for Different X (X = F, Cl, Br, I, NO₂, CN)

X	occupation			
	Pt fragment		CH ₄ fragment	
	1b ₁	2a ₁	σ_{CH}	σ_{CH}^*
F	1.78	0.28	1.70	0.21
Cl	1.89	0.21	1.76	0.09
Br	1.90	0.21	1.77	0.09
I	1.92	0.19	1.79	0.07
NO ₂	1.94	0.13	1.83	0.05
CN	1.96	0.09	1.84	0.02

Table 7. Pt–C, Pt–H, and C–H Bond Lengths^a in *trans*-PtCl₂X(CH₄)⁻ with Different X (X = F, Cl, Br, I, NO₂, and CN)

X	$r(\text{Pt–C})$	$r(\text{Pt–H})$	$r(\text{C–H})$
	1 ^b	2 ^b	3 ^b
F	2.27	1.63	1.27
Cl	2.36	1.72	1.19
Br	2.38	1.74	1.19
I	2.41	1.77	1.18
NO ₂	2.53	1.85	1.15
CN	2.53	1.98	1.13

^a Bond lengths in Å. ^b See geometry 1 in Scheme 1.

Table 8. Further Decomposition of $\Delta E_{\text{orbital}}$ in *trans*-PtCl₂X(CH₄)⁻ Complexes

Pt/C ^a	orbital interaction energy ($\Delta E_{\text{orbital}}^{\text{f}}$)					
	F	Cl	Br	I	NO ₂	CN
All/All ^b	-81.6	-55.2	-52.8	-46.7	-40.2	-36.4
1/2 ^c	-48.8	-28.6	-26.9	-23.2	-18.4	-12.6
0/2 ^d	-18.1	-9.2	-8.5	-7.0	-5.0	-3.3
1/0 ^e	-21.4	-15.9	-15.2	-14.0	-12.2	-8.8

^a Pt/C means Pt fragment/C (methane) fragment. ^b All/All means all virtual orbitals are kept for both fragments in calculation. ^c 1/2 means one virtual orbital is kept for Pt fragment and two virtual orbitals are kept for CH₄ fragment. ^d 0/2 means no virtual orbitals for Pt fragment but two virtual orbitals are kept for CH₄ fragment. ^e 1/0 means one virtual orbital is kept for Pt fragment and no virtual orbitals for CH₄ fragment. ^f Energy in kcal/mol.

donation and back-donation will result in an elongation of the C–H bond, Table 7, and an increase in the distortion energy, especially for CH₄.

The order of the interaction energy F > Cl > Br > I > NO₂ > CN is also reflected in the steric interaction term ΔE_{steric} . The larger the interaction energy is in absolute terms, the shorter the Pt–C bond length, Table 7, and the larger the steric interaction energy between the two fragments, Table 5. It is important to note that the complex with the strongest Pt–CH₄ bond (X = F) has both the largest destabilizing and (in absolute terms) stabilizing contributions to the total interaction energy $\Delta E_{\text{bond}}^{\text{int}}$. However, if we consider the ratio

$$\chi = \frac{|\Delta E_{\text{orbital}}|}{(\Delta E_{\text{dist}} + \Delta E_{\text{steric}})} \quad (10)$$

we find that it decreases through the series F > Cl > Br > I > NO₂ > CN. Thus we gain more stabilization per unit of destabilization energy ($\Delta E_{\text{dist}} + \Delta E_{\text{steric}}$) at the beginning of the series than at the end.

