

Contents lists available at ScienceDirect

[Thermochimic](dx.doi.org/10.1016/j.tca.2011.02.004)a Acta

journal homepage: www.elsevier.com/locate/tca

Thermodynamic and acoustic [properties](http://www.elsevier.com/locate/tca) [of](http://www.elsevier.com/locate/tca) [binary](http://www.elsevier.com/locate/tca) [m](http://www.elsevier.com/locate/tca)ixtures of ethers. IV. Diisopropyl ether or oxolane with N,N-dimethylaniline or N,N-diethylaniline at 303.15, 313.15 and 323.15 K

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article info

Article history: Received 28 October 2010 Received in revised form 29 January 2011 Accepted 2 February 2011 Available online 19 February 2011

Keywords: Density Speed of sound Ether Binary mixture Dialkyl aniline Excess property

ABSTRACT

Densities ρ and speeds of sound u, of binary mixtures formed by N,N-dimethylaniline (DMA) or N,Ndiethylaniline (DEA) with diisopropyl ether (DIPE) or oxolane have been measured over the entire range of composition at a temperature of (303.15, 313.15 and 323.15 K) and atmospheric pressure. The ρ and u values were used to calculate isentropic compressibilities κ_S , Rao's molar sound functions R, intermolecular free lengths L_f , specific acoustic impedances Z, excess molar volumes V_{m}^{E} , excess isentropic compressibilities $\kappa^{\rm E}_{\rm S}$, excess intermolecular free lengths $L^{\rm E}_{\rm f}$ and excess specific acoustic impedances $Z^{\rm E}$. The results have been used to investigate molecular interactions and structural effects in these mixtures. The speeds of sound in present mixtures have been estimated using several empirical and theoretical models to determine their relative predicting ability in terms of pure component properties.

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1. Introduction

The thermodynamic, acoustic and transport properties of nonelectrolyte liquid–liquid mixtures provide information about type and extent of molecular interactions, and can be used for the development of molecular models for describing the behav[iour of](#page-10-0) solutions [1–5].

We are engaged in systematic investigations of binary mixtures of aromatic amine with ether at different temperatures. In the previous papers [6–8] we have reported speeds of sound, isentropic compressibilities, excess molar volumes and various other [relate](#page-10-0)d properties of binary mixtures of diisopropyl ether (D[IPE\)](#page-10-0) [or](#page-10-0) oxolane with aniline, N-methylaniline, N-ethylaniline, o-toluidine and m-toluidine (primary and secondary arylamines) at 303.15, 313.15 a[nd](#page-10-0) [323](#page-10-0).15 K. It was observed that there exist specific interactions through hydrogen bonding and dipole–dipole interactions between primary/secondary arylamine and ether molecules. It would be advantageous to investigate and interpret the volumetric and acoustic results when a tertiary amine is mixed with ether molecules, because in contradiction to primary and secondary arylamines, the tertiary ones cannot form hydrogen bonds with ethers [9–13]. Therefore, in this paper, we extend our studies to the binary mixtures formed from DIPE or oxolane with N,N-dimethylaniline (DMA) or N,N-diethylaniline (DEA). Speeds of sound at different temperatures in the investigated binary mixtures have also been estimated through Nomoto model (NM) [14], Van Dael model (VM) [15], Ernst et al. model (EM) [16], impedance model (IM) [17], Schaaffs' collision factor theory (CFT) [18], Jacobson's free length theory (FLT) [19], and Prigogine–Flory–Patterson–Oswal theory (PFPOT) [20–23].

2. Experimental

[All](#page-10-0) chemicals used in this study were of analytical grade obtained from s.d. fine-chem., Ltd. The claimed mass fraction purity for the chemicals was >0.995. These liquids were dried over 4\AA molecular sieves and partially degassed prior to use. The purity of these experimental liquids was checked by comparing the observed densities and speeds of sound with those reported in the literature. The measured values are presented in Table 1 along with the available literature values.

The densities of pure liquids and the binary mixtures were measured by calibrated single stem capillary pycnometer, having a bulb volume of about 10 cm^3 and capillary of internal diameter of 0.7 rnm. The speeds of s[ound](#page-1-0) [in](#page-1-0) [p](#page-1-0)ure liquids and in their

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Table 1 Densities (ρ), speeds of sound (u), thermal expansion coefficients (α), and heat capacities (C_p^o) for pure liquids.

