

Low-temperature heat capacity and thermodynamic properties of four boron nitride modifications

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Abstract

A short review of the low-temperature heat capacity studies of four boron nitride (BN) modifications is presented. C_p versus (T) dependences of high-ordered hexagonal and disordered (turbostratic) modifications were studied by adiabatic calorimetry. The influence of disordering on the heat capacity of boron nitride is shown. The values of thermodynamic properties (heat capacity, entropy, enthalpy, and formation enthalpy) of four BN modifications is reported.

INTRODUCTION

Boron nitride, first synthesized 150 years ago by W.H. Balmain, is now successfully used in many fields of science and technology. Electronics and nuclear engineering materials, superhard construction materials, lubricants, and refractory materials are only a few of the applications of different boron nitride modifications.

THE STRUCTURE OF THE BORON NITRIDE POLYMORPHOUS MODIFICATION

Under normal conditions, boron nitride exists in four polymorphous modifications: two graphite-like forms, hexagonal (hBN) (Fig. 1) and uncommon rhombohedral (rBN) (Fig. 2), and two dense forms, wurtzitic

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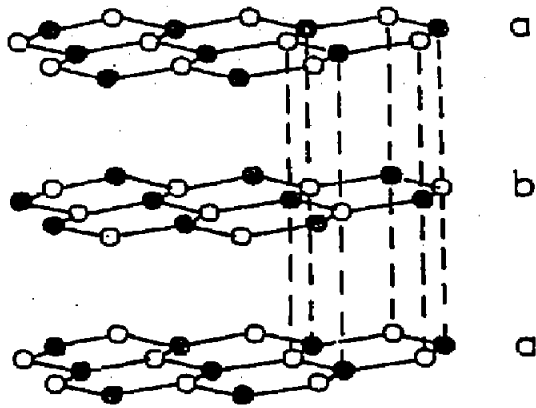


Fig. 1. Crystal structure of graphite-like BN modification (a and b are types of layers).

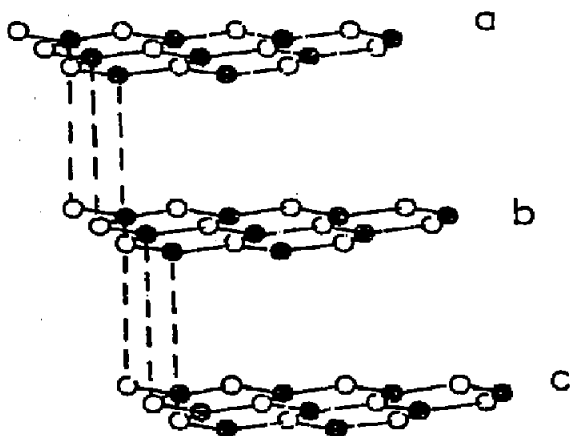


Fig. 2. Crystal structure of rhombohedral BN modification (a, b and c are types of layers)

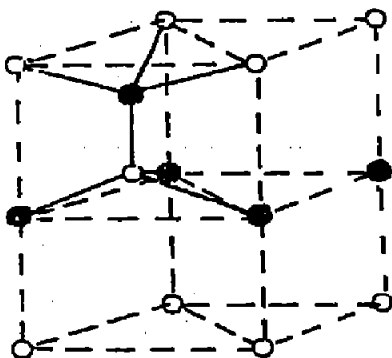


Fig. 3. Crystal structure of wurtzitic BN modification.

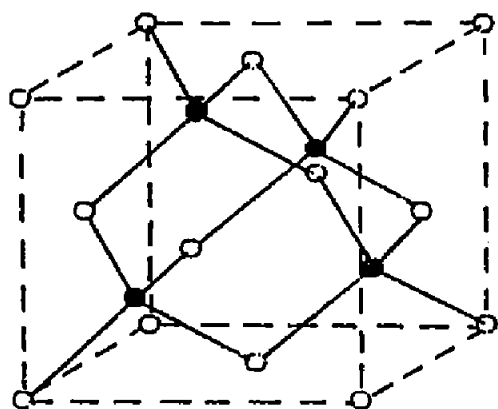


Fig. 4. Crystal structure of cubic BN modification.

