

energy and perhaps the correct monopole matrix element with the ground state. It would be interesting to see the same method applied in detail to C^{12} .

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Probable Absence of K Capture in the Decay of Lead-205†

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A number of experiments designed to detect the K -capture decay of Pb^{205} are discussed. No K x-ray activity which could be ascribed to Pb^{205} was observed in deuteron-bombarded natural lead or neutron-irradiated 25.7% Pb^{204} . From these results, the L -capture half-life of Pb^{205} , and its absence in natural lead, it is concluded that Pb^{205} does not decay by K capture with a half-life in the range 2 seconds $< t_{1/2} < 10^{10}$ years.

IN view of the recent determination¹ of the L -capture half-life of Pb^{205} , it seemed appropriate to present evidence for the lack of K capture. These data extend the K -capture half-life limits given in previous communications.^{2,3}

The abundance of Pb^{205} occurring naturally has been reported to be less than 0.001%⁴ of natural lead. The age of the solid shell of the earth has been estimated to be 4.5×10^9 years.⁵ The abundance of Pb^{204} is reported to be $1.45 \pm 0.007\%$.⁶ If it is assumed that the primordial abundances of Pb^{204} and Pb^{205} were equal, one calculates from the present upper limit of 1 part Pb^{205} per 1450 parts of Pb^{204} that the half-life of Pb^{205} is $\leq 5 \times 10^8$ years.

A K -capture half-life in the region 5 days to 10^{10} years was excluded on the basis of two experiments. In the first of these, 0.400 g of lead, enriched to 25.7% in Pb^{204} ,⁷ was exposed to an integrated flux of 3×10^{20} thermal neutrons/cm² in the Materials Testing Reactor (MTR). If the thermal-neutron capture cross section of Pb^{204} is taken to be 0.9 barns,⁸ one calculates that

8.0×10^{16} atoms of Pb^{205} were present at the end of the irradiation if no decay occurred during the irradiation. After suitable radiochemical purification, the sample in the form of $PbSO_4$ was counted on a NaI scintillation spectrometer. The energy scale was calibrated with the 74.7-keV Pb x-ray of 6.4-day Bi^{206} . No activity (< 1 count/min) was detected in this energy region (the Tl K x-ray energy is 72.7 keV). Taking 10 disintegrations/min as the upper limit of the Pb^{205} activity, one calculates the K -capture half-life to be longer than 10^{10} years. Since the irradiated sample was counted approximately one month after the irradiation, a K -capture half-life of the order of 5 days would have been detectable.

The second experiment involved the radiochemical separation of bismuth isotopes from a 0.005-in. lead foil bombarded with 1500 microampere-hours of 21-MeV deuterons. Using the same procedure for milking lead from the separated bismuth sample as was used in the L -capture experiment,¹ we found less than 0.2 count/min of activity due to Pb^{205} in the Tl K x-ray region, corresponding to a K -capture half-life longer than 5×10^8 years.

K -capture half-lives in the region of a few seconds to several days were investigated in two separate experiments. In the first of these, two samples of lead, one of natural abundance and the other enriched to 25.7% in Pb^{204} , were irradiated in the thermal flux region of the Brookhaven reactor. The two samples were counted within 4 minutes after the end of the irradiation in a single-channel pulse analyzer—NaI(Tl) arrangement which had been preset to detect Tl x-rays. No activity was observed that satisfied the 1.45 to 25.7 abundance of Pb^{204} in the two samples.

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¹ J. R. Huizenga and J. Wing, *Phys. Rev.* **102**, 926 (1956).

² Sugihara, Herber, Bennett, and Coryell, *Phys. Rev.* **95**, 298 (1954).

³ J. R. Huizenga and C. M. Stevens, *Phys. Rev.* **96**, 548 (1954).

⁴ White, Collins, and Rourke, *Phys. Rev.* **101**, 1786 (1956).

⁵ Patterson, Brown, Tilton, and Inghram, *Science* **121**, 69 (1955).

⁶ Farquhar, Palmer, and Aitken, *Nature* **172**, 860 (1953).

⁷ Obtained from the Stable Isotopes Division, U. S. Atomic Energy Commission.

⁸ H. Pomerance, *Phys. Rev.* **88**, 412 (1952).

In the second experiment, the bismuth activities resulting from the $\text{Pb}(d, xn)\text{Bi}$ reactions produced by 15-Mev deuterons in the MIT cyclotron were adsorbed on Dowex-1 anion resin and the lead was quickly eluted to look for a very short half-life. The eluate flowed rapidly into the well of a NaI scintillation counter. An increase in counting rate was observed, but the activity had a relatively long half-life ($t_{1/2}$ ~several hours) and could be attributed to the equilibrium concentration of Bi activities in the eluate. A half-life as short as 2 seconds could have been detected. No Pb activity, the initial concentration of which fell with the 14-day half-life of the Bi^{205} parent, could be observed in washings of the Dowex-1 column which would have removed any Pb decay daughters.⁹⁻¹¹

⁹ E. C. Campbell and F. Nelson, Phys. Rev. **91**, 499 (1953).