Liquid	T(K)	ρ (kg m ⁻³)		u (m s ⁻¹)		α^{o} a (kK ⁻¹)	$C_{\rm p}^{\rm o}$ (J mol ⁻¹ K ⁻¹)
		Expt.	Lit.	Expt.	Lit.		
DIPE	303.15	712.9	712.96 ^b 713.06 ^d 713.1 ^f 713.20g	976	976.04 ^b 974.9^e	1.468	218 ^c
	313.15	702.9	702.80 ^f 702.59g	932	932.17 ^b	1.489	$222^{h,i}$
	323.15	691.9	691.533 ^b 691.86 ^g	889	888.61 ^b	1.512	$226^{h,i}$
Oxolane	303.15	876.9	877.0^{j} 876.7 ¹	1254	1254 1255 ¹ 1256 ^m	1.248	125 ^k
	313.15	865.7	865.4^{k} 870.5°	1207	1209 ⁿ 1206.4^{k}	1.265	128 ^k
	323.15	854.8		1165		1.291	130 ^k
DMA	303.15	948.4	948.0P 948.33 ^q	1468		0.861	231 ⁱ
	313.15	939.7	939.8 ^p 939.72 ^q	1431		0.868	235 ⁱ
	323.15	931.5		1393		0.876	239 ⁱ
DEA	303.15 313.15 323.15	926.1 917.5 964.6	925.9 ^p 918.0P	1434 1396 1358		0.862 0.869 0.877	291 ⁱ 296^{i} 301 ⁱ

^a Derived from density data Ref. [24–30].

^f Ref. [26].

^h Ref. [28].

- ⁱ From group contribution method [33].
- ^j Ref. [34].
- ^k Ref. [28]. ^l Ref. [35].

^m Ref. [36].

 n Ref. [37].

^o Ref. [38].

- ^p Ref. [30].
- ^q Ref. [39].

[b](#page-10-0)inary mixtures were measured using single-crystal variable-path [u](#page-10-0)ltrasonic interferometer (Mittal Enterprises, New Delhi, Model: [M](#page-10-0)-82) operating at 2 MHz. For all the measurements, temperatures were controlled by circulating water through an ultra thermostat JULABO F-25 (made in Germany) keeping temperature fluctuations within \pm 0.02 K. The details of the experimental procedure have been described elsewhere [40]. The binary mixtures were prepared by mixing known masses of pure liquids in air tight stoppered bottles taking due precautions to minimize evaporation losses. All the mass measurements were performed on an electronic balance (Mettler-AE 240, Switzerland) accurate to \pm 0.05 mg. The samples for the measu[remen](#page-10-0)t were immediately used after preparation. The uncertainty in density and speed of sound measurements was within 0.1 kg m^{-3} and 1 m s⁻¹.

3. Results and discussion

The results for the densities, speeds of sound, isentropic compressibilities, Rao's molar sound functions [41], specific acoustic impedances and intermolecular free lengths, excess molar volumes, and excess isentropic compressibilities for four binary mixtures of DIPE or oxolane with DMA or DEA at 303.15, 313.15, and 323.15 K are given in Tables 2–5.

The isentropic compressibilities κ_S [,](#page-10-0) [R](#page-10-0)ao's molar sound functions R, specific acoustic impedances Z and intermolecular free lengths

 $L_{\rm f}$ were obtained from the experimental densities ρ and speeds of sound u using the following relations:

$$
\kappa_{\rm s} = \left(\rho u^2\right)^{-1} \tag{1}
$$

$$
R = u^{1/3}V\tag{2}
$$

$$
Z = u\rho \tag{3}
$$

$$
L_{\rm f} = K_{\rm Jac}/(\mu \rho^{1/2}) = K_{\rm Jac} \kappa_{\rm S}^{1/2} \tag{4}
$$

where molar volume $V = \sum (x_i M_i)/\rho$, in which x_i and M_i are the mole fraction and molecular mass of component i. K_{lac} = (91.368 + 0.3565T) × 10⁻⁸ is temperature dependent Jacobson's constant [19]. The values of κ_S , R, and Z were uncertain within \pm 0.2% while L_f within 1%.

Excess molar volume $V_{\mathrm{m}}^{\mathrm{E}}$, excess isentropic compressibility $\kappa_{\mathrm{S}}^{\mathrm{E}}$, excess intermolecular free length L_f^E and excess specific acoustic [imped](#page-10-0)ance Z^E in each mixture were calculated with following expressions:

$$
Y^{\rm E} = Y - Y^{\rm id} \tag{5}
$$

where, Y represent either V, $\kappa_{\mathsf{S}}, L_{\mathsf{f}},$ and Z. The $V^{\text{id}}_{\mathsf{m}}$ for an ideal mixture was calculated from usual relation

$$
V_{\rm m}^{\rm id} = \sum x_i V_i^{\rm o} = \sum \frac{x_i M_i}{\rho_i} \tag{6}
$$

^b Ref. [24].

^c Ref. [31].

 d Ref. [25].

^e Ref. [32].

^g Ref. [27].

^a Unit = m^{10/3} s^{-1/3} mol⁻¹.

