

$T=3/2$ Level in F^{19} at 9.07 Mev

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Contrary to previous reports, we have detected α particles emitted from the 9.07-Mev state of F^{19} formed by the reaction $O^{18}(p,\alpha)N^{15}$. The α -particle reduced width indicates that the isotopic spin of this state is $T=\frac{3}{2}$ with an admixture of about 15% of $T=\frac{1}{2}$. The spin of the state is determined to be $\frac{7}{2}$ from the α -particle angular distribution, while odd parity seems to be favored.

THE reactions $O^{18}(p,\gamma)F^{19}$ and $O^{18}(p,\alpha)N^{15}$ are both resonant at a series of bombarding energies up to 3 Mev.^{1,2} The absence of a (p,α) resonance at the energy 1.169 Mev where the (p,γ) reaction is resonant, has been interpreted¹ as evidence for the $T=\frac{3}{2}$ character of the 9.07-Mev state of F^{19} . Considering the level density at this energy, it is surprising that the level should be purely $T=\frac{3}{2}$ and not mixed with a $T=\frac{1}{2}$ component by the Coulomb forces. We therefore searched for and found α particles from this resonance with the results given in Table I.

This reaction was studied with the 2-Mev Van de Graaff generator of the Ecole Normale Supérieure. Targets were 20- μ g/cm² self-supporting Al_2O_3 foils prepared³ by anodic oxidation of aluminum in an electrolyte containing 90% of O^{18} . Alpha particles were detected with p - n junctions prepared in n -type 75-ohm cm silicon.⁴ The solid angle subtended at the target by the detector was 1.4×10^{-3} steradian. The resonant excitation curve was traced by varying the voltage applied directly to the target by means of a ± 5 -kev calibrated supply, in steps of 250 ev. The beam energy furnished by the accelerator itself was constant to better than 250 ev during a 20-minute run. By studying the reaction $Al^{27}(p,\gamma)$ at 993 kev ($\Gamma \leq 140$ ev), the energy profile of the incident beam was found to be closely Gaussian with a half-width of 1 kev. The same resonance of $Al^{27}(p,\gamma)$ established the absolute energy scale, combined with the threshold of the $Li^7(p,n)$ reaction (1881 kev). The thickness of the target, which corresponds to a total stopping power of 10.8×10^{-15} ev cm² per atom of O^{18} at resonance, was measured to be 2.75 kev. By comparison with the excitation curve of $Al^{27}(p,\gamma)$ at 993 kev we determined the width of this resonance of $O^{18}(p,\alpha)$ to be $\Gamma \leq 500$ ev. We measured simultaneously the γ -ray yield to verify that the same resonance is involved for the γ rays and α particles. The

α -particle angular distribution was measured from 40° to 170° at 15° intervals, and the results normalized to a monitor p - n junction detecting the α particles at 170°.

It was very difficult to measure this angular distribution since the resonance is superposed on the large tails of neighboring resonances. For forward angles the resonance contributed 50% of the observed counting rate, and near 90° only 20%. However, the least-square analysis of the data gave a coefficient of P_6 equal to 0.57 ± 0.13 .

A simple explanation of the narrow width of the resonance can result from high-orbital angular momenta for proton and α particle. The observed angular distribution is consistent with a spin of $\frac{7}{2}$ and parity odd or even within the experimental errors. Odd parity is preferred since it leads to a reasonable proton reduced width. This choice is further supported by the preferential γ -ray de-excitation¹ to the 2.78-Mev ($9/2^+$; $7/2^+$), the 1.35-Mev ($5/2^-$), and the 0.198-Mev ($5/2^+$) excited states of F^{19} . The angular distribution of the γ rays¹ leads to the choice of $9/2^+$ for the 2.78-Mev level of F^{19} , consistent with the collective model interpretation of this level scheme.⁵

By comparison with the systematics of α -particle reduced widths in light nuclei which range from 10^0 to 10^{-2} , we see that the present value of 4×10^{-4} indicates some hindrance of the α -particle transition. Probably, therefore, the state is mainly $T=\frac{3}{2}$ but with an amplitude of at most 15% of $T=\frac{1}{2}$.

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TABLE I. The parameters of the resonance in the reaction $O^{18}(p,\alpha)N^{15}$ measured in this experiment.

Energy at resonance	Γ_{res}	$\sigma_{\text{max}}\Gamma$	θ_{p^2}	θ_{α^2}	Spin scheme	Angular distribution ^b
1165.2 ± 1.0 kev	≤ 500 ev	6 ± 1.5 ev-barn	$\leq 10^{-3a}$	$\leq 4 \times 10^{-4}$	$\frac{1}{2}(s^2)\frac{5}{2}(s^2)\frac{1}{2}$ $\frac{1}{2}(s^2)\frac{7}{2}(s^2)\frac{1}{2}$	$1 + (8/7)P_2(\cos\theta) + (6/7)P_4(\cos\theta)$ $1 + (25/21)P_2(\cos\theta) + (81/77)P_4(\cos\theta) + (25/33)P_6(\cos\theta)$

^a Fraction of Wigner limit evaluated for f -wave proton; g -wave proton leads to a reduced width exceeding this limit.

^b Theoretical angular distribution. The presence of the $P_6(\cos\theta)$ term in the measured distribution leads to the choice of $7/2$ for the level spin.

¹ J. W. Butler and H. D. Holmgren, Phys. Rev. **116**, 1485 (1959).

² H. A. Hill and S. M. Blair, Phys. Rev. **104**, 198 (1956).

³ The procedure will be published by G. Amsel and D. Samuel.

⁴ G. Amsel, P. Baruch, and O. Smulkowski, Nuclear Instr. **8**, 92 (1960); IRE Trans. on Nuclear Sci. NS-8, 21 (1961).

⁵ G. Rakavy, Nuclear Phys. **4**, 375 (1957).