

Elastic Scattering of Heavy Ions by Pb^{206} , Pb^{207} , and $\text{Pb}^{208}\dagger$

D. D. KERLEE

Seattle Pacific College Institute for Research, Seattle, Washington

AND

H. L. REYNOLDS AND E. GOLDBERG

Lawrence Radiation Laboratory, University of California, Livermore, California

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The angular distributions of C^{12} , N^{14} , O^{16} , and Ne^{20} elastically scattered by Pb^{206} , Pb^{207} , and Pb^{208} have been measured at laboratory energies of approximately 10.4 MeV/nucleon. The elastically scattered ions were recorded in photographic emulsions at laboratory angles from 17° to 175° . In general, measurements were extended only to angles where the ratio of the cross section to the Coulomb cross section, σ/σ_c , was greater than 0.2. In one case, the measurement was extended to a region where $\sigma/\sigma_c \leq 1.2 \times 10^{-5}$. The cross sections all exhibited a behavior similar to that previously reported for the same ions on Au^{197} and Bi^{209} . Excellent agreement with the semiclassical strong absorption model was obtained. Nuclear interaction distances calculated by fitting the strong absorption calculations are consistent with $r_0 = 1.45$ F, where $R = r_0(A_1^{1/3} + A_2^{1/3})$.

INTRODUCTION

ELASTIC scattering experiments conducted by Rutherford and subsequent workers¹ formed one of the important bases upon which the nuclear atom was conceived and established. Later, alpha-particle elastic-scattering experiments²⁻⁴ and more recent experiments involving the elastic scattering of heavy ions⁵⁻¹⁰

have yielded quantitative information on the respective interaction radii, cutoff values of angular momentum, and values of r_0 . Less quantitative information concerning nuclear surface properties has also been obtained. Analysis of the data by means of the semiclassical strong-absorption model, or "sharp cutoff" model, by Blair^{3,11,12} has been quite successful in

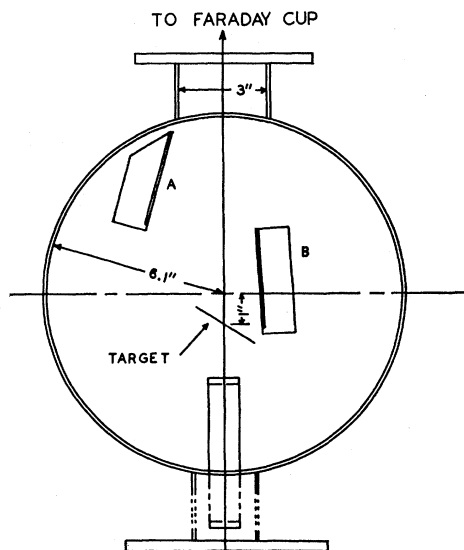


FIG. 1. Positions of emulsions in scattering chamber. Elastic scattering was detected primarily on plate A.

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¹ E. Rutherford and J. Chadwick, *Phil. Mag.* **50**, 889 (1925).

² G. W. Farwell and H. E. Wegner, *Phys. Rev.* **95**, 1212 (1954).

³ D. D. Kerlee, J. S. Blair, and G. W. Farwell, *Phys. Rev.* **107**, 1343 (1957).

⁴ Other early alpha particle work is given in reference 3.

⁵ H. L. Reynolds and A. Zucker, *Phys. Rev.* **102**, 1378 (1956).

⁶ E. Goldberg and H. L. Reynolds, *Phys. Rev.* **112**, 1981 (1958).

⁷ H. L. Reynolds, E. Goldberg, and D. D. Kerlee, *Phys. Rev.* **119**, 2009 (1960).

⁸ J. A. McIntyre, S. D. Baker, and T. L. Watts, *Phys. Rev.* **116**, 1212 (1960).

⁹ M. L. Halbert and A. Zucker, *Phys. Rev.* **115**, 1635 (1960).

¹⁰ A. Zucker, *Bull. Am. Phys. Soc.* **6**, 285 (1961).