 $\kappa^{\rm id}_{\rm S}$ for an ideal mixture were calculated from the relation recommended by Benson and Kiyohara [42]

The $V^{\rm E}_{\rm m}$, $\kappa^{\rm E}_{\rm S}$, $L^{\rm E}_{\rm f}$ and $Z^{\rm E}$ were correlated by Redlich–Kister polynomial equation

$$
\kappa_{\rm S}^{\rm id} = \sum \phi_{\rm i} \left\{ \frac{\kappa_{\rm S,i}^{\rm o} + \text{TV}_{\rm i}^{\rm o}(\alpha_{\rm i}^{\rm o})^2}{C_{\rm p,i}^{\rm o}} \right\} - T \left(\sum x_{\rm i} V_{\rm i}^{\rm o} \right) \frac{\sum \phi_{\rm i} \alpha_{\rm i}^{\rm o \, 2}}{\sum x_{\rm i} C_{\rm p,i}^{\rm o}} \tag{7}
$$

 $L_{\rm f}^{\rm id}$ and $Z^{\rm id}$ for an ideal mixture were obtained as

$$
L_{\rm f}^{\rm id} = K_{\rm Jac} (\kappa_{\rm S}^{\rm id})^{1/2} \tag{8}
$$

$$
Zid = \rhoid uid = \left(\frac{\rhoid}{\kappa_Sid}\right)^{1/2}
$$
 (9)

in which the M_i , $V_i^{\rm o}$, $\alpha_i^{\rm o}$ and $C_{\rm p,i}^{\rm o}$ are respectively, the molar mass, molar volume, isobaric thermal expansion coefficient and molar isobaric heat capacity of pure component i, $\phi_i = x_iV_i/\sum x_iV_i$ is the volume fraction of i in the mixture, stated in terms of the unmixed components and $\rho^{\rm id} = \sum \phi_{\rm i} \rho_{\rm i}^{\rm o}$. The values of $\alpha^{\rm o}$ and $C_{\rm p}^{\rm o}$ used for these calculations are included in Table 1.

For compact and smooth representation, the values of u and κ_S were fitted to a polynomial equation of the form

$$
u \text{ or } \kappa_{\mathsf{S}} \text{ or } R = \sum_{i=0}^{m} A_i X_1^i \tag{10}
$$

$$
Y^{E} = x_{1}(1 - x_{1}) \sum_{i=0}^{n} B_{i}(1 - 2x_{1})^{i}
$$
 (11)

where $Y^{\rm E} = V_{\rm m}^{\rm E}$, $\kappa_{\rm S}^{\rm E}$, $L_{\rm f}^{\rm E}$, and $Z^{\rm E}$.

The coefficients A_i of Eq. (10) and B_i of Eq. (11) and corresponding standard deviations σ obtained from a least-squares fit with equal weights assigned to each point are given in Tables 6 and 7.

It is observed from Tables 2–5, that the values of speeds of sound, specific acoustic impedances decrease with mole fraction of DIPE or oxolane and also with the rise in temperature for all the mixtures studied. The values of isentropic compressibility and intermolecular free length increase with th[e](#page-5-0) [mole](#page-5-0) [fraction](#page-5-0) [o](#page-5-0)f DIPE or oxolane and with the rise in temperature from 303.15 to 323.15 K. The Figs. 1–4 show the dependence of $V_{\rm m}^{\rm E}$, $\kappa_{\rm S}^{\rm E}$, $L_{\rm f}^{\rm E}$ and $Z^{\rm E}$ on composition. The $V_{\rm m}^{\rm E}$, $\kappa_{\rm S}^{\rm E}$ and $L_{\rm f}^{\rm E}$ are negative for all the four studied mixtures over the entire mole fraction range, and their magnitude increases with the rise in temperature from 303.15 to 323.15 K. Further, the values of $V_{\rm m}^{\rm E}$ for both the DIPE mixtures are much more negative compared to both oxolane mixtures. Other very large negative values of $V_{\mathrm{m}}^{\mathrm{E}}$ reported in the literature include 2methoxyethanol + tert-butylamine (-2.4 to -2.6×10^{-6} m³ mol⁻¹ [43]) and DMSO + cyclohexylamine (-3.4×10^{-6} m³ mol⁻¹) [44].

The observed values of $V_{\mathrm{m}}^{\mathrm{E}}$ for presently investigated DIPE + DMA and +DEA mixtures are less negative than DIPE + aniline [7], +toluidines [8]. Similarly, $V_{\mathrm{m}}^{\mathrm{E}}$ for presently investigated oxolane + DMA Densities (ρ), speeds of sound (u), isentropic compressibilities (κ_S), Rao's molar sound functions (R), specific acoustic impedances (Z) and intermolecular free lengths (L_f), excess molar volumes (V $_{\rm m}^{\rm E}$) and excess isentropic compressibilities ($\kappa_{\rm S}^{\rm E}$) for DIPE+DEA mixture at 303.15, 313.15 and 323.15 K.

