

Phase Shifts for Nucleon-Nucleon Scattering at 280 Mev*

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RECENT results on the polarization of nucleons as measured by double scattering have given more specific indications regarding phase shifts¹ than have been previously available. In particular, the more recent data² at 280 Mev indicate³ the presence of phase shifts for orbital angular momenta $L > 1$ at this energy. A closer examination shows that even more specific conclusions can be drawn.

If the analysis of $(P\sigma)_{p-p}$ is confined to orbital angular momenta $L < 4$, it may be shown that even if one assumes the presence of coupling between 3P_2 and 3F_2 , the terms in $\sin\theta \cos\varphi \cos^5\theta$ can be present only if there is a 3F_4 phase shift δ_4^F . The presence of terms of above type in $(P\sigma)_{p-p}$ indicates, therefore, that there are effects caused by F -wave phase shifts δ^F originating directly rather than through coupling to states of lower orbital angular momenta. The conclusion regarding the presence of direct δ^F effects rests heavily on the presence of the term $0.116(\sin\theta) \text{ mb/sterad}$ in the analysis of $(P\sigma)_{p-p}$. According to the authors,² the presence of this term is not quite certain since the data can be represented reasonably well without it. Considering the Fourier analysis of the experimental curve by means of Fourier's formula, the coefficient 0.116 is a sum of two positive integrals in the intervals from $(0^\circ, 30^\circ)$ and $(60^\circ, 90^\circ)$ and one negative integral over $(30^\circ, 60^\circ)$. Each of the positive integrals is ~ 3 times the sum of the three so that the accuracy of the result is rather poor. The most uncertain of the three is the integral over $(0^\circ, 30^\circ)$. The knowledge of the curve in this region would admit an error of 10 percent in the integral over it corresponding to an error of perhaps 30 percent in the number 0.116. A qualitatively incorrect value of this coefficient thus appears unlikely but is not altogether excluded since both for $p-p$ and $p-n$ the experimental points at $\theta \cong 18^\circ$ are systematically off the curves. If the data could be made more accurate in this respect, clarification regarding F waves would result.

Since there is little doubt regarding the presence of noncentral forces at low energies, as indicated by the quadrupole moment of the deuteron, it appears likely that there are differences between different δ^D and between different δ^F . The coupling of 3S_1 to 3D_1 is very likely for the same reason. Calculation shows that if $\delta_2^D = \delta_3^D = 0$ and if all phase shifts but δ^D , δ^S are neglected then $(P\sigma)_{p-n}$ contains no term in $\sin\theta \cos\theta P_2 \times (\cos\theta)$ the only dependence being on $\sin\theta \cos\theta$. While it is probable that the differences in the δ^D and the coupling of 3S_1 to 3D_1 are connected through the interaction energy, one can separate phenomenologically the

coupling from the direct occurrence of δ_2^D and δ_3^D . In this sense the 3D_1 state caused by 3S_1 does not produce $\sin\theta \cos\theta P_2(\cos\theta)$ terms, without the aid of cross products involving the pairs S - G , D - D , P - F , etc., the sum of the L 's in a product being even if an odd power of $\cos\theta$ multiplying the factor $\sin\theta$ is desired.

Among the D - D combinations the 3D_1 - 3D_2 term does not occur with $\sin\theta \cos^3\theta$ similarly to the absence of 3F_3 - 3F_2 combinations with $\sin\theta \cos^5\theta$. Presence of coupling between 3D_1 and 3S_1 or between 3F_2 and 3P_2 does not interfere with this rule. The conclusion regarding the role of 3S_1 - 3D_1 coupling differs only in emphasis from that of Fried, the possibility of ${}^3D_2 = {}^3D_3 = 0$ in the presence of the coupling being specifically considered here.

The 3F phase shifts required by the data are of the order of 15° , an amount sufficiently smaller than the 3P phase shifts to make it conceivable that phase shifts for $L > 3$ are not important in the analysis of $p-p$ data at this energy. The repulsive character of potentials expected for triplets with odd L in $p-p$ scattering on the symmetric theory may be expected to contribute to the smallness of these phase shifts.

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¹ A. Garren, Phys. Rev. **92**, 213 (1953); de Carvalho, Heiberg, Marshall, and Marshall, Phys. Rev. **94**, 1796 (1954).

² Chamberlain, Donaldson, Segrè, Tripp, Wiegand, and Ypsilantis, Phys. Rev. **95**, 850 (1954).

³ B. D. Fried, Phys. Rev. **95**, 851 (1954). In the paper by Thaler, Bengtson, and Breit, Phys. Rev. **94**, 683 (1954) quoted by Fried, Eq. (9) contains a misprint, $<$ appearing for $>$, in contradiction to the text immediately to the right of Eq. (9).

Polarization of High-Energy Protons in Elastic Scattering on Helium and Carbon*

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WE have previously reported¹ results of experiments concerning polarization of proton beams by scattering from complex nuclei. Other authors² have also dealt with this problem. Particular interest has been attached to polarization by elastic scattering and several publications³⁻⁵ have treated the theory for this case.

As previously, we have attempted to isolate elastic scattering by insertion of an absorber into the counter telescope used to detect the scattered protons. Owing to range straggling in the absorber and to inhomogeneity of the beam, it is impossible to exclude with certainty all protons resulting from inelastic scattering. For example, we mention the scattering by carbon in which the lowest excited state is at 4.4 Mev, while our counting arrangement could be guaranteed to reject scatter-