We display finally in Table 8 a further analysis of ΔE_{int} as a function of X. In the first entry “All/All”, we have the full interaction energy with all virtual orbitals on the two fragments included. The entry “1/0” stems from a calculation in which 2a₁ is the only virtual orbital included. The corresponding interaction energy thus represents the stabilization from donation

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Table 9. Decomposition of the Bonding Energy^{a,d} between the *trans*-PtCl₂X⁻ Fragment and H₂O in *trans*-PtCl₂X(H₂O)⁻ for Different X (X = F, Cl, Br, I, NO₂, and CN)

X	ΔE_{dist}	$\Delta E_{\text{steric}}^c$	$\Delta E_{\text{orbital}}$	χ^b	$\Delta E_{\text{bond}}^{\text{int},f}$	ΔE_{solv}	$\Delta E_{\text{bond}}^{\text{d},e}$
F	4.5	16.9	-43.1	2.02	-21.7	9.3	-12.4
Cl	1.6	16.9	-36.1	1.95	-17.6	9.3	-8.3
Br	1.6	17.9	-35.8	1.84	-16.3	8.2	-8.1
I	1.1	18.9	-35.0	1.74	-15.0	7.4	-7.6
NO ₂	4.2	18.9	-33.1	1.43	-10.0	3.3	-6.7
CN	1.4	12.7	-28.0	1.99	-14.0	15.6	1.6

^a Energy in kcal/mol. ^b $\chi = |\Delta E_{\text{orbital}}| / (\Delta E_{\text{dist}} + \Delta E_{\text{steric}})$. ^c $\Delta E_{\text{steric}} = \Delta E_{\text{elstat}} + \Delta E_{\text{Pauli}}$. ^d $\Delta E_{\text{bond}} = \Delta E_{\text{bond}}^{\text{int}} + \Delta E_{\text{solv}}$. ^e Bond energy in solution. ^f Bond energy in the gas phase, see eq 9.

into 2a₁. In the same way, the entry “0/2” represents the stabilization due to donation into σ_{CH}^* . Finally, 1/2 has both donations to 2a₁ as well as σ_{CH}^* . We note that the interaction energy for 1/2 in absolute terms exceeds the sum of 1/0 + 0/2, thus testifying to some synergism. On the other hand 1/0 + 0/2 is much smaller than All/All in absolute terms. The difference represents intrafragment polarizations to compensate for the steric interaction energy ΔE_{steric} .

In order to obtain the total bond strength ΔE_{bond} , we have to add to the gas-phase bond energy $\Delta E_{\text{bond}}^{\text{int}}$ the contribution from the solvation energy ΔE_{solv} . The change in solvation energy ΔE_{solv} from PtCl₂X⁻ + CH₄ to PtCl₂X(CH₄)⁻ is positive, as PtCl₂X(CH₄)⁻ has a larger volume than PtCl₂X⁻. After the inclusion of solvation effects, the order of the “trans influence” is F < Cl ~ Br ~ I < NO₂ < CN in PtCl₂X(CH₄)⁻.

3.4. Energy Decomposition Analysis of Bond between H₂O and *trans*-PtCl₂X⁻ in *trans*-PtCl₂X(H₂O)⁻. The most favorable pathway of methane uptake for *trans*-PtCl₂X(H₂O)⁻ is the substitution of H₂O by CH₄, and its energy barrier is largely determined by the dissociation energy of H₂O, as demonstrated in a previous investigation.³ The energy decomposition of the bond between H₂O and *trans*-PtCl₂X⁻ in PtCl₂X(H₂O)⁻ is shown in Table 9.

The bonding energy in *trans*-PtCl₂X(H₂O)⁻ between *trans*-PtCl₂X⁻ and H₂O decreases in the order F > Cl > Br > I > NO₂ > CN, see $\Delta E_{\text{bond}}^{\text{int}}$, which is in accordance with the prediction by the trans-directing ability of the trans ligand X. The above trend is reflected in χ , the increase of $|\Delta E_{\text{orbital}}|$ per unit energy of $(\Delta E_{\text{dist}} + \Delta E_{\text{steric}})$, as χ is largest for F and smallest for NO₂. Behind this trend is the reduction in the overlap between the *trans*-PtCl₂X⁻ and the H₂O fragments, as X becomes less electronegative just as in the interaction between *trans*-PtCl₂X⁻ and CH₄; see Sections 3.1 and 3.2. Solvation energy does not alter the trend in ΔE_{bond} , which is determined by electronic effects.