^a Unit = $m^{10/3}$ s^{-1/3} mol⁻¹.

Table 3

and +DEA mixtures are less negative than oxolane + aniline [6] +monoalkylaniline [6], +toluidines [8]. This clearly indicates influence of tertiary amine over primary and secondary amines.

The over all behaviour of $V_{\rm m}^{\rm E}$ and $\kappa_{\rm S}^{\rm E}$ for the present mixtures can be envisaged as a resultant of opposite effects: (i) disruption of ether–ether interactions; (ii) break [down](#page-10-0) of dipolar ord[er](#page-10-0) [in](#page-10-0) ether an[d](#page-10-0) [ter](#page-10-0)tiary arylamine; (iii) dipole–dipole interactions in tertiary aryl amine and ether; (iv) free volume effect; and (v) characteristic pressure P^* effect [45,46]. The free volume ($V_{\rm fs}^{\rm E}$), characteristic pressure ($V_{\rm pk}^{\rm E}$) and interaction
($V_{\rm int}^{\rm E}$) contributions to the total $V_{\rm total}^{\rm E}$ can be estimated from the
Prigogine–Flory–Patterson (PFP) theory [26–28]. Table marizes these contributions at equimol[ar compo](#page-10-0)sition for the present mixtures. The estimated sum of free volume and P^* contributions $(V_{fv}^E + V_{P*}^E)$ to the V_{total}^E as per the PFP theory is in the range fro[m](#page-7-0) -1.821 to -2.291×10^{-6} -2.291×10^{-6} -2.291×10^{-6} m³ mol⁻¹ DIPE + DMA or DEA and -0.152 to -0.296×10^{-6} m³ mol⁻¹ oxolane + DMA or DEA, while interaction contribution is in the range from −0.114 to -0.968×10^{-6} m³ mol⁻¹ DIPE + DMA or DEA and -0.213 to -0.536×10^{-6} m³ mol⁻¹ oxolane + DMA or DEA. This points to the fact that in very large negative values of $V_{\mathrm{m}}^{\mathrm{E}}$ for DIPE+DMA or DEA the major contribution is because of the free volume and P^* effects, which include size, shape and conformation of the molecules involved. In case of oxolane mixtures, the combined contribution due to the free volume and P^* effects is much smaller as compared to the DIPE mixtures. This difference can be explained by considering the difference in thermal expansion coefficient α or Flory's reduced volume \tilde{V} and characteristic pressure P^* of unlike components (listed in Tables 1 and 9). The free volume (V_{fv}^E) term is proportional to $-(\tilde{V}_1 - \tilde{V}_2)^2$ i.e. $-(\alpha_1 - \alpha_2)^2$ [22,45] and is always negative and its magnitude depends upon difference in thermal expansion coefficients or Flory's reduced volumes of involv[ed components.](#page-1-0) The difference between α or \tilde{V} of DIPE and DMA/DEA is relatively very large as compared to the α or \tilde{V} of oxolane and DMA/DEA, t[hat](#page-10-0) [is](#page-10-0) [the](#page-10-0) reason free volume (V_{fv}^E) term for DIPE+DMA/MEA is very large negative compared to the oxolane + DMA/DEA mixtures. The characteristic pressure (V_{P*}^E) term is proportional to $(P_1^*-P_2^*)(\tilde{V}_1-\tilde{V}_2)$ and can have either sign depending upon the magnitude of P^* and α (or \tilde{V}) of unlike components [22,45]. For DIPE + DMA/MEA the $(P_1^* - P_2^*)$ is large negative while $(\tilde{V}_1 (\tilde{V}_2)$ large positive, making (V_{P*}^E) term large and negative. For oxolane+DMA, $(P_1^*-P_2^*)$ is small negative and $(\tilde{V}_1-\tilde{V}_2)$ positive, with product $(P_1^* - P_2^*)(\tilde{V}_1 - \tilde{V}_2)$ i.e. $(V_{P_*}^E)$ term [for oxola](#page-10-0)ne + DMA is small negative. Contrary to this for oxolane+DEA, $(P_1^* - P_2^*)$ is small positive and $(\tilde{V}_1 - \tilde{V}_2)$ positive, as a result V_{P*}^{E} is small and positive.