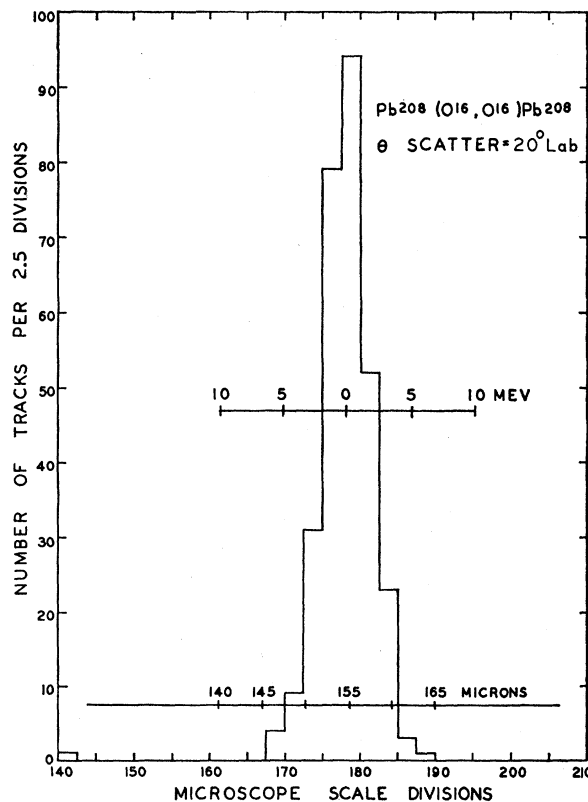


FIG. 2. A typical range spectrum in emulsion for O^{16} ions elastically scattered from Pb^{208} at a scattering angle of 20° in the laboratory. The micron scale has been corrected for track dip angle in the emulsion. The emulsion density is 3.976 g/cc.

¹¹ J. S. Blair, *Phys. Rev.* **95**, 1218 (1954).

¹² J. S. Blair, *Phys. Rev.* **108**, 827 (1957).

TABLE I. Measured ranges and energies for the incident particles.

Reaction	Range ^a (μ)	Energy at emulsion surface (MeV)	Reaction energy (MeV)	E/M (HILAC) (MeV)
$C^{12}+Pb^{206}$	198.6	122.8	124.2	10.38
$N^{14}+Pb^{206}$	175.8	145.0	146.8	10.52
$O^{16}+Pb^{206}$	156.3	164.7	167.1	10.48
$Ne^{20}+Pb^{206}$	130.6	203.6	207.2	10.41
$C^{12}+Pb^{207}$	197.5	122.4	123.7	10.33
$O^{16}+Pb^{207}$	155.1	163.9	166.1	10.41
$Ne^{20}+Pb^{207}$	131.7	204.7	208.1	10.44
$C^{12}+Pb^{208}$	195.1	121.6	122.8	10.26
$N^{14}+Pb^{208}$	176.9	145.7	147.2	10.53
$O^{16}+Pb^{208}$	154.8	163.7	166.0	10.41
$Ne^{20}+Pb^{208}$	129.9	202.8	206.2	10.34

^a The emulsion density was 3.976 g/cc.

regions where the "classical" parameter $\eta = zz'e^2/\hbar v$ is large compared to unity.

In the present work, the elastic scattering of C^{12} , N^{14} , O^{16} , and Ne^{20} ions from the lead isotopes, Pb^{206} , Pb^{207} , and Pb^{208} , has been studied using nuclear emulsions.

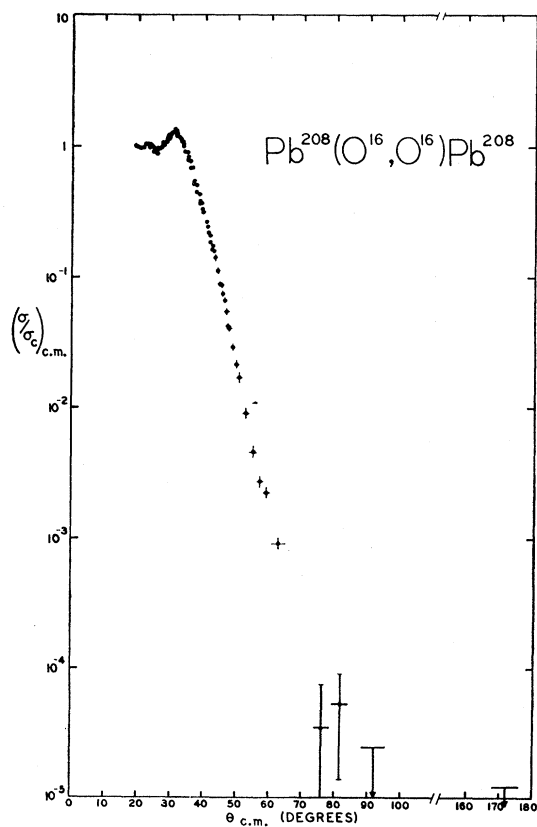


FIG. 3. Ratio of the differential elastic scattering cross section to the Coulomb cross section vs the center-of-mass scattering angle for O^{16} scattered by Pb^{208} . The laboratory energy before scatter is 166.0 MeV. Upper limits to the ratio are indicated at 92° and 172° c.m.