3.5. Energy Decomposition Analysis of Interaction between CH₄ and PtCl₂X⁻ in C–H Activation Transition State 3. After the formation of the methane complex PtCl₂X(CH₄)⁻, the system might proceed with the C–H activation through the transition state 3 of Scheme 1. We note that we were unable to find a transition state for X = F. In this case, PtCl₂F(CH₄)⁻ proceeds directly to the C–H activation product without an energy barrier.

Table 10 displays key geometrical parameters of the transition state for C–H activation of CH₄ by *trans*-PtCl₂X⁻; see complex 3 of Scheme 1. It is clearly seen that the C–H bond is lengthened while the Pt–C and Pt–H bonds are shortened, compared to the PtCl₂X(CH₄)⁻ complex 1 of Scheme 1. The geometry of 3 is very similar for different X. Thus, the bond lengths are the same to within 0.22 Å for Pt–C, 0.04 Å for

Table 10. Bond Lengths^a in Transition State for C–H Activation of CH₄ by PtCl₂X⁻ with Different X (X = F, Cl, Br, I, NO₂, and CN)

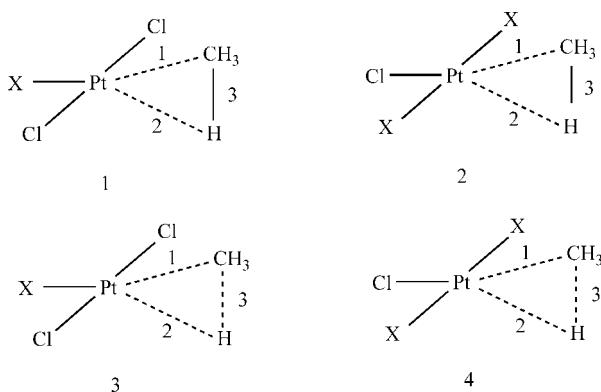
X	$r(\text{Pt}-\text{C})$ 1 ^b	$r(\text{Pt}-\text{H})$ 2 ^b	$r(\text{C}-\text{H})$ 3 ^b
F			
Cl	2.14	1.59	1.36
Br	2.13	1.60	1.34
I	2.05	1.63	1.25
NO ₂	1.93	1.63	1.27
CN	2.15	1.61	1.31

^a Bond lengths in Å. ^b See geometry 3 in Scheme 1.

Table 11. Decomposition of the Bonding Energy^{a,d} between the *trans*-PtCl₂X⁻ Fragment and the CH₄ Fragment in *trans*-PtCl₂X(CH₄)⁻ for Different X (X = F, Cl, Br, I, NO₂, and CN) at the Transition State

X	ΔE_{dist}	$\Delta E_{\text{steric}}^c$	$\Delta E_{\text{orbital}}$	χ^b	$\Delta E_{\text{bond}}^{\text{int},f}$	ΔE_{solv}	$\Delta E_{\text{bond}}^{\text{d},e}$
F							
Cl	29.7	60.2	-91.5	1.02	-1.6	13.6	12.0
Br	34.2	64.6	-96.5	0.98	2.3	12.6	14.9
I	30.8	66.9	-96.6	0.99	1.1	11.6	12.7
NO ₂	28.8	86.5	-98.5	0.85	16.8	10.8	27.6
CN	33.2	105.8	-111.7	0.80	27.3	25.9	53.2

^a Energy in kcal/mol. ^b $\chi = |\Delta E_{\text{orbital}}| / (\Delta E_{\text{dist}} + \Delta E_{\text{steric}})$. ^c $\Delta E_{\text{steric}} = \Delta E_{\text{elstat}} + \Delta E_{\text{Pauli}}$. ^d $\Delta E_{\text{bond}} = \Delta E_{\text{bond}}^{\text{int}} + \Delta E_{\text{solv}}$. ^e Bond energy in solution. ^f Bond energy in the gas phase, see eq 9.