(wBN) (Fig. 3) and cubic (cBN) (Fig. 4). The characteristics of these boron nitride modifications are given in Table 1.

LOW-TEMPERATURE HEAT CAPACITY

Graphite-like hexagonal BN

The low-temperature heat capacity studies of hBN were reported in refs. 5–7. Dworkin et al. [5] determined the C_p versus (T) dependence in the 20–300 K range using adiabatic calorimetry. The studied sample of hBN

TABLE 1

Crystallochemical characteristics of BN polymorphous modifications

| Space group | hBN P6 ₃ /mmc | rBN R3m | cBN F43m | wBN P6 ₃ mc |
|---|-----------------------------|---------------|---------------|---------------------------|
| Lattice period (Å) | | | | |
| <i>a</i> | 2.504 [1] | 2.504(5) [2] | 3.6160(3) [3] | 2.551(0) [4] |
| <i>c</i> | 6.6612 [1] | 10.000(2) [2] | – | 4.214(8) [4] |
| Density ^a (from X-ray data) (g cm ⁻³) | 2.279 | 2.276 | 3.486 | 3.470 |
| Shortest interatomic distance B–N (Å) ^a | 1.446 | 1.450 | 1.565 | 1.568 |

^a Calculated from the indicated lattice parameters.

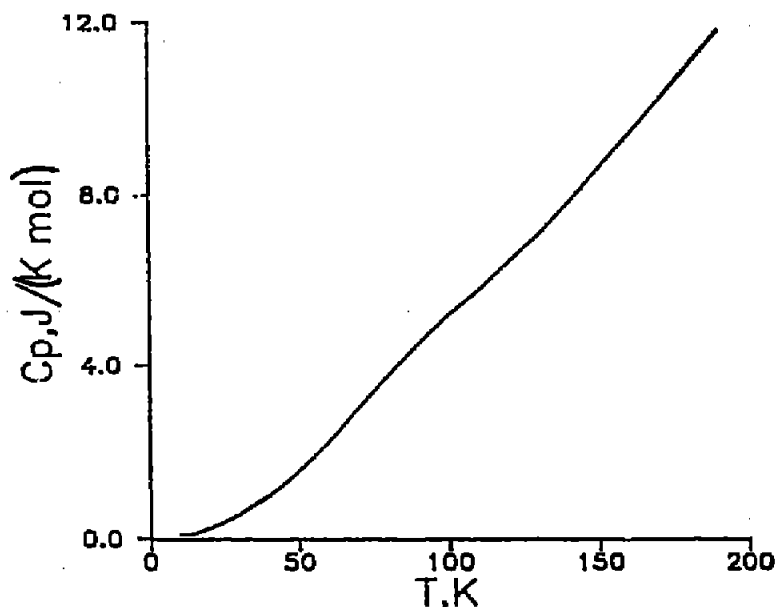


Fig. 5. C_p versus (T) curve for graphite-like BN [7] in the region of a slightly convex anomaly.

contained 0.15% of iron and boron oxides impurities. Later, Sirota et al. [6] investigated the heat capacity of this modification, but their sample contained larger amounts of impurities ($\approx 2.8\%$) and their C_p versus (T) data had a large deviation of experimental points (3–10%). In ref. 7, the C_p versus (T) dependence of hBN was determined in the 10–300 K range in an adiabatic calorimeter. The content of impurities in the sample was less than 0.15%; here [7], the authors took into account the influence of impurities on the C_p versus (T) data.

A comparison of these results [5–7] show that they differ slightly (in the order of percents). A slightly convex anomaly is a common feature for all the experimental curves [5–7] in the range of 60–140 K (Fig. 5). It is interesting that the C_p versus (T) dependence for graphite has no anomaly in this region [8].

Cubic BN

The low-temperature heat capacity studies of cubic boron nitride are reported in refs. 9–12. The cBN sample studied by Sirota and Kofman [9] contained up to 4% of impurities, whereas samples studied in refs. 10–12 contained 99.8–99.9% of the main material. Moreover, the methods of heat capacity measurement used in ref. 9 did not provide accurate C_p versus (T) data. The results of refs. 10–12 are in good agreement with one another (deviations of about 1%), unlike those of ref. 9 which are sufficiently lower

(dozens of percent). Therefore, only the more reliable results of refs. 10-12 were used for the estimation of cBN thermodynamic properties.