TABLE II. Target thicknesses and isotopic concentrations.

Target	Thickness (mg/cm ²)	Isotopic concentration (%)
Pb^{206}	0.97	64.9
Pb^{207}	0.82	71.5
Pb^{208}	0.90	88.9

Values of η range from approximately 25 to 40. Angular distributions were obtained for angular intervals from 17° (lab) to an angle such that $\sigma/\sigma_c(\theta) \approx 0.2$. In the case of O^{16} on Pb^{208} , an interval of from 17° to 175° was studied. Upper limits to the value of σ/σ_c were obtained at large angles. At 172° c.m., $\sigma/\sigma_c < 1.2 \times 10^{-5}$.

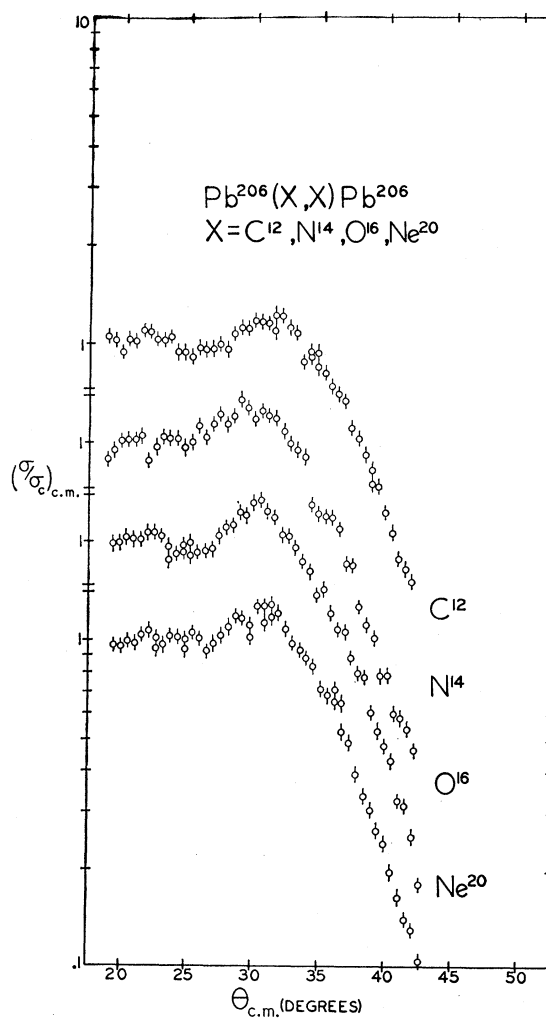


FIG. 4. Ratio of the differential cross section for elastic scattering to the Coulomb cross section vs the center-of-mass angle for C^{12} , N^{14} , O^{16} , and Ne^{20} scattered by Pb^{206} . Relative positions are arbitrary. The curves have been normalized to unity at small angles. Three hundred tracks were counted for each point.

EXPERIMENTAL PROCEDURE

The experimental arrangement has been described in previous papers.^{6,7} These papers will be referred to for description of the experimental procedure in the presentation of this further work.

The targets were exposed to heavy-ion projectiles of approximately 10.4 MeV nucleon from the Berkeley heavy-ion linear accelerator (HILAC). The beams from the HILAC were deflected by a bending magnet before they approached the target box (Fig. 1). The distance from the bending magnet to the first collimator slit was approximately $3\frac{1}{2}$ m. The beams passed through two $\frac{1}{8}$ -in. diam collimator slits, placed 7 in. apart, before striking the thin targets of isotopic lead. The scattered particles were detected by nuclear emulsions in positions *A* and *B*. Position *A* was loaded with Ilford E-1 emulsions, primarily for detection of elastically scattered particles. Ilford K-2 emulsions in positions *B*, *C*, and *D* (*C* and *D* not shown) were used in heavy-ion-