Scheme 1. Structures of *trans*-PtCl₂X(CH₄)⁻ at Reactant (1) and Transition State (3) and of *trans*-PtClX₂(CH₄)⁻ at Reactant (2) and Transition State (4)

Pt–H, and 0.11 Å for the C–H bond. The similar structure of CH₄ in 3 is also reflected in the modest variation in ΔE_{dist} , Table 11.

Table 11 displays the interaction energy $\Delta E_{\text{bond}}^{\text{int}}$ decomposed according to eq 9 for the transition state 3. We note that $\Delta E_{\text{bond}}^{\text{int}}$ becomes increasingly positive and destabilizing through the series Cl < Br < I < NO₂ < CN. This trend is set by the ratio χ defined in eq 10. Thus the amount of stabilizing orbital interaction obtained in absolute terms per unit of steric interaction and distortion is largest at the start of the series, as the overlaps here are larger between 2a₁ and σ_{CH}^* as well as 1b₁ and σ_{CH}^* ; see Table 4. The term ΔE_{bond} exhibits the same trend, but the inclusion of solvation energy enlarges the energy gap between the halogen ligands and NO₂ and CN.

3.6. Energy Decomposition Analysis of the Bond between *trans*-PtClX₂⁻ and CH₄ in *trans*-PtClX₂(CH₄)⁻ at the Reactant and Transition State. After having analyzed how the trans ligand X affects the C–H activation of *trans*-PtCl₂X(CH₄)⁻, we now turn to the influence of the cis ligand X on C–H activation, taking as our example the *trans*-PtClX₂(CH₄)⁻, complex 2 of Scheme 1.

Table 12 displays the energy decomposition of the bond between *trans*-PtClX₂⁻ and CH₄ in *trans*-PtClX₂(CH₄)⁻. In

Table 12. Decomposition of the Bonding Energy^{a,d} between the *trans*-PtClX₂⁻ Fragment and CH₄ in *trans*-PtClX₂(CH₄)⁻ for Different X (X = F, Cl, Br, I, NO₂ and CN)

X	ΔE_{dist}	$\Delta E_{\text{steric}}^c$	$\Delta E_{\text{orbital}}$	χ^b	$\Delta E_{\text{bond}}^{\text{int},f}$	ΔE_{solv}	$\Delta E_{\text{bond}}^{d,e}$
F	10.1	32.7	-52.9	1.24	-10.1	17.9	7.8
Cl	9.1	35.1	-51.6	1.17	-7.4	14.2	6.8
Br	10.0	35.2	-52.7	1.17	-7.5	11.5	4.0
I	10.3	37.6	-56.2	1.17	-8.3	8.8	0.5
NO ₂	20.6	28.0	-47.9	0.98	0.7	4.1	4.8
CN	7.7	27.9	-48.1	1.35	-12.5	7.5	-5.0

^a Energy in kcal/mol. ^b $\chi = |\Delta E_{\text{orbital}}| / (\Delta E_{\text{dist}} + \Delta E_{\text{steric}})$. ^c $\Delta E_{\text{steric}} = \Delta E_{\text{elstat}} + \Delta E_{\text{Pauli}}$. ^d $\Delta E_{\text{bond}} = \Delta E_{\text{bond}}^{\text{int}} + \Delta E_{\text{solv}}$. ^e Bond energy in solution. ^f Bond energy in the gas phase, see eq 9.