Wurtzitic BN

The low-temperature heat capacity studies of boron nitride wurtzitic modification are reported in refs. 7 and 13. The sample studied in ref. 13 contained only about 92% wBN. The results of this research had a considerable deviation of experimental points (up to 5%). Gorbunov et al. [7] studied the heat capacity of wBN in the 6-300 K range using adiabatic calorimetry. The sample contained $\approx 99.7\%$ of the main material. A small anomaly with a maximum at 21 K was found on the C_p versus (T) curve of wBN. The nature of this anomaly was not explained by the authors. A comparison of the above results has shown that the C_p versus (T) curve in ref. 13 lies considerably higher than the curve in ref. 7 (up to 25%) and even higher than the curve for hBN. This demonstrates the low reliability of data in ref. 13.

Rhombohedral BN

The rhombohedral modification of BN usually occurs as an admixture in the graphite-like hexagonal modification; its properties have been insufficiently studied. Recently, Solozhenko et al. [14] proposed a method of rBN synthesis involving the interaction of boron and ammonia at high pressure and high temperature.

The low-temperature heat capacity of the synthesized material has been measured by adiabatic calorimetry [15]. The studied sample contained 2% of hBN and 0.8% of boron oxide, and carbide and magnesium nitride impurities. The degree of three-dimensional ordering was 0.93. The authors took into consideration the influence of impurities while calculating the thermodynamic properties. No anomaly of the heat capacity was observed, unlike hBN.

THE INFLUENCE OF DISORDERING ON THE HEAT CAPACITY OF BORON NITRIDE

A small convex anomaly of graphite-like hexagonal boron nitride, reported in ref. 7, attracted our attention. We have analysed the possible factors which may result in the appearance of an anomaly on the C_p versus (T) curve of hBN. This anomaly can probably be explained as the consequence of (1) the influence of impurities; (2) the high dispersion of the sample; (3) the increased distance between layers in the crystal structure.

The literature contains conflicting information on the influence of impurities on the heat capacity. An impurity may either penetrate into the crystal lattice of the studied material (as, for example, on doping ammonium chloride with heavy metals [16]), or intercalate in the layered structures (as described for PbI_2 [17]). Both cases featured a slightly convex anomaly of heat capacity in the region below 40 K, caused by the interaction of atoms or molecules of intercalants. But for the hBN reported in ref. 7 (content of impurities $\approx 0.15\%$), the interaction of admixtural atoms is improbable, and we assume that this does not occur.

Because the graphite-like hexagonal boron nitride studied in ref. 7 was a very dispersed powder with a particle size of $\geq 1 \mu\text{m}$, an anomaly of hBN heat capacity may be attributed to the influence of the size-factor on the C_p values, because it is known that the surplus heat capacity of ultra-dispersed samples yields a curve with the maximum at 100 K [18, 19].

However, we believe that the most probable reason for the appearance of a slightly convex anomaly of the C_p of hBN [5–7] was the occurrence of a low-frequency mode in the vibrational spectrum of this modification because of the increased distance between layers caused by the disordering of the layers alternating along the c -axis. The disordering of graphite-like hexagonal boron nitride due to the displacement of layers in arbitrary positions by translation or turning in the basal plane, has been investigated earlier [20] and was named turbostratic defects of packing.

For the experimental examination of our hypothesis, we studied two samples of graphite-like boron nitride: high-ordered hBN and fully disordered (turbostratic) BN. The particles sizes were nearly equal to minimize the size-factor influence on the heat capacity.

For the heat capacity investigation, a sample of turbostratic BN (tBN) synthesized according to ref. 21, and a sample of Fluka AG fine-fraction hBN were studied. The starting powders were subjected to thermal treatment: tBN at 1400 K in a vacuum of 0.1 Pa for 3 h; hBN at 2100 K in nitrogen atmosphere (oxygen content less than 0.003%) for 5 h. All subsequent operations with the samples were carried out in an atmosphere of dry inert gas.