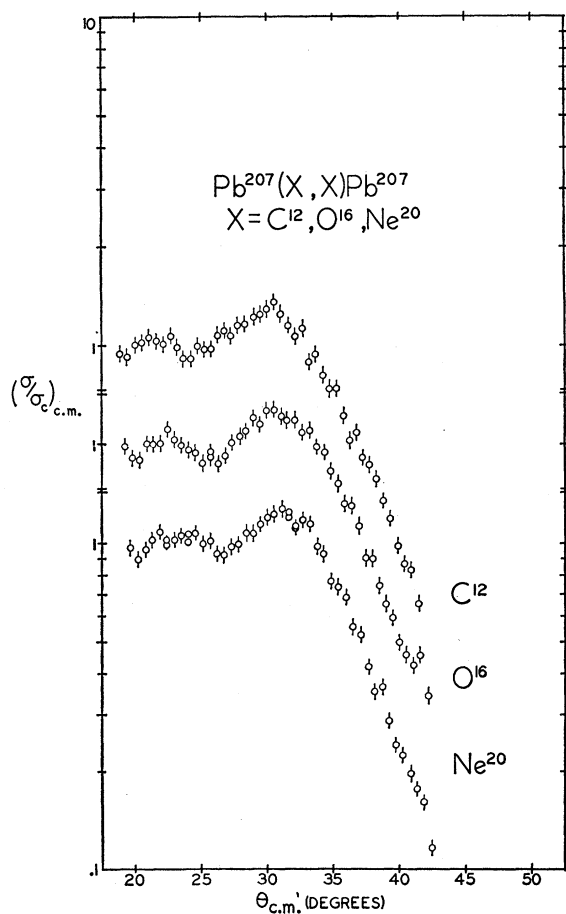


FIG. 5. Ratio of the differential cross section for elastic scattering to the Coulomb cross section vs the center-of-mass angle for C^{12} , O^{16} , and Ne^{20} scattered by Pb^{207} . Relative positions are arbitrary. The curves have been normalized to unity at small angles. Three hundred tracks were counted for each point.

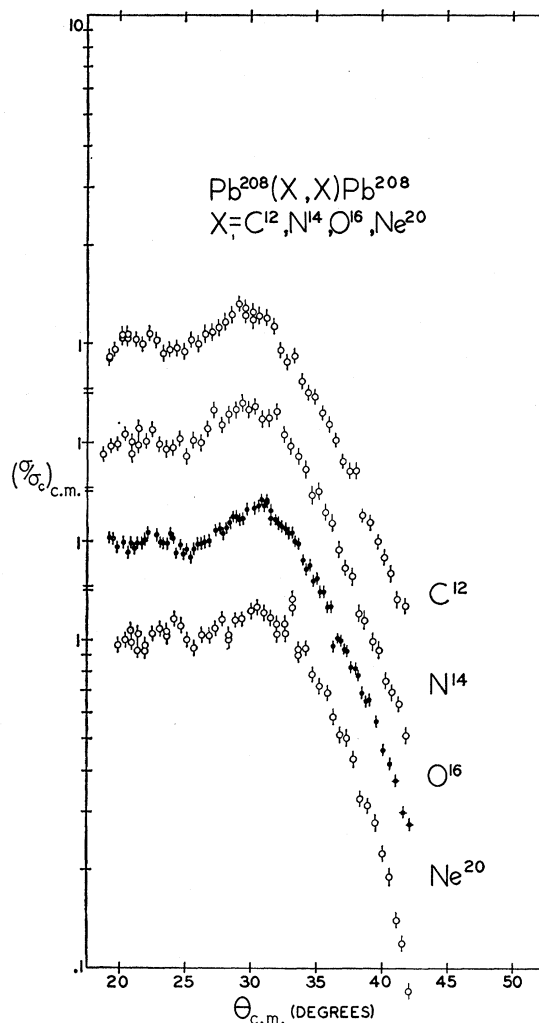


FIG. 6. Ratio of the differential cross section for elastic scattering to the Coulomb cross section vs the center-of-mass angle for C^{12} , N^{14} , O^{16} , and Ne^{20} scattered by Pb^{208} . Relative positions are arbitrary. The curves have been normalized to unity at small angles. One thousand tracks were counted for each point of the O^{16} , Pb^{208} reaction. Three hundred tracks were counted for each point of the other curves.

induced fission experiments¹³ performed simultaneously with the elastic experiments. The K-2 emulsions recorded elastically scattered ions sufficiently well at end of range to obtain elastic angular distributions at position *B* as well as at *A*, thus providing elastic angular distribution measurements on either side of the beam axis. The characteristic rapid decrease of the elastic differential scattering cross section with increasing angle allowed a determination of symmetry with the beam axis, thus eliminating all but minor beam-alignment uncertainties.

