Composite Particle Model for the Nucleon and the (3,3) Resonance, J. S. BALL AND D. Y. WONG [Phys. Rev. **133**, B179 (1964)]. An error has been discovered in the computer program used to calculate the results given in Table I. After correction Table I should read as follows:

	γ_{33}	γ_1	m	g^2	Γ_{33}	a_1	a_3	c_{11}	c_{31}	c_{13}	c_{33}	d_{13}	d_{33}	W_c
1	0.05	−1.0	7.26	24.0	1.01	−0.501	−0.192	−0.620	−0.056	−0.048	0.298	−0.0006	0.0008	18.6
2	0.05	−0.5	7.22	21.0	1.00	−0.451	−0.185	−0.484	−0.051	−0.051	0.297	−0.0011	0.0011	18.2
3	0.05	0	7.52	14.0	1.02	−0.369	−0.194	−1.06	−0.041	−0.059	0.303	−0.0014	0.0013	18.6
4	0.06	−1.0	7.47	23.4	1.00	−0.477	−0.168	−1.20	−0.056	−0.050	0.288	−0.0007	0.0009	18.9
5	0.06	−0.5	7.50	16.9	1.01	−0.425	−0.167	−1.17	−0.050	−0.055	0.302	−0.0012	0.0010	18.4
6	0.06	0	7.82	12.0	1.05	−0.338	−0.195	1.44	−0.036	−0.064	0.304	−0.0015	0.0012	18.6
7	0.07	−1.0	7.85	15.9	0.96	−0.400	−0.155	+1.51	−0.053	−0.057	0.267	−0.0007	0.0009	20.5
8	0.07	−0.5	7.7	20.6	1.04	−0.367	−0.154	⋯	−0.050	−0.059	0.297	−0.0012	0.0010	18.5
9	0.07	0	8.14	14.2	1.05	−0.296	−0.197	0.482	−0.030	−0.069	0.307	−0.0015	0.0012	18.7

The only qualitative change in the results is that the nucleon is not sufficiently bound over the entire range of coupling constants investigated.

We would like to thank G. L. Shaw and P. W. Coulter of Stanford University for bringing this error to our attention.

Pion Exchange Currents in Deuteron Photodisintegration Dispersion Theory, MALCOLM H. SKOLNICK [Phys. Rev. **136**, B1493 (1964)]. Equation (4.33) should read

$$A(p)=e^{\delta(p)}b\left\{A_0(p)+r(p)+\frac{2\bar{R}M}{r_s(\gamma^2-\alpha_+^2)}\left(\frac{p^2+\alpha_+^2}{p^2+\alpha_-^2}\right)^{1/2}\left[\frac{p^2+\gamma^2}{l_c^2-\alpha_+^2}\left(\frac{\alpha_+}{p^2+\alpha_+^2}+\frac{l_c(\gamma^2-\alpha_+^2)}{(p_c^2+\gamma^2)(p^2+l_c^2)}\right)\right]-\frac{\gamma}{p_c^2+\gamma^2}\right\}. \quad (4.33)$$

Similarly Eqs. (4.39) and (4.41) should read

$$A_{\text{pole}}=\frac{\bar{R}M}{p^2+l_c^2}+\frac{2\bar{R}M}{r_s(\gamma^2-\alpha_+^2)}\left(\frac{p^2+\alpha_+^2}{p^2+\alpha_-^2}\right)^{1/2}\left\{\left(\frac{p^2+\gamma^2}{l_c^2-\alpha_+^2}\right)\left[\frac{\alpha_+}{p^2+\alpha_+^2}+\frac{l_c(\gamma^2-\alpha_+^2)}{(p_c^2+\gamma^2)(p^2+l_c^2)}\right]-\frac{\gamma}{p_c^2+\gamma^2}\right\}, \quad (4.39)$$

$$a(\nu_c)=\frac{1}{\nu_c+B}+K\left\{\left(\frac{\gamma^2}{\nu_c+L/M}\right)\left[(1/\alpha_+)+\frac{L/M}{\nu_c(M\nu_c+\gamma^2)^{1/2}}\right]-\gamma/\nu_c\right\}. \quad (4.41)$$

Electron Spin-Echo Envelope Modulation, L. G. ROWAN, E. L. HAHN, AND W. B. MIMS [Phys. Rev. **137**, A61 (1965)]. Directly following Eq. (10), the next line gives a relationship $2\hbar\omega_1=g\beta H_{\text{rf}}$. The symbol g should be replaced by $g'=(g_1g_{11}/g)$ for the case where H_{rf} is perpendicular to the plane containing H_0 and the crystalline c axis. In the next line below, $(2g\beta)$ should be replaced by $(g'\beta/2)$. The symbol g appearing in Eq. (15) and everywhere in the Appendix should be replaced by g'.