Table 13. Decomposition of the Bonding Energy^{a,d} between the *trans*-PtClX₂⁻ Fragment and the CH₄ Fragment in *trans*-PtClX₂(CH₄)⁻ for Different X (X = F, Cl, Br, I, NO₂, and CN) at the Transition State

X	ΔE_{dist}	$\Delta E_{\text{steric}}^c$	$\Delta E_{\text{orbital}}$	χ^b	$\Delta E_{\text{bond}}^{\text{int},f}$	ΔE_{solv}	$\Delta E_{\text{bond}}^{d,e}$
F	22.9	50.2	-80.8	1.11	-7.7	17.4	9.7
Cl	26.2	60.2	-91.5	1.06	-5.1	13.6	8.5
Br	28.9	60.6	-94.9	1.06	-5.4	11.1	5.7
I	24.8	58.9	-90.4	1.08	-6.7	8.2	1.5
NO ₂	31.2	51.7	-79.4	0.96	3.5	3.3	6.8
CN	27.2	61.8	-96.7	1.09	-7.7	7.8	0.1

^a Energy in kcal/mol. ^b $\chi = |\Delta E_{\text{orbital}}| / (\Delta E_{\text{dist}} + \Delta E_{\text{steric}})$. ^c $\Delta E_{\text{steric}} = \Delta E_{\text{elstat}} + \Delta E_{\text{Pauli}}$. ^d $\Delta E_{\text{bond}} = \Delta E_{\text{bond}}^{\text{int}} + \Delta E_{\text{solv}}$. ^e Bond energy in solution. ^f Bond energy in the gas phase, see eq 9.

Table 14. Decomposition of the Energy Barrier^a of C–H Activation of *trans*-PtCl₂X(CH₄)⁻ for Different X (X = F, Cl, Br, I, NO₂, and CN)

X	$\Delta E_{\text{dist}}^{\ddagger}$	$\Delta E_{\text{steric}}^{\ddagger}$	$\Delta E_{\text{orbital}}^{\ddagger}$	χ^{\ddagger}	$\Delta E_{\text{g}}^{\ddagger}$	$\Delta E_{\text{solv}}^{\ddagger}$	$\Delta E_{\text{s}}^{\ddagger}$
F							
Cl	20.6	25.2	-40.0	0.87	5.8	-0.7	5.1
Br	25.6	30.4	-47.5	0.85	8.5	-0.3	8.2
I	24.0	34.9	-52.8	0.90	6.1	0.0	6.1
NO ₂	23.4	58.1	-65.3	0.80	16.2	-1.4	14.8
CN	29.5	85.4	-85.8	0.75	29.1	1.8	30.9

^a Energy in kcal/mol. ^b $\Delta E_{\text{dist}}^{\ddagger} = \Delta E_{\text{dist}}^{\text{TS}} - \Delta E_{\text{dist}}^{\text{G}}$. ^c $\Delta E_{\text{steric}}^{\ddagger} = \Delta E_{\text{steric}}^{\text{TS}} - \Delta E_{\text{steric}}^{\text{G}}$. ^d $\Delta E_{\text{orbital}}^{\ddagger} = \Delta E_{\text{orbital}}^{\text{TS}} - \Delta E_{\text{orbital}}^{\text{G}}$. ^e $\chi^{\ddagger} = |\Delta E_{\text{orbital}}^{\ddagger}| / (\Delta E_{\text{dist}}^{\ddagger} + \Delta E_{\text{steric}}^{\ddagger})$. ^f $\Delta E_{\text{solv}}^{\ddagger} = \Delta E_{\text{solv}}^{\text{TS}} - \Delta E_{\text{solv}}^{\text{G}}$. ^g $\Delta E_{\text{g}}^{\ddagger} = \Delta E_{\text{elstat}} + \Delta E_{\text{Pauli}}$. ^h $\Delta E_{\text{s}}^{\ddagger} = \Delta E_{\text{bond}}^{\text{int,TS}} - \Delta E_{\text{bond}}^{\text{int,G}}$, see eq 9 for $\Delta E_{\text{bond}}^{\text{int}}$. ⁱ $\Delta E_{\text{s}}^{\ddagger} = \Delta E_{\text{bond}}^{\text{TS}} - \Delta E_{\text{bond}}^{\text{G}}$ in solution, $\Delta E_{\text{bond}} = \Delta E_{\text{bond}}^{\text{int}} + \Delta E_{\text{solv}}$.

