

NON-ISOTHERMAL KINETICS WITH NON-LINEAR TEMPERATURE PROGRAMME. V. NON-LINEAR HEATING PROGRAMMES WITH EXACT SOLUTIONS FOR THE TEMPERATURE INTEGRAL

E. SEGAL

Polytechnical Institute of Bucharest, Department of Physical Chemistry and Electrochemical Technology, Bucharest (Romania)

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ABSTRACT

This paper deals with some general non-linear heating programmes which allow the temperature integral to be solved exactly. Some possible uses of these programmes are discussed.

One of the most important problems connected with non-isothermal kinetics involves obtaining exact solutions for the temperature integral. Indeed, for the heating of a sample at a constant rate a ($a = dT/dt$), the integral kinetic equation has the form

$$F(\alpha) = \int_0^\alpha \frac{d\alpha}{f(\alpha)} = \frac{A}{a} \int_{T_0}^T e^{-E/RT} dT \quad (1)$$

or, neglecting as usual the lower limit, T_0

$$F(\alpha) = \frac{A}{a} \int_0^T e^{-E/RT} dT \quad (2)$$

The right-hand side of eqns. (1) and (2) is called the temperature integral. For the given linear programme of heating this integral cannot be solved exactly [1].

In order to obtain exact solutions for the temperature integral, hyperbolic and parabolic programmes of heating were used. These programmes are given by

$$dT/dt = a_2 T^2 \quad (3)$$

$$1/T = b - a_2 t \quad (4)$$

for the hyperbolic programme [2], and

$$\frac{dT}{dt} = \frac{1}{b_1} \frac{E}{2RT + E} \quad (5)$$

$$b_1 \frac{RT^2}{E} + b_1 T + c = t \quad (6)$$

for the parabolic programme [3] (a_2 , b and c are constants for a given programme). These equations may be used as starting points for working some more general non-linear heating programmes. Thus, if instead of the value 2 for the temperature exponent, one takes positive integers p , so as $p \geq 2$, eqn. (3) becomes

$$\frac{dT}{dt} = a_p T^p \quad (7)$$

which defines a family of non-linear heating programmes for different values of p .

Equation (2), with the heating rate given by eqn. (7), takes the form

$$F(\alpha) = \frac{A}{a_p} \int_0^T \frac{e^{-E/RT}}{T^p} dT \quad (8)$$

where the temperature integral can now be solved exactly. In order to obtain the solution, one has to operate the usual change of variable

$$x = -E/RT \quad (9)$$

In terms of the variable x , eqn. (8) can be written as

$$F(\alpha) = (-1)^p \frac{A}{a_p} \left(\frac{R}{E} \right)^{p-1} \int_{-\infty}^x x^{p-2} e^x dx \quad (10)$$

The integral

$$I = \int x^{p-2} e^x dx$$

can be calculated exactly through repeated integration by parts, with the result [4]

$$I = e^x \left[x^{(p-2)} - (p-2)x^{p-3} + (p-2)(p-3)x^{p-4} \dots \right. \\ \left. + (-1)^{(p-3)}(p-2)!x + (-1)^{(p-2)}(p-2)! \right] \quad (11)$$

Returning to the old variable by means of eqn. (9), eqn. (10) becomes

$$F(\alpha) = (-1)^p \frac{A}{a_p} \left(\frac{R}{E} \right)^{(p-1)} e^{-E/RT} \left[\left(\frac{E}{RT} \right)^{(p-2)} - (p-2) \left(\frac{E}{RT} \right)^{(p-3)} \right. \\ \left. + (p-2)(p-3) \left(\frac{E}{RT} \right)^{(p-4)} \dots + (-1)^{(p-3)}(p-2)! \frac{E}{RT} + (-1)^{(p-2)}(p-2)! \right] \quad (12)$$

Obviously, for $p = 2$, one gets

$$F(\alpha) = \frac{A}{a_2} \frac{R}{E} e^{-E/RT} \quad (13)$$

which is the particular form of eqn. (12) for the hyperbolic heating programme [2]. Equation (12) could be used in working integral methods for non-isothermal kinetics, like the method given by Coats and Redfern [5].

Another family of non-linear heating programmes which allows the temperature integral to be solved exactly may be obtained from the general form of eqn. (2)

$$F(\alpha) = A \int_0^T f(T) e^{-E/RT} dT \quad (14)$$

where

$$f(T) = \frac{1}{dT/dt} \quad (15)$$

The problem involves the proper choice of a function $f_1(T)$ which has to fulfil the condition [3]

$$\int f(T) e^{-E/RT} dT = f_1(T) e^{-E/RT} \quad (16)$$

where the right-hand side is implied by the form of the left-hand integrand. Operating the derivative of both members in eqn. (16), one gets:

$$f(T) = f_1'(T) + \frac{E}{RT^2} f_1(T) \quad (17)$$

a condition which $f_1(T)$ has equally to fulfil. If f_1 is chosen as

$$f_1 = \frac{R}{a_2 E} \quad (18)$$

from relationship (17), one gets eqn. (3) for the hyperbolic programme of heating. For $f_1(T)$ given by

$$f_1(T) = \frac{RT}{bE} \quad (19)$$

eqn. (17) leads to

$$f(T) = \frac{R}{bE} + \frac{1}{bT} \quad (20)$$

or, taking into account eqn. (15) and integrating eqn. (20), it turns out that

$$t = b_1 \frac{RT}{E} + b_1 \ln T + c \quad (21)$$

where $b_1 = 1/b$.