¹⁰ F. Nelson and K. A. Kraus, J. Am. Chem. Soc. **76**, 5916 (1954).

¹¹ Note added in proof.—From the negative results of this experiment, nothing can be concluded regarding the existence of

These data indicate that Pb^{205} does not decay by K capture with a half-life in the range of 2 sec to 10^{10} years. A partial K -capture half-life of less than 2 sec can be ruled out on the basis of the observed¹ L -capture half-life of 5×10^7 years. Hence it is concluded that if Pb^{205} decays by K capture at all the half-life for this decay mode must be greater than 10^{10} years.

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a short-lived Pb^{205} isomer, except that Bi^{205} electron capture does not heavily populate the $i_{13/2}$ level in Pb^{205} . In the analogous Bi^{207} decay, 87% of the e.c. (electron capture) transitions go to the $i_{13/2}$ state of the 0.8 sec Pb^{207m} . The half-life of this $h_{9/2}$ to $i_{13/2}$ transition is 9.2 years ($8 \gamma/0.87$) and the competing $h_{9/2}$ to $f_{7/2}$ transition in Bi^{207} has an energy of only 50 ± 40 kev. In the decay of Bi^{205} , this latter transition probably has considerably more energy, and may be the transition most influencing the half-life. The $h_{9/2}$ to $i_{13/2}$ decay may be roughly estimated by the ratio $14d$ to 9 years or 0.004, and this low rate of population of the $i_{13/2}$ level of Pb^{205m} may well be the reason that we do not observe this isomer.

Nucleon-Hole Interaction in jj Coupling*

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A theorem connecting the energy levels in jj coupling of a nucleon-nucleon and nucleon-hole system is derived, and applied in particular to Cl^{38} and K^{40} .

THE interaction of a nucleon with a hole in a closed shell has been considered by Racah and others¹ by an application of tensor algebraic methods. An alternate approach, independent of the nature of the two-nucleon interaction, is possible for the jj -coupling case. It is based on the property that the coefficients of fractional parentage (c.f.p.) connecting the states of one- and two-hole systems have a particularly simple analytical form.² This approach yields a theorem which gives the energy levels of a nucleon-hole system directly in terms of the energy levels of the corresponding nucleon-nucleon system.

Consider a proton (or a neutron) hole in shell j , and a neutron (or a proton) in shell j' coupled to a resultant spin J . The $2j$ protons are antisymmetrized with respect to each other, but not with respect to the neutron. (Note, however, that if, as we assume, all the lower-lying proton and neutron subshells are filled, this state does have a specified isotopic spin.)

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¹ G. Racah, Phys. Rev. **62**, 438 (1942); Marty, Nataf, and Prentki, J. phys. radium **15**, 134 (1954).

² G. Racah, Phys. Rev. **63**, 367 (1943), Eq. (19); C. Schwartz and A. de-Shalit, Phys. Rev. **94**, 1257 (1954).

The interaction Hamiltonian for the system, $\sum_{i < j} H_{ij}$, may then be written as a sum of two parts: (i) The interactions among the identical nucleons giving the energy of the hole, $E[j^{-1}]$, independent of J , and (ii) the neutron-proton interactions. To evaluate the second part, c.f.p. are used to write the state of the hole $|j^{-1}\rangle$ in terms of $2j-1$ antisymmetrized protons coupled to a single proton; then the neutron and the proton are recoupled by means of a Racah coefficient. Now, using the explicit analytical form for the c.f.p., the sum over the parent states of the hole can easily be carried out, leaving only a sum over the states J_0 , of a proton in shell j coupled to a neutron in shell j' . Finally, we obtain for the relative energies $E_J[j^{-1}j']$ of the neutron-proton hole configuration,³

$$E_J[j^{-1}j'] = - \sum_{J_0} (2J_0+1) W(jj'j'j; JJ_0) E_{J_0}[jj'], \quad (1)$$

and the reciprocal relation,

$$E_J[jj'] = - \sum_{J_0} (2J_0+1) W(jj'j'j; JJ_0) E_{J_0}[j^{-1}j']. \quad (1')$$

³ For detailed derivation of (1) and (1') see S. P. Pandya, U. S. Atomic Energy Commission Report NYO-7590 (unpublished), Appendix.