if we use  $\Gamma_{\eta}(2\gamma) = 192 \text{ eV}$ , and  $r \approx 1.7$  when we employ  $\Gamma_{\eta}(2\gamma) = 25 \text{ eV}$ . The ratio 2.4 is consistent with experiment. Estimates for  $\Gamma_{\eta}(2\gamma)$  made on a unitary symmetry model also lie between 25 eV and 192 eV and, hence, for these estimates *r* will lie between 1.7 and 2.4. Hori *et al.* as well as Gell-Mann *et al.* ....."

If  $(\delta g_{\pi NN}/g_{\pi NN})$  is taken to be 0.7%, then the values (9) and (10) for  $\Gamma_{\eta}(\pi^+\pi^-\pi^0)$  and  $\Gamma_{\eta}(3\pi^0)$  are unchanged provided that  $g_{\eta NN^2}/4\pi \approx 1$ . However, then Eq. (11) gives  $\Gamma_{\eta}(2\gamma) \approx 12$  eV so that r lies between 1.6 and 2.4 according as  $\Gamma_n(2\gamma)$  lies between 12 eV and 192 eV. Again with  $(\delta g_{\pi NN}/g_{\pi NN})$  $\approx 0.7\%$  and  $g_{\eta NN^2}/4\pi \approx 2$ ,  $\Gamma_{\eta}(\pi^+\pi^-\pi^0)$  and  $\Gamma_{\eta}(3\pi^0)$ become, respectively, 112 eV and 179 eV. These values are consistent with the estimates of Barret and Barton<sup>1</sup> who estimated the  $\eta^0 \rightarrow \pi^0$  vertex by a quite different approach based on unitary symmetry. In this case r lies between 1.8 and 3.3 according as  $\Gamma_{\eta}(2\gamma)$  lies between 25 eV and 192 eV, consistent with experiment. In this case  $\Gamma_{\eta}(2\gamma)$  can be equal to  $\Gamma_{\eta}(3\pi^0)$  depending on what value we take between 25 eV and 192 eV for  $\Gamma_{\eta}(2\gamma)$ .

In the last paragraph but one, if we take  $\lambda/16\pi = -0.15$ , the pseudoscalar coupling constant  $g_{\Sigma N \kappa^2}/4\pi$  has the values 48 to 24 according as  $R(K_{2^0} \rightarrow \pi^+\pi^-\pi^0) = 1.5 \times 10^6 \text{ sec}^{-1}$  or  $3 \times 10^6 \text{ sec}^{-1}$ . For scalar  $K\Sigma N$  coupling,

$$g_{\Sigma NK}^2/4\pi \approx 0.48.$$

Thus, for pseudoscalar coupling,  $g_{\Sigma NK}^2/4\pi$  comes out to be quite large, showing that the K pole in  $\Sigma^- \rightarrow n + \pi^-$  does not dominate and that one has to consider other contributions also.

We are grateful to Barbara Barrett for pointing out the error in our paper.

<sup>1</sup> Barbara Barret and G. Barton (to be published).

Angular Distribution of Muons in  $\pi$ - $\psi$  Decay at Rest, H. HULUBEI, J. S. AUSLÄNDER, E. M. FRIED-LÄNDER, AND Ş. ȚIȚEICA [Phys. Rev. 129, 2789 (1963)]. 1. In Table I, column headed "Sample size," row " $\Omega^*$ ": instead of 19126<sup>b</sup> read (19126)<sup>b</sup>. This figure does not represent an *actual* sample size, but a *fictitious* one. 2. In Fig. 8, (a) and (b) must be interchanged in order to obtain agreement between (i) drawings and (ii) figure caption and text.

Relaxation-Time Measurements in Ruby by a dc Magnetization Technique, SHIH-YU FENG AND N. BLOEMBERGEN [Phys. Rev. 130, 531 (1963)]. In the caption of Fig. 4 and in the line of the text immediately following this figure, it is erroneously stated that " $H_{de}$ =2990 G." This should read " $H_{de}$ = 1580±20 G." The value originally quoted belongs to another transition at 0° orientation. A check of our experimental records revealed the correct value, although the precision is rather poor. An accurate machine solution of the spin Hamiltonian at the frequency used in the experiment gives the following result for the harmonic point:

	Resonance	
Orientation	$H_{ m dc}$ (G)	Ratio $\nu_{34}/\nu_{13}$
21°	1585	2.92:2
22°	1674	3.01:2

We wish to thank Dr. W. Grant for calling our attention to this error.

**Partial-Wave Bethe-Salpeter Equation,** NOBORU NAKANISHI [Phys. Rev. **130**, 1230 (1963)]. In the denominators of Eqs. (3.21) and (4.5), and in the argument of the  $\delta$  function of Eq. (4.9), x and (1-x) should be interchanged.

Branching Ratios of  $\pi^-$  Mesons Stopped in Hydrogen and Deuterium, JAMES W. RYAN [Phys. Rev. 130, 1554 (1963)]. Add:

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