The negative values of interactional contribution ($V_{\rm int}^{\rm E}$) to $V_{\rm m}^{\rm E}$ for the present mixtures suggest that dipole-dipole interactions (μ for DIPE/oxolane = 1.34/1.75 D, and for DMA/DEA \approx 1.61 D), between unlike molecules dominate over the dispersive ether–ether interactions, breaking of dipolar order of aromatic amines as well as of DIPE/oxolane. The negative $V^{\rm E}_{\rm m}$ and $\kappa^{\rm E}_{\rm S}$ values at equimolar compo-

^a Unit=m^{10/3} s^{-1/3} mol^{-1.}

Table 5

Densities (ρ), speeds of sound (u), isentropic compressibilities (κ_S), Rao's molar sound functions (R), specific acoustic impedances (Z) and intermolecular free lengths (L_f),
excess molar volumes (V $_{\rm m}^{\rm E}$) and e

Table 5 (Continued)

x_1	ρ (kg m ⁻³⁾	u (ms ⁻¹)	κ_S (TPa ⁻¹)	10^6 R^a	10^3Z (kg m ⁻⁻² s ⁻¹)	$10L_f$ nm	10^6 V_m^E (m ³ mol ⁻¹)	κ_c^{E} (TPa ⁻¹)
0.2899	0.9028	1330	626	1546	1201	0.515	-0.474	-26
0.3879	0.8995	1315	643	1453	1183	0.522	-0.543	-31
0.4892	0.8953	1296	665	1358	1160	0.531	-0.565	-34
0.5937	0.8904	1275	691	1260	1135	0.541	-0.594	-36
0.6904	0.8849	1252	721	1169	1108	0.553	-0.586	-35
0.7948	0.8778	1226	758	1072	1076	0.567	-0.552	-31
0.8935	0.8687	1200	799	983	1042	0.582	-0.391	-24
	0.8548	1165	862	888	996	0.604		

^a Unit = m^{10/3} s^{-1/3} mol⁻¹.

sitions for common amine follow the order: DIPE > oxolane, while there is a marginal difference between the values of DMA or DEA.

As expected, the trend of negative L_f^E values (Fig. 3 and Table 7) is similar to κ_S^E in the studied mixtures. The negative values of L_f^E
are generally observed in the mixtures with specific interactions between unlike molecules. The negative $L_{\rm f}^{\rm E}$ values at equ[imolar](#page-9-0) compositions for common amine follow the order: DIPE > oxolane. Fig. 4 shows that Z^E is positive for all the studied mixtures. Positive trends in Z^E have also been reported for binary mixtures of methanol + piperidine [47]. Specific acoustic impedance is a quantity, which depends on the molecular packing of the systems. The positive values of Z^E are in accordance with the presence of struc-

 a Unit = $(m^{10/3} s^{-1/3} mol^{-1})$.

tural effect and specific interactions between the unlike molecules [6–8].

The dependence of Rao's molar sound functions R on mole fractions x_1 is shown in Fig. 5. It is observed that there are slight negative deviations from the linear dependence on x_1 for DIPE binary mixtures while almost linear for oxolane mixtures. As expected the R are independent of temperature (Tables 2–5). According to Rao's approach [41], the speed of sound is directly related to the strength of inter[molecu](#page-9-0)lar energy. If the intermolecular energy were simply a consequence of additive non-bonding interaction, the Rao's molar sound function [should predic](#page-2-0)t additive behaviour. The fact that is not so implies some additional interactions in the mixtures compared with those in the simple liquids [48].

4. Estimation of speeds of sound

Recently Gayol et al.[49] and [Kham](#page-10-0)mer and Shaw [50] described and tested several predictive methods for speed of sound in alkanol + n-alkanes mixtures and discussed the combing rules used in different models. Glinski [51] discussed the additivity of sound velocity in twenty four randomly selected binary mixtures and found Nom[oto](#page-10-0) [mo](#page-10-0)del [14] based on Rao's [hypo](#page-10-0)thesis [41] provides

Fig. 1. Dependence of excess molar volumes on mole fraction of DIPE or oxolane, (a) DIPE + DMA, (b) DIPE + DEA, (c) oxolane + DMA and (d) oxolane + DEA at T = 303.15 K (●), 313.15 K (\blacksquare), and 323.15 K (\blacktriangle). Solid lines have been drawn from Eq. (11) using coefficient given in Table 7.

results similar to those of Ernst et al. model [16] while the Van Dael model [15] often fails. For all three models large deviations were observed for mixtures with associating [cons](#page-2-0)tituents. Here, the speeds of sound in the investigated binary mixtures have been estimated at 303.15, 313.15 and 323.15 [K from](#page-10-0) the Nomoto model [14] Van Dael model [15], Ernst et al. model [16], impedance model [17], c[ollisio](#page-10-0)n factor theory [18], intermolecular free length theory [19], and Prigogine–Flory–Patterson–Oswal theory [20–23]. The pertinent relations in these calculations and their theoretical basis have been [outli](#page-10-0)ned several time[s and](#page-10-0) will not be repe[ated h](#page-10-0)ere. PFPO model is the only model amongst the considered here which has interaction parameter X_{12} . Normally X_{12} is obtained from excess [mola](#page-6-0)r enthalpy data [21]. In absence of excess molar enthalpy data needed to determine interaction parameter X_{12} , we have determined speed of sound using $X_{12} = 0$ as well as X_{12} derived from equation [21]

$$
X_{12} = P_1^* \left[1 - \left(\frac{P_2^*}{P_1^*} \right)^{1/2} \left(\frac{V_2^*}{V_1^*} \right)^{1/6} \right]^2 \tag{12}
$$

Table 8

Free volume contribution $V_{\rm fv}^{\rm E}$ P^* contribution $V_{\rm pz}^{\rm E}$ and interaction contribution to $V_{\rm int}^{\rm E}$ to total excess molar volume $V_{\rm total}^{\rm E}$ at equimolar composition as per the PFP theory.