Factors involving angular resolution have not

¹³ E. Goldberg, H. L. Reynolds, and D. D. Kerlee, *Proceedings of the Second Conference on Reactions Between Complex Nuclei* (John Wiley & Sons, Inc., New York, 1960), paper D-6, p. 230.

changed appreciably from the earlier report.⁷ Contributions to the spread in scattering angle due to initial beam divergence, collimator slit widths, multiple scattering in the target, swath length scanned in the emulsion, and loss of angular resolution due to beam energy, result in a spread in scattering angle of less than $\pm 0.8^\circ$ at half-maximum.

Errors in target box alignment were corrected to within 0.1° by comparison of elastic scattering in plates *A* and *B*. The error in scattering angle due to target box alignment, measurement of the internal geometry of the box, and data processing was $\pm 0.2^\circ$. The above discussion of errors in angle and angular resolution pertains to a range of scattering of from 20° to 40° and applies to particles detected on *A* plate.

The energy of the heavy ions was determined by measuring range spectra and applying the range-energy measurements of Heckman *et al.*¹⁴ The ranges and

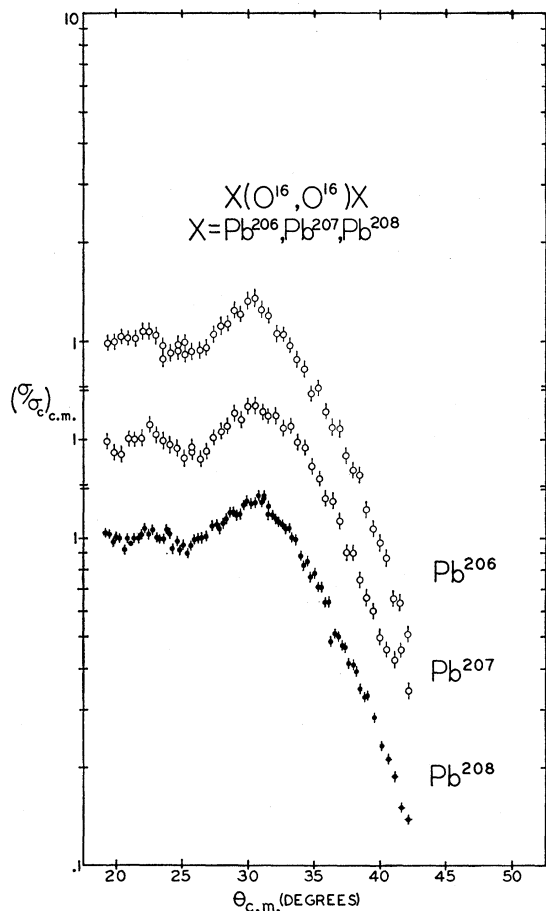


FIG. 7. Ratio of the differential cross section for elastic scattering to the Coulomb cross section vs the center-of-mass angle for O^{16} scattered by Pb^{206} , Pb^{207} , and Pb^{208} . Relative positions are arbitrary. The curves have been normalized to unity at small angles.

¹⁴ H. H. Heckman, B. L. Perkins, W. G. Simon, F. M. Smith, and W. H. Barkas, *Phys. Rev.* **117**, 544 (1960).

