so

$$N_i = S^{-1}{}_{i\mu}\bar{N}_\mu\,,$$
$$N_i^0 = S^{-1}{}_{i\mu}\bar{N}_\mu{}^0 \quad \text{for} \quad i=1,\cdots,n. \qquad \text{(III.5')}$$

Equations (III.4') can be inverted by the Wiener-Hopf method to give:

$$\bar{N}_i \underset{s>s_M}{=} O_i\bar{N}_i{}^0 \quad \text{for} \quad i=1,\cdots,n. \qquad \text{(III.6')}$$

Thus we have:

$$N_i(s) = B_i(s) + \int_{s_\mu}^{s_M} U_{i\mu}(s,s')N_\mu(s')ds' + S^{-1}{}_{\mu\nu}\int_{s_M}^{\sigma}(U_{i\mu}O_\nu)(s,s')\bar{N}_\nu{}^0(s')ds'$$
$$\text{for} \quad s<s_M \quad i=1,\cdots,n \qquad \text{(III.7'a)}$$

where

$$(U_{i\mu}O_\nu)(s,s') = \int_{s_M}^{\sigma} ds'' U_{i\mu}(s,s'')O_\nu(s'',s')\,,$$

$$N_i{}^0(s) = B_i(s) + \int_{s_\mu}^{s_M} U_{i\mu}(s,s')N_\mu(s')ds' + \int_{s_M}^{\sigma} K_{i\mu}(s,s')N_\mu(s')ds'$$

or

$$\bar{N}_i{}^0(s) = \bar{B}_i(s) + S_{i\mu}\int_{s_\mu}^{s_M} U_{\mu\nu}(s,s')N_\nu(s')ds' + S_{i\mu}S^{-1}{}_{\nu\rho}\int_{s_M}^{\sigma}(K_{\mu\nu}O_\rho)(s,s')\bar{N}_\rho{}^0(s')ds'$$
$$\text{for} \quad s_M<s<\sigma; \quad i=1,\cdots,n \qquad \text{(III.7'b)}$$

where

$$\bar{B}_i(s) = S_{i\mu}B_\mu(s)\,,$$

$$(K_{\mu\nu}O_\rho)(s,s') = \int_{s_M}^{\sigma} ds'' K_{\mu\nu}(s,s'')O_\rho(s'',s')\,.$$

Equations (III.7') are a system of coupled integral equations with $N_i(s)$ for $s_1<s<s_M$ and $\bar{N}_i{}^0(s)$ for $s_M<s<\sigma$ as unknown functions. The functions $U_{ij}(s,s')$ are square integrable for $s_1<s<s_M$, $s_1<s'<\sigma$, and the functions $K_{ij}(s,s')$ are square integrable for $s_M<s$, $s'<\sigma$. The singular function $O_\nu(s'',s')$ no longer appears alone but only folded with these nonsingular functions, so the difficulty explained at the beginning of this erratum will not arise. The functions $(U_{i\mu}O_\nu)(s,s')$ can easily be shown to be square integrable for $s_1<s<s_M$, $s_1<s'<\sigma$ and likewise the functions $(K_{\mu\nu}O_\rho)(s,s')$ for $s_M<s$, $s'<\sigma$. Thus, we have achieved a system of Fredholm equations.

Electric Polarization of the Deuteron by a Point Charge, RICHARD J. DRACHMAN [Phys. Rev. **132**, 374 (1963)]. An algebraic error was responsible for several incorrect results. In Eq. (17), the coefficient of the last term should be 3, rather than $\frac{3}{2}$. Two subsequent equations should now read as follows:

$$V_2{}^{(1)}(t) = \frac{2Z^2Me^4}{3\hbar^2t^4}\{-\tfrac{3}{8}+e^{-2t}[-2t^5+t^4-t^3+3t^2-\tfrac{3}{2}t+\tfrac{3}{2}] - 9e^{-4t}/8 - 4t^6\,\mathrm{Ei}(-2t)\}\,, \qquad \text{(18)}$$

$$V_2{}^{(1)}(t) \sim -\frac{Z^2Me^4}{4\hbar^2t^4}\left\{1-e^{-2t}(4t^2+4t-16)+3e^{-4t} - \frac{16}{3}t^5e^{-2t}\sum_{n=6}^{N}\frac{n!(-1)^n}{(2t)^n}\right\}. \qquad \text{(20)}$$

Using the exact equation (18) instead of the asymptotic equation (20), one obtains new and smaller entries in Table I:

TABLE I. The deviation of the dipole polarization potential from t^{-4} form: $V_2{}^{(1)}(t) = -(Z^2Me^4/4\hbar^2t^4)(1+D)$, where $x=t/2\gamma=2.2t$ F.

t	D
3	−0.103
4	−0.024
5	−0.005
6	−0.001

Finally, the limiting value discussed in the first two sentences of p. 377 should read $V_2{}^{(1)}(0)=0$. Delete the remainder of the paragraph.

Decay of Pd^{100} to Odd-Odd Rh^{100}, J. S. EVANS AND R. A. NAUMANN [Phys. Rev. **138**, B1017 (1965)]. We have incorrectly used the ambiguous word "lifetime" in our paper. If "lifetime" be replaced by "half-life" everywhere, then the calculations and conclusions are numerically correct.

Doubly Excited States in Lithium, J. D. GARCIA AND J. E. MACK [Phys. Rev. **138**, A987 (1965)]. Parts of Table IV and the accompanying discussion are incorrect, partly in view of the work of Werner[13] and of Toresson and Edlén,[14] to which we should have referred, and partly on account of a printer's misalignment. In Table IV, first column, for line (5) "1420 Å (8.72 eV)" should be deleted, and instead of "9572 Å (1.30 eV)," which was aligned with line (7), there should be "8517 Å (1.45 eV)," placed between lines (6) and (7) to indicate ambiguity; actually the line may be a transition between (unlisted) doublet terms. The terms associated with the four new assignments should have been listed in a new Table V, which must be considered approximate and incomplete pending further studies of the fine and hyperfine structure and isotope shift