According to granulometric and electron microscopic analyses, the size of the tBN particles varied within the limits 0.1–0.2 μm ; the specific surface of the studied sample, calculated by the BET method from the isotherms of nitrogen absorption at 77 K, was $210 \pm 20 \text{ m}^2 \text{ g}^{-1}$. However, the value of the specific surface area may be slightly high due to the possible sorption of nitrogen in the sample volume. The degree of three-dimensional ordering of the tBN crystal lattice measured by the method in ref. 22 using the X-ray data, was equal to 0; the distance between layers d_{002} was 3.56 Å and the closest B–N interatomic distance was 1.45 Å, i.e. the same as for high-ordered hBN.

The mean value of hBN particle size according to granulometric

TABLE 2

Impurity contents of BN samples (mass%)

| | C | O | Si | Ca | Al | Mg | Fe | Cl |
|-----|-----|-----|------|--------|--------|--------|--------|------|
| tBN | 0.1 | 0.1 | 0.01 | 0.0002 | 0.0004 | 0.0001 | 0.0003 | 0.01 |
| hBN | 0.4 | 0.1 | 0.03 | 0.02 | 0.015 | 0.0006 | 0.0025 | – |

measurements was $0.4 \mu\text{m}$ and the specific surface area was $60 \pm 10 \text{ m}^2 \text{ g}^{-1}$. The degree of three-dimensional ordering of the crystal lattice was 0.98 ± 0.2 , and the distance between layers d_{002} was $3.330(4) \text{ \AA}$.

The contents of impurities in the studied samples determined by mass spectrometric investigation and X-ray microanalysis, are given in Table 2.

TABLE 3

Experimental heat capacity of high-ordered hBN

| T (K) | C_p ($\text{J K}^{-1} \text{ mol}^{-1}$) | T (K) | C_p ($\text{J K}^{-1} \text{ mol}^{-1}$) | T (K) | C_p ($\text{J K}^{-1} \text{ mol}^{-1}$) |
|------------|---|------------|---|------------|---|
| 15.04 | 0.1654 | 82.71 | 3.698 | 182.79 | 11.01 |
| 16.04 | 0.1798 | 85.42 | 3.847 | 186.75 | 11.46 |
| 17.28 | 0.1983 | 88.14 | 4.099 | 190.67 | 11.63 |
| 18.80 | 0.2259 | 91.34 | 4.327 | 194.34 | 12.02 |
| 20.54 | 0.2777 | 94.99 | 4.571 | 198.19 | 12.30 |
| 22.47 | 0.2996 | 98.62 | 4.819 | 203.23 | 12.68 |
| 24.50 | 0.3534 | 102.42 | 5.066 | 208.37 | 13.05 |
| 26.80 | 0.4224 | 105.29 | 5.261 | 213.54 | 13.44 |
| 29.19 | 0.5458 | 107.69 | 5.459 | 218.53 | 13.83 |
| 31.85 | 0.6129 | 111.36 | 5.722 | 223.96 | 14.26 |
| 34.78 | 0.7394 | 115.16 | 5.987 | 229.04 | 14.65 |
| 37.53 | 0.8604 | 119.23 | 6.275 | 234.01 | 15.03 |
| 40.46 | 0.9911 | 123.28 | 6.594 | 238.99 | 15.46 |
| 43.72 | 1.160 | 127.13 | 6.866 | 243.97 | 15.79 |
| 46.59 | 1.322 | 131.05 | 7.133 | 248.96 | 16.18 |
| 49.68 | 1.525 | 135.11 | 7.431 | 254.21 | 16.61 |
| 52.68 | 1.706 | 139.00 | 7.737 | 259.40 | 16.96 |
| 55.92 | 1.899 | 142.68 | 8.033 | 264.38 | 17.36 |
| 58.91 | 2.069 | 146.56 | 8.308 | 269.43 | 17.76 |
| 61.88 | 2.228 | 150.51 | 8.617 | 274.48 | 18.14 |
| 64.74 | 2.432 | 154.41 | 8.911 | 279.54 | 18.51 |
| 67.59 | 2.660 | 158.16 | 9.216 | 284.57 | 18.87 |
| 70.36 | 2.840 | 161.93 | 9.480 | 289.55 | 19.23 |
| 73.09 | 3.040 | 167.03 | 9.854 | 294.65 | 19.59 |
| 75.85 | 3.244 | 171.07 | 10.16 | 299.82 | 19.95 |
| 78.57 | 3.444 | 174.96 | 10.47 | 304.81 | 20.30 |
| 81.27 | 3.640 | 178.84 | 10.83 | | |

