Errata

Properties of Liquid He³ at Low Temperature, K. A. BRUECKNER AND J. L. GAMMEL [Phys. Rev. **109**, 1040 (1958)]. The numerical values determining the magnetic susceptibility according to Eq. (32) are given incorrectly. The correct values are:

$$M/M^* = 0.543,$$

$$\frac{Mk_F\Omega}{2\pi^2}[a_0(k_F,k_F)_{\rm av} - a_e(k_F,k_F)_{\rm av}] = -0.780,$$

$$\frac{2Mk_F\Omega}{\pi^2}\left\{\int\frac{d\mathbf{k}_i}{4\pi k_F{}^3}k_F\frac{\partial}{\partial k_F}[a_0(\mathbf{k}_i,\mathbf{k}_j)_{k_j=k_F} + \frac{1}{8}\int\frac{d\mathbf{k}_i}{4\pi k_F{}^3}\int\frac{d\mathbf{k}_j}{4\pi k_F{}^3}k_F\frac{\partial}{\partial k_F} \times[5a_0(\mathbf{k}_i,\mathbf{k}_j)+a_e(\mathbf{k}_i,\mathbf{k}_j)]\right\} = 0.320.$$

These give $E_s/E_s(F) = 0.083$ and $\chi = 12.0\,\chi_F$, where χ_F is the susceptibility of an ideal Fermi gas.

Study of (d,α) Reactions on Some Light Nuclei, G. E. FISCHER AND V. K. FISCHER [Phys. Rev. **114**, 533 (1959)]. The data in Fig. 12 were incorrectly transcribed from Freemantle *et al.*[6] and the ordinate scale unit should be 0.1 mb/steradian rather than 0.01 mb/steradian. Similarly, the integrated experimental cross section for $O^{16}(d,\alpha_0)N^{14}$ given in Table I should be 2.2 mb rather than 0.22 mb. In Table II, the following values for $O^{16}(d,\alpha_0)N^{14}$ should be changed: $(d\sigma/d\Omega)_{\rm exp}$ should be 0.40 mb rather than 0.04 mb and $(\gamma_0)_{\rm normalized}$ should be 0.08 rather than 0.008. These corrections invalidate the conclusion that the $O^{16}(d,\alpha_0)N^{14}$ reaction can best be described by compound-nucleus theory.

Scattering of 200-Mev Positrons by Electrons, J. A. POIRIER, D. M. BERNSTEIN, AND JEROME PINE [Phys. Rev. **117**, 557 (1960)]. The result quoted in line 6 of the abstract is incorrect, and should be $(13\pm 9)\%$.

Calculation of the Magnetic Hyperfine Structure Coupling Constants of NO, HÉLÈNE LEFEBVRE-BRION AND C. M. MOSER [Phys. Rev. **118**, 675 (1960)]. The configuration interaction function on p. 677 should be:

$$\Psi = 0.95958\psi_0 + 0.00719\psi_1 + 0.06902\psi_2 + 0.06050\psi_2 - 0.16466\psi_3 - 0.15233\psi_4 + 0.11617\psi_4 - 0.07637\psi_7 + 0.00577\psi_8 + 0.03247\psi_9.$$

The values of the constants obtained from this function thus will be slightly changed from the values given in the paper.

	$\frac{1}{2}a$	$\frac{1}{3}d$	$\frac{1}{3}c$
Table IV, line 2	14.22	12.47	-7.03 $(-9)^{\rm b}$.
Table VII, line 2	A	125.7.	

Page 678, $\psi^2(0) = 0.096$ a.u.
Page 680, line 5 from bottom, $q = 1.559 - 1.342 = 0.217$ a.u.

Perturbation Theory Applied to the Nuclear Many-Body Problem, J. S. LEVINGER, M. RAZAVY, O. ROJO, AND N. WEBRE [Phys. Rev. **119**, 230 (1960)]. The left-hand column in Table II should be titled "$2x^2$."

