totic form for Eq. (A11):

$$\tan\delta_l^{\pm}\sim\alpha b_l^{\pm}x^{2l+1}[(2l+1)!!]^{-2}[1-\alpha b_l^{\pm}(2l+1)^{-1}]^{-1}.$$

This same expression is obtained with the approximation $\tan\delta_l^0=0$ in Eq. (A11) (for $x\ll l$). The numerator of this expression is just the Born approximation for the phase shifts when $x\ll l$, while the denominator is a correction term resulting from the distortion of the outgoing wave by the spin-orbit interaction. For l near to or less than X the $\tan\delta_l^0$ term in Eq. (A11) must generally be treated in detail [as in Eq. (35)]. However, near a resonance in $\tan\delta_l^0$ Eq. (A11) gives for $x\ll l$

$$\tan\delta_l^{\pm}\sim -x^{2l+1}[1+\alpha b_l^{\pm}(2l+1)^{-1}]\times\{\alpha b_l^{\pm}[(2l-1)!!]^2\}^{-1},$$

and the polarization can be appreciable. On the other hand, a resonance in $\tan\delta_l^{\pm}$ occurs for $x\ll l$ when

$$\tan\delta_l^0=-x^{2l+1}[1-\alpha b_l^{\pm}(2l+1)^{-1}]\times\{\alpha b_l^{\pm}[(2l-1)!!]^2\}^{-1},$$

which splits the resonance energies for $\tan\delta_l^{\pm}$.

New Measurements of β-Circularly-Polarized γ Angular-Correlation Asymmetry Parameters in Allowed β Decay, L. G. MANN, D. C. CAMP, J. A. MISKEL, AND R. J. NAGLE [Phys. Rev. **137**, B1 (1965)]. A recent re-examination of our Co^{58} source showed Co^{60} impurity which amounted to about 2% of the total activity at the time of the original measurements. Another recently acquired source for which the specifications indicated $<0.01\%$ Co^{60} shows a similar 2% contamination. Since the polarimeter is much more sensitive to the radiations from Co^{60} than to those from Co^{58}, this impurity introduced about a 10% error in our measurement of the β-γ (circularly-polarized) asymmetry parameter of Co^{58}. Correcting for the Co^{60}, the result in Table V should read $A=-0.185\pm0.011$ instead of -0.213 ± 0.012 for Co^{58}. The corresponding values for $|M_F|$ and the isospin impurity coefficient α are $(0.9\pm0.6)\times10^{-3}$ and $(0.34\pm0.21)\times10^{-3}$, respectively. These results now agree very well with the nuclear alignment data.

Consistency Conditions on the Strong Interactions Implied by a Partially Conserved Axial-Vector Current, STEPHEN L. ADLER [Phys. Rev. **137**, B1022 (1965)]. In Eqs. (16) and (23),

$$[(p_{10}/M)(p_{20}/M)2k_0]^{1/2}$$

should be

$$[(p_{10}/M)(p_{20}/M)2q_0]^{1/2}.$$

Theory of the Electronic Thermal Conductivity of Superconductors Containing Paramagnetic Impurities, LEONARD W. GRUENBERG [Phys. Rev. **138**, A78 (1965)]. In Eq. (II.24), the right-hand side of each equation should be multiplied by $\frac{1}{4}$.

Equation (III.9) should read

$$\kappa_s/\kappa_n=\frac{\omega_0}{\pi^2kT}(\alpha^{-2/3}-1)\times\{1+(\Gamma_{tr}^{+})^{-1}[2\Gamma^M(\alpha^{-2/3}-1)+2\Gamma^{M'}]\}^{-1}e^{-\beta\omega_0},$$

where $\alpha=2\Gamma^M/\Delta$.

In Eq. (III.10), τ^+ should be replaced by τ_{tr}^{+}.

The denominator of the first term on the right-hand side of Eq. (III.4) should be squared.

Taking note of the identity $T_2^2+T_3T_4=0$, we see that our Eq. (II.26) is identical to Eq. (2.47) of Ambegaokar and Griffin.[1] On the other hand, our results in Sec. IIIC differ strikingly from theirs. This is because we have looked at the limiting case $\Gamma^N\approx0$ while they assume $\Gamma^N\gg\Gamma^M$. We believe that for the case of Gd in very pure La, $\Gamma^N\ll\Gamma^M$ and hence our results apply.

If $\Gamma^N>kT_c$, the Boltzmann equation does not give the correct limiting value for κ_s/κ_n for temperatures near T_c.

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[1] V. Ambegaokar and A. Griffin, Phys. Rev. **137**, A1151 (1965).

Further Evidence for Pignotti's *R* Trajectory, ROGER J. N. PHILLIPS AND WILLIAM RARITA [Phys. Rev. **138**, B723 (1965)]. In Eq. (1) interchange the right-hand sides of lines (3) and (4).

Ground-State Properties and Low-Lying States of the N^{14} Nucleus, NAZAKAT ULLAH AND R. K. NESBET [Phys. Rev. **134**, B308 (1964)]. In carrying out a detailed configuration interaction calculation, an error was found in the computer program used to compute radial integrals. On repeating the numerical calculation with the corrected program, it turns out that the approximate Hartree-Fock binding energy is only one-third of its experimental value. The rms radius is, as before, in fair agreement with its experimental value, the quadrupole moment has the correct sign, but the magnetic dipole moment is too large by a factor of 2. The excitation energy of the low-lying state $J^{\pi}=1^+$, $T=0$ is now in good agreement with its experimental value.

The corrected Tables III and IV, giving the new parameters and results of the approximate Hartree-