considering ΔE_{bond} of Table 12, we note that the stability of the Pt–CH₄ bond in solution follows the order F < Cl < Br < NO₂ < I < CN, which, except for the relative position of NO₂, I, is the opposite of that of the trans influence in PtCl₂X(CH₄)⁻. The trend is set by the solvation energies ΔE_{solv} , which follows the trend F > Cl > Br > NO₂ > I > CN. For the majority of the systems ΔE_{solv} is positive, as the product *trans*-PtX₂Cl(CH₄)⁻ has a larger volume, and consequently less solvation stabilization, than the reactant *trans*-PtX₂Cl⁻. The change in volume is largest at the start of the series.

For the transition state 4 in Scheme 1, ΔE_{bond} of Table 13 follows the same trend as in the adduct PtX₂Cl(CH₄)⁻. The trend is again set by ΔE_{solv} . It is perhaps not so surprising that electronic factors have only a modest influence on the trend in the Pt–CH₄ bond strength as a function of X both in the adduct PtX₂Cl(CH₄)⁻ 2 and the transition state 4 since all systems have the same ligand trans to the Pt–CH₄ bond.

3.7. Comparison of the C–H Cleavage Energy Barrier of *trans*-PtCl₂X(CH₄)⁻ and *trans*-PtClX₂(CH₄)⁻. Tables 14 and 15 display the energy barriers ΔE^{\ddagger} both in the gas phase and in solution for C–H cleavage involving *trans*-PtCl₂X(CH₄)⁻ and *trans*-PtClX₂(CH₄)⁻, respectively. ΔE^{\ddagger} is further decomposed into its activation components $\Delta E_{\text{dist}}^{\ddagger}$, $\Delta E_{\text{steric}}^{\ddagger}$, $\Delta E_{\text{orbital}}^{\ddagger}$,

Table 15. Decomposition of the Energy Barrier^a of C–H Activation of *trans*-PtClX₂(CH₄)⁻ for Different X (X = F, Cl, Br, I, NO₂, and CN)

X	$\Delta E_{\text{dist}}^{\ddagger}$	$\Delta E_{\text{steric}}^{\ddagger}$	$\Delta E_{\text{orbital}}^{\ddagger}$	χ^{\ddagger}	$\Delta E_{\text{g}}^{\ddagger}$	$\Delta E_{\text{solv}}^{\ddagger}$	$\Delta E_{\text{s}}^{\ddagger}$
F	12.8	17.5	-28.0	0.92	2.3	-0.5	1.8
Cl	17.2	25.2	-40.0	0.94	2.4	-0.7	1.7
Br	18.9	25.5	-42.2	0.95	2.2	-0.4	1.8
I	14.5	21.3	-34.2	0.95	1.6	-0.6	1.0
NO ₂	10.6	23.7	-31.6	0.92	2.7	-0.9	1.8
CN	19.5	34.8	-48.6	0.90	5.7	0.3	6.0