TABLE 4

Experimental heat capacity of turbostratic BN

| T (K) | C_p (J K ⁻¹ mol ⁻¹) | T (K) | C_p (J K ⁻¹ mol ⁻¹) | T (K) | C_p (J K ⁻¹ mol ⁻¹) |
|------------|---|------------|---|------------|---|
| 11.59 | 0.1635 | 76.38 | 4.210 | 178.80 | 12.04 |
| 12.87 | 0.1713 | 79.32 | 4.449 | 183.71 | 12.43 |
| 14.01 | 0.2645 | 82.29 | 4.672 | 188.65 | 12.82 |
| 15.02 | 0.3472 | 85.15 | 4.886 | 193.47 | 13.21 |
| 15.99 | 0.3709 | 87.95 | 5.069 | 198.44 | 13.62 |
| 17.11 | 0.4429 | 91.24 | 5.329 | 203.47 | 14.02 |
| 18.59 | 0.5261 | 95.19 | 5.571 | 209.01 | 14.47 |
| 20.42 | 0.6173 | 99.05 | 5.868 | 213.91 | 14.87 |
| 22.42 | 0.7323 | 102.76 | 6.110 | 218.88 | 15.30 |
| 24.49 | 0.8495 | 106.51 | 6.363 | 223.85 | 15.68 |
| 26.58 | 0.9678 | 110.24 | 6.614 | 229.11 | 16.10 |
| 28.44 | 1.076 | 113.94 | 6.889 | 234.38 | 16.50 |
| 30.82 | 1.187 | 117.71 | 7.190 | 238.30 | 16.83 |
| 33.76 | 1.338 | 121.51 | 7.456 | 242.82 | 17.23 |
| 36.86 | 1.501 | 125.20 | 7.725 | 246.91 | 17.57 |
| 40.26 | 1.663 | 130.36 | 8.093 | 250.99 | 17.95 |
| 43.84 | 1.852 | 135.11 | 8.478 | 255.04 | 18.22 |
| 47.32 | 2.128 | 139.92 | 8.849 | 259.58 | 18.60 |
| 50.63 | 2.341 | 144.55 | 9.230 | 264.59 | 18.96 |
| 54.33 | 2.604 | 149.05 | 9.599 | 269.59 | 19.39 |
| 56.26 | 2.629 | 149.25 | 9.603 | 274.64 | 19.79 |
| 59.22 | 2.949 | 151.89 | 9.847 | 280.07 | 20.22 |
| 62.22 | 3.129 | 154.80 | 10.10 | 285.37 | 20.59 |
| 65.07 | 3.344 | 158.76 | 10.38 | 290.25 | 21.03 |
| 67.90 | 3.568 | 163.57 | 10.80 | 295.23 | 21.41 |
| 70.74 | 3.785 | 168.43 | 11.20 | 300.06 | 21.66 |
| 73.55 | 4.004 | 173.66 | 11.62 | 305.11 | 22.06 |

The heat capacities of hBN and tBN were determined by adiabatic calorimetry in the temperature range of 15–300 K. The experimental C_p versus (T) data are given in Tables 3 and 4. The smoothness of the C_p versus (T) curves was performed by LSM method in the normal polynomial. Values of the thermodynamic functions of hBN and tBN are given in Tables 5 and 6, respectively.