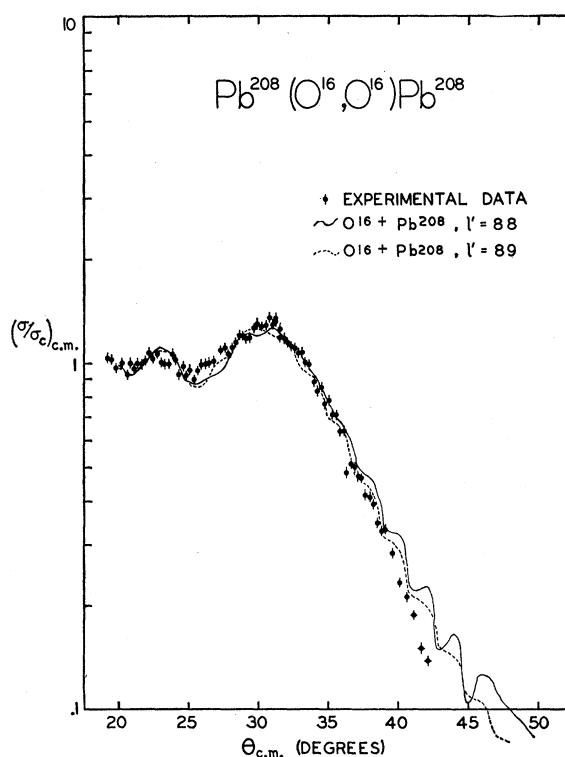


FIG. 8. Ratio of the differential cross section for elastic scattering to the Coulomb cross section vs the center-of-mass angle for O^{16} scattered by Pb^{208} . Sharp-cutoff calculations are shown for $l' = 88$, and $l' = 89$.

laboratory energies for scattering from Pb^{206} , Pb^{207} , and Pb^{208} are given in Table I. The incident energy of the projectiles in the laboratory system was approximately 10.4 MeV/nucleon. A magnetic electron suppressor similar to that described earlier⁶ was used in these experiments. No rubbing of the emulsion⁷ was necessary.

A typical range spectrum of O^{16} on Pb^{208} at a laboratory angle of 20° is shown in Fig. 2. The width is consistent with a range straggling of 1.0% and a standard deviation of the incident-beam energy spread⁷ of 1.0%.

The isotopic targets were self-supporting foils of Pb^{206} , Pb^{207} , and Pb^{208} . Isotopic concentrations and target thicknesses are given in Table II. The energy loss of the O^{16} ions in the Pb^{208} target was approximately 1.1 MeV.

RESULTS

In Fig. 3 the experimental angular distribution obtained for O^{16} projectiles incident on Pb^{208} is plotted as the ratio of the differential cross section to the Coulomb cross section, σ/σ_c , vs the center-of-mass angle. It has been assumed that the ratio approaches unity as the scattering angle decreases. At small angles, each point represents 1000 tracks with a resulting

standard deviation of approximately 3%. Beyond 40° c.m., fewer tracks were measured for each point.

The acceptance angle was broadened for these points to obtain improved sensitivity. The larger standard deviations and angles of acceptance are indicated. No tracks identifiable as elastic, by range and density measurements, were found beyond 82°. Areas scanned are indicated at 92° and 172°. Upper limits to the ratio, σ/σ_e , are indicated on the basis of the areas scanned.

The scattered particle range spectrum has a normal shape up to values of $\sigma/\sigma_e \geq 0.2$. Beyond this a broadening of the elastic peak occurs. Consequently, points where $\sigma/\sigma_e < 0.2$ have been determined by accepting only those tracks with the correct range and density for normal elastic scattering. However, it is possible that these points may still contain some nonelastic events. The value of σ/σ_e has an upper limit of 1.2×10^{-5} at 172° c.m.

Composite graphs of the individual experiments are shown in Figs. 4 to 7. With the exception of O¹⁶ on Pb²⁰⁸, each point represents 300 tracks with a resulting standard deviation of approximately 6%. Figure 4 shows the experimental angular distributions obtained for C¹², N¹⁴, O¹⁶, and Ne²⁰ on Pb²⁰⁶ plotted as σ/σ_e vs the center-of-mass angle. Figure 5 shows C¹², O¹⁶, and Ne²⁰ on Pb²⁰⁷. Figure 6 shows the four projectiles C¹², N¹⁴, O¹⁶, and Ne²⁰ on Pb²⁰⁸. The O¹⁶, Pb²⁰⁸ experiment is with improved statistics on each point as described above. Figure 7 is a composite of O¹⁶ on the three lead isotopes.

DISCUSSION

Cutoff values of angular momentum, and values of the nuclear interaction distance have been determined from the Blair model. A representative fit to the experimental data is shown for O¹⁶ on Pb²⁰⁸ in Fig. 8.