^a Energy in kcal/mol. ^b $\Delta E_{\text{dist}}^{\ddagger} = \Delta E_{\text{dist}}^{\text{TS}} - \Delta E_{\text{dist}}^{\text{G}}$. ^c $\Delta E_{\text{steric}}^{\ddagger} = \Delta E_{\text{steric}}^{\text{TS}} - \Delta E_{\text{steric}}^{\text{G}}$. ^d $\Delta E_{\text{orbital}}^{\ddagger} = \Delta E_{\text{orbital}}^{\text{TS}} - \Delta E_{\text{orbital}}^{\text{G}}$. ^e $\chi^{\ddagger} = |\Delta E_{\text{orbital}}^{\ddagger}| / (\Delta E_{\text{dist}}^{\ddagger} + \Delta E_{\text{steric}}^{\ddagger})$. ^f $\Delta E_{\text{solv}}^{\ddagger} = \Delta E_{\text{solv}}^{\text{TS}} - \Delta E_{\text{solv}}^{\text{G}}$. ^g $\Delta E_{\text{g}}^{\ddagger} = \Delta E_{\text{elstat}} + \Delta E_{\text{Pauli}}$. ^h $\Delta E_{\text{s}}^{\ddagger} = \Delta E_{\text{bond}}^{\text{int,TS}} - \Delta E_{\text{bond}}^{\text{int,G}}$, see eq 9 for $\Delta E_{\text{bond}}^{\text{int}}$. ⁱ $\Delta E_{\text{s}}^{\ddagger} = \Delta E_{\text{bond}}^{\text{TS}} - \Delta E_{\text{bond}}^{\text{G}}$ in solution, $\Delta E_{\text{bond}} = \Delta E_{\text{bond}}^{\text{int}} + \Delta E_{\text{solv}}$.

and $\Delta E_{\text{solv}}^{\ddagger}$. Here $\Delta E_{\text{dist}}^{\ddagger} = \Delta E_{\text{dist}}^{\text{TS}} - \Delta E_{\text{dist}}^{\text{G}}$ with similar definitions for the other components; see Tables 14 and 15. The term χ^{\ddagger} , the increase in $\Delta E_{\text{orbital}}^{\ddagger}$ as per unit energy of $\Delta E_{\text{dist}}^{\ddagger}$ and $\Delta E_{\text{steric}}^{\ddagger}$, is also calculated for both Pt species.

It is clearly seen from Table 14 that the energy barrier for C–H activation in the case of *trans*-PtCl₂X(CH₄)⁻ increases in the order Cl ~ Br ~ I < NO₂ < CN both in the gas phase and in solution. Among the different components, $\Delta E_{\text{dist}}^{\ddagger}$ varies modestly among Cl, Br, I, and NO₂, but increases substantially for X = CN. This is understandable since CH₄ for X = CN is only modestly distorted in the ground state of PtCl₂X(CH₄)⁻. It must then undergo the largest distortion to reach the transition state where CH₄ is equally distorted for all systems. Both $\Delta E_{\text{steric}}^{\ddagger}$ and $\Delta E_{\text{orbital}}^{\ddagger}$ increase continuously in absolute terms from Cl to CN.

We can understand this trend by observing that the system with the strongest trans-directing ligand has the longest Pt–CH₄ distance in the ground state, whereas the Pt–CH₄ distances are quite similar for all systems in the transition state. Thus the system with the most trans-directing ligand must undergo the largest change in the Pt–CH₄ bond length and consequently the largest variation in $\Delta E_{\text{steric}}^{\ddagger}$ and $\Delta E_{\text{orbital}}^{\ddagger}$. It is finally worth noting that the activation correlates with χ^{\ddagger} , the stabilization gained in $-\Delta E_{\text{orbital}}^{\ddagger}$ for each energy unit spent of $\Delta E_{\text{steric}}^{\ddagger} + \Delta E_{\text{dist}}^{\ddagger}$, Table 14. The systems with the larger χ^{\ddagger} are those with the lower barrier and the higher overlap of $\langle \sigma_{\text{CH}} | 2a_1 \rangle$ and $\langle \sigma_{\text{CH}}^* | 1b_1 \rangle$.

Table 15 exhibits the energy barrier and its components for C–H activation in the case of *trans*-PtClX₂(CH₄)⁻. Due to the similarities in geometry of the Pt–CH₄ frame in *trans*-PtClX₂(CH₄)⁻, both at the ground state and the transition state with different X, $\Delta E_{\text{steric}}^{\ddagger}$, $\Delta E_{\text{dist}}^{\ddagger}$, and $\Delta E_{\text{orbital}}^{\ddagger}$ vary only slightly from F to CN. Also $\Delta E_{\text{solv}}^{\ddagger}$ is modest as the solvation energy changes little from adduct to transition state. However, a jump in ΔE^{\ddagger} is seen for X = CN.