Equations (20) and (21) define a new heating programme which allows an exact solution of the temperature integral. Taking into account eqns. (19) and (16), eqn. (14) takes the form

$$F(\alpha) = \frac{AR}{bE} T e^{-E/RT} \quad (22)$$

whose logarithmic form could be used to determine the activation energy and the pre-exponential factor. These two kinetic parameters are easy to evaluate from the slope and the intercept of the straight line ($\log F(\alpha)/T, 1/T$).

Equation (17) with

$$f_1(T) = \frac{RT^2}{bE} \quad (23)$$

leads to [3]

$$f(T) = \frac{2RT}{bE} + \frac{1}{b} \quad (24)$$

and

$$\frac{RT^2}{E} b_1 + b_1 T + c = t \quad (25)$$

where $b_1 = 1/b$, i.e. to a parabolic heating programme. For this programme, the integral kinetic equation (14) takes the particular form

$$F(\alpha) = \frac{AR}{bE} T^2 e^{-E/RT} \quad (26)$$

This allows evaluation of E and A from the slope and the intercept of the straight line ($\log F(\alpha)/T^2, 1/T$).

Equation (17) with the general function $f_1(T)$ given by

$$f_1(T) = \frac{RT^m}{bE} \quad (27)$$

where m is a positive integer ($m \geq 2$) becomes

$$f(T) = \frac{mRT^{m-1}}{bE} + \frac{T^{m-2}}{b} \quad (28)$$

By integration of eqn. (28), one gets

$$t = b_1 \frac{RT^m}{E} + b_1 \frac{T^{m-1}}{m-1} + c \quad (29)$$

where $b_1 = 1/b$. Equations (28) and (29) define another family of non-linear programmes which lead to exact solutions of the temperature integral. In this case, from eqns. (14), (16) and (17), it turns out that

$$F(\alpha) = A \frac{RT^m}{bE} e^{-E/RT} \quad (30)$$

the logarithmic form of which could be used to evaluate E and A in the usual way.

One more family of non-linear programmes which allow an exact solution of the temperature integral can be obtained for a function $f_1(T)$ given by:

$$f_1(T) = T^q e^{kT} \quad (31)$$

where q is a positive integer and k is a constant. Taking into account eqn. (31), eqn. (17) becomes

$$f(T) = qT^{q-1} e^{kT} + kT^q e^{kT} + \frac{E}{R} T^{q-2} e^{kT} \quad (32)$$

By integration with respect to temperature, the first two terms give $T^q e^{kT}$. As for the integral of the last term, this is $(E/R) I$, where I is given by eqn. (11). Thus, the integration of eqn. (32) leads to

$$t = e^{kT} \left[T^q + \frac{E}{R} \left(\frac{T^{q-2}}{k} - \frac{(q-2)T^{q-3}}{k^2} + \frac{(q-2)(q-3)}{k^3} T^{q-4} \dots \right. \right. \\ \left. \left. + (-1)^{q-3} \frac{(q-2)!}{k^{q-2}} + \frac{(-1)^{q-2} (q-2)!}{k^{q-1}} \right) \right] \quad (33)$$

with $q \geq 2$.

For $f_1(T)$ given by eqn. (31), eqn. (14) takes the form

$$F(\alpha) = AT^q e^{kT} e^{-E/RT} \quad (34)$$

Although the three families of non-linear heating programmes mentioned above are mainly suitable for the integral methods in non-isothermal kinetics, they can also be used to work out differential methods. Thus, for the heating rate (7), through generalization of the formulae given for $p = 2^{1,2}$, one gets

$$\frac{d^2\alpha}{d\alpha dT} = -\frac{p}{T} + \frac{E}{RT_{\max}^2} + \frac{f'(\alpha)}{f(\alpha)} \frac{d\alpha}{dT} \quad (35)$$

$$\left(\frac{d}{dT}\right)_{\max} = \frac{pRT_{\max} - E}{RT_{\max}^2} \frac{f(\alpha_{\max})}{f'(\alpha_{\max})} \quad (36)$$

$$\frac{d \ln a_p T_{\max}^{p-2}}{d\left(\frac{1}{T_{\max}}\right)} = -\frac{E}{R} \quad (37)$$

All these equations lead to variants of the classical differential methods for evaluating the kinetic parameters under non-isothermal conditions [1,6].

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