It is obvious from the C_p versus (T) curves that the heat capacity of high-ordered hBN has no anomaly and is in good agreement with that given in ref. 7, with the exception of the temperature range 100–200 K (Fig. 6). The investigation of the low-temperature heat capacity of turbostratic graphite-like boron nitride shows that its heat capacity is considerably higher than the corresponding hBN. In contrast to the C_p versus (T) curve for tBN, a convex anomaly is clearly featured. This demonstrates that the size

TABLE 5

Thermodynamic properties of high-ordered hBN

| T (K) | C_p^\ominus (J K ⁻¹ mol ⁻¹) | S^\ominus (J K ⁻¹ mol ⁻¹) | $H^\ominus(T) - H^\ominus(0)$ (J mol ⁻¹) | $-[G^\ominus(T) - H^\ominus(0)]/T$ (J K ⁻¹ mol ⁻¹) |
|------------|---|---|---|--|
| 16 | 0.1603 | 0.0530 | 0.641 | 0.0129 |
| 20 | 0.2501 | 0.0981 | 1.462 | 0.0249 |
| 25 | 0.3914 | 0.1685 | 3.066 | 0.0458 |
| 30 | 0.5632 | 0.2545 | 5.452 | 0.0728 |
| 35 | 0.7635 | 0.3560 | 8.769 | 0.1055 |
| 40 | 0.9906 | 0.4724 | 13.15 | 0.1436 |
| 45 | 1.242 | 0.6033 | 18.74 | 0.1870 |
| 50 | 1.517 | 0.7482 | 25.63 | 0.2355 |
| 60 | 2.129 | 1.077 | 43.86 | 0.3463 |
| 70 | 2.813 | 1.456 | 68.57 | 0.4761 |
| 80 | 3.556 | 1.879 | 100.4 | 0.6236 |
| 90 | 4.221 | 2.336 | 139.3 | 0.7878 |
| 100 | 4.900 | 2.815 | 184.9 | 0.9660 |
| 110 | 5.602 | 3.315 | 237.4 | 1.156 |
| 120 | 6.323 | 3.833 | 297.0 | 1.357 |
| 130 | 7.060 | 4.368 | 363.9 | 1.568 |
| 140 | 7.809 | 4.918 | 438.3 | 1.788 |
| 150 | 8.567 | 5.483 | 520.2 | 2.015 |
| 160 | 9.332 | 6.060 | 609.7 | 2.249 |
| 170 | 10.10 | 6.649 | 706.8 | 2.491 |
| 180 | 10.88 | 7.248 | 811.7 | 2.738 |
| 190 | 11.65 | 7.857 | 924.4 | 2.992 |
| 200 | 12.43 | 8.474 | 1045 | 3.250 |
| 210 | 13.20 | 9.099 | 1172 | 3.514 |
| 220 | 13.97 | 9.731 | 1308 | 3.782 |
| 230 | 14.74 | 10.37 | 1452 | 4.054 |
| 240 | 15.50 | 11.01 | 1603 | 4.331 |
| 250 | 16.26 | 11.66 | 1762 | 4.611 |
| 260 | 17.02 | 12.32 | 1929 | 4.894 |
| 270 | 17.77 | 12.97 | 2103 | 5.181 |
| 273.15 | 18.02 | 13.18 | 2159 | 5.272 |
| 280 | 18.51 | 13.63 | 2284 | 5.471 |
| 290 | 19.25 | 14.29 | 2473 | 5.764 |
| 298.15 | 19.85 | 14.84 | 2632 | 6.012 |
| 300 | 19.99 | 14.96 | 2669 | 6.063 |

of the particles is not the reason for the appearance of the convex anomaly of hBN heat capacity. Differences in the heat capacities of turbostratic and high-ordered hBN (1), and hBN [7] and high-ordered hBN (2) in the range of 15–300 K are given in Fig. 7. We believe that the similarity of these curves testifies to the partial disordering of the sample studied in ref. 7.

