

Errata

Properties of Liquid He³ at Low Temperature, K. A. BRUECKNER AND J. L. GAMMEL [Phys. Rev. **109**, 1040 (1958)]. The numerical values determining the magnetic susceptibility according to Eq. (32) are given incorrectly. The correct values are:

$$\begin{aligned} M/M^* &= 0.543, \\ \frac{Mk_F\Omega}{2\pi^2} [a_0(k_F, k_F)_{av} - a_e(k_F, k_F)_{av}] &= -0.780, \\ \frac{2Mk_F\Omega}{\pi^2} \left\{ \int \frac{d\mathbf{k}_i}{4\pi k_F^3} \frac{\partial}{\partial k_F} [a_0(\mathbf{k}_i, \mathbf{k}_j)_{k_j=k_F} \right. \\ &\quad \left. + \frac{1}{8} \int \frac{d\mathbf{k}_i}{4\pi k_F^3} \int \frac{d\mathbf{k}_j}{4\pi k_F^3} \frac{\partial}{\partial k_F} \right. \\ &\quad \left. \times [5a_0(\mathbf{k}_i, \mathbf{k}_j) + a_e(\mathbf{k}_i, \mathbf{k}_j)] \right\} = 0.320. \end{aligned}$$

These give $E_s/E_s(F) = 0.083$ and $\chi = 12.0 \chi_F$, where χ_F is the susceptibility of an ideal Fermi gas.

Study of (d, α) Reactions on Some Light Nuclei, G. E. FISCHER AND V. K. FISCHER [Phys. Rev. **114**, 533 (1959)]. The data in Fig. 12 were incorrectly transcribed from Freemantle *et al.*⁶ and the ordinate scale unit should be 0.1 mb/steradian rather than 0.01 mb/steradian. Similarly, the integrated experimental cross section for $O^{16}(d, \alpha_0)N^{14}$ given in Table I should be 2.2 mb rather than 0.22 mb. In Table II, the following values for $O^{16}(d, \alpha_0)N^{14}$ should be changed: $(d\sigma/d\Omega)_{exp}$ should be 0.40 mb rather than 0.04 mb and $(\gamma_0)_{normalized}$ should be 0.08 rather than 0.008. These corrections invalidate the conclusion that the $O^{16}(d, \alpha_0)N^{14}$ reaction can best be described by compound-nucleus theory.

Scattering of 200-Mev Positrons by Electrons, J. A. POIRIER, D. M. BERNSTEIN, AND JEROME PINE [Phys. Rev. **117**, 557 (1960)]. The result quoted in line 6 of the abstract is incorrect, and should be $(13 \pm 9)\%$.

Calculation of the Magnetic Hyperfine Structure Coupling Constants of NO, HÉLÈNE LEFEBVRE-BRION AND C. M. MOSER [Phys. Rev. **118**, 675 (1960)]. The configuration interaction function on p. 677 should be:

$$\begin{aligned} \Psi &= 0.95958\psi_0 + 0.00719\psi_1 + 0.06902\psi_2 + 0.06050\psi_2 \\ &\quad - 0.16466\psi_3 - 0.15233\psi_4 + 0.11617\psi_4 - 0.07637\psi_7 \\ &\quad + 0.00577\psi_8 + 0.03247\psi_9. \end{aligned}$$

The values of the constants obtained from this

function thus will be slightly changed from the values given in the paper.

Table IV, line 2 $\frac{1}{2}a$ 14.22 $\frac{1}{3}d$ 12.47 $\frac{1}{3}c$ $-7.03(-9)^b$.
Table VII, line 2 A 125.7.
Page 678, $\psi^2(0) = 0.096$ a.u.
Page 680, line 5 from bottom, $q = 1.559 - 1.342 = 0.217$ a.u.

Perturbation Theory Applied to the Nuclear Many-Body Problem, J. S. LEVINGER, M. RAZAVY, O. ROJO, AND N. WEBRE [Phys. Rev. **119**, 230 (1960)]. The left-hand column in Table II should be titled " $2x^2$."

