

# Symmetry of the $ND^{-1}$ Solutions for Coupled Scattering Amplitudes

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We prove that the  $ND^{-1}$  matrix solutions for coupled scattering amplitudes are symmetric provided the given discontinuity of the scattering matrix across the unphysical cut is symmetric.

A USEFUL method developed by Omnes<sup>1</sup> and Chew and Mandelstam<sup>2</sup> to construct single-channel partial-wave amplitudes satisfying the requirements of analyticity and unitarity has been extended by Mandelstam<sup>3</sup> and Bjorken<sup>4</sup> to the multiple-channel case. The method consists in writing the scattering matrix  $G$  as a product of a matrix  $N$  which contains the unphysical (dynamical) singularities of  $G$  and the inverse of a matrix  $D$  which contains the physical (unitarity) singularities of  $G$ . A well-known requirement from time-reversal invariance is that  $G$  be a symmetric matrix. We prove that  $G$ , written in the form

$$G = ND^{-1}, \quad (1)$$

satisfies this condition, provided only that the given discontinuity of  $G$  across the unphysical cuts is a symmetric matrix.

By assumption, the matrices  $N$  and  $D$  satisfy dispersion relations

$$N = -\frac{1}{\pi} \int_{-\infty}^{s_L} \frac{\text{Im}N(s')}{s' - s} ds', \quad (2)$$

$$D = 1 + \frac{1}{\pi} \int_{-\infty}^{+\infty} \frac{\text{Im}D(s')}{s' - s} ds', \quad (3)$$

where  $s_L$  and  $s_R$  are the end points of the unphysical ( $-\infty < s < s_L$ ) and the physical ( $s_R < s < \infty$ ) cuts. For simplicity we restrict the unphysical singularities to the real axis. From Eqs. (1)–(3) it follows that

$$\begin{aligned} \text{Im}N &= (\text{Im}G)D \quad \text{when } s < s_L \\ &= 0 \quad \text{when } s_L < s, \\ \text{Im}D &= \text{Im}(G^{-1})N \quad \text{when } s_R < s \\ &= 0 \quad \text{when } s < s_R. \end{aligned} \quad (4)$$

The unitarity of the  $S$  matrix requires that

$$\text{Im}G^{-1} = -\rho \quad \text{when } s_R < s, \quad (5)$$

where  $\rho$  is the diagonal density of states matrix, and is therefore known. On the other hand,  $\text{Im}G$  along the unphysical cut depends on the form of the interaction. The only assumption we make is that  $\text{Im}G$  is a symmetric matrix.

Consider now

$$G - G^T = ND^{-1} - (D^{-1})^T N^T, \quad (6)$$

where the superscript  $T$  indicates the transposed matrix. Multiplying Eq. (6) on the right-hand side by  $D$  and on the left-hand side by  $D^T$ , taking the imaginary part, and substituting Eqs. (4) and (5), we obtain

$$\begin{aligned} \text{Im}[D^T(G - G^T)D] \\ &= D^T(\text{Im}G)D - D^T(\text{Im}G)^T D \quad \text{when } s < s_L \\ &= -N^T \rho^T N + N^T \rho N \quad \text{when } s_R < s. \end{aligned} \quad (7)$$

Since  $\rho$  is diagonal and  $\text{Im}G = (\text{Im}G)^T$ , the right-hand side of Eq. (7) is everywhere zero. It follows that the function  $D^T(G - G^T)D$  is analytic everywhere in the  $s$  plane, and vanishes at  $\infty$ , Eqs. (2)–(3). Hence it is identically zero and we obtain

$$G = G^T. \quad (8)$$

If  $D$  has zeros, consideration of the matrix  $(\det D)G$  leads to the same conclusion, Eq. (8).

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<sup>1</sup> R. Omnes, *Nuovo cimento* **8**, 316 (1958).

<sup>2</sup> G. Chew and S. Mandelstam, *Phys. Rev.* **119**, 467 (1960).

<sup>3</sup> G. Chew (private communication); M. Nauenberg, thesis, Cornell University, 1960 (unpublished); F. Ferrari, M. Nauenberg, and M. Pusterla (to be published).

<sup>4</sup> J. Bjorken, *Phys. Rev. Letters* **4**, 473 (1960).