This change does not effect the physical results of echo envelope determinations because the over-all amplitude, determined by the g' factor associated with off-diagonal transition matrix elements, is normalized to arbitrary units of amplitude.

Internal Field in General Dipole Lattices, F. W. DE WETTE AND G. E. SCHACHER [Phys. Rev. **137**, A78 (1965)]. In Eq. (23), for $(j_2+j_2\xi_2)$, read $(j_2+j_3\xi_2)$. In Eqs. (44) and (45), replace $\Omega_{j_3}(\mu_1,\mu_2)$ by $\tilde{\Omega}_{j_3}(\mu_1,\mu_2)$. The definition of $\tilde{\Omega}_{j_3}$ is obtained from Eq. (24) by placing a minus sign in front of the first term inside the curly brackets. In Eqs. (51) and (52), replace $\Omega_0(\mu_1,\mu_2)$ by $\tilde{\Omega}_0(\mu_1,\mu_2)$. The definition of $\tilde{\Omega}_0$ is obtained from Eq. (25) by placing a minus sign in front of the first term inside the curly brackets. The exponent in the right-hand side of Eq. (72) should read

$$-2\pi i\{[j_1+(\lambda_3+j_3)\xi_1]\mu_1 +[j_2+(\lambda_3+j_3)\xi_2]\mu_2+k_3\lambda_3\}.$$

The authors are indebted to Dr. L. T. Klauder for pointing out a number of mistakes in the original manuscript.

Pressure Dependence of the Emission from Ga($As_{1-x}P_x$) Electroluminescent Diodes, G. E. FENNER [Phys. Rev. **137**, A1000 (1965)]. The expression $\eta(x)=\eta(0)\exp -x/\Lambda$ following Eq. (1) should read $\eta(x)=\eta(0)-x/\Lambda$. The factor $\eta(0)$ in Eq. (2) should be dropped and the calculated curves in Fig. 5 shifted by the appropriate factor $\eta(0)$.

Weak Interactions and Self-Consistent Theories, MAHIKO SUZUKI [Phys. Rev. **136**, B769 (1964)]. The statement made at the end of the paper, "present arguments lead to the contradictions independently of the behavior of the weak vertices near the light cone" is erroneous. In fact, if one admits a singular behavior like $1/q^2$, implying a massless boson, the theorem in the paper fails and therefore one can avoid the claim that self-consistent theories

encounter a disagreement with experiment in weak interactions. Then, one must instead overcome difficulties associated with the massless boson. I am indebted to Dr. Th. A. J. Maris for pointing this out.

π^--p Interactions at 683 MeV/c, R. A. BURNSTEIN, G. R. CHARLTON, T. P. DAY, G. QUARENI, A. QUARENI-VIGNUDELLI, G. B. YODH, AND I. NADELHAFT [Phys. Rev. **137**, B1044 (1965)]. There were some omissions in Ref. 7 of this paper, which should read: [7]M. Olsson and G. B. Yodh, Phys. Rev. Letters **10**, 353 (1963); Bull. Am. Phys. Soc. **9**, 27 (1964); University of Maryland Technical Report No. 358, 1964 (unpublished) and M. Olsson (thesis), University of Maryland Technical Report No. 379, 1964 (unpublished). The model herein is based upon the isobar model of S. J. Lindenbaum and R. M. Sternheimer, Phys. Rev. **105**, 1874 (1957); **106**, 1107 (1957); R. M. Sternheimer and S. J. Lindenbaum, Phys. Rev. **109**, 1723 (1958); **123**, 333(1961).

Frequency Dependence of the Two-Magnon Ferrimagnetic Resonance Linewidth, P. E. SEIDEN AND M. SPARKS [Phys. Rev. **137**, A1278 (1965)]. Reference 1 should read: E. Schlömann, in Proceedings of the Conference on Magnetism and Magnetic Materials, 1956, Boston, Massachusetts (unpublished). In the sentence after Eq. (14) on p. A1280 the factor $\cos\theta_k$ should be $\cos\theta_u$.