TABLE 6

Thermodynamic properties of turbostratic BN

| T (K) | C_p^\ominus (J K ⁻¹ mol ⁻¹) | S^\ominus (J K ⁻¹ mol ⁻¹) | $H^\ominus(T) - H^\ominus(0)$ (J mol ⁻¹) | $-[G^\ominus(T) - H^\ominus(0)]/T$ (J K ⁻¹ mol ⁻¹) |
|------------|---|---|---|--|
| 12 | 0.211 | 0.0703 | 0.632 | 0.0176 |
| 15 | 0.340 | 0.1306 | 1.458 | 0.0334 |
| 20 | 0.597 | 0.2619 | 3.799 | 0.0719 |
| 25 | 0.877 | 0.4242 | 7.484 | 0.1248 |
| 30 | 1.152 | 0.6079 | 12.56 | 0.1893 |
| 35 | 1.405 | 0.8042 | 18.95 | 0.2683 |
| 40 | 1.648 | 1.008 | 26.58 | 0.3431 |
| 45 | 1.962 | 1.220 | 35.61 | 0.4283 |
| 50 | 2.308 | 1.444 | 46.28 | 0.5184 |
| 60 | 3.022 | 1.927 | 72.93 | 0.7111 |
| 70 | 3.751 | 2.446 | 106.5 | 0.9207 |
| 80 | 4.481 | 2.994 | 148.0 | 1.145 |
| 90 | 5.203 | 3.564 | 196.4 | 1.382 |
| 100 | 5.918 | 4.148 | 252.0 | 1.629 |
| 110 | 6.628 | 4.746 | 314.7 | 1.885 |
| 120 | 7.344 | 5.353 | 384.6 | 2.148 |
| 130 | 8.080 | 5.970 | 461.7 | 2.418 |
| 140 | 8.853 | 6.597 | 546.4 | 2.694 |
| 150 | 9.684 | 7.236 | 639.0 | 2.975 |
| 160 | 10.49 | 7.886 | 739.9 | 3.262 |
| 170 | 11.30 | 8.546 | 848.9 | 3.553 |
| 180 | 12.11 | 9.215 | 965.9 | 3.849 |
| 190 | 12.93 | 9.892 | 1091 | 4.149 |
| 200 | 13.75 | 10.58 | 1224 | 4.453 |
| 210 | 14.57 | 11.27 | 1366 | 4.761 |
| 220 | 15.38 | 11.96 | 1515 | 5.073 |
| 230 | 16.20 | 12.66 | 1674 | 5.387 |
| 240 | 17.01 | 13.37 | 1839 | 5.705 |
| 250 | 17.81 | 14.08 | 2013 | 6.026 |
| 260 | 18.61 | 14.80 | 2196 | 6.349 |
| 270 | 19.40 | 15.51 | 2386 | 6.676 |
| 273.15 | 19.66 | 15.74 | 2447 | 6.779 |
| 280 | 20.19 | 16.24 | 2584 | 7.004 |
| 290 | 20.96 | 16.95 | 2789 | 7.335 |
| 298.15 | 21.58 | 17.54 | 2963 | 7.606 |
| 300 | 21.72 | 17.68 | 3003 | 7.667 |

THERMODYNAMIC PROPERTIES OF BORON NITRIDE POLYMORPHOUS MODIFICATIONS

Based on the expert evaluation of the low-temperature heat capacities of cubic, wurtzitic and rhombohedral BN modifications, indicated in this paper, as well as on the heat capacity values of the ordered-graphite-like