Polarization of Protons Scattered from C¹², T. A. TOMBRELLO, R. BARLOUTAUD, AND G. C. PHILLIPS [Phys. Rev. **119**, 761 (1960)]. The expression for the polarization on page 762 should read

$$\mathbf{P}(\theta) = \frac{2\,{\rm Im}(f_c f_i^*)}{|f_c|^2 + |f_i|^2}(\hat{k}_2 \times \hat{k}_1).$$

The authors wish to express their appreciation to Professor H. H. Barschall for pointing out this error.

Magnetic Scattering of Neutrons by Exchange-Coupled Lattices, A. W. SÁENZ [Phys. Rev. **119**, 1542 (1960)]. In line 1 of Eq. (2.7), ϵ should read ϵ'. In Eqs. (2.9a), the lines reading $-if(\mathbf{e}\cdot\boldsymbol{\lambda}) \times(\mathbf{e}\cdot[\mathbf{S}_i(0)\times\mathbf{S}_j(t)])$, $+\alpha f\{([\mathbf{e}\times\boldsymbol{\lambda}]\cdot[\mathbf{e}\times\mathbf{S}_i(0)]) \times([\mathbf{e}\times\boldsymbol{\lambda}']$, and $\times([\mathbf{e}\times\mathbf{S}_i(0)]\cdot[\mathbf{e}\times\mathbf{S}_j(t)])+i(\mathbf{e}\cdot\boldsymbol{\lambda}')$ should read $+if(\mathbf{e}\cdot\boldsymbol{\lambda})(e\cdot[\mathbf{S}_i(0)\times\mathbf{S}_j(t)])$, $+\alpha\{f \times[([\mathbf{e}\times\boldsymbol{\lambda}]\cdot[\mathbf{e}\times\mathbf{S}_i(0)])([\mathbf{e}\times\boldsymbol{\lambda}']$, and $\times([\mathbf{e}\times\mathbf{S}_i(0)] \cdot[\mathbf{e}\times\mathbf{S}_j(t)])]-i(\mathbf{e}\cdot\boldsymbol{\lambda}')$, respectively. In the first of Eqs. (2.9b), $\exp[-2W_0(\mathbf{q}_0)]$ and a_l should read $\exp[-W_0(\mathbf{q}_0)]$ and $a_l\exp[-W_l(\mathbf{q}_0)]$, respectively. In the definition of $\psi_1(\mathbf{e};\alpha)$ in Eqs. (4.5), -2 should read 2. The terms $-2f(\mathbf{e}\cdot\boldsymbol{\lambda})(\mathbf{e}\cdot\boldsymbol{\mu})\mathfrak{N}(\epsilon,\mathbf{q})$ in (4.12), $+2\eta f(\mathbf{e}\cdot\boldsymbol{\lambda})(\mathbf{e}\cdot\boldsymbol{\mu})$ in (4.16a), and $+2\eta f(\mathbf{e}_0\cdot\boldsymbol{\lambda})(\mathbf{e}_0\cdot\boldsymbol{\mu})$ in (4.17) should read $+2f(\mathbf{e}\cdot\boldsymbol{\lambda})(\mathbf{e}\cdot\boldsymbol{\mu})\mathfrak{N}(\epsilon,\mathbf{q})$, $-2\eta f(\mathbf{e}\cdot\boldsymbol{\lambda})(\mathbf{e}\cdot\boldsymbol{\mu})$, and $-2\eta f(\mathbf{e}_0\cdot\boldsymbol{\lambda})(\mathbf{e}_0\cdot\boldsymbol{\mu})$, respectively. In line 14 after Eq. (4.18), $d\theta>0(<0)$ should read $d\theta<0(>0)$.

Pseudodipolar Anisotropy in Cubic Ferromagnets at Low Temperatures, S. H. CHARAP AND P. R. WEISS [Phys. Rev. **116**, 1372 (1959)]. In the solution [Eq. (41)] of Eq. (40), certain terms peculiar to the face-centered cubic lattice and characterized by the parameter B (see Appendix A) have been omitted. The solution which has been given is, for the fcc, but the leading term in an expansion of the exact solution in powers of $B/4\pi(1-\gamma_j)$ $(\approx 7\times 10^{-3})$. The remaining terms are those which arise in the orders of perturbation theory beyond