Electromagnetic Structure of the Giant Resonance in Oxygen-16, F. H. LEWIS, JR. [Phys. Rev. **134**, B331 (1964)]. Because of a numerical error in the original calculations, the graphs shown below should replace Figs. 1, 2, and 4 of the paper, as indicated. The discussion of the results should remain unchanged except for the remarks in the last paragraph starting on p. B335 concerning the state at 17.3 MeV. The theoretical cross section for this

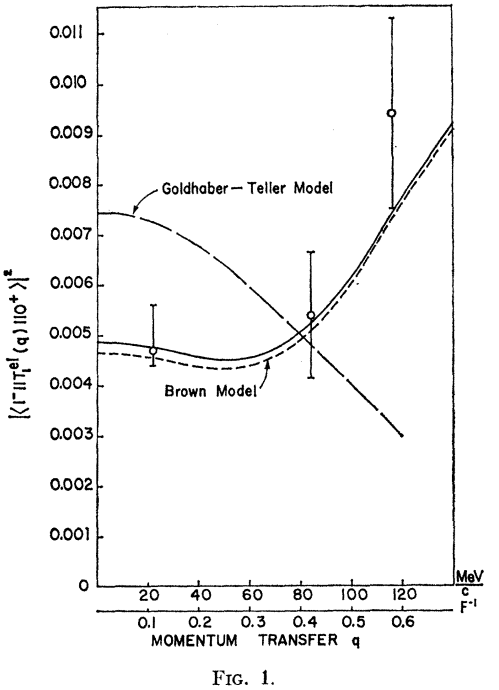

FIG. 1.

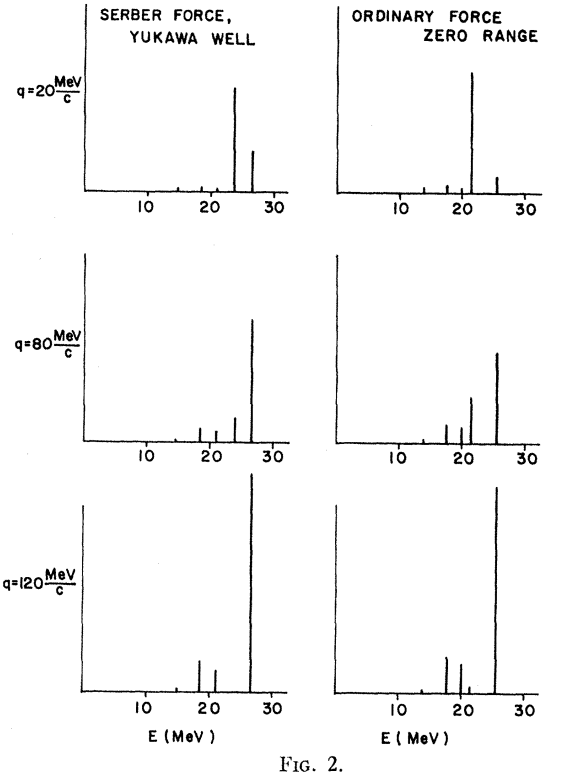

FIG. 2.

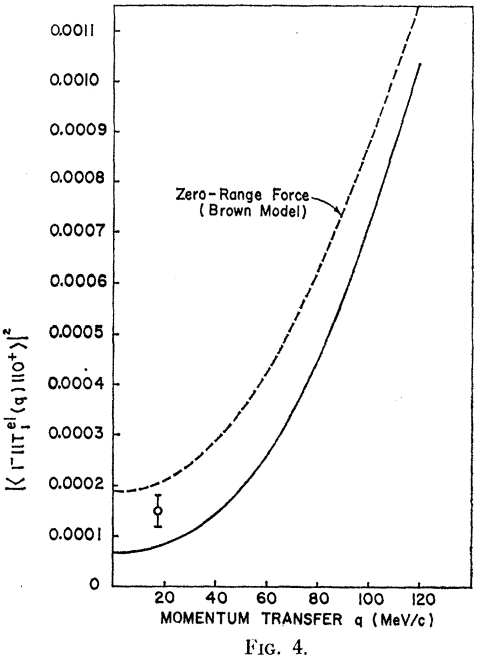

FIG. 4.