T(K)	10^6 $V_{\text{f}_{\text{V}}}^{\text{E}}$ (m ³ mol ⁻¹)	$10^6 V_{p_{*}}^E$ (m ³ mol ⁻¹)	$10^6 (V_{\text{fv}}^E + V_{P*}^E)$ (m ³ mol ⁻¹)	10^6 V_{int}^E (m ³ mol ⁻¹)	$10^6 V_{total}^E$ (m ³ mol ⁻¹)
DIPE + DMA					
303.15	-0.558	-1.288	-1.846	-0.101	-1.947
313.15	-0.603	-1.417	-2.020	-0.188	-2.208
323.15	-0.681	-1.610	-2.291	-0.327	-2.618
$DIPE + DEA$					
303.15	-0.597	-1.224	-1.821	-0.224	-2.045
313.15	-0.643	-1.344	-1.987	-0.483	-2.470
323.15	-0.692	-1.500	-2.192	-0.970	-3.162
THF + DMA					
303.15	-0.181	-0.042	-0.223	-0.233	-0.456
313.15	-0.201	-0.053	-0.254	-0.383	-0.637
323.15	-0.236	-0.060	-0.296	-0.536	-0.832
THF+DEA					
303.15	-0.199	0.047	-0.152	-0.213	-0.365
313.15	-0.221	0.047	-0.174	-0.283	-0.457
323.15	-0.243	0.050	-0.193	-0.385	-0.578

Table 9 Flory's characteristic and reduced parameters of neat components.

Component	T(K)	Ũ	10^6V^* (m ³ mol ⁻¹)	$10^{-6}P^*$ (J m ⁻³)
DIPE	303.15	1.3407	106.85	419
	313.15	1.3529	107.45	408
	323.15	1.3655	108.15	394
Oxolane	303.15	1.3004	63.21	617
	313.15	1.3122	63.45	612
	323.15	1.3242	63.67	608
DMA	303.15	1.2216	104.60	634
	313.15	1.2293	104.90	632
	323.15	1.2343	105.40	629
DEA	303.15	1.2218	131.88	600
	313.15	1.2295	132.28	596
	323.15	1.2375	132.64	592

where P^* and V^* are Flory's reduction parameters [21]. In order to determine relative predicting ability of various models, the standard percentage deviation σ % was calculated by considering percentage deviations for all mole fractions using the relation (13) and recorded in Table 10.

$$
\sigma \mathscr{E}\left[\sum \left\{\frac{100(u_{\exp}-u_{\text{cal}})}{u_{\exp}}\right\}/n\right]^{1/2} \tag{13}
$$

where *n* represents number of data points.

The study of Table 10 reveals that the values of σ % for the NM, VM, EM, IM, CFT, FLT, and PFPOT are in the range from 0.1 to 2.8, 5.7 to 10.6, 0.3 to 5.2, 0.4 to 2.7, 0.4 to 1.3, 0.6 to 3.3 and 0.2 to 1.7 (0.1–1.2), respectively for the four binary mixtures at investigated three different temperatures. The average values of sta[ndard](#page-10-0) [dev](#page-10-0)iations $\langle \sigma \rangle$ 1.2%, 1.5%, 0.9%, 1.5%, and 0.9% for the NM, IM, CFT, FLT and PFPOT, respectively are comparatively less than those found for other estimations. In EM and VM average σ % are 2.3 and 7.9%. Further it is observed, the estimated results by all the models in oxolane mixtures are better than those for the mixtures involving DIPE. This is attributed to the

Fig. 2. Dependence of excess isentropic compressibilities on mole fraction of DIPE or oxolane, (a) DIPE + DMA, (b) DIPE + DEA, (c) oxolane + DMA and (d) oxolane + DEA at T=303.15 K (\bullet), 313.15 K (\blacksquare), 323.15 K (\blacktriangle). Solid lines have been drawn from Eq. (11) using coefficient given in Table 7.

Fig. 3. Dependence of excess intermolecular free length on mole fraction of DIPE or oxolane at T=303.15 K. (a) DIPE+DMA (●) and oxolane+DMA (▲). (b) DIPE+DEA (■) and oxolane + DEA (\Diamond), Solid lines have been drawn from Eq. (11) using coefficient given in Table 7.