TABLE III. Cutoff l' values, nuclear interaction distances, and reaction cross sections.^a

Reaction	Cutoff l'	Interaction distance R^b (Fermis)	r_0 (Fermis)	σ_r (barns)
C ¹² +Pb ²⁰⁶	64.0±1.1	11.8±0.2	1.44±0.02	2.1
N ¹⁴ +Pb ²⁰⁶	78.0±1.1	12.1±0.2	1.46±0.02	2.3
O ¹⁶ +Pb ²⁰⁶	89.0±1.7	12.2±0.2	1.45±0.03	2.3
Ne ²⁰ +Pb ²⁰⁶	111.5±2.2	12.6±0.2	1.46±0.02	2.4
C ¹² +Pb ²⁰⁷	65.5±1.1	12.0±0.3	1.46±0.04	2.2
O ¹⁶ +Pb ²⁰⁷	87.5±1.7	12.2±0.2	1.44±0.02	2.3
Ne ²⁰ +Pb ²⁰⁷	112.0±2.2	12.5±0.2	1.45±0.03	2.4
C ¹² +Pb ²⁰⁸	66.0±1.1	12.1±0.2	1.47±0.02	2.2
N ¹⁴ +Pb ²⁰⁸	78.0±1.1	12.1±0.2	1.45±0.02	2.3
O ¹⁶ +Pb ²⁰⁸	88.5±1.7	12.3±0.2	1.45±0.02	2.3
Ne ²⁰ +Pb ²⁰⁸	111.5±2.2	12.6±0.2	1.45±0.02	2.4

^a All errors are probable errors. Uncertainty in the scattering angle due to chamber geometry is included in the cutoff l' -value error.

^b The value of $E_{c.m.}$ used in the calculation of the interaction distance was the nonrelativistic energy corresponding to the measured (slightly relativistic) reaction energy ($E_{c.m.R}$), where $E_{c.m.} = 0.984 E_{c.m.R}$. Also the reduced mass was determined from the rest mass of the particles.⁷

The method of analysis is described in previous papers.^{6,7} After finding the best-fit l' values, nuclear interaction distances R were obtained from the equation,

$$E_{c.m.} = \frac{Z_1 Z_2 e^2}{R} + \frac{\hbar^2 l'(l'+1)}{2\mu R^2},$$

where μ is the reduced mass. The results are shown in Table III. Values of r_0 were obtained from the relation $R = r_0(A_1^{1/3} + A_2^{1/3})$. The results are consistent with an r_0 of 1.45 F for all of the reactions. Interaction cross sections, computed from the relation

$$\sigma_r = \pi \lambda^2 \sum_0^{l'} (2l+1),$$

are also given in Table III.

The error in the radius was determined by the relation⁷

$$|dR|^2 = |adE_{c.m.}|^2 + |bdl'|^2.$$

The constants a and b for the lead isotopes are shown in Table IV.

Comparisons of the elastic scattering data for O¹⁶ on Pb²⁰⁶, Pb²⁰⁷, and Pb²⁰⁸ indicate no essential difference in the behavior of the rise. This is in contrast to the trend observed with alpha-particle scattering⁸ on these isotopes. It was observed with alpha particles that the rise increased as the neutron number N increased from 124 to shell closure at 126. The heavy-ion measurements with larger values of the classical parameter η do not show this trend. Further, a study of the four projectiles on Pb²⁰⁸ shows that few distinctive features are revealed by the scan across projectiles. The earlier alpha-particle work indicated a decrease in the rise in the rare-earth region. McIntyre¹⁵ has shown that differences are de-

TABLE IV. Values of constants in error determination for interaction distance.

Reaction	a (Fermis/MeV)	b (Fermis)
C ¹² +Pb ²⁰⁶	-0.091	0.118
N ¹⁴ +Pb ²⁰⁶	-0.080	0.103
O ¹⁶ +Pb ²⁰⁶	-0.073	0.090
Ne ²⁰ +Pb ²⁰⁶	-0.060	0.073
C ¹² +Pb ²⁰⁷	-0.093	0.119
O ¹⁶ +Pb ²⁰⁷	-0.075	0.091
Ne ²⁰ +Pb ²⁰⁷	-0.060	0.074
C ¹² +Pb ²⁰⁸	-0.093	0.119
N ¹⁴ +Pb ²⁰⁸	-0.080	0.103
O ¹⁶ +Pb ²⁰⁸	-0.077	0.091
Ne ²⁰ +Pb ²⁰⁸	-0.060	0.073

¹⁵ See reference 13, J. A. McIntyre, S. D. Baker, and K. H. Wang, paper C-8, p. 180.

tectable in the rise and in the oscillations precursory to the rise by comparing the reactions $\text{Pb}^{208}(\text{O}^{16}, \text{O}^{16})\text{Pb}^{208}$, and $\text{Tb}^{159}(\text{F}^{19}, \text{F}^{19})\text{Tb}^{159}$. In the present work there is an indication of similar trends in the scan across projectiles (Fig. 6). However, experimental angular resolution and statistics do not allow quantitative evaluation of this effect. Further work in the rare-earth region would be of interest.