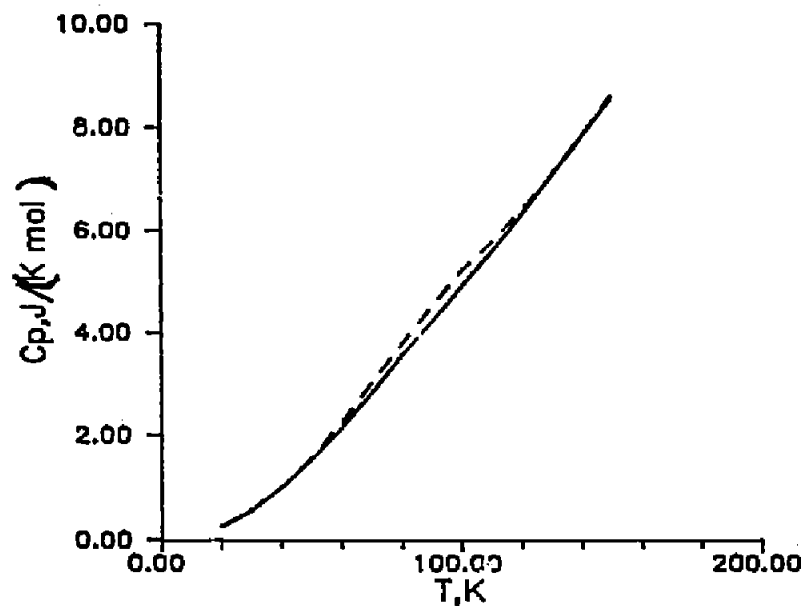


Fig. 5. C_p versus (T) curves for hBN: —, present work; ---, ref. 7.

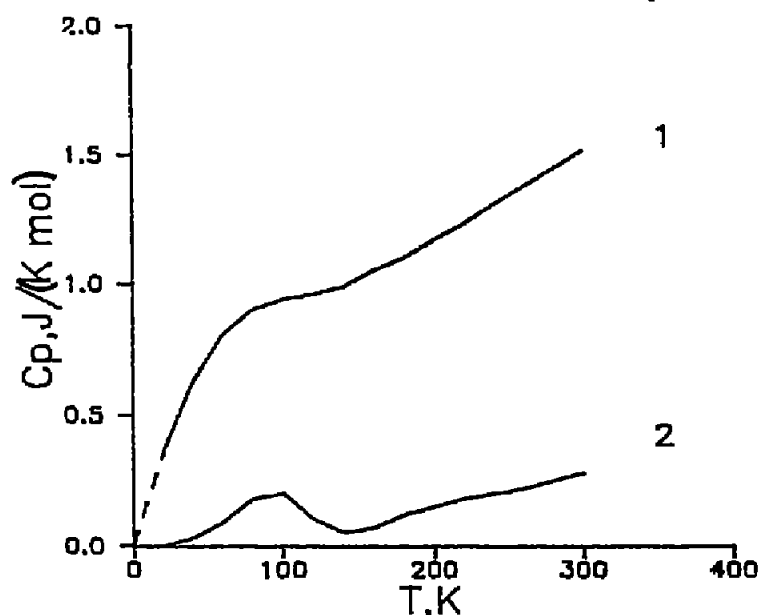


Fig. 7. Differences in the heat capacities of (curve 1) turbostratic and high-ordered hBN; (curve 2) hBN [7] and high-ordered hBN.

TABLE 7

Thermodynamic properties of boron nitride polymorphous modifications

| | hBN ^a | rBN [15] | cBN [10-12] | wBN [7] |
|---|---------------------|---------------------|---------------------|---------------------|
| $C_p^\ominus(298.15 \text{ K})$ ($\text{J K}^{-1} \text{ mol}^{-1}$) | 19.85 ± 0.06 | 20.63 ± 0.06 | 15.95 ± 0.10 | 16.45 ± 0.05 |
| $S^\ominus(298.15 \text{ K})$ ($\text{J K}^{-1} \text{ mol}^{-1}$) | 14.90 ± 0.10 | 15.83 ± 0.10 | 6.710 ± 0.06 | 7.340 ± 0.05 |
| $H^\ominus(298.15 \text{ K}) - H^\ominus(0)$ (J mol^{-1}) | 2632 ± 25 | 2805 ± 25 | 1457 ± 15 | 1541 ± 15 |

^a This work.

TABLE 8

Formation enthalpies for hBN, cBN and wBN

| | hBN [23] | cBN [24] | wBN [25] |
|--------------------------------------|------------------|------------------|------------------|
| $\Delta H^\ominus(298.15 \text{ K})$ | -250.6 ± 2.1 | -266.8 ± 2.2 | -263.6 ± 2.3 |

hexagonal modification, we propose the following data as reference thermodynamic values for the four boron nitride modifications, see Table 7.

The standard formation enthalpies of three boron nitride modifications are given in Table 8. The corresponding data for the rhombohedral modification of boron nitride will be published in the near future. And thus the determination of the key standard thermodynamic values for boron nitride will be complete.

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