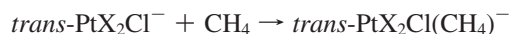
4. Conclusion

We have studied the influence of different ligands X (X = F, Cl, Br, I, NO₂, CN) on the Pt–CH₄ bond and C–H activation barriers in *trans*-PtCl₂X(CH₄)⁻, **1**, and *trans*-PtX₂Cl(CH₄)⁻, **2**, where X is respectively trans (**1**) and cis (**2**) to CH₄. For **1** with X trans to CH₄, we find that the Pt–CH₄ bond energy decreases as F > Cl > Br > I > NO₂ > CN, in agreement with the experimentally established “trans influence” of these ligands.

The Pt–CH₄ bond in complex **1** of Scheme 1 is established through the bonding interaction between the empty 2a₁ σ orbital on *trans*-PtCl₂X⁻ and σ_{CH} on CH₄ as well as between the occupied 1b₁ orbital on the metal fragment and σ_{CH}^* on CH₄. Here 2a₁ is an out-of-phase combination between d_z on the metal

and σ_x on X. Likewise, $1b_1$ is an out-of-phase combination between d_{xz} on Pt and π_x on the ligand X. Through the series X = F, Cl, Br, I, NO₂, CN, the orbital energies of σ_x and π_x increase due to the decrease of the electronegativity of the elements involved in the orbitals. As a result, the contribution to $2a_1$ and $1b_1$ increases, whereas the d contribution declines. The decline in d- contribution will in turn lead to a decrease in the bonding overlaps $\langle \sigma_{CH} | 2a_1 \rangle$ and $\langle \sigma_{CH}^* | 1b_1 \rangle$, as only the d component in $2a_1$ and $1b_1$ contributes to these overlaps. The decline in the bonding overlaps will in turn weaken the Pt–CH₄ bond. Thus the experimentally well-established “trans influence” for the series of ligands investigated here can be related to their electronegativity. The trans effect is even more pronounced in the transition state with the result that the barrier of C–H activation with CH₄ trans to X follows the order F < Cl < Br < I < NO₂ < CN.

When X is situated in the cis position, the stability of *trans*-PtX₂Cl(CH₄)[−] in both the ground state and transition state follows the order F < Cl < Br < NO₂ < I < CN. The order is set by solvent effects that favor the association process



for ligands X with a large volume.

Not only C–H activation, but also the rate-determining methane uptake step is influenced by the trans ligand. There are many reports on ligand displacement rather than C–H bond

Table 16. Comparison of the Activation Energy of Methane Uptake and C–H Activation with Different Trans Ligand in *trans*-PtCl₂X(H₂O)[−]

ligand	activation energy (ΔG_s^\ddagger) ^a	
	methane uptake ^b	C–H activation ^b
F	31.3	0
Cl	23.3	1.1
Br	23.4	2.1
I	22.0	3.6
NO ₂	20.9	16.6
CN	24.9	28.6

^a Energy in kcal/mol. ^b Data are from ref 5.

cleavage as the rate-determining step with various metal centers, including Pt.^{25–27} As one might expect from our analysis, they all involve weakly trans-directing ligands. However, as the trans-directing ability of X increases, the uptake barrier decreases as the Pt–L bond for the leaving group L is weakened, whereas the C–H activation barrier increases. This trend was illustrated by calculations on PtCl₂X(H₂O)[−]; see Table 16. It is thus clear that the optimal trans ligand X in the Shilov systems is one that labilizes the leaving ligand sufficiently without causing a substantial C–H activation barrier.

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