Fig. 4. Dependence of excess specific acoustic impedance on mole fraction of DIPE or oxolane at T=303.15 K. (a) DIPE + DMA (\bullet) and oxolane + DMA (\blacktriangle). (b) DIPE + DEA (\blacksquare) and oxolane + DEA (\Diamond), Solid lines have been drawn from Eq. (11) using coefficient given in Table 7.

fact that in oxolane mixtures, the i[ntera](#page-2-0)ctions are not so [com](#page-6-0)plex as in the case of DIPE mixtures. Considering the average values of standard deviations, the estimation ability of speed of sound in the presently investigated mixtures follows the sequence PFPOT > CFT > NM > FLT ≅ IM > EM > VM. This indicates PFPOT provides best estimates while VM least estimates of speed of sound in present mixtures. The PFPOT is the only model which involves the interaction parameter. It is also observed that speeds of sound in IDPE + amines mixtures estimated using X_{12} obtained by the equation [12] in the PFPO theory are better than that when $X_{12} = 0$ was taken, while reverse is the case for oxolane + amines mixtures.

Fig. 5. Dependence of Rao's molar sound function on mole fraction of DIPE or oxolane at T=303.15 K. (a) DIPE + DMA (\bullet) and oxolane + DMA (\blacktriangle). (b) DIPE + DEA (\blacksquare) and oxolane + DEA (\Diamond) .

Table 10 Standard percentage deviation ((%)) between experimental and estimated speed of sound using empirical or theoretical relations for binary mixtures.

^a $X_{12} = 0$; $X_{12} =$ from Eq. (12).

5. Conclusions

The m[agnitu](#page-7-0)de of negative values of V_m^E , $\frac{E}{S}$, and L_f^E and positive values of Z^E for binary mixtures of DIPE with DMA or DEA is much larger than that observed for oxolane mixtures at 303.15, 313.15, and 323.15 K. This difference is attributed to the effect due to the difference in free volume and internal pressure of involved components. The speci"c interactions between unlike molecules through dipole…dipole interactions exist and dominate over the dispersion interactions. The overall estimation ability of speed of sound in presently investigated mixtures follows the sequence: PFPOT > CFT > NM > FLE IM > EM > VM.

Acknowledgement

The authors are thankful to Dr. P. Sriyutha Murthy, Scienti"c Of"cer D, Biofouling & Bio"lm Processes Section, Water & Steam Chemistry Division, BARC Facilities IGCAR Campus, Kalpakkam 60 102, Tamil Nadu, India for his constant encouragement throughout the study.

References

- [1] J.M. Prausnitz, R.N. Lichtenthaler, E.G. Azevedo, Molecular Thermodynamic of Fluid-Phase Equilibria, 2nd Ed., Prentice-Hall, Inc., 1986.
- [2] S.L. Oswal, N.B. Patel, J. Chem. Eng. Data 40 (1995) 840...844.
- [3] S.L. Oswal, P. Oswal, J.P. Dave, J. Mol. Liq. 94 (2001) 203…219.
- [4] S.L. Oswal, K.D. Prajapati, N.Y. Ghael, S.P. Ijardar, Fluid Phase Equilib. 218 (2004) 131…140.
- [5] S.L. Oswal, R.L. Gardas, R.P. Phalak, J. Mol. Liq. 116 (2005) 109…118.
- [6] S.L. Oswal, V. Pandiyan, B. Krishnakumar, P. Vasantharani, Thermochim. Acta 507…508 (2010) 27…34.
- [7] V. Pandiyan, P. Vasantharani, S.L. Oswal, A.N. Kannappan, J. Chem. Eng. Data 56 (2011) 269…277.
- [8] V. Pandiyan, S.L. Oswal, P. Vasantharani, Thermochim. Acta (2011), doi:10.1016/j.tca.2011.01.019.
- [9] J.A. Gonzalez, I. Mozo, I. Garcia de la Fuente, J.C. Cobos, Can. J. Chem. 83 (2005) 1812…1825.
- [10] N. Riesco, S. Villa, J.A. González, I. Garcša de la Fuente, J.C. Cobos, Fluid Phase Equilib. 202 (2002) 345…358.
- [11] S.L. Oswal, S.P. Ijardar, J. Solution Chem. 38 (2009) 321…344.
- [12] T.M. Letcher, U. Domanska, J. Chem. Thermodyn. 26 (1994) 553…560.
- [13] T.M. Letcher, U. Domanska, J. Chem. Thermodyn. 26 (1994) 1241…1247.
- [14] O. Nomoto, J. Phys. Soc. Jpn. 13 (1958) 1528…1532.

[15] W. Vand Dael, Thermodynamic Properties and Velocity of Sound, Btterworths, London, 1975.

- [16] S. Ernst, J. Glinski, B. Jezowska-Trzebiatowaska, Acta. Phys. Polond A55 (1979) 501…516.
- [17] S. Kalidoss, T. Saivasamoorthy, J. Pure Appl. Ultrason, 19 (1997) 9...15.
- [18] W. Schaaffs, K.-H. Molekularakustik (Eds.), Hellwege, Springer-Verlag, Berlin,
- 1975.
- [19] B. Jacobson, Acta Chem. Scand. 6 (1952) 1485...1498.
- [20] Molecular Theory of Solution, North-Holland, Amsterdam, 1957.
- [21] P.J. Flory, J. Am. Chem. Soc. 87 (1965) 1833…1838.
- [22] D. Patterson, G. Delmas, Disc. Faraday Soc. 49 (1970) 98…105.
- [23] S.L. Oswal, Acoust. Lett. 14 (1990) 17…25.
- [24] R. Gonzalez-Olmos, M. Iglesias, B.M.R.P. Santos, S. Mattedi, Phys. Chem. Liq. 46 (3) (2008) 223…237.
- [25] U.P. Govender, T.M. Letcher, S.K. Garg, J.C. Ahluwalia, J. Chem. Eng. Data 41 (1996) 147…150.
- [26] X. Meng, J. Wu, Z. Liu, J. Chem. Eng. Data 54 (2009) 2353…2358.
- [27] P. Venkatesu, M.J. Lee, H.M. Lin, J. Chem. Thermodyn. 37 (2005) 996…1002.
- [28] B. Giner, B. Oliver, I. Giner, G. Pera, C. Lafuente, J. Solution Chem. 36 (2007) 375…386.
- [29] J.A. Riddick, W.B. Bunger, T.K. Sakano, Organic Solvents Physical Properties and Methods of Puri"cation, 4th ed., Wiley Interscience, New York, 1986.
- [30] R. Palepu, J. Oliver, D. Campbell, J. Chem. Eng. Data 30 (1985) 355…360.
- [31] J. George, N.V. Sastry, Int. J. Thermophys. 24 (2003) 1697…1719.
- [32] K.V.N. Suresh Reddy, G. Sankara Reddy, A. Krishnaiah, Thermochim. Acta 440 (2006) 43…50.
- [33] T.M. Reid, J.M. Prausnitz, B.E. Poling, The Properties of Gases and Liquids, 4th ed., McGraw Hill, New York, 1987.
- [34] M. Gupta, I. Vibhu, J.P. Shukla, Fluid Phase Equilib. 244 (2006) 26…32.
- [35] S.L. Oswal, R.L. Gardas, R.P. Phalak, Thermochim. Acta 426 (2005) 199…206.
- [36] A. Ali, A.K. Nain, Pramana J. Phys. 58 (2002) 695…701.
- [37] A.K. Nain, Phys. Chem. Liq. 45 (2007) 371…388.
- [38] TRC Tables. Selected Values of Properties of Chemical Compounds. Thermodynamic Research Center Data Project, Texas A&M University, College Station, 971, 49, 2605…2609.
- [39] A.G. Oskoei, N. Safaei, J. Ghasemi, J. Chem. Eng. Data 53 (2008) 343…349.
- [40] P. Vasantharani, L. Balu, R. Ezhil Pavai, S. Shailajha, Glob. J. Mol. Sci. 4 (1) (2009) 42…48.
- [41] M.R. Rao, J. Chem. Phys. 9 (1941) 682…685.
- [42] G.C. Benson, O. Kiyohara, J. Chem. Thermodyn. 11 (1979) 1061…1064.
- [43] C.M. Kinart, W.J. Kinart, D.A. Checinska-Majak, Cwiklinska, J. Mol. Liq. 109 (2004) 19…22.
- [44] A. Ali, A.K. Nain, D. Chand, B. Lal, Phys. Chem. Liq. 45 (2007) 79…91.
- [45] H.T. Van, D. Patterson, J. Solution Chem. 11 (1982) 793…805.
- [46] M. Costas, D. Patterson, J. Solution Chem. 11 (1982) 807…821.
- [47] G.V. Rama Rao, A.V. Sarma, D. Ramachandran, C. Rambabu, Ind. J. Chem. 46A (2007) 1972…1978.
- [48] A.M. Awwad, R.A. Pethrick, J. Chem. Thermodyn. 16 (1984) 131…136.
- [49] A. Gayol, M. Iglesias, J.M. Goenaga, R.G. Concha, J.M. Resa, J. Mol. Liq. 135 (2007) 105…114.
- [50] M. Khammar, J.M. Shaw, Fluid Phase Equilib. 288 (2010) 145…154.
- [51] J. Glinski, J. Solution Chem. 31 (2002) 59…69.