Errata

Self-Consistent Three-Body Calculation of Pion-Nucleon Scattering, RONALD AARON [Phys. Rev. **151**, 1293 (1966)]. In Eq. (2) the quantity $\int_0^\infty dq$ should be replaced by $\int_0^\infty q^2 dq$. In Eq. (3) the quantities $\omega_k^{1/4}$ and $\omega_{k'}^{1/4}$ should be replaced by $\omega_k^{1/2}$ and $\omega_{k'}^{1/2}$, respectively.

Three-Photon Decay of Positronium ^1S State as a Test of Charge Conjugation Invariance, JOSEPH SCHECHTER [Phys. Rev. **132**, 841 (1963)]. In the matrix element of Eq. (7) $\mathbf{K}_i \cdot \mathbf{K}_j$ should be replaced by $K_i \cdot K_j$. Then it vanishes, so that a form with more powers of the momentum must be used. I am grateful to Dr. H. Weisberg and Dr. S. Berko for pointing this out. See also A. Dolgov and L. Ponomarev, Moscow Report No. N433, 1966 (unpublished).

General Analysis of Nucleon-Nucleon Scattering: Critical Tests for Regge-Pole Theory, ELLIOT LEADER AND RICHARD C. SLANSKY [Phys. Rev. **148**, 1491 (1966)]. We are very grateful to Dr. J. K. Perring for pointing out that some of our formulas are not correct in the relativistic region. The following corrections should be made:

(1) In Eqs. (2.11g)–(2.11j), (4.8f)–(4.8i), (4.33), and (4.39), replace $\frac{1}{2}\theta$ by θ_L, the laboratory angle of scattering.

(2) Replace Eq. (2.11e) by

$$\begin{aligned}
I_0 C_{KP} = {} & \tfrac{1}{4}[\,|\phi_3-\phi_4|^2 - |\phi_1+\phi_2|^2]\cos(\theta_L-\phi_L) \\
& +\mathrm{Re}\{\phi_5^*(\phi_1+\phi_2-\phi_3+\phi_4)\sin(\theta_L-\phi_L)\} \\
& -\tfrac{1}{4}\,\mathrm{Re}\{(\phi_1-\phi_2+\phi_3+\phi_4)^*(-\phi_1+\phi_2+\phi_3+\phi_4)\} \\
& \qquad\qquad \times\cos(\theta_L+\phi_L) \\
= {} & s^{-1}\,\mathrm{Re}\{(g_4-g_5)^*[2\beta(\theta_L-\phi_L)g_3 \\
& +\alpha(\theta_L-\phi_L)(g_1+g_2)] - (g_4+g_5)^*(g_1-g_2) \\
& \qquad\qquad \times\cos(\theta_L+\phi_L)\},
\end{aligned}$$

where ϕ_L is the laboratory angle of the recoil nucleon.

It has also been pointed out that the use of D_t for the polarization transfer parameter is not yet common practice. It is defined by

$$I_0 D_t = (\mathbf{N}0;\, 0\mathbf{N})$$

in the notation of, e.g., M. J. Moravcsik [*The Two-Nucleon Interaction* (Clarendon Press, Oxford, England, 1963)].

All relativistic experimental parameters used conform to the definitions of D. W. L. Sprung, Phys. Rev. **121**, 925 (1961).

Introduction to the N-Quantum Approximation in Quantum Theory, O. W. GREENBERG [Phys. Rev. **139**, B1038 (1965)]. It has been pointed out by A. Halprin that the mass renormalization terms diverge in fifth order of the coupling constant. To correct this difficulty, the following changes should be made: The paragraph below Fig. 6, p. B1045, should read: "Next we study the two self-mass diagrams in the bracket in Fig. 7. The first term is linear in f, while the second is quadratic. Here, we linearize the second term, replacing the f in the mass renormalization loop by the Born approximation f_0 ...." The open circles in the mass renormalization loop in Fig. 11 should be closed circles. The first term on the second line of Eq. (9) should read:

$$(2\pi)^{-3} g(1-L)\int dq\, \delta_m(q)[f_0(k_1+k_2-q) - f_0(k_1+k_2-q)\,|_{(k_1+k_2)^2=m^2}]f(k_1+k_2).$$

The first term in curly brackets in the equation for K on p. B1046 should read:

$$(2\pi)^{-3} g\frac{\partial}{\partial k^2}\int dq\, \delta_m(q) f_0(k-q).$$

The terms in $\mathfrak{M}$ in Eq. (10) should read $\mathfrak{M}\{k;\, f_0\}$. The equation above Fig. 9 should read

$$g_0 = g[1-g^{-1}M_1\{f_0\} + (2g)^{-1}F_0\{f_0+\phi\}]^{-1}.$$

The equation on p. B1047 and the remainder of Sec. 3 should read

$$\phi_1(k) = (m^2-k^2)^{-1}[\Phi\{k;\, f_0\} + \mathfrak{M}\{k;\, f_0\}].$$

"A careful but straightforward analysis of the iterative solution to Eq. (10) shows that all iterations are finite." Therefore, with the changes given above, the difficulty pointed out by Halprin is removed, and the power-series-expansion solution of Eq. (10) for the first N-quantum approximation to the model with $\mathfrak{L}_I = gA^3$ is finite, term by term, in all orders of the coupling constant.

We are happy to thank Dr. Richard Brandt for helpful conversations.