Polarization of Protons Scattered from C¹², T. A. TOMBRELLO, R. BARLOUTAUD, AND G. C. PHILLIPS [Phys. Rev. **119**, 761 (1960)]. The expression for the polarization on page 762 should read

$$P(\theta) = \frac{2 \operatorname{Im}(f_e f_i^*)}{|f_e|^2 + |f_i|^2} (\hat{k}_2 \times \hat{k}_1).$$

The authors wish to express their appreciation to Professor H. H. Barschall for pointing out this error.

Magnetic Scattering of Neutrons by Exchange-Coupled Lattices, A. W. SÁENZ [Phys. Rev. **119**, 1542 (1960)]. In line 1 of Eq. (2.7), ϵ should read ϵ' . In Eqs. (2.9a), the lines reading $-if(\mathbf{e} \cdot \boldsymbol{\lambda}) \times (\mathbf{e} \cdot [\mathbf{S}_i(0) \times \mathbf{S}_j(t)])$, $+\alpha f\{([\mathbf{e} \times \boldsymbol{\lambda}] \cdot [\mathbf{e} \times \mathbf{S}_i(0)]) \times ([\mathbf{e} \times \boldsymbol{\lambda}'] \cdot [\mathbf{e} \times \mathbf{S}_j(t)]) + i(\mathbf{e} \cdot \boldsymbol{\lambda}') \times ([\mathbf{e} \times \mathbf{S}_i(0)] \cdot [\mathbf{e} \times \mathbf{S}_j(t)])\}$ should read $+if(\mathbf{e} \cdot \boldsymbol{\lambda})(\mathbf{e} \cdot [\mathbf{S}_i(0) \times \mathbf{S}_j(t)])$, $+\alpha f \times [([\mathbf{e} \times \boldsymbol{\lambda}] \cdot [\mathbf{e} \times \mathbf{S}_i(0)])([\mathbf{e} \times \boldsymbol{\lambda}'] \cdot [\mathbf{e} \times \mathbf{S}_j(t)]) - i(\mathbf{e} \cdot \boldsymbol{\lambda}') \cdot [\mathbf{e} \times \mathbf{S}_j(t)]] - i(\mathbf{e} \cdot \boldsymbol{\lambda}') \cdot [\mathbf{e} \times \mathbf{S}_i(0)]$, respectively. In the first of Eqs. (2.9b), $\exp[-2W_0(\mathbf{q}_0)]$ and a_l should read $\exp[-W_0(\mathbf{q}_0)]$ and $a_l \exp[-W_l(\mathbf{q}_0)]$, respectively. In the definition of $\psi_1(\mathbf{e}; \alpha)$ in Eqs. (4.5), -2 should read 2 . The terms $-2f(\mathbf{e} \cdot \boldsymbol{\lambda})(\mathbf{e} \cdot \mathbf{y})\Re(\epsilon, \mathbf{q})$ in (4.12), $+2\eta f(\mathbf{e} \cdot \boldsymbol{\lambda})(\mathbf{e} \cdot \mathbf{y})$ in (4.16a), and $+2\eta f(\mathbf{e}_0 \cdot \boldsymbol{\lambda})(\mathbf{e}_0 \cdot \mathbf{y})$ in (4.17) should read $+2f(\mathbf{e} \cdot \boldsymbol{\lambda})(\mathbf{e} \cdot \mathbf{y})\Re(\epsilon, \mathbf{q})$, $-2\eta f(\mathbf{e} \cdot \boldsymbol{\lambda})(\mathbf{e} \cdot \mathbf{y})$, and $-2\eta f(\mathbf{e}_0 \cdot \boldsymbol{\lambda})(\mathbf{e}_0 \cdot \mathbf{y})$, respectively. In line 14 after Eq. (4.18), $d\theta > 0 (< 0)$ should read $d\theta < 0 (> 0)$.

Pseudodipolar Anisotropy in Cubic Ferromagnets at Low Temperatures, S. H. CHARAP AND P. R. WEISS [Phys. Rev. **116**, 1372 (1959)]. In the solution [Eq. (41)] of Eq. (40), certain terms peculiar to the face-centered cubic lattice and characterized by the parameter B (see Appendix A) have been omitted. The solution which has been given is, for the fcc, but the leading term in an expansion of the exact solution in powers of $B/4\pi(1 - \gamma_j)$ ($\approx 7 \times 10^{-3}$). The remaining terms are those which arise in the orders of perturbation theory beyond