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Measurement of the $\text{O}^{16}(n, p)\text{N}^{16}$ Cross Section from 11 to 19 MeV

J. A. DEJUREN* AND R. W. STOOKSBERRY

Bettis Atomic Power Laboratory,† Westinghouse Electric Corporation, Pittsburgh, Pennsylvania

AND

M. WALLIS

Los Alamos Scientific Laboratory, Los Alamos, New Mexico

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The present experiment was undertaken to improve upon the accuracy and energy resolution of the $\text{O}^{16}(n, p)\text{N}^{16}$ cross section and to extend the neutron energy range covered. The oxygen was contained in a dioxane liquid scintillator which was irradiated by accelerator-produced neutrons. Shaped holders were used at appropriate angles from the target to obtain good energy resolution for neutrons produced by both $\text{T}(d, n)\text{He}^4$ and $\text{D}(t, n)\text{He}^4$ reactions. After the irradiation the scintillator was automatically transferred to a 4π β -counting system. The resulting cross section data exhibit a previously undetermined resonance near 11.8 MeV. Integration of the experimental cross sections over U^{235} fission spectrum formulations yield values of 16 to 19 μb per neutron.

I. INTRODUCTION

AN earlier measurement of the $\text{O}^{16}(n, p)\text{N}^{16}$ cross section from 12.4 to 18 MeV was made by Martin¹ in 1953. The energy resolution below 15 MeV was ± 1 MeV and the absolute value near 14.7 MeV was a factor of 2 higher than later results obtained by DeJuren and Stooksberry,² Kantele,³ and Seemann and Moore.⁴ The present experiment was undertaken to improve the energy resolution and to extend the energy range covered. A cooperative experiment was undertaken by Bettis and Los Alamos personnel utilizing the Los Alamos 3.0-MeV Van de Graaff accelerators and the $\text{T}(d, n)\text{He}^4$ and $\text{D}(t, n)\text{He}^4$ reactions to cover the energy range from 11 to 19.4 MeV. As in the 14.7-MeV measurement² at Bettis Atomic Power Laboratory, the oxygen was contained in a dioxane-based liquid scintillator and was irradiated in various geometrically shaped holders at different angles with respect to the beam direction to obtain the different

neutron energies. After an irradiation the liquid was automatically transferred to a Vycor cell between two EMI 7064 photomultipliers and the β rays were counted in a 4π system. The system used in the 14.7-MeV experiment was again used but without photomultiplier cooling. Beta particles above 20 keV were counted, with a detection beta efficiency of 99.5%.²

II. EXPERIMENTAL PROCEDURE

Two 3-MeV Van de Graaff accelerators were used to provide sources of monoenergetic neutrons by means of the $\text{D}(t, n)\text{He}^4$ and the $\text{T}(d, n)\text{He}^4$ reactions. The respective energy ranges available were 11 to 19.4 MeV and 12 to 19 MeV at accessible laboratory angles. Even though the first reaction provided the complete range of neutron energies covered in this experiment, of the two accelerators, only one could be used to accelerate tritons and the time that this machine could be used was limited. Therefore, some of the intermediate neutron energies were obtained using the $\text{T}(d, n)\text{He}^4$ reaction on the second machine. Backgrounds were present due to the t -T reaction⁵ when the t -D source was being used and the $\text{D}(d, n)\text{He}^3$ reaction when the d -T source was being used. These two backgrounds

* Now at Atomics International, Canoga Park, California.

† Operated for the U. S. Atomic Energy Commission by Westinghouse Electric Corporation.

¹ H. C. Martin, Phys. Rev. **93**, 498 (1954).

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⁵ S. J. Bame, Jr., and W. T. Leland, Phys. Rev